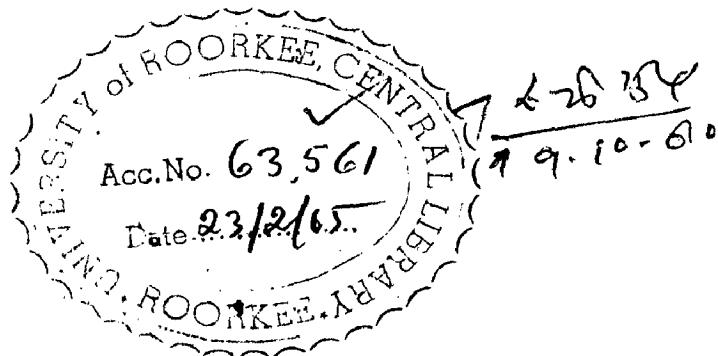


**ELECTRICAL ANALOGUE COMPUTATION  
OF PERIODIC HEAT FLOW  
THROUGH BUILDINGS**

THESIS SUBMITTED FOR THE AWARD OF THE DEGREE  
OF DOCTOR OF PHILOSOPHY  
IN APPLIED PHYSICS



BY  
**K. RAMAMOHAN RAO, M.Sc.**

**UNIVERSITY OF ROORKEE  
ROORKEE**

**1963**

ANALOGUE COMPUTATION  
OF PERIODIC HEAT FLOW  
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the Degree of Doctor of Philosophy  
in Applied Physics

by

K. ASWATHAN RAO, M.Sc.

UNIVERSITY OF ROORKEE

1963

1063

C O M P I L A T I O N

Certified that the Thesis entitled " Electrical Analogue Computation of Heat Flow through Buildings " which is being submitted by Shri K. Ramachan Rao for the award of the Degree of Doctor of Philosophy in applied Physics, of the University of Poona, is his own work carried out under my supervision and guidance. The matter embodied in this Thesis has not been submitted for the award of any other Degree or Diploma of any University.

This is to further certify that he has worked for a period of two and a half years (from June 1961 to December 1963) on this problem.

P.V. Indiresan

( P.V. Indiresan )  
Associate Professor,  
Electrical Engineering Deptt.,  
University of Poona.

ROONKSH

Date : 25.12.1963

A B S T R A C T

O F

T H E T H E o I S

## ABSTRACT

This thesis contains a record of research work carried out on periodic heat flow characteristics of building elements and the prediction of indoor air temperatures of buildings by Electrical Analogue method. Such a study was found essential for evolving building designs to suit diverse climatic conditions of India.

A building element is represented on this analogue by a suitable combination of resistance - capacitance network. The type and the number of lumps required to represent a building fabric is decided by the thickness and physical properties of the material. The surface heat transfer coefficients are represented by suitable resistances.

A steady state sinusoidal excitation has been adopted to simulate the diurnal solar temperature variations. The main advantages of this approach over other simulation techniques are (i) eliminates the need of complicated electronic equipment like function generators and recording devices, (ii) makes possible to separate the effects of climatic factors from the thermal characteristics of the building fabric and (iii) only two quantities viz., amplitude decrement and phase lag angle are sufficient to measure on the analogue network.

To facilitate phase measurements, a direct reading audio phasor meter capable of reading upto  $\pm 1^\circ$  has been specially designed and fabricated.

The number of lumps required to represent any building element adequately on the analogue has been established both theoretically and experimentally. Computed data, on the transfer matrix coefficients, for lumped and distributed systems, required for evaluating the lumping errors in the estimation of temperatures and heat fluxes, are also included.

To define completely the transient behaviour of a building component a set of three analogous system functions (one transfer and two driving point) are required. Relationships between these functions and the transfer matrix coefficients, for commonly occurring boundary conditions have been derived in this thesis.

A detailed study of the type and thickness of each layer and the effect of the order of arrangement on the overall thermal functions of the composite construction, has also been made.

A simple and flexible computational method for the prediction of the indoor air temperatures, in terms of thermal system functions, 'U' values and sol-air temperatures, has been evolved. As this method utilises

the pretabulated thermal system function data, it eliminates the necessity of representing a full room or a building with all its structural elements and input generators, on the analogue, for each specific study.

This method, when used with the data presented in this thesis, leads to a quantitative evaluation of the relative thermal efficiencies of the various possible structural combinations and other design factors. A few applications of the data are also illustrated.

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## **C H A P T E R   1**

### **INTRODUCTION**

## CHAPTER I

### INTRODUCTION

Accurate prediction, of the thermal behaviour of building under periodic variations of temperatures and heat flow, is of considerable practical importance in building design. Although a good deal of work has been done in this direction and new methods and techniques evolved, several problems of practical interest, still await solutions. This owes mainly to the multiplicity of the variables involved and the lack of a simple and flexible method for studying the transient heat flow, through building fabrics.

A purely analytical solution of this problem is not impossible, but it is very laborious and time consuming, even with the availability of such efficient methods, as matrix analysis, and symbolic calculus. If more than a few variables are to be investigated simultaneously, field measurements are not practicable. This imposes a serious limitation on those results which do not lead to any generalised conclusions. Further, to determine the influence of each of the several variables that one normally comes across in heat transfer problems, a large number of test houses have to be built in different climatic zones and the results analysed.

statistically. Such a procedure is time consuming and uneconomical.

Electrical analogue methods have been found to be more suitable for such problems and yield quick and economical results.

The electrical analogue simulates thermal problems in terms of electrical quantities. Within the limits set by the accuracy of the analogue and the basic assumptions made in the theory, the analogue method may be regarded as a means of carrying out controlled experiments on a variety of building elements, under identical conditions.

Most of the investigations employing electrical analogue methods were mainly restricted to the estimation of fabric cooling loads for air conditioning purposes. Though these techniques can be extended to unconditioned buildings, these will have only a limited application. The simulation of a room or a building with all its structural components, and the associated input generators, make the network too complex for experimental work. Even the initial cost of the equipment becomes exorbitant. Moreover the efforts involved in simulating the several variables (climatic and design) and their possible combinations on such an analogue make the method cumbersome and unsuitable for design problems.

It is, therefore, apparent that there is a genuine need for the development of a simple and flexible method to obtain a quick evaluation of thermal efficiencies of buildings and building elements. The primary requisite of such a method is that it should be capable of taking into account the influence of various factors, other than the constructional variations, viz., ventilation, orientation, surface treatments, internal heat sources and sinks, regional and seasonal variations of temperatures and solar radiation etc.

A very profitable line of approach would be i) to obtain the basic reference data on transient thermal characteristics of a large number of building sections, that one might come across in practice, by the electrical analogue method, and ii) to devise a simple computational method to predict the indoor air temperatures, utilising the above pretabulated data, for any combination of building sections under any climatic conditions.

This thesis is mainly directed towards the development of such a method and the data.

In addition several other problems of interest such as i) the analogue representation of the building sections by lumped cascaded networks, ii) influence of surface heat transfer coefficients on the lumping errors and thermal system functions, iii) the effect of different

combinations of layers of materials on the thermal behaviour of the individual layers and on the composite section as a whole, have also been studied in this thesis.

Such a study will lead to a better understanding of the transient behaviour of composite constructions and will finally result in the appropriate design of composite constructional elements, best suited to the climatic conditions of a given place.

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## **C H A P T E R 2**

### **PERIODIC HEAT FLOW THROUGH BUILDINGS**

## CHAPTER 2

### PERIODIC HEAT FLOW THROUGH BUILDINGS

#### 2.1 Introduction

In the steady state heat flow analysis the heat storage capacity of the structure does not come into the picture. In climates where, there are large periodic variations of air temperature and solar radiation, the thermal capacity of external and internal building elements play a great role in modifying the temperature distribution and the rate of heat flow, with respect to both space and time.

The mechanism of heat flow through a building element, say a wall, under transient conditions may be visualised in the following manner. An elementary thickness of wall just beneath the outside surface, receives a net heat gain by convection and radiation. A portion of this energy is spent in heating up the element (in question) while the balance is conducted on to the next elementary thickness of the wall and so on. As a definite interval of time is required for the incident heat to raise the temperature of the element (in question) there is a phase difference between the temperature or heat waves at the outer and the inner surfaces. Further, as some of the

input energy is used up in heating the wall, the amplitude of the wave at the inner surface is considerably reduced.

The main object of the investigations on periodic heat flow through the building elements is to establish a relationship between the thermal conditions of the outdoor and indoor environments and the properties of the intervening structure. This necessitated a development of the theory of non-steady state heat flow, and solutions to the Fourier heat conduction equation under natural boundary conditions. The differential equation of heat flow in solids which gives the temperature distribution with time was given by Fourier as

$$\frac{\partial t}{\partial T} = \alpha \left( \frac{\partial^2 t}{\partial x^2} + \frac{\partial^2 t}{\partial y^2} + \frac{\partial^2 t}{\partial z^2} \right) \quad \dots (1)$$

where  $t$  = temperature at the point  $(x, y, z)$  at time  $T$ .

$\alpha = K/\rho s$  - thermal diffusivity in  $\text{Ft}^2/\text{Hr.}$

$K$  = thermal conductivity in  $\text{Btu}/\text{Ft}/\text{Hr}/{}^\circ\text{F.}$

$s$  = specific heat  $\text{Btu}/\text{Lb}/{}^\circ\text{F}$

$\rho$  = density  $\text{Lb}/\text{Ft}^3.$

In special conditions (parallel bounding surfaces with large area compared to the thickness) the heat flow can be considered as one dimensional. Then the above equation reduces to

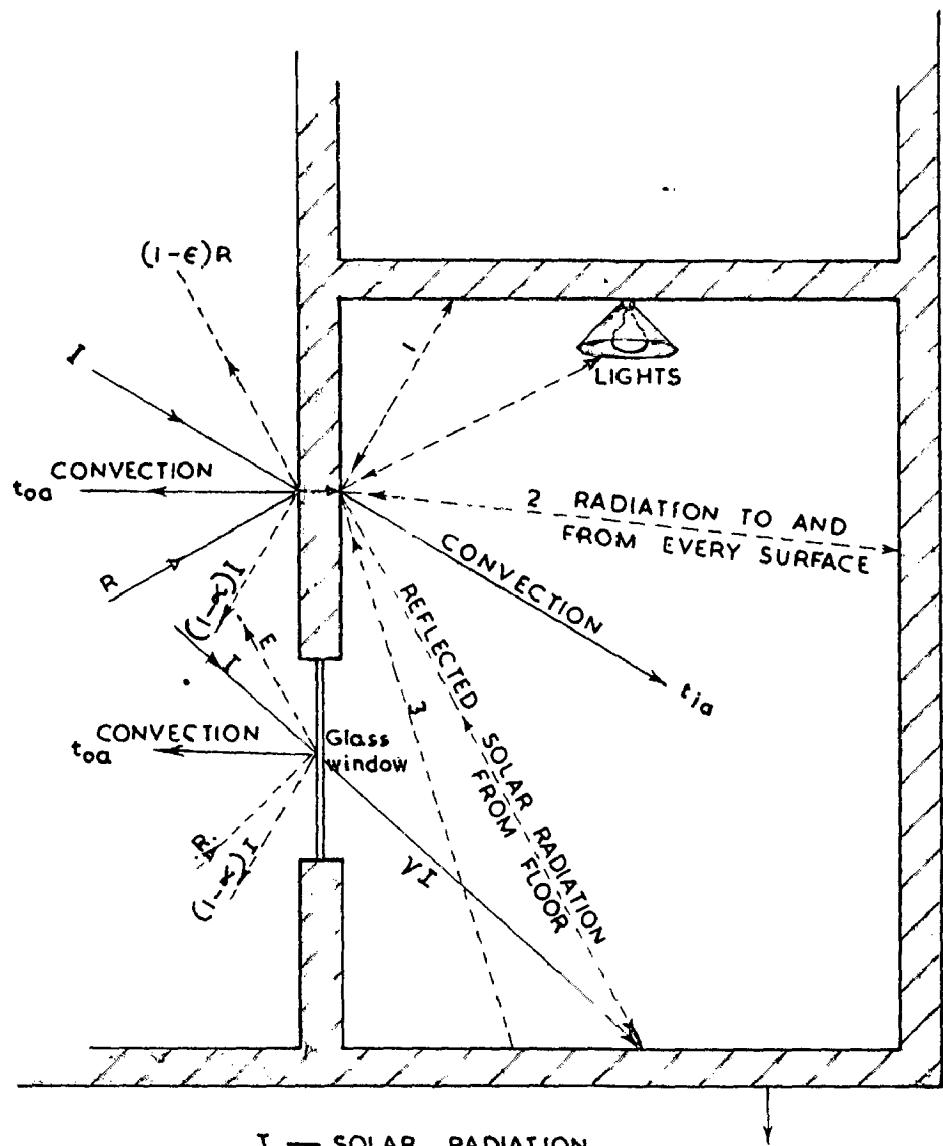
$$\frac{\partial t}{\partial T} = \alpha \frac{\partial^2 t}{\partial x^2} \quad \dots (2)$$

For all practical purposes the building components such as walls, roof, floor etc., satisfy the condition of unidirectional heat flow. Hence this equation (2) is to be solved, for a given building element, with the appropriate external and internal boundary conditions, which of course vary from situation to situation. The heat exchange taking place at the outer and inner surfaces form the outside and inside boundary conditions respectively. Hence the expressions for these heat exchanges are of fundamental importance and should be expressed in the simplest possible form consistent with the requirements, to a satisfactory degree of accuracy.

### 2.2 Heat Exchange at the Outside Surface of a Building Element

The phenomena of heat exchange at the outside surface of a building element as illustrated in Fig. (2.1) can be resolved into the following components :-

- i) Heat transfer by convection between the surface and the outdoor air.
- ii) Total solar radiation (Direct + Diffuse) absorbed by the surface.
- iii) Heat transfer by low temperature radiation exchange between the surface and the outdoor environment.
- and iv) The rate of heat transfer by conduction into the building element.



I — SOLAR RADIATION

$\gamma$  — SOLAR TRANSMISSIVITY OF WINDOW

$t_{os}$  — OUTSIDE SURFACE TEMP.

$t_{is}$  — INSIDE SURFACE TEMP.

$q$  — HEAT FLOW FROM OUTSIDE SURFACE TO INSIDE

$t_{oa}$  — AMBIENT TEMPERATURE

$t_{ia}$  — ROOM AIR TEMPERATURE

$\alpha$  — SOLAR ABSORPTIVITY OF THE SURFACE

R — LONG WAVE RADIATION FROM THE SURROUNDINGS

E — RADIATION EMITTED FROM THE SURFACE

$\epsilon$  — EMISSIVITY OF SURFACE

## 2.1 HEAT BALANCE AT THE OUTSIDE AND INSIDE SURFACES OF AN EXPOSED WALL OF AN ENCLOSURE

The heat balance equation at the outside surface is given by

$$q_{os} + h_{os} (t_{oa} - t_{os}) + \alpha I_T \pm I_L = 0 \quad \dots (3)$$

where

$q_{os}$  = instantaneous heat flux ontoing the building element at the outside surface per unit area.

$h_{os}$  = outside surface convection heat transfer coefficient.

$t_{oa}$  = outside air temperature.

$t_{os}$  = outside surface temperature.

$I_T$  = intensity of the total solar radiation incident on the surface (Direct + Diffuse).

$\alpha$  = absorptivity of surface for solar radiation.

$I_L$  = net long wave radiation exchange between the surface and its surroundings.

All those component parts of the equation (3) depend upon the weather elements and the surface properties. Seasonal and diurnal variations of these climatic factors will have to be taken into consideration. Combining all these factors into a single parameter will simplify the computations to a great extent.

### 2.3 Sol-air Temperature

Huckley and Wright (1) have introduced the concept of sol-air temperature, which combines the air temperature and solar radiation into a single equivalent temperature. As per the definition of the sol-air

temperature, the external boundary condition is given by

$$q_{os} = h_{oc} (t_{sa} - t_{os}) \quad \dots (4)$$

where  $t_{sa}$  is the sol-air temperature. In their original work, these authors have not taken the low temperature radiation, into account and the sol-air temperature is given by the equation

$$t_{sa} = t_{oa} + \frac{\alpha I_T}{h_{oc}} \quad \dots (5)$$

Neglecting the low temperature radiation is not justifiable and the above formula for sol-air temperature requires modification. The modified equation according to Roux (2) may be written as

$$t_{sa} = t_{oa} + \frac{\alpha I_T \pm I_{LT}}{h_{oc}} \quad \dots (6)$$

To compute the sol-air temperature, a knowledge of the incident solar radiation intensity, surface absorption coefficient, low temperature radiation surface convection heat transfer coefficient, are required, apart from the outdoor air temperatures data. The intensity of solar radiation incident upon a surface, depends upon many factors such as orientation, latitude, day and time, atmospheric conditions etc. Hence for a given locality, though all the surfaces of a building are subjected to the same ambient air temperature, the sol-air temperature

will be different for differently oriented surfaces. The solar radiation intensity data (measured as well as computed) for a number of places and differently oriented surfaces is available in the literature (3). Preparation of design curves of diurnal variations of solar-air temperature for summer and winter seasons for various places, facilitate the calculations of heat transfer through building elements and the design of buildings to suit a given climate.

The surface absorption coefficient for solar radiation depends upon the surface colour and the angle of incidence. Considerable work has been done on this aspect (4, 5) and large amount of reference data is available in the literature (3, 7).

#### 3.4 Low Temperature Radiation

The low temperature radiation received by a building element, depends upon the mean radiant temperature of the surroundings, ground and atmosphere. It is usually assumed that the outdoor surroundings emit long wave radiation as a black body at air temperature (3). Hence the net radiation exchange between the surface and its surroundings is expressed as

$$I_u = \epsilon \sigma (T_{os}^4 - T_{ou}^4) \quad \dots (7)$$

where

$\epsilon$  = absorptivity and emissivity of the surface for long-wave thermal radiation.

$\sigma$  = Stephan Boltzman constant.

$T_{os}$  = absolute surface temperature, and

$T_{oa}$  = absolute air temperature.

Burnt (9) has pointed out that the outdoor surroundings by no means radiate as a black body at air temperature and correlated the long wave radiation emitted by cloudless atmosphere with water vapour pressure present in the atmosphere based on Dines and Dines measurements (10).

The correlation is given by

$$R = R_b ( a + b \sqrt{P_w} ) \quad \dots (8)$$

where

$R_b$  = low temperature radiation emitted by a blackbody at outdoor air temperature (i.e.,  $\sigma T_{oa}^4$ )

$R$  = low temperature radiation emitted by the atmosphere.

$P_w$  = the water vapour pressure in inches of mercury.

'a' and 'b' are constants and given as 0.55 and 0.33 respectively for horizontal surfaces and as 0.3 and 0.165 for vertical surfaces.

Roux (2) pointed out, that while the Burnt's equation (8) gives an accurate estimate of low temperature radiation during night times, is not strictly valid for conditions for day times when solar radiation intensity is high. However, the equation (8) is very useful as a working

formula in the absence of more reliable data. Under cloudy conditions the low temperature radiation of the atmosphere increases and approaches the black body radiation at air temperature for an overcast sky condition. Substituting the value of  $I_b$  in the sol-air temperature equation (5) we get

$$t_{oa} = t_{ou} + \frac{\alpha L_T + (\epsilon \sigma T_{ou}^4 (a + b \sqrt{R_b}) - \epsilon \sigma T_{ou}^4)}{h_{oc}} \dots (8)$$

This expression includes a term of surface temperature, which is a function of the physical properties of the building element as well. As the sol-air temperature is to be computed exclusively from the weather data, the term  $t_{oc}$  from the above expression is to be eliminated. For this purpose a new term "artificial surface radiative heat transfer coefficient" has been defined (2) as

$$h_{or} = \frac{\epsilon \sigma (T_{os}^4 - T_{ou}^4)}{(t_{os} - t_{ou})} \dots (10)$$

With the help of the equation (10) we obtain the sol-air temperature as

$$t_{os} = t_{ou} + \frac{\alpha L_T - (\epsilon \sigma T_{ou}^4 - a)}{(h_{oc} \pm h_{or})} \dots (11)$$

The terms  $(h_{oc} \pm h_{or})$  are combined into a single term

which is called, the outside surface heat transfer coefficient ( $h_o$ )

## 2.5 Outside Surface Heat Transfer Coefficient ( $h_o$ )

As seen from the above equation ' $h_o$ ' consists of two parts, namely, convective and radiative coefficients. The convective coefficient is made up of two components, namely the forced convection ( $h_{oc}$ ) and natural convection ( $h_{onc}$ ). Several investigators (11 - 15) have studied the nature of the forced and natural convective heat transfer and given empirical equations relating several factors involved in the phenomena. Both the convective and radiative coefficients have been found to vary considerably during day and night, as they are dependent on the air temperature and wind velocity (16, 17).

However, it is combined value of  $h_{oc}$  and  $h_{or}$  rather than their separate individual values that is of interest. Fortunately, this combined value  $h_o$  (i.e.,  $h_{oc} + h_{or}$ ) appears to remain fairly constant (13), although the individual variations of  $h_{oc}$  and  $h_{or}$  are large. This approximation of taking the combined heat transfer coefficient ( $h_o$ ), as constant, greatly simplifies the computations of heat transfer through building elements. Errors due to this approximation are found to be small (13).

With these simplifications the boundary condition at the outside surface may be written as

$$-K \left( \frac{\partial t}{\partial x} \right)_o = h_o (t_{sa} - t_{os}) \quad \dots (12)$$

where

$$-K \left( \frac{\partial t}{\partial x} \right)_o = q_{os}$$

### 3.6 Heat Exchange at the Inside Surface

Various heat transfer processes that take place at the inside surface, are illustrated in Fig.(3.1). At this surface the heat transfer takes place to the inside air by convection and by radiative exchange between all the interior surfaces. The heat balance equation for an inside surface can be expressed as

$$q_{is} + h_{ic} (t_{ia} - t_{is}) + \sum_{n=1}^{n=M} s^F n \sigma (T_n^4 - T_s^4) = 0 \quad \dots (13)$$

where

$q_{is}$  = the heat flow through the wall at the inside surface.

$h_{ic}$  = convective heat transfer coefficient at the inside surface.

$t_{ia}$  = inside air temperature.

$t_{is}$  = inside surface temperature.

$n$  = subscript indicating one of the surfaces enclosing.

$M$  = number of surfaces forming the enclosure.

$s^F n$  = overall interchange factor for radiant energy exchange between surfaces.

As in the case of outside surface coefficient, if the inside surface, convective ( $h_{ic}$ ) and radiative ( $h_{ir}$ ) heat transfer coefficients are combined into a single term ( $h_i$ ) and assumed constant, the expression for the internal boundary condition simply becomes

$$- k \left( \frac{\partial t}{\partial x} \right)_i = h_i (t_{is} - t_{ia}) \quad \dots (14)$$

where

$$- k \left( \frac{\partial t}{\partial x} \right)_i = q_{is}$$

and

$$h_i = (h_{ic} + h_{ir})$$

The major advantage of combining the convection and radiation heat exchange at the inside surface, is that the overall radiation interchange factors need not be known. This saves the difficulty of finding the geometric view factors between the surfaces. Another significant simplification achieved by this assumption is that the heat flow paths to the room air, can be considered separately, without reference to any other surface.

## 2.7 Types of Problems met with in Practice

The types of non-steady state heat flow problems related to buildings can be broadly divided into the following three categories :-

- 1) Air conditioned buildings - where the indoor

air temperature is maintained constant by artificial cooling or heating. For this class of buildings the periodic heat flow through all the components of the fabric, due to external air temperature and solar radiation variations, are to be precisely calculated.

- ii) Air conditioned buildings with insufficient plant capacity - where the indoor air temperatures, which vary to a certain extent, are to be calculated.
- and iii) Un-conditioned buildings - where the calculation of indoor air temperatures under various climatic conditions is the main objective.

Once these outside and inside boundary conditions are established, for any given situation the solution for the Fourier heat conduction equation can be obtained by several established methods such as Analytical (18,20), Numerical (21,22), Matrix (23,24,25), Symbolic Calculus (23,27), Influence and Weighting functions (23,29,30), Direct Analogous (31 - 40) and (Electronic Analogue and Digital) Computers (41 - 46) etc. The relative merits of the various methods employed for calculating heat conduction for the determination of cooling loads for air conditioned

buildings, have been discussed by Stephenson (46).

### 2.8 Assumptions made in the Theory

In all the methods, certain simplifying assumptions are made for obtaining solutions of the practical design problems. The usual ones are :-

- i) Heat flow through the building element is unidirectional.
- ii) The differential equation describing the heat flow is linear, which implies that the physical properties of the materials ( $K$ ,  $\rho$ , and  $c$ ) are independent of temperature.
- iii) Heat exchange phenomena at the external surface of a building element can be expressed by a single parameter (vol-air temperature).
- iv) The heat transfer by convection and radiation at the inside and outside surface can be represented by a combined film conductance which does not vary with time and temperature.
- v) The absorptivity of the outside surface of building elements, to solar radiation is independent of the angle of incidence.
- vi) The variations of the outdoor air temperature and solar radiation intensity with time are identical on successive days.

Most of the methods enumerated above were developed for estimating the fabric cooling loads, for air conditioning purposes. However, some of them are extended for the unconditioned buildings too. As pointed out earlier (Chapter 1) though these methods provide solutions for any specific situation with good accuracy, the complexity and the effort involved is so great, that for generalised design problems of buildings, for various climates, these become rather unsuitable. Simpler methods, utilizing maximum amount of pretabulated, thermal characteristic data, are to be evolved.

---

## **C H A P T E R 3**

### **ANALOGUE SETUP AND MEASUREMENTS**

## CHAPTER 3

### ANALOGUE SETUP AND MEASUREMENTS

#### 3.1 Introduction

Electrical Analogue techniques for the study of transient heat flow problems in different fields are well established. Although the thermal electrical analogy is based on the similarity of mathematical formulation of heat flow through a material with that of the electric current flow in a transmission line

$$\frac{\partial t}{\partial T} = \alpha \frac{\partial^2 t}{\partial x^2}$$

$$\frac{\partial E}{\partial T} = \frac{1}{RC} \frac{\partial^2 E}{\partial x^2}$$

In practice the analogue networks may be of different designs (31,34,36,47). The main difference between the earlier R-C models and the recent ones centres round the choice of the time scale, which has a direct bearing on the size of the components and the methods of generating and recording the input and output wave forms.

Paschkis (31) and Boukou (32) who are the pioneers of the direct analogue method have used

"slow time" scale and studied the thermal transients over long durations, of the order of several minutes, with ordinary galvanometers and potentiometric type recorders. The major drawback in such "slow time" analogues is the requirement of large time constants which are obtained by the use of capacitors of several hundreds of Microfarads and resistors of several Megohms. Cost of those components of the necessary quality is prohibitive, while leakage current present in the large capacitors, constitute serious source of error.

Burnand (43) has shown that for the study of periodic heat flow problems "fast time" analogues can be employed with advantage. A repetition frequency of 50 c/s was successfully used by them. Lator-Lawson and McGuire (33) described a device, for the study of transients of only a few milliseconds duration. Robertson and Gross (34) have used much faster time analogues with a time compression of the order of  $10^6$  (1 sec of electrical is equal to  $10^6$  sec of thermal) in their study of fire resistance characteristics of building elements. The fast time scaling has the advantage of requiring circuits of smaller time constants, permitting thereby the use of capacitors and resistors of smaller value and size. The leakage current problem is almost absent in this case. Owing precision

capacitors and resistors in these ranges is neither difficult nor expensive. However, the use of fast time analogue calls for more complicated electronic devices for the generation and recording of the special wave forms. The use of cathode ray oscilloscopes for measuring and recording the wave forms has become a common feature in such studies.

In all these studies, an electrical signal representing the outside temperature variation with time is to be generated. This is done electronically in several ways. Robertson and Gross (34) have used a photoformer type function generator (40) whereas Burnard employed special wave shaping devices (41). Input devices like diode function generators commonly used in analogue computers, can also be used with these direct analogues. In this study a fast time scale (a frequency of 100 c/s) is used. A brief description of the analogue setup and the experimental procedure adopted is given below.

### 3.2 Resistance - Capacitance Analogue

A building element, whose transient thermal characteristics are to be studied, is often represented on the electrical analogue model as a Resistance - Capacitance lumped system consisting of a cascaded Y, L  $\pi$  networks. The actual value of the electrical

resistance and capacitance of each section of the network lump depends on the thermal conductivity, specific heat, density of the material and the thickness of the building element, the choice of the scaling factor and the number of lumps used to represent the element.

In order to make the analogue representation more flexible, so that any material or combination of materials (with their usual thicknesses), can be represented as an R-C network, panels of resistances and capacitances to provide any desired value of R and C, are to be constructed. Sufficient number of such panels should be available in order that an analogue lumped representation, of dense as well as thick materials is feasible. For this study, twelve panels, each consisting of fixed resistors permanently connected on proper insulating supports having values from 100 ohms to 1 Megohm, and capacitors having values from 10 pF to 20,000 pF, were constructed. The resistors and capacitors were so connected that any desired resistance between 100 ohms to 2.1111 Megohms and any capacitor between 10 pF to 61,110 pF, can easily be selected in a short time by a suitable plugging arrangement.

If the required resistance and capacitor values fall beyond the range provided by the panels,

either the resistance scaling factor or the capacitance scaling factor may be suitably altered without affecting the time scaling factor (frequency) or the time scaling factor itself may be varied, thus automatically affecting both the resistance and capacitance scaling factors.

The ease with which the selection of the resistance capacitance values, is one of the salient features of the panel design. This arrangement not only facilitates the construction of analogue models on the panels, but also permits the use of a given model to simulate a whole range of structural thicknesses, simply by a suitable change of the time scale (i.e., frequency).

All the resistors used were of precision cracked carbon type with 1% tolerance and 1 watt rating. All the capacitors are precision polystyrene type with 2% tolerance.

All the individual resistors and capacitors were checked for their rated values with a precision G.I. type 1650 A impedance bridge before their installation in the panels. The front and back views of the panels are shown in plates 1 and 2 respectively.

The arrangement of the resistors and capacitors in a single unit is diagrammatically shown in Fig. (3.1). As all the resistances in this unit are

RESISTANCE CAPACITANCE PANELS

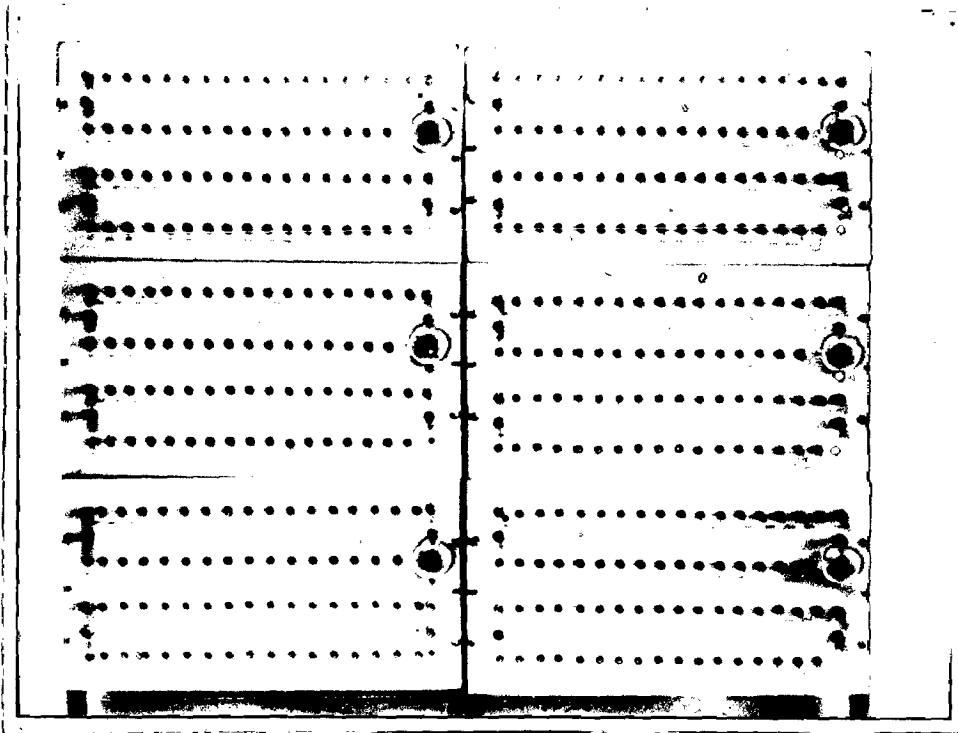


PLATE 1 Front View

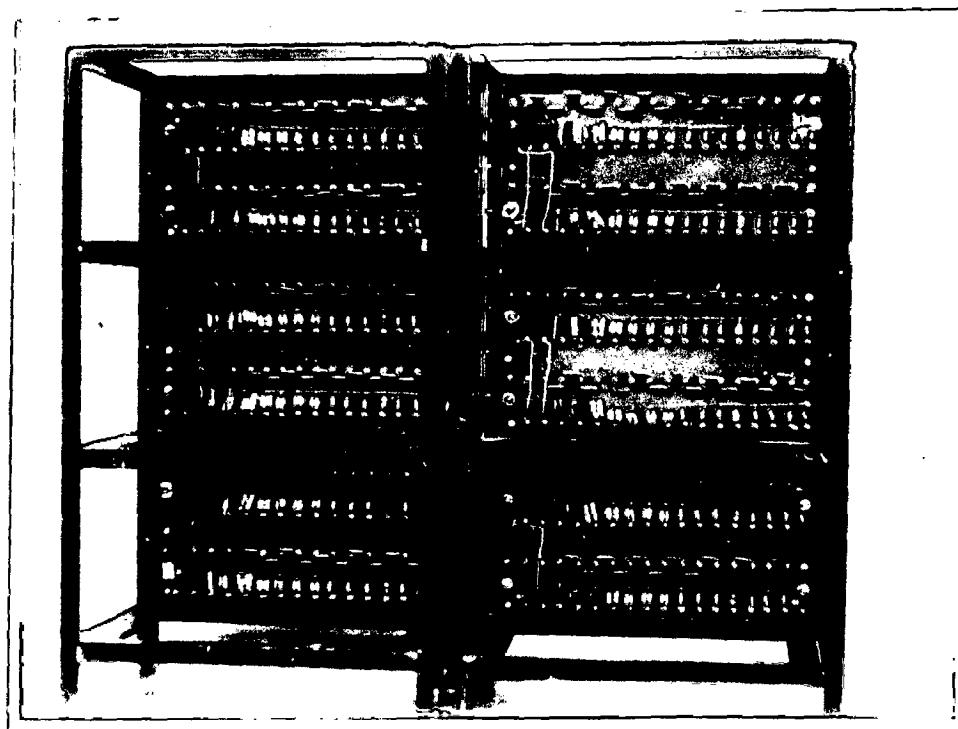


PLATE 2 Back View

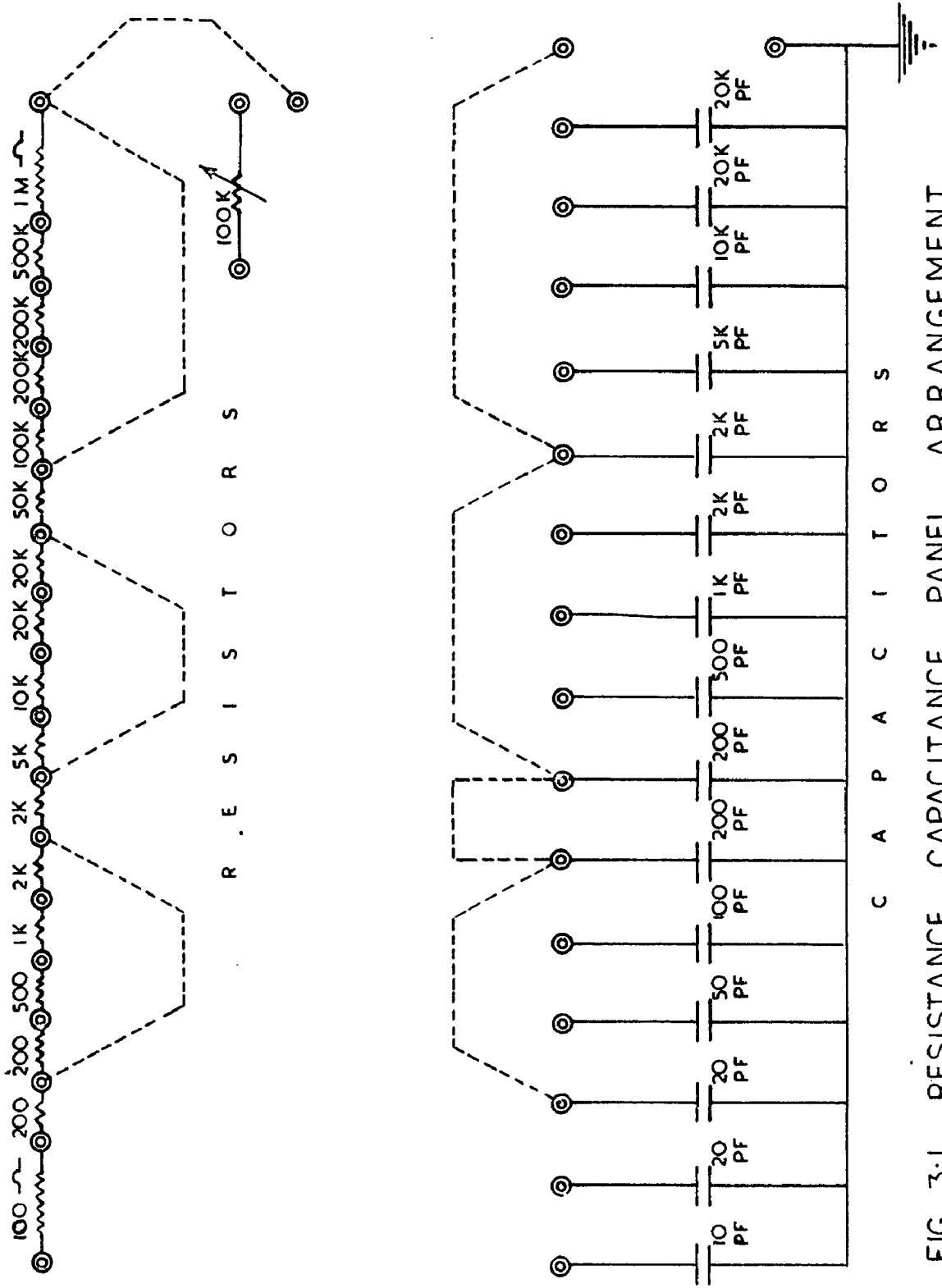


FIG. 3·1. RESISTANCE CAPACITANCE PANEL ARRANGEMENT

permanently connected in series, those resistances between the plugged points are shorted. The plugged in terminals of the capacitor units, come in parallel and the total capacitance of the circuit will be additive. Finally at the output terminals the required resistance and capacitance values are obtained. As an example, the dotted lines indicate the external plugging connectors to obtain a value of 62.3 Kilo ohms resistance and 2420 pF capacitance. In this manner any desired value of 'R' and 'C' can easily be selected from each unit quickly.

Apart from those twelve units of fixed resistors, six units of linear tan turn helical potentiometers (Heliopots) each of 100 K ohms resistance, are also installed on the panels to represent purely resistive circuit elements like surface film resistances and enclosed "Air spaces".

### 3.3 Scaling Factors

The relevant thermal quantities of a building element that are to be determined and converted into equivalent electrical quantities, (in order to represent it on the Electrical analogue) are its Thermal resistance and Thermal capacity.

The thermal resistance and the thermal

capacity per unit area depend upon the thickness (l) of the material and the thermal conductivity (K) specific heat (s) and density (P).

Thermal resistance  $R_T$  is given by  $\frac{L}{K}$

Thermal capacity  $C_T$  is given by  $L \rho s$ .

where

L = thickness in Ft.

K = thermal conductivity in Btu/(Ft.)(hr)(°F).

s = specific heat in Btu/(lb)(°F).

P = density in lb/Ft.<sup>3</sup>

Once the thermal resistance and thermal capacity are calculated from the properties of the material, their corresponding electrical resistance and capacitance can be determined with use of suitable scaling factors. Scaling factors actually correlate the thermal quantities to electrical quantities. There are three independent scaling factors, namely :

- 1) Potential scaling factor

$$n_p = \frac{\text{Unit of potential in the thermal}}{\text{Unit of potential in the electrical}} = \frac{\text{O.F.}}{\text{Volt}}$$

- 2) Time scaling factor

$$n_t = \frac{\text{Unit of time in thermal}}{\text{Unit of time in electrical}} = \frac{\text{hr.}}{\text{sec.}}$$

3) Resistance scaling factor

$$\begin{aligned} \eta_R &= \frac{\text{Unit of resistance in thermal}}{\text{Unit of resistance in electrical}} \\ &= \frac{(^\circ\text{F})(\text{Fr})(\text{Ft})/\text{Btu}}{\text{ohms}} \end{aligned}$$

Other scaling factors, for capacitance and current are interpolated with the above three independent factors. The product of resistance and capacitance is called time constant of the circuit and is of dimensions of time. Hence the relation between the time scaling factor and the resistance and capacitance scaling factors is given by

$$\eta_T = \eta_R \times \eta_C$$

Similarly the relation between the Flux scaling factor and the potential and resistance scaling factors is given by

$$\eta_\Phi = \frac{\eta_V}{\eta_R}$$

In the present study a thermal cycle of 24 hour period has been represented by an electrical generator with a frequency of 100 c/s. This gives a time scaling factor of 1/2400 sec of electrical time i.e., one hour of thermal time corresponds to factor  $\eta_T$  of 1/2400 of electrical time. By this a time compression of  $8.64 \times 10^6$  is obtained.

Once the time scaling factor has been fixed there is a choice in fixing either the capacitance or the resistance scaling factor. The capacitance scaling factor was chosen as  $N_C = 10^8$  so that 1 Btu/ $^{\circ}$ F corresponds to 10000 pF. This automatically fixed the resistance scaling factor as  $N_R = 41.66 \times 10^3$  i.e., one unit of thermal resistance corresponds to  $41.66 \times 10^3$  ohms.

### 3.4 Representation of Electrical Model

Using these scaling factors the total equivalent electrical resistance and capacity of the building element to be represented on the Analogue model are easily obtained. As mentioned earlier, to represent the building component accurately on the analogue as a lumped R.C. network, the number of lumps necessary have to be known. This problem of lumping is dealt in Chapter (4) in detail. Let us suppose that it is required to use 'N' lumps of 'T' circuit elements. The resistance value of each arm of the 'T' network will be equal to  $R/2N$  and the shunt capacitance will be equal to  $C/N$ . These values of  $R/2N$  and  $C/N$  are to be selected from each unit in the manner described earlier. As an example of the manner in which an electrical model may be constructed, a brick wall, 9" thick, is considered below. It is assumed that the

model is to have 8 sections of 'T' type cascaded circuit. This is shown in Fig. (3.2). The thermal properties of the brick wall considered are :-

Thermal conductivity  $K = 6.3 \text{ Btu/Ft}^2/\text{hr}/{}^\circ\text{F}$  per inch.

Specific heat 's' = 0.200 Btu/Lb/{}^\circ\text{F}

Density 'ρ' = 120 Lb/Cu.Ft.

The total thermal resistance of the wall ( $R_T$ )

$$= \frac{9}{6.3} = 1.423({}^\circ\text{F})(\text{hr})(\text{Ft})^2/\text{Btu}.$$

$$\text{Thermal resistance per lump (R/l)} = \frac{1.423}{8} = 0.1785$$

Thermal resistance for each arm of the 'T' circuit

$$R/2l = 0.03935$$

$$\text{The total thermal capacity} = \frac{9 \times 0.200 \times 120}{12}$$

$$C_T = 18 \text{ Btu/}{}^\circ\text{F}$$

$$\text{Thermal capacity per lump} = (C/l) = \frac{18}{8} = 2.25 \text{ Btu/}{}^\circ\text{F}$$

The electrical equivalents are obtained as

$$R_0/2l = 0.03935 \times 41.66 \times 10^3 \text{ ohms}$$

$$C_0/l = 2.25 \times 10^4 = 22.5 \times 10^3 \text{ pF.}$$

These can be easily represented on the panels. Then the equivalent resistance values of the exterior and interior wall surface heat transfer coefficients are to be calculated and added on either side of the lumped LC network. This completes the representation of the wall panel with its boundary conditions on the analogue.

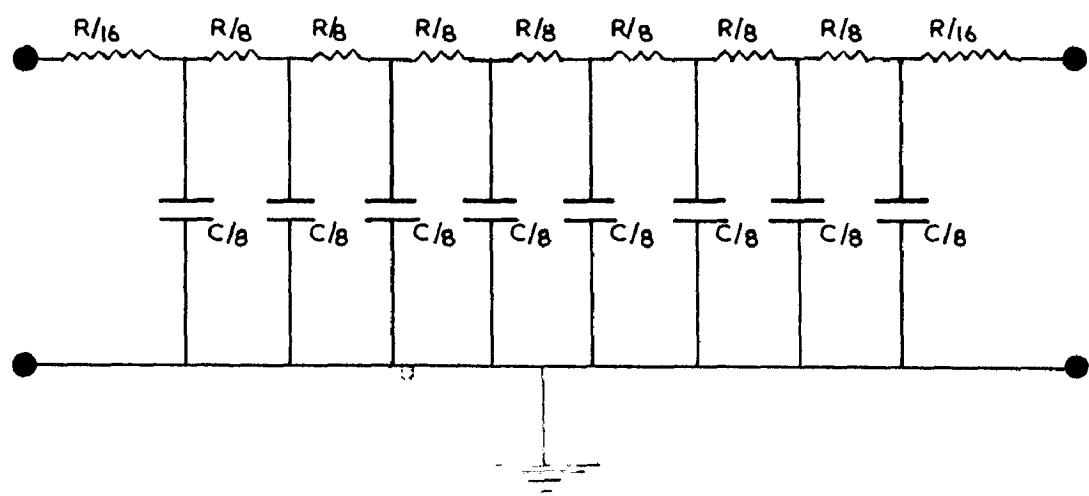
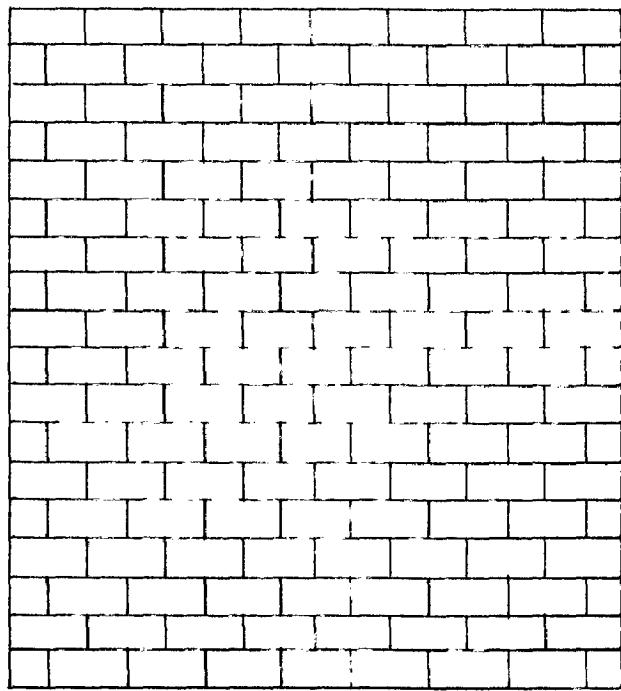


FIG. 3·2. R-C LUMPED T NETWORK REPRESENTATION  
OF A BRICK WALL

### 3.5 signal Generation

The next requirement is to generate an electrical signal to simulate (a) the external air temperature variation, (b) the solar radiation, and (c) the low temperature radiation exchange at the external surface. Mackay (1) has introduced the concept of sol-air temperature, which combines all the above phenomena into a single temperature variation. This approach has the advantage of requiring only one potential generator instead of three separate ones representing each of these three phenomena. The sol-air temperature depends on many factors, such as, latitude, time of the year, orientation of the surface, surface colour, wind speed, outdoor air temperature, dust content and precipitable water vapour present in the atmosphere etc. All these quantities are either available from Meteorological data or can be computed to a fair degree of accuracy. The sol-air temperature variation with time is then computed for any given situation. An electrical wave form representing the sol-air temperature is to be generated. As mentioned earlier, special function generators which demand complicated electronic equipment are to be employed for this purpose.

### 3.6 Steady State Sinusoidal Method

An alternative approach is to analyse the sol-air temperature wave form and represent as a Fourier trigonometric series of harmonic waves of different amplitudes and phases, super imposed over a steady level.

$$t_{sa}(T) = t_{sa}(\text{mean}) + \sum_{n=0}^{\infty} t_{sa,n} \cos(\omega_n T - \gamma_n)$$

For all practical purposes, three or four harmonics are sufficient for a good representation of any periodic wave form. This approach makes the problem of input signal generation very simple. Any stable commercially available Audio Oscillator which can generate 100 c/s to 1000 c/s will do the job. This obviously eliminates the need for complicated expensive equipment. The amplitude decrement ( $\lambda$ ) and phase lag ( $\Phi$ ) characteristics of the network model for the fundamental and three or more higher harmonics have to be determined once only. Then the inside surface temperature equation can readily be obtained as

$$t_{is}(T) = t_{is}(\text{mean}) + \sum_{n=0}^{\infty} \lambda_n t_{sa,n} \cos(\omega_n T - \gamma_n - \varphi_n)$$

Another major advantage of this approach is apparent. The amplitude decrement and phase lag

characteristics of a given R-C network depend only on the frequency but not on the amplitude of the input wave form. Hence, the inside surface temperature variation under any external climatic conditions can easily be obtained once  $\lambda$  and  $\Phi$  are determined. This is possible because even though the sol-air temperature profile is influenced by the variations of external climatic factors, (which in turn depend upon latitude, time of the year, orientation etc.), it can still be represented by the same Fourier series, with variations only in amplitude and initial phase angles. Instead, when the sol-air time temperature variations are impressed to the model directly, in its complex form, as for example by a photoform function generator, the output wave form will be true only for that condition. Hence, for each variation in the climatic factor, and the consequent variation in sol-air temperature, a new wave form should be generated and the output wave form recorded. This makes the material behaviour and the climatic factors inseparable. Another disadvantage with this method is that for measuring the amplitude variations and phase lag characteristics of such complex wave forms, cathode ray oscillographic recording has to be employed. This method of recording will not yield accurate measurement.

For sinusoidal excitation, there is no need

to record the wave forms, since for linear networks with such excitation, the response will also be sinusoidal with a lesser amplitude and a phase shift. Hence, only the amplitude of the output wave form and the phase shift introduced by the network have to be measured. Precise measurements of amplitude may be made with inexpensive equipment, like a Vacuum-tube Voltmeter. Accurate phase shift measurements, however, need special consideration.

### 3.7 Input Signal Generation

In the light of above advantages, the steady state sinusoidal method has been adopted in these studies. The only drawback of this approach is that it involves a little extra mathematical operation in carrying out the Fourier analysis of the sol-air temperature wave form. However, the advantages of this method over weigh this drawback.

A Hewlett Packard (H.P) type 201 C Audio Oscillator having a frequency range of 100 c/s to 10 Mc/s has been used as an input signal source. The fundamental frequency used is 100 c/s. The frequency calibration and stability of the oscillator was checked against a precision electronic frequency counter, (H.P type 523 B) which can read upto 0.1 c/s. The oscillator was found to be highly stable, the frequency

drift was only of the order of a cycle over a prolonged period of several hours. It was also found, that the voltage output was also constant over long periods and had a flat response within the full frequency range. The output impedance of the oscillator was quite low (600 ohms) and hence the voltage output was not affected to any significant extent by the network.

### 3.8 Measuring Equipment

The amplitude (peak to peak) of the voltages of the input and output wave forms of the networks were measured by an R.C.A. Senior Voltammist vacuum tube voltmeter type HV CB A. This V.T.V.M. had an input impedance of 0.33 megohms shunted by 70 pF for A.C. measurements. Though this is a high input impedance compared to the most of the network simulations, in certain cases, especially while studying the insulating materials, the network output impedances were of the order a few hundred kilo ohms when the loading of the network by the measuring meter became noticeable. To prevent this loading effect a cathode follower unit was designed and employed as a low impedance matching unit. This unit has an input impedance of greater than 60 meg-ohms and the output impedance of 600 ohms. The cathode follower will have gain of less than unity and for this unit it was found to be 0.9, and was very stable in

operation. In order to get the correct output of the network the actual measured amplitude at the cathode follower terminals is to be multiplied by a factor (10/9). The amplitude decrement factor ( $\gamma$ ) is defined as a ratio of the output to input amplitude. As the attenuation of the amplitude for some networks and especially for higher harmonics is sufficiently high, if an input voltage of unit amplitude is fed to the network the output may be of the order of few millivolts. The lowest voltage that can be read in the V.T.V.M. was 0.1 V (peak to peak). Hence, in determining the amplitude decrement factors, an input amplitude of 60 V (peak to peak) was applied and the output amplitudes are mostly within the measurable range. By dividing the output voltage (after correcting for the Cathode follower) by the input voltage (60 V) the amplitude decrement factor is obtained.

In spite of this procedure, in a few cases the output was found to be below 0.1 V. For such measurements a commercially available Electronic Millivoltmeter reading upto 0.1 mv was used after calibration.

### 3.9 Phase Measurement

Accurate measurement of phase lag angles posed a problem. Commercially made equipment were not readily available. Initially it was measured on a

double Beam oscilloscope. Solartron type CD 711 ~, by simultaneously displaying the input and output wave forms and the phase shift was noted on the calibrated expanded time base. Though this method enabled one to make quantitative measurements, the accuracy was not high, especially for small phase shifts. The other disadvantage was that the measuring procedure was time consuming and tedious.

To facilitate phase measurements, a direct reading audio phasemeter was designed and fabricated. The instrument reads with an accuracy of  $\pm 1^\circ$ , throughout the audio frequency range, over a wide range of input amplitudes 0.3 to 100 V and irrespective of the audio signal frequency under measurement. Photographs of the phase meter are shown in plates 3 and 4.

#### .10 Design Principles of Phase Meter

Various techniques have been developed for the phase angle measurements of periodic signals. These employed cathode ray tube method (50), vector addition of voltages (51) and null methods (52). But these techniques lack in precision and are not capable of measuring phase angles directly in degrees. In some methods the amplitudes of the two signals must be made equal and the calibration is frequency-sensitive.

# AUDIO PHASE METER



PLATE 3 Front View

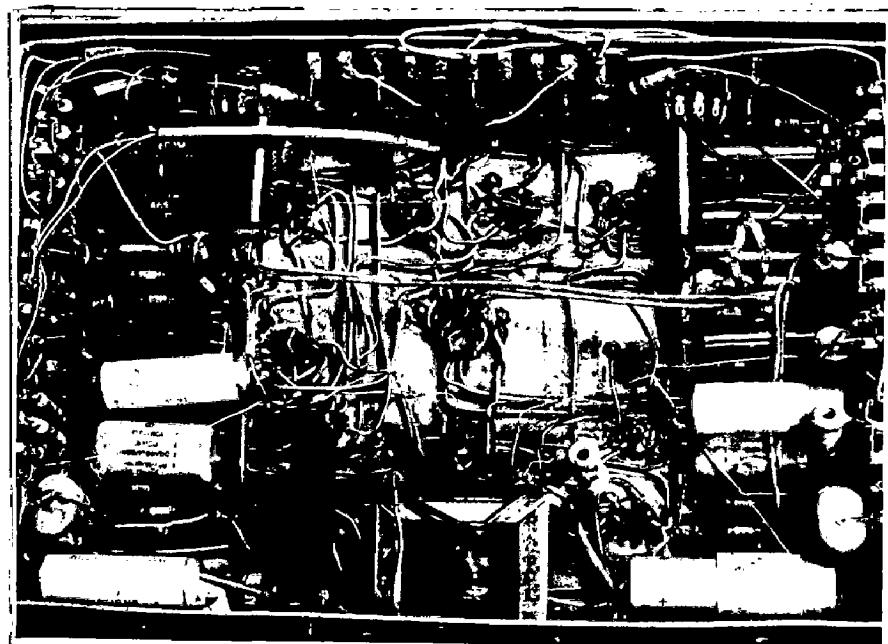


PLATE 4 Bottom View

## AUDIO PHASE METER

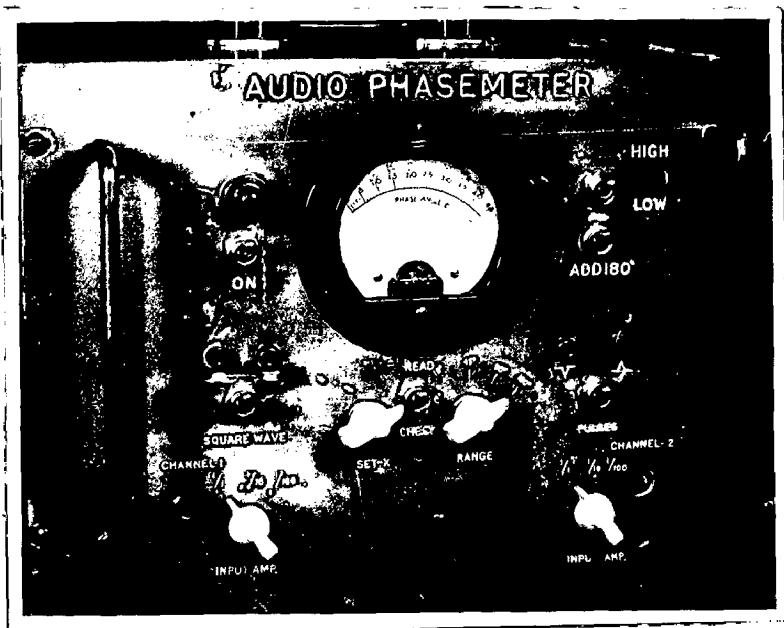


PLATE 3 Front View

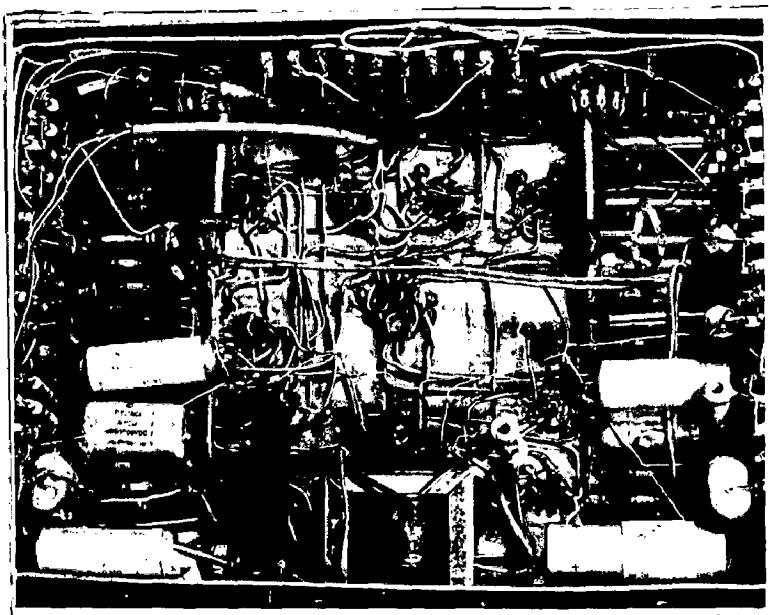


PLATE 4 Bottom View

Use of the foregoing methods have been superseded by direct reading phasemeters (53, 54). The basic principle of this type of phase meter is sensing the phase shifts as a relative time delay between the zero-axis crossings of the reference and phase delayed wave forms.

The direct phase measuring techniques consist of two operations namely conversion of input sine waves into square waves, preserving their times of zero-axis crossings, and an electronic device whose indications are linear functions of zero crossing-time intervals. Though the principle appears to be simple, complexities of circuit design are encountered in maintaining zero-axis crossings of input and output waves, throughout and in reducing the instrumental errors. The block diagram of the phase meter is shown in Fig.(3.3).

The instrument consists of two identical channels. The first stage of each channel is a cathode follower with a high input impedance (about 70 megohms) and prevents the loading of the network under study. The output of the cathode follower is clipped symmetrically at 1.5 volts before it is fed to the grid of the next high gain amplifier. Signal amplitudes whose peak value exceeds 1.5 volts are clipped off at this level, while smaller signals were clipped in the subsequent stages of the amplifier. With the aid of an attenuator,

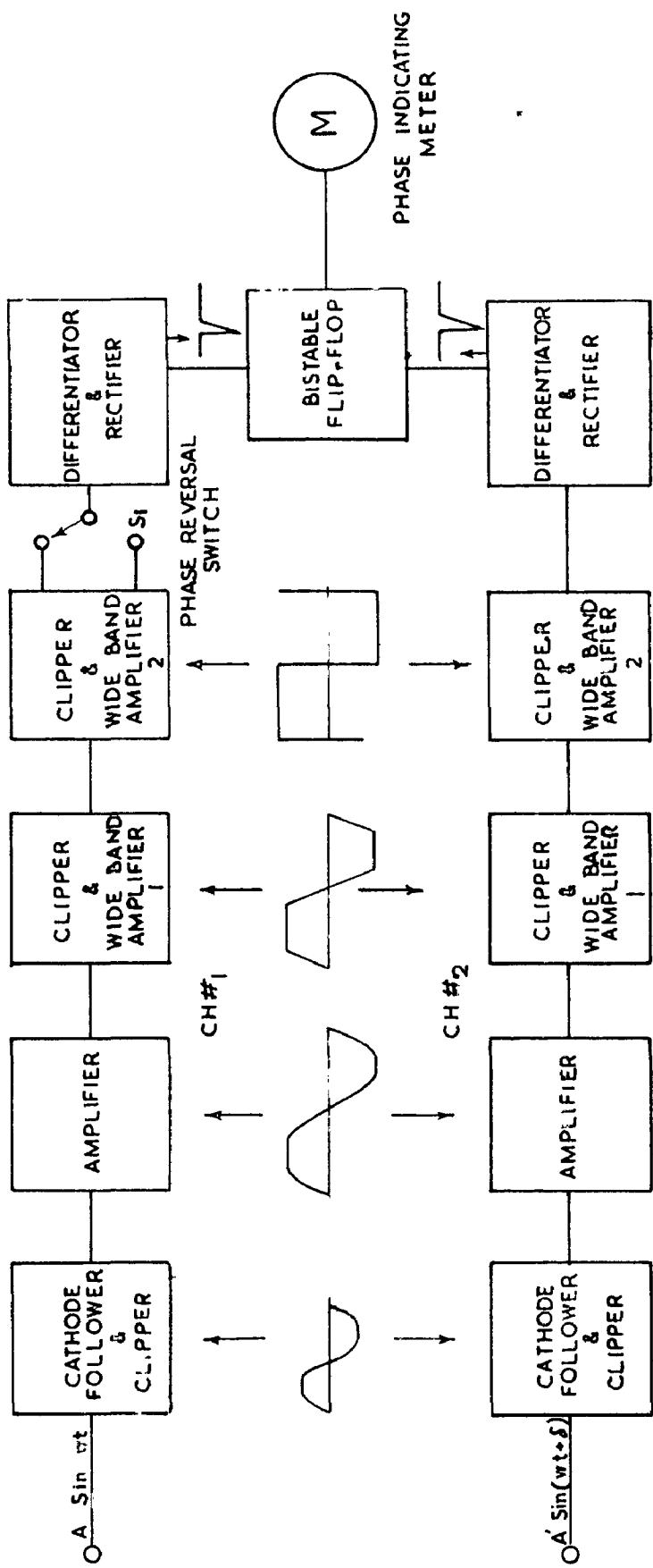


FIG.3.3 BLOCK DIAGRAM OF PHASE METER

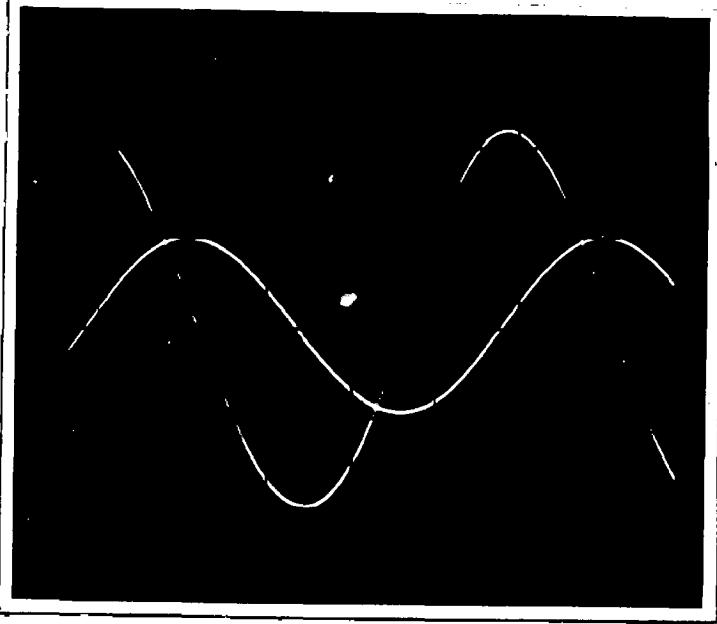


PLATE 5  
Input to the two channels  
(excitation & response of  
the R.C. Network)

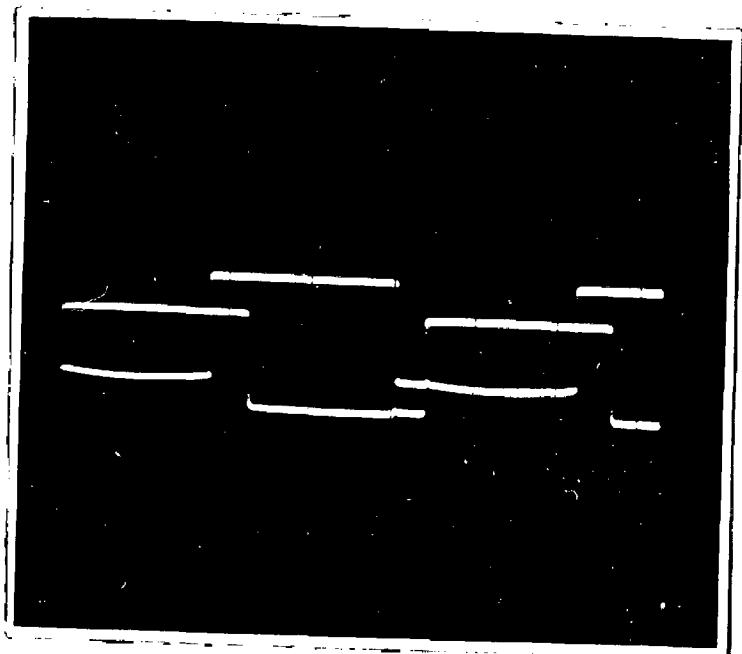


PLATE 6  
Symmetrical Squarewave  
output of the two  
channels.

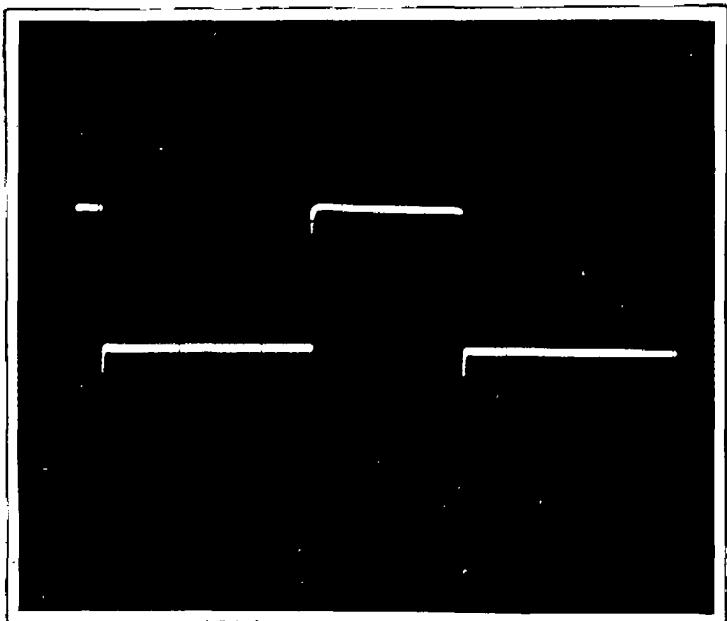


PLATE 7  
Bistable output  
corresponding to the  
phase difference

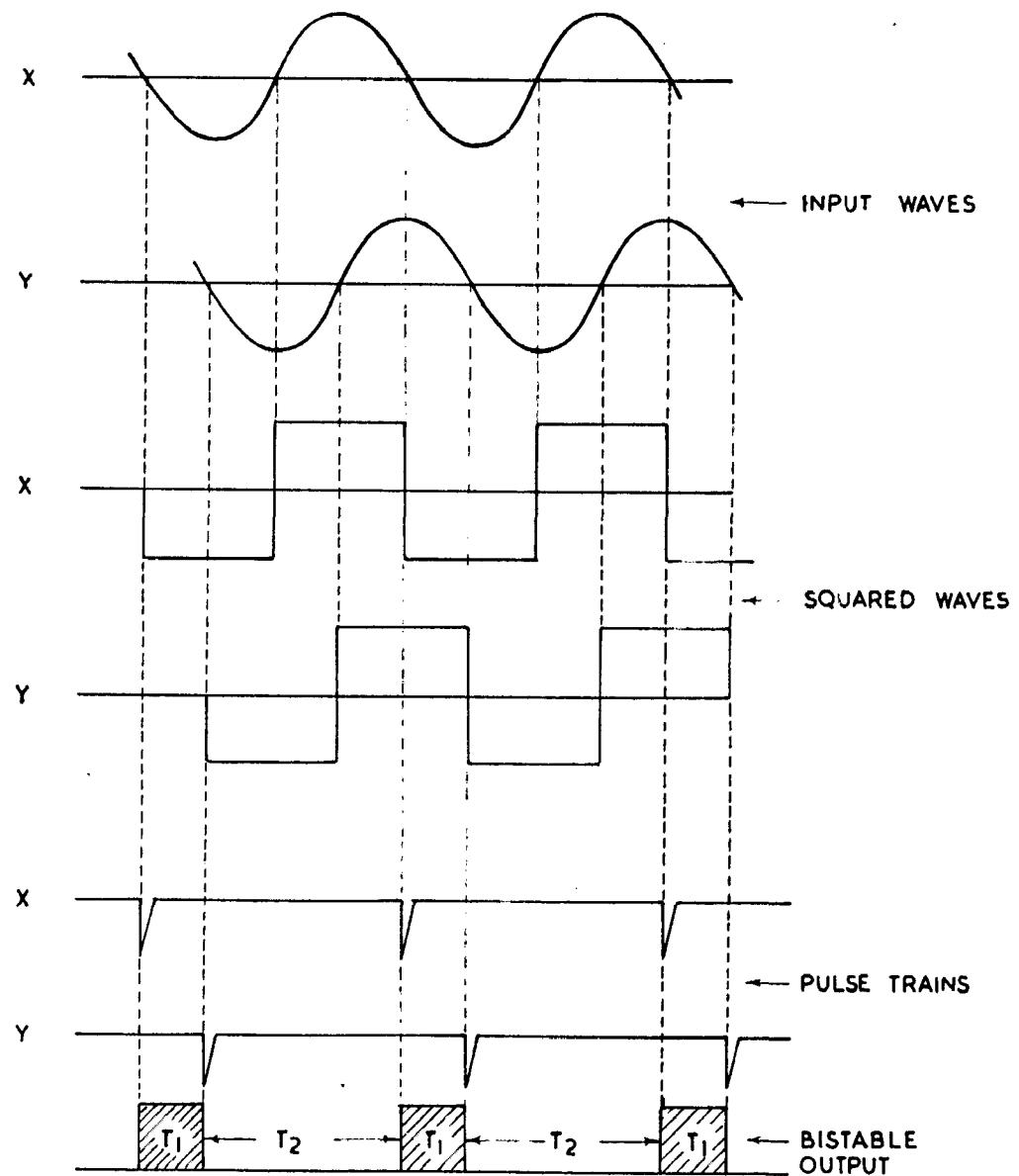


FIG. 3.4 PRINCIPLE OF PHASE MEASUREMENTS

again restoring (the conduction to A) the initial conditions and maintains it for a time interval  $T_2$ . The action is repetitive. It is obvious that greater the phase lag between the reference and the test signals, the larger is the fraction of a cycle for which B conducts and hence will be higher in the D.C. average indicated by the microammeter put in series with its anode. The phase difference is directly proportional to the fraction  $T_1 / (T_1 + T_2)$  and hence to the average current through the anode circuit of B. A typical bistable output with phase delayed signals is shown in plate 7.

$$\text{Phase angle } \phi \propto \frac{T_1}{(T_1 + T_2)}$$

$$\phi = c \frac{T_1}{T_1 + T_2} = c I_b$$

where  $c$  is a constant and  $I_b$  is the average current through the anode circuit of B indicated by the meter. The meter dial was calibrated directly in terms of phase angles (degrees) linearly. The phase angles from 0-360° degrees in six ranges 0-45°, 0-90°, 0-180°, 180-225°, 180-270° and 180-360° were obtained with suitable shunts to the microammeter.

The phase inverter incorporated in channel 1 was used to produce 180° phase shift. This arrangement

enables measurement of phase angles beyond  $180^\circ$  ( $180-360^\circ$ ) by transposing them into the lower ranges and hence without loss of measuring accuracy. The circuit diagram of the phasometer is shown in Fig. (3.8).

With this specially designed audio phase meter, the problem of accurate determination of phase lag angles has been simplified considerably and being a direct reading type, quick measurements are made possible, compared to the cathode ray oscillographic method.

The complete experimental setup for the determination of thermal system functions of building elements by the electrical analogue method is shown in plate 8 and illustrated by the block diagram Fig.(3.6). With this setup it is possible to simulate any type of building component, wall, roof or floor, etc. homogeneous or composite (multi layer), with little effort. This system is particularly suitable for the study of the effect of various parameters, especially, the influence of surface heat transfer coefficients on the overall thermal behaviour under periodic heat flow conditions.

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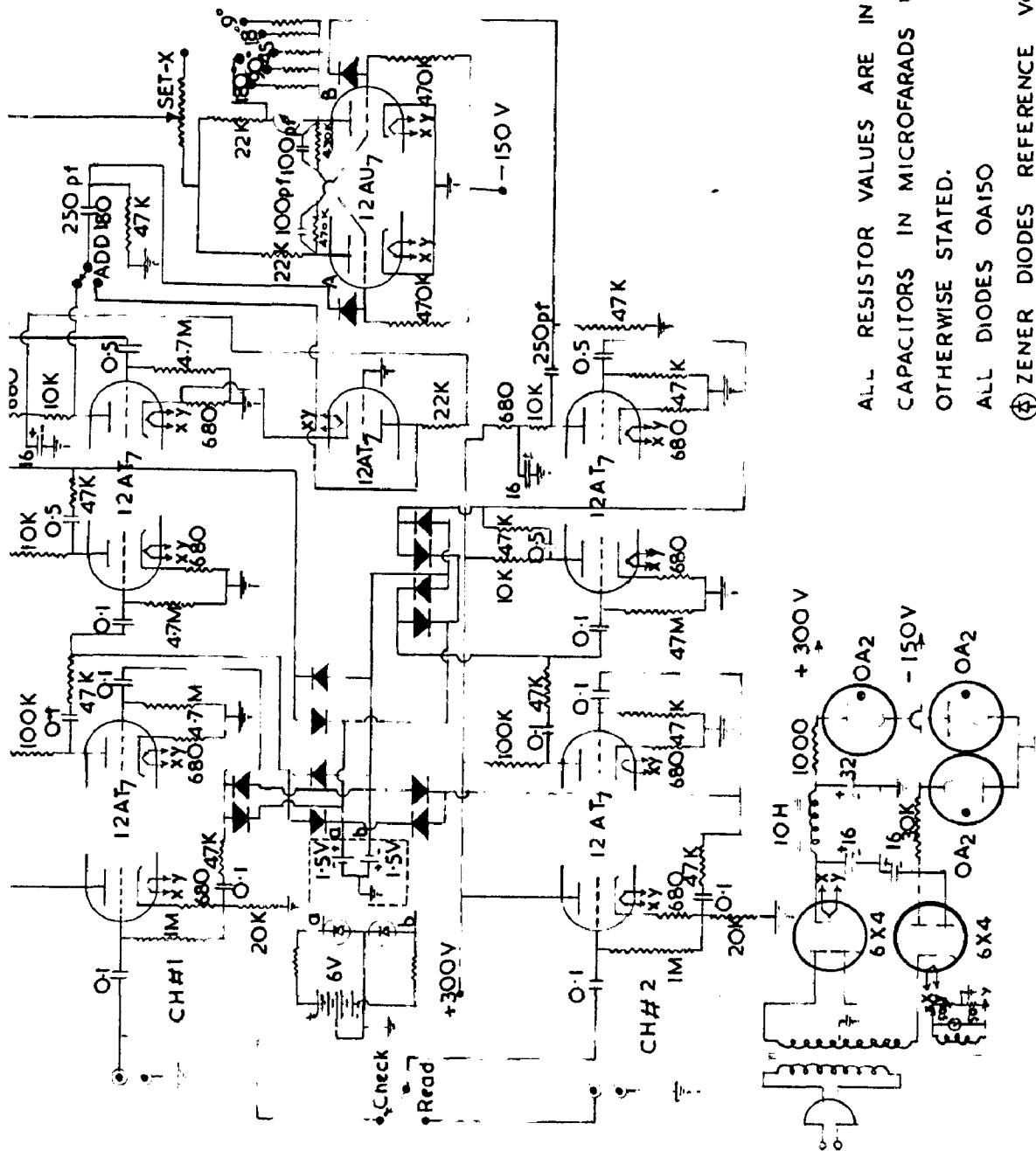


FIG. 3-5 SCHEMATIC CIRCUIT DIAGRAM OF THE AUDIO PHASE METER

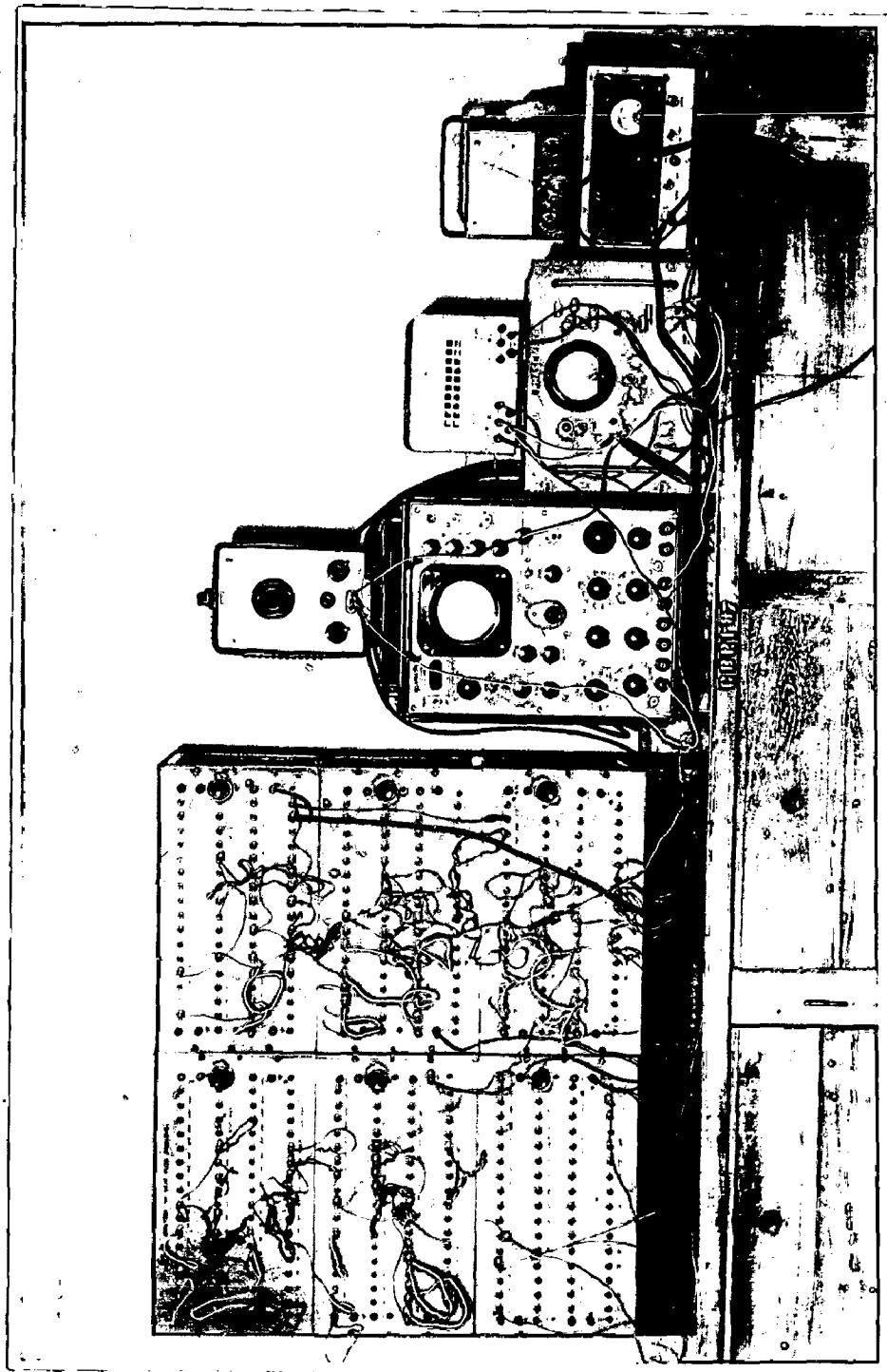


PLATE 8 Photograph of the Experimental Setup

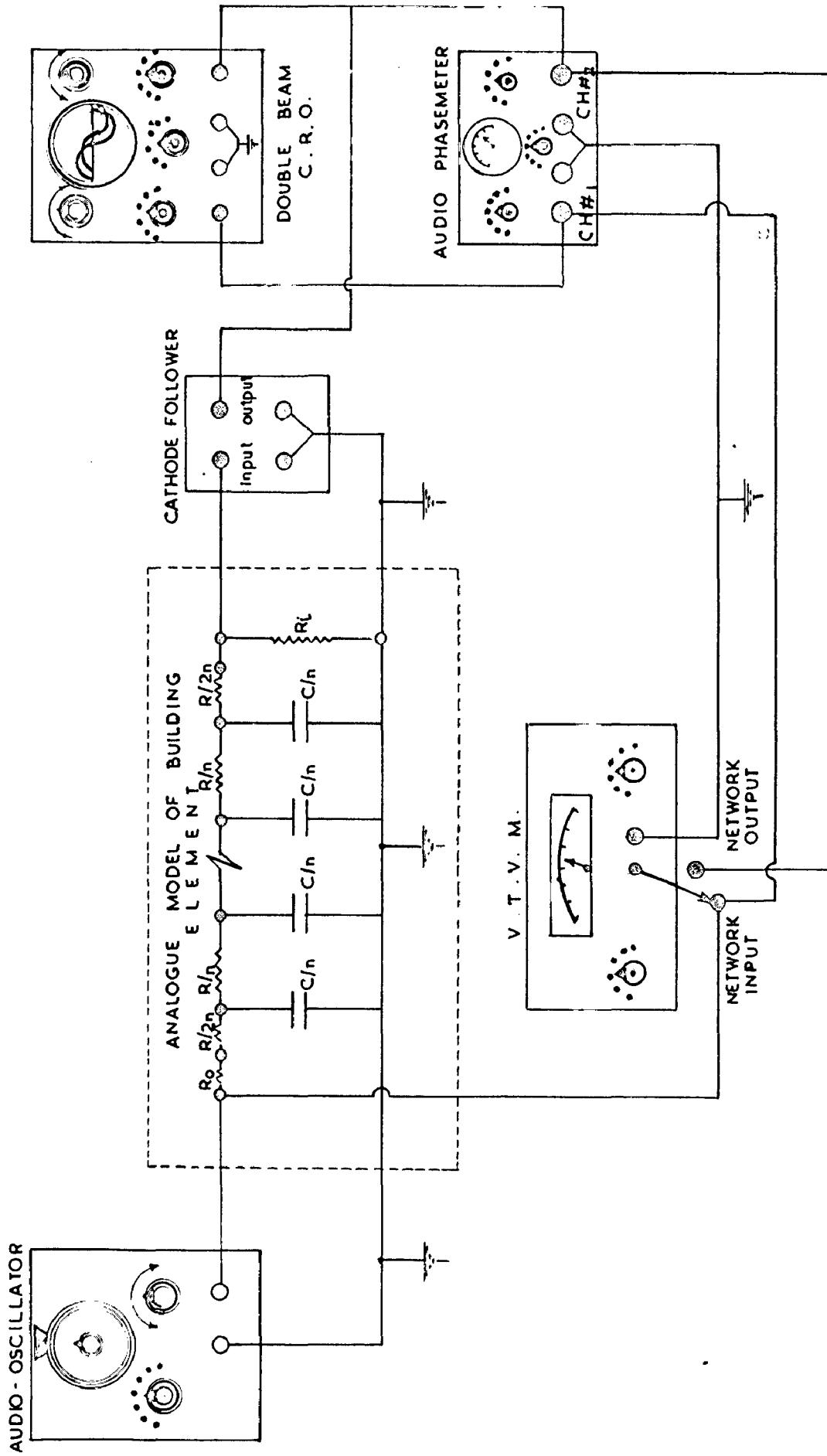


FIG. 3.6 BLOCK DIAGRAM OF THE EXPERIMENTAL SET UP

**C H A P T E R 4**

**LUMPING ERRORS**

C H A P T E R 4LUMPING METHODS4.1 Introduction

The heat flow through a body is analogous to the current flow through a non-inductive transmission line with distributed resistance and capacitances. To represent a building element in the electrical model by a transmission line of reasonable length is difficult. Hence, in the study of the problems of heat transfer through building elements, approximate representations by lumped circuits are used. Mathematically "lumping" is the representation of a partial differential equation by an approximate finite difference equation. The behaviour of the lumped circuit approaches the distributed system as the number of lumps are increased. However, there are certain practical limitations due to the component qualities and cost, which restrict the number of lumps used in the network. So the first problem is to determine the optimum number of lumps required to represent a thermal system for obtaining reasonably accurate results. The errors introduced due to lumping depend on many factors such as thickness and properties of the material, the type of lumped circuit ( $T$ ,  $L$  or  $\Pi$  networks),

number of lumps used, boundary conditions, surface film resistances, frequency of the input wave form and the location of the point of interest in the network.

Several investigators have studied this problem under specified boundary conditions. Puschkis and Neissler (55) have investigated the inaccuracies of ac lumped circuits at different depths and time intervals using a step function excitation. The influence of film resistances on the errors were also studied. In their studies, 13 lumps were taken to be the equivalent of the distributed value. The voltage variations (which corresponds to the temperature rise) with time at the junction points of the various lumped circuit elements were compared with those of the 13 lump values. Their results indicated that (i) five lumps gave sufficient accuracy except for relatively small values of  $t/\tau_C$  (ratio of the time at which measurements were taken to the time constant of the circuit), (ii) for high ratios of the film resistance to the circuit resistance the number of lumps used had no significant influence, (iii) inside film resistance had smaller influence than the external film resistance, and (iv) lumped circuit representation of the building element without the inclusion of film resistances, would present maximum errors. Lawson and McGuire (33) have

shown that for a semi infinite slab, with a step function excitation, the responses are identical after the first section for T, L and  $\pi$  sections. Klein, et.al. (56) have analysed the distribution errors in a finite bar, subjected to a step function excitation at one end, the other maintained at constant potential. They have found that the errors, at the centre of the bar have shown a high positive departure during the early part of the transient and reduced to negative errors - 6 per cent for 5 lumps and - 2 per cent for 11 lumps. Clarko (57) has analysed the error in a finite bar represented by different number of L sections one end of which is supplied w with a constant heat flux, the other end being perfectly insulated. It was found that the deviation of the lumped system behaviour from that of the distributed one, at the insulated end of the bar was always positive but decreased with time. Friedman (58) has estimated an upper limit of the error between the exact solution and its lumped parameter analogue for the case examined by Klein et.al. The solution for the lumped system of 'T' sections was given by Jaeger (59) and the exact analytical solution was provided by Carslaw and Jaeger (60). The errors were greatest at the first section and decreased rapidly for points farther off from the source. Robertson and Gross (34) have compared the percentage errors for open

and short circuit terminations for step function excitation, and concluded that the changes in the circuit termination do not result in appreciable change in the errors within the lumped circuit.

Most of the above investigators have considered the transient response for step function excitation. However, as mentioned earlier, the external boundary conditions for a building element can be represented by the soil-air temperature which can be expressed with sufficient accuracy as a Fourier series with three or four harmonic components. Hence, it is sufficient to consider the lumping errors for the fundamental and two or three higher harmonic sinusoidal inputs. Designing of lumped circuits for a limited frequency range leads to lesser number of sections than are needed for a transient response for a step function. Burnand (43) has suggested that division of a commonly used wall or roof element into 3 or 4 sections would be a sufficient approximation to the true analogy.

Mason (61) presented an approximate method of selecting the number of lumps, based on filter network analysis in terms of the wave length of the temperature wave through the medium. He has shown that if the length of such lumped section 'L' is  $\delta/8$ , the errors would be within 8 per cent of the distributed

one. Nakagawa and Paraboloo (36) have applied this one-eighth wave length representation to the periodic heat conduction through building elements and have derived an expression between the thickness of the material representing each lump and the thermal diffusivity of the material ( $a$ ), and the frequency of the thermal cycle ( $f$ ) as

$$L = 0.44 \sqrt{\frac{a}{f}} \quad \dots(1)$$

Draho et.al. (62) had pointed out that the above criterion may be useful for homogeneous structures, but do not hold good for composite constructions. They have also shown that the fraction of wave length needed for 0.5 accuracy varied between 1/23 to 1/12 for composite constructions. This indicates that in a composite construction the manner in which a given material is located in relation to other materials, influences the size of lump to be used for obtaining a specified accuracy. Stephenson and Nitulescu (63) have compared the theoretical frequency responses of distributed and lumped networks employing Matrix method. They have attempted to present a rational method of designing active as well as passive networks to represent homogeneous slabs, to any desired degree of accuracy. Recently Murray and Lendle (64) have derived matrix equations for estimating the lumping

errors of any parameter.

The determination of errors introduced due to lumping requires a comparison of the network solution with an exact analytical solution for the distributed system. The analytical solution is a difficult one particularly for composite constructions. Instead Drake et.al. (62) have adopted an experimental procedure in which the number of lumps are increased to an extent, where further increase of lumps does not significantly alter the transfer functions.

Though a good deal of theoretical as well as experimental studies were carried out on this problem of lumping, no definite generalisations regarding the choice of the numbers of lumps based on the knowledge of the thermal properties (Resistance and Capacitance) are available. Hence, a systematic study of the effect of various parameters on the lumping errors was carried out theoretically (Matrix method) and experimentally. The total errors in analogue studies are, due to lumping, component tolerances and measurements. In order to separate the lumping errors from other errors, theoretical calculation of lumping errors have to be made. The procedure followed here was to calculate the transfer functions for both the distributed and lumped circuit representations, starting with one lump and increasing to 10 lumps, for different

thermal time constants (RC) ranging from 0.5 to 200 and then to compare their moduli and arguments.

The following aspects were included in the study :-

1. Selection of the type of network (T, L or  $\Pi$ ) configuration, for obtaining a specified accuracy with minimum number of sections.
2. Determination of the variation of percentage errors in transfer function (amplitude decrements and phase lags) as a function of the thermal time constants (RC) and the number of lumps, for the best type of network (obtained from the results of (1)).
3. Evaluation of the influence of film resistances on the lumping errors of transfer functions.
4. Determination of the effect of input frequency on the lumping errors.

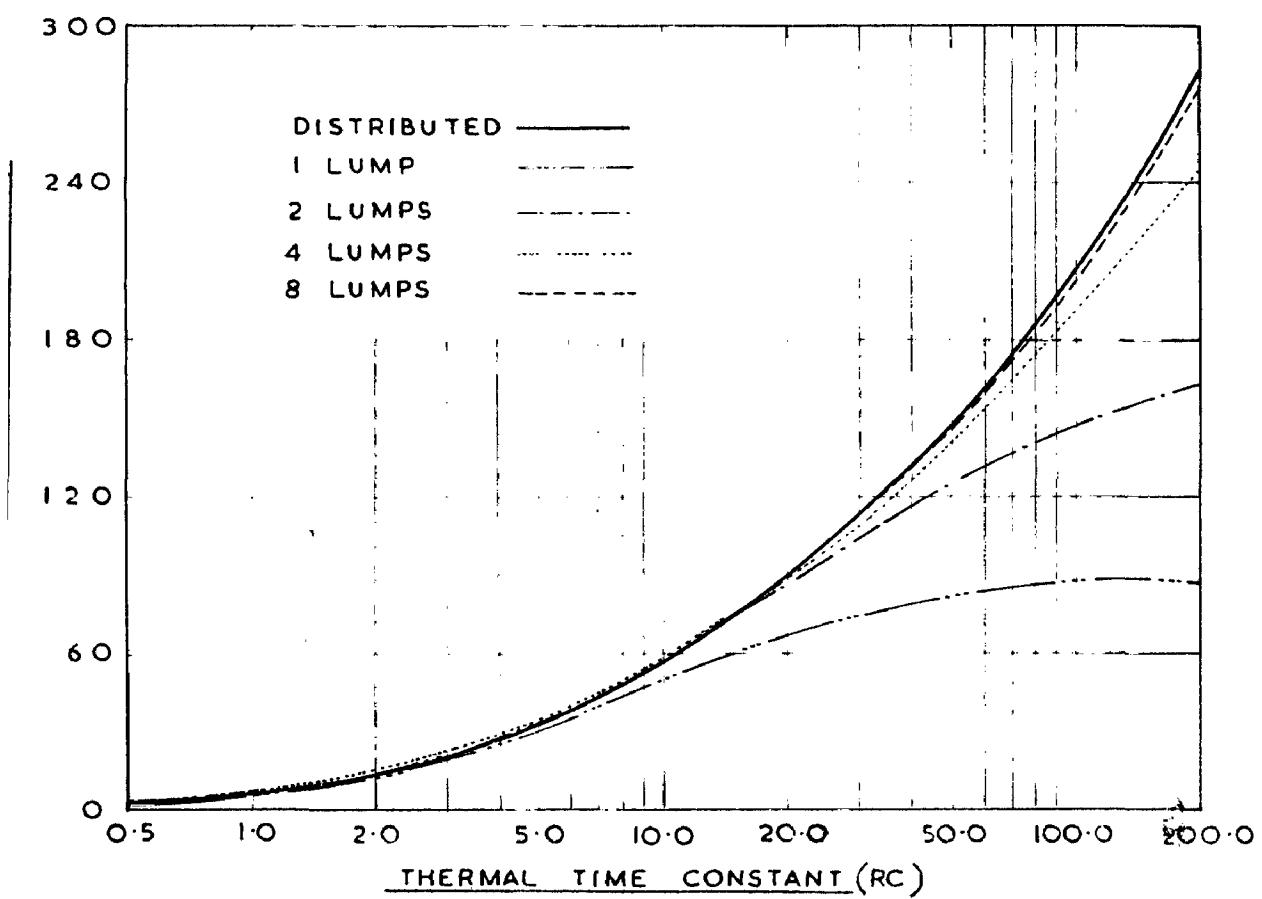
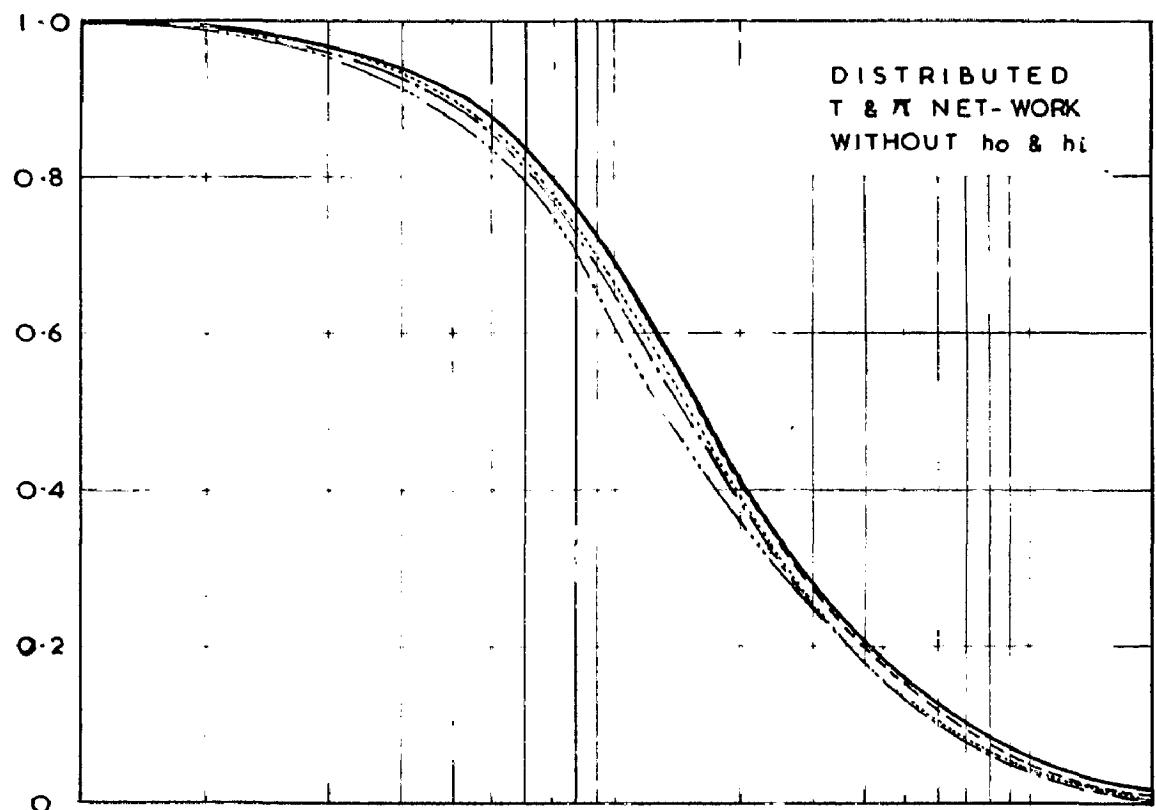
#### 4.2 Choice of the Network

To determine which type of network (T, L or  $\Pi$ ) will yield best results with lesser number of lumps for any building element of a given thermal time constant (RC), the transmission matrices of the lumped and distributed systems were computed and compared.

For a lumped circuit of 'N' network elements, the overall transfer matrix is obtained by matrix multiplication of the transfer matrix of each lump 'N' times. The overall transmission matrices for 1, 2, 4, G, 8, and 10 lumps, for all the three network configurations are calculated. The calculated lumped matrix elements are for RC values ranging from 0.1 to 200 and the distributed matrix elements are included in Appendix (I). This data will enable a quick estimation of the lumping errors, in the determination of the surface temperatures, heat fluxes and the thermal system functions.

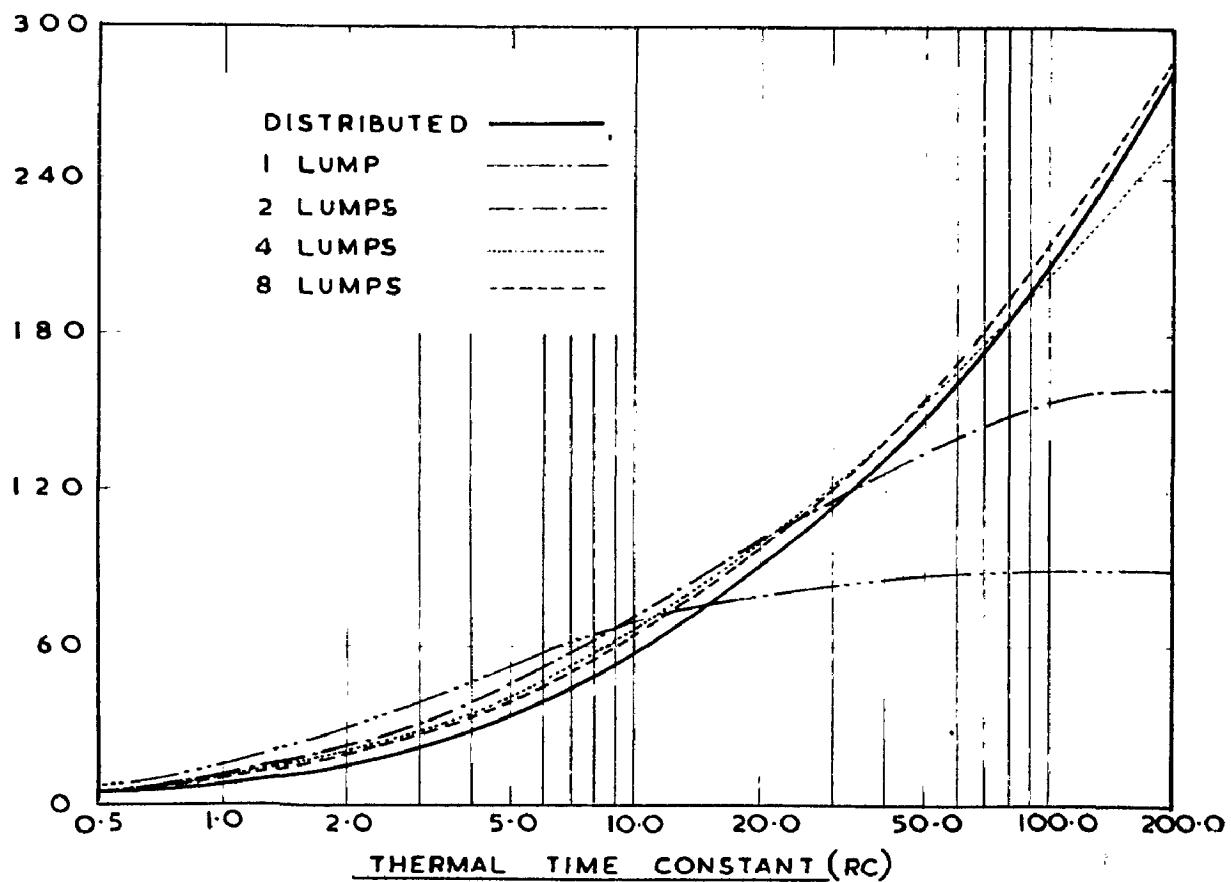
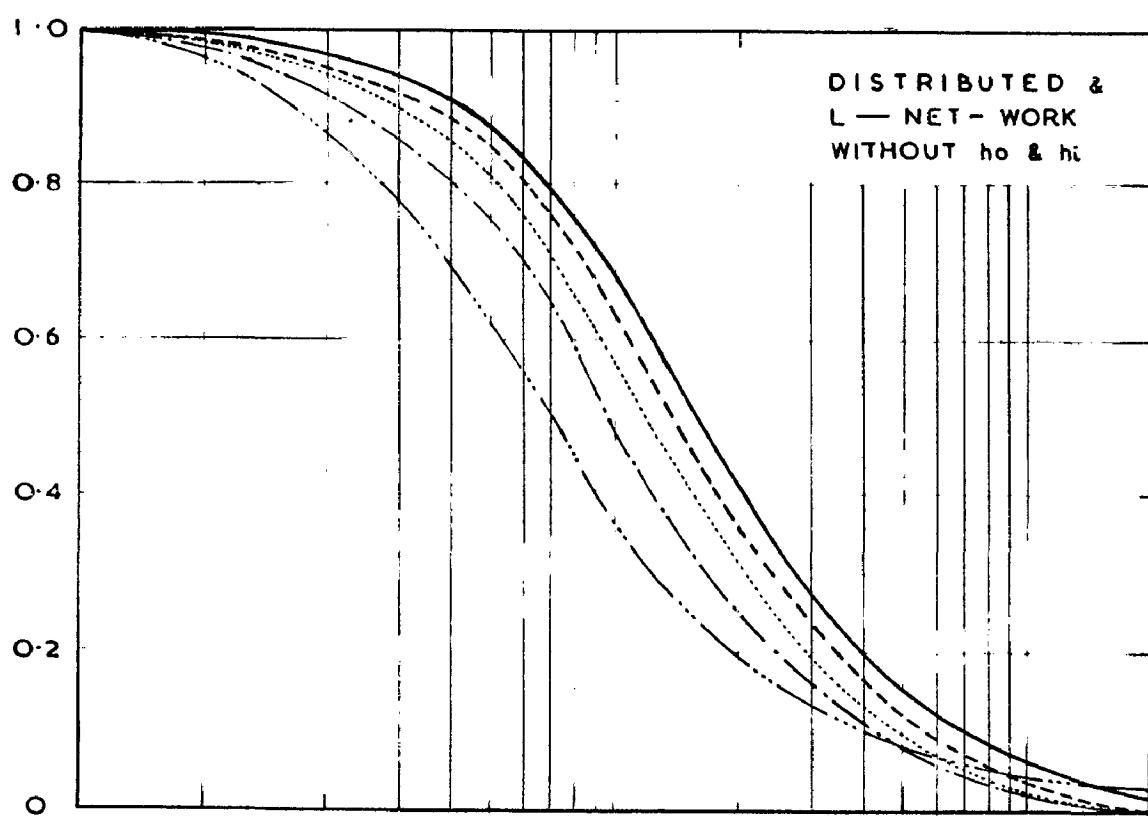
For surface to surface transmission with open circuit termination (considering the material alone without surface heat transfer coefficients) the transfer function is given by the reciprocal of the element 'A' of the transmission matrix i.e.,  $\frac{1}{A}$ . Where A is a complex number. In order to estimate the lumping errors, the modulus and arguments of the transfer functions for distributed and the lumped T and  $\pi$  networks of 1, 2, 4 and 8 lumps for RC values from 0.5 to 60, have been determined and compared in Figs. (4.1 and 4.2). The matrix element 'A' is the same for T and  $\pi$  networks.

The results have shown that the lumped T and  $\pi$  circuit configurations give less deviation



COMPARISON OF LUMPED (T &  $\Pi$  CIRCUIT) AND DISTRIBUTED TRANSFER FUNCTIONS

FIG. 4.1



COMPARISON OF LUMPED (L-CIRCUIT) AND DISTRIBUTED  
TRANSFER FUNCTIONS

FIG. 4.2

from the distributed values than the L type. In both the errors in amplitude as well as phase decreased with increasing number of lumps for any given RC value. For large RC values less than 4 lumps are totally inadequate. Based on those results, the 'T' type of network was finally chosen for further studies on the effect of various factors on the lumping errors.

#### 4.3 Thermal Time Constant Versus Lumping Errors

The transfer functions are computed, for distributed and lumped 'T' circuits of 1, 2, 4, 6, 8 and 10 lumps for RC values from 0.1 to 200 are given in Table (4.1). The percentage errors due to lumping and its reduction with the number of lumps were determined for matrix elements A, B and C (both modulus and argument) and are presented in Figs. (4.3, 4.4 & 4.5 respectively). These comparisons were made for the fundamental frequency which is most significant. The results of those studies are broadly summarised in Table (4.2)

TABLE (4.2)

Thermal time constant RC	No. of lumps required
Upto 1	1 lump
Between 1 and 6	2 lumps
" 6 and 20	4 lumps
" 20 and 50	8 lumps
" 50 and 100	10 lumps
Above 100	more than 10 lumps

TABLE II  
Effect of Surface Coefficients on the Critical Current Density of a Superconductor

(Independent of surface coefficients)

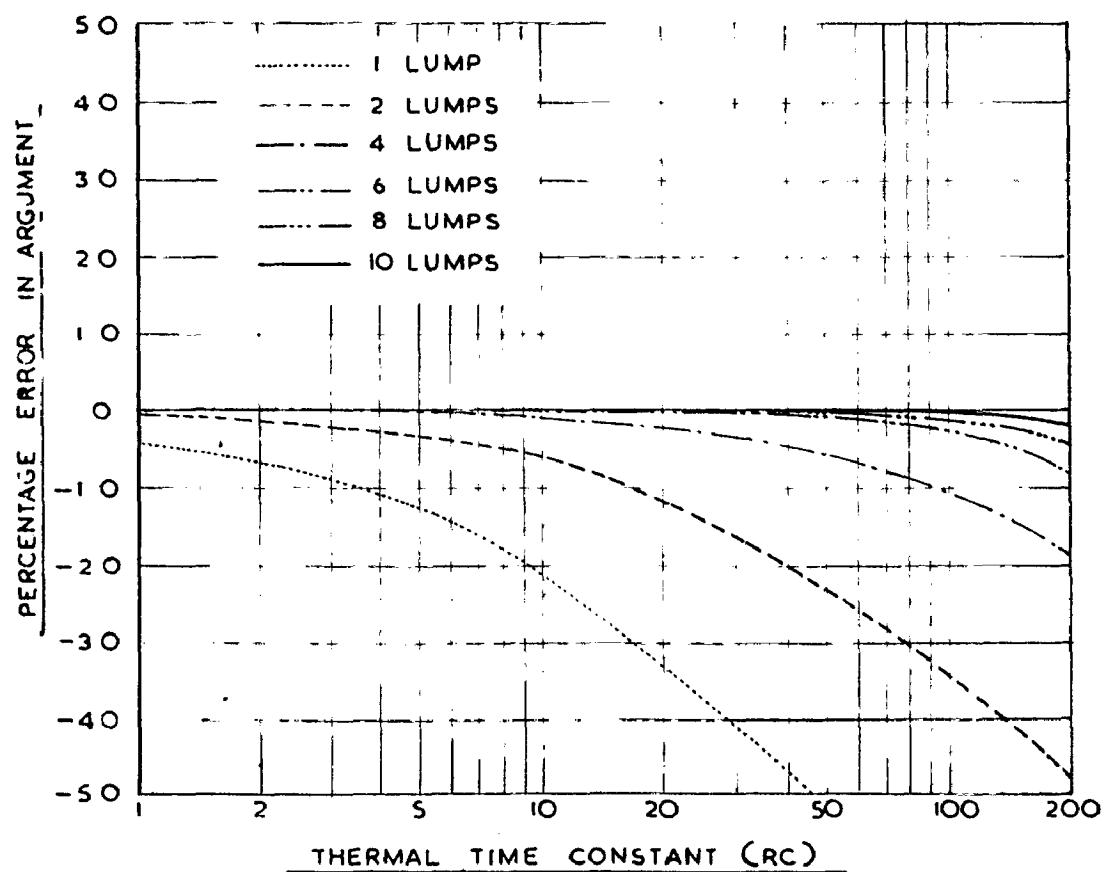
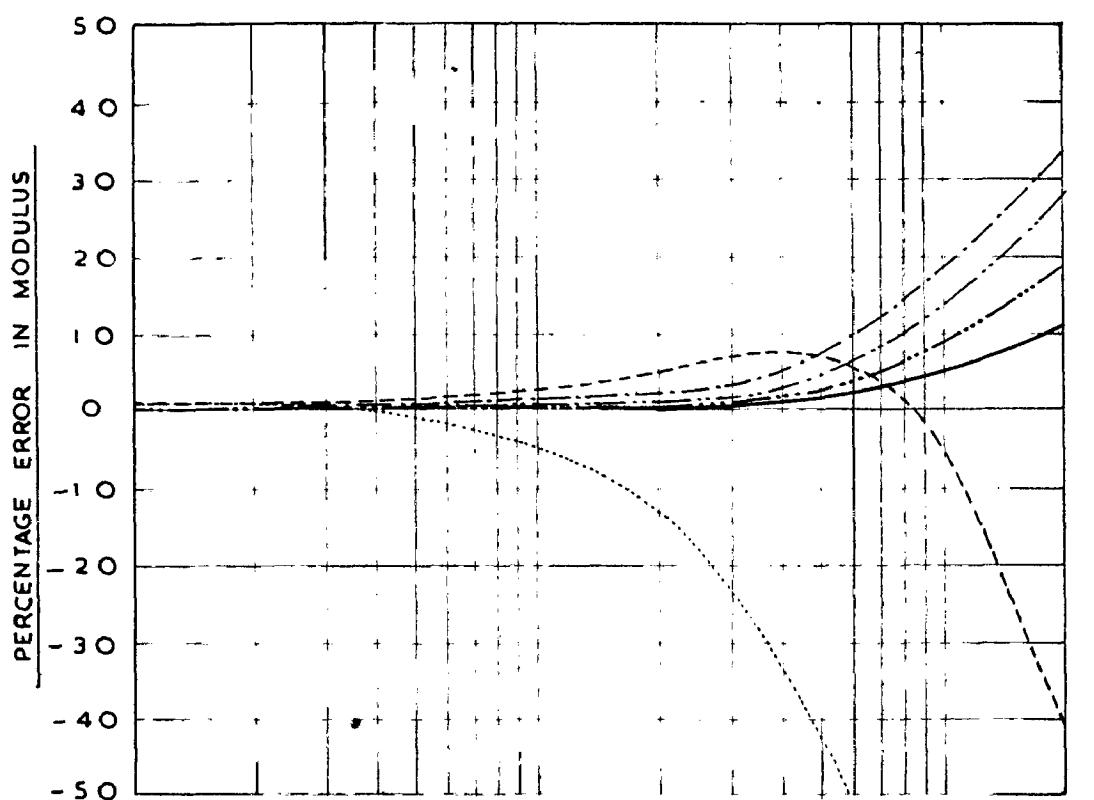


FIG.4.5. LUMPING ERRORS OF TRANSMISSION MATRIX  
ELEMENT - C  
(T - NETWORK)

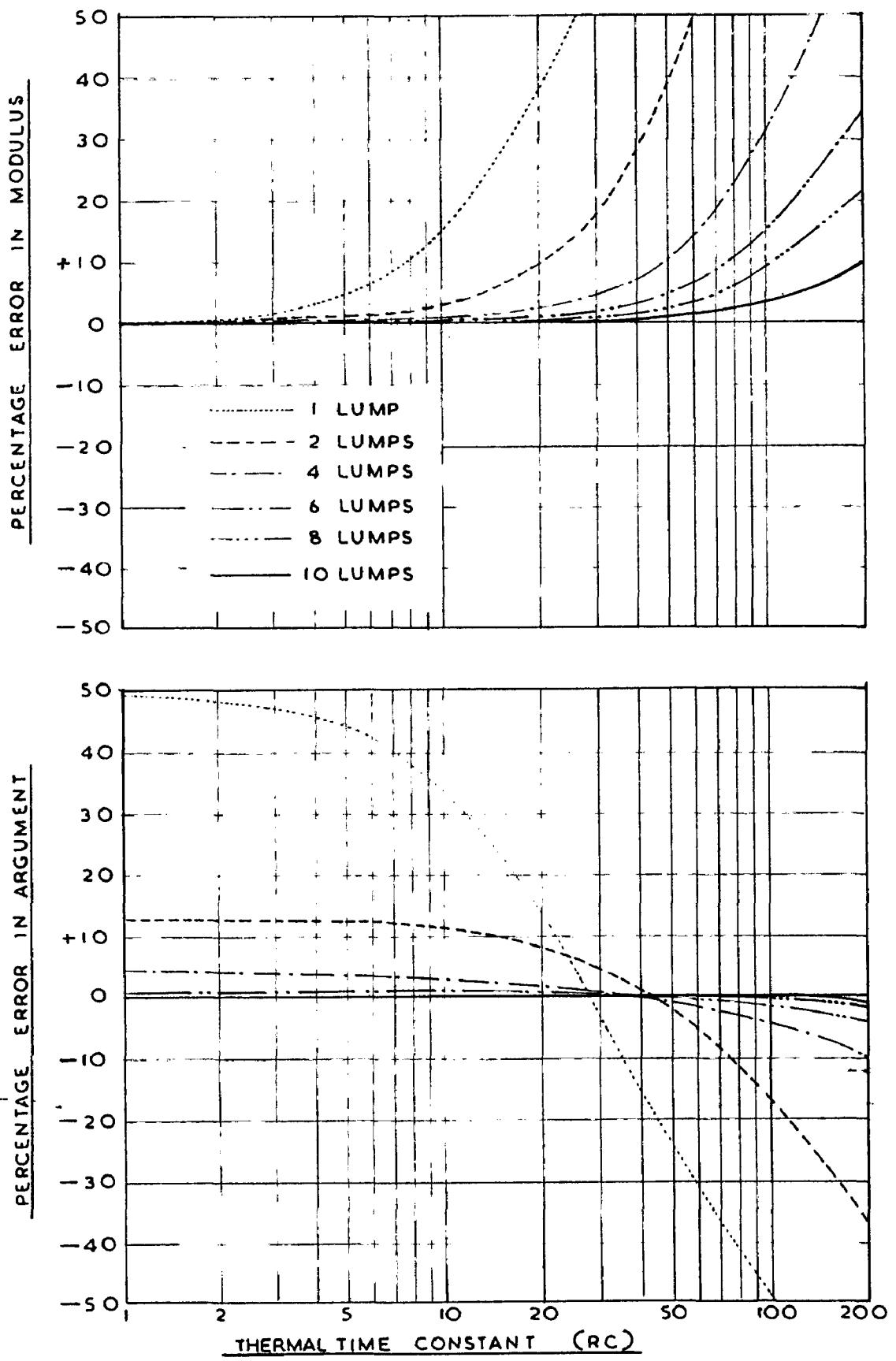


FIG. 4.4 LUMPING ERRORS OF TRANSMISSION MATRIX

ELEMENT T - B  
(T - NETWORK)

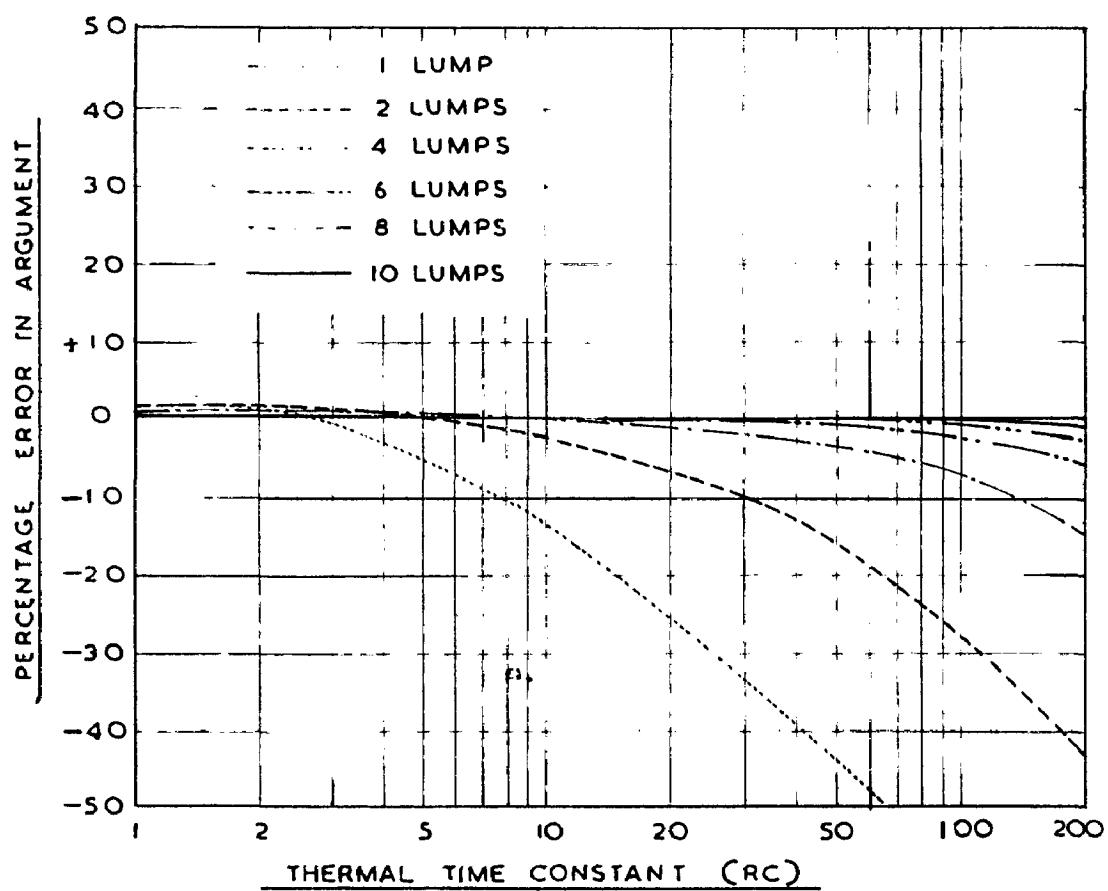
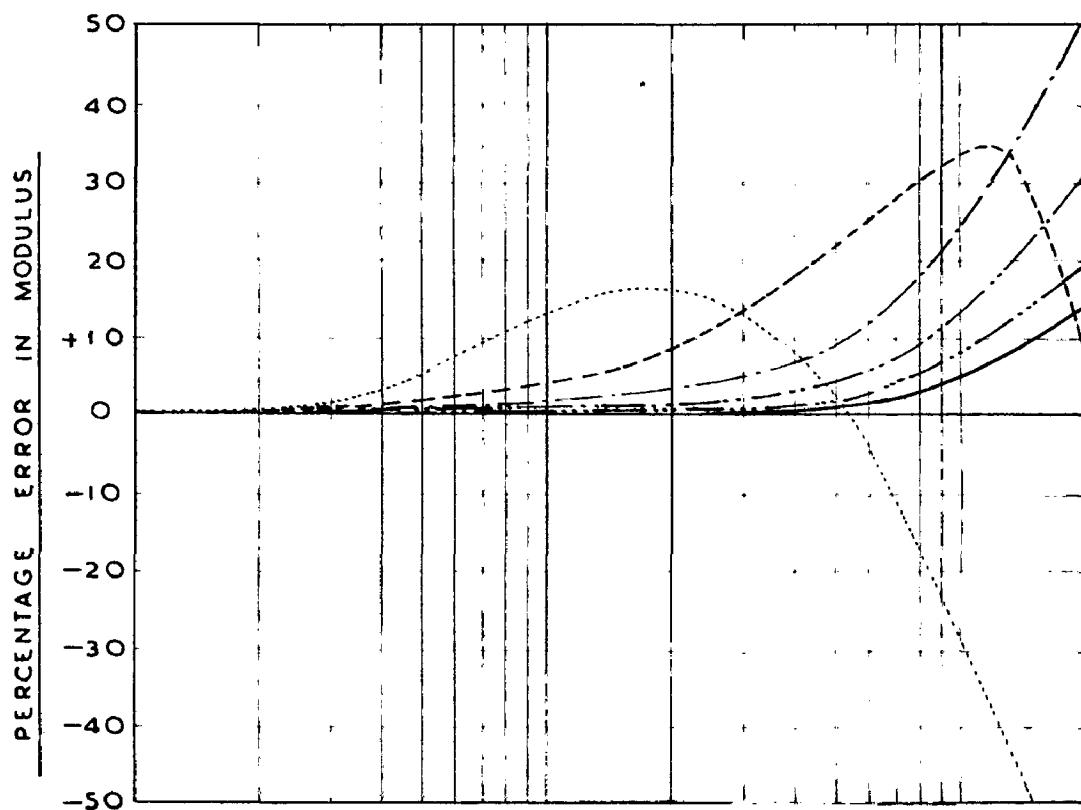


FIG.4.3. LUMPING ERRORS OF TRANSMISSION MATRIX ELEMENT - A  
(T-NETWORK)

This indicates that the number of lumps are proportional to the square root of  $\cdot RC$ . This provides a basis for the direct selection of the number of lumps required to represent any building element of known thermal time constant ( $RC$ ), ensuring an accuracy of 2 to 3 % in the transfer functions. The above criterion of lumping was employed for further studies by the Analogue method. The transfer functions were determined by the Analogue method, for building elements having different time constants, for the fundamental and three higher harmonics. These are compared with the corresponding distributed system transfer functions (computed) in Fig. (4.6). Since experimental errors were of the order of 2 to 3 %, any attempt to obtain better precision by increasing the number of lumps than those recommended here is not justifiable.

#### 4.4 Effect of Film Resistances on Lumping Errors

It has been pointed out (49, 55) that the boundary conditions influence the lumping errors but quantitative variation of these errors with surface film resistances have not so far been fully evaluated for sinusoidal inputs.

Inclusion of the surface film resistances on either side of the building element will alter the overall transmission matrix. This transfer matrix is



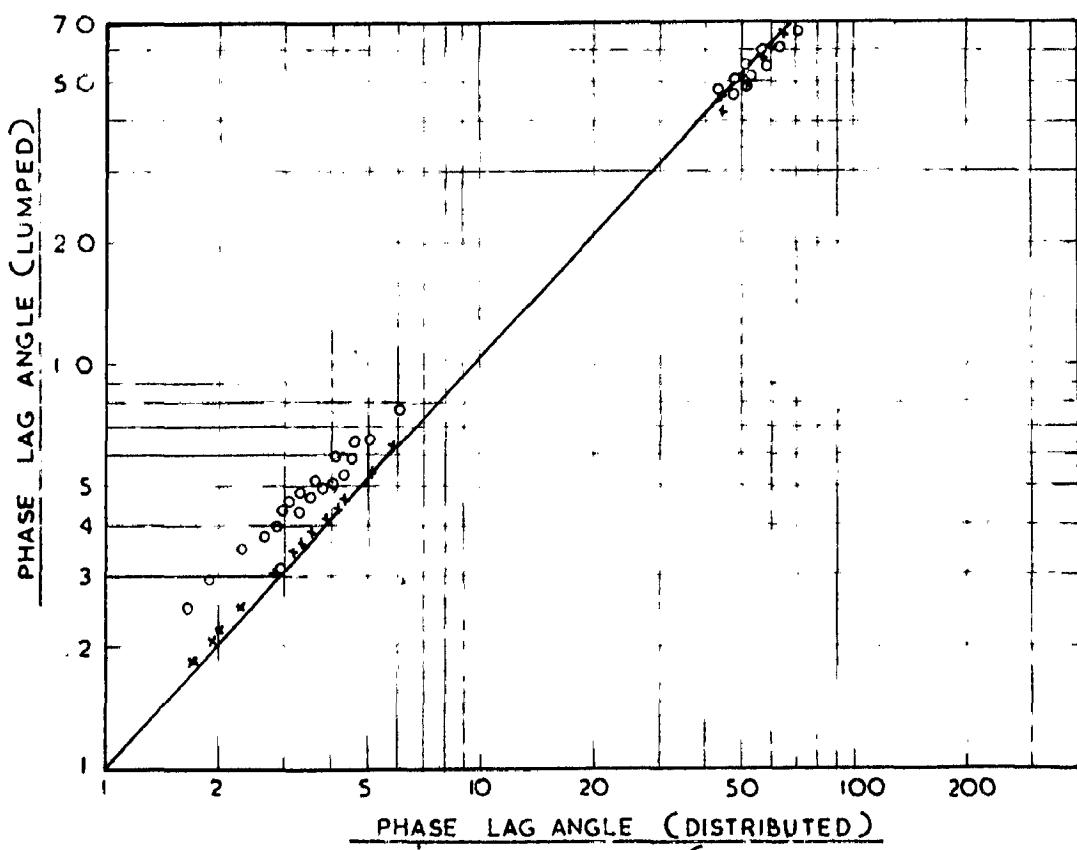
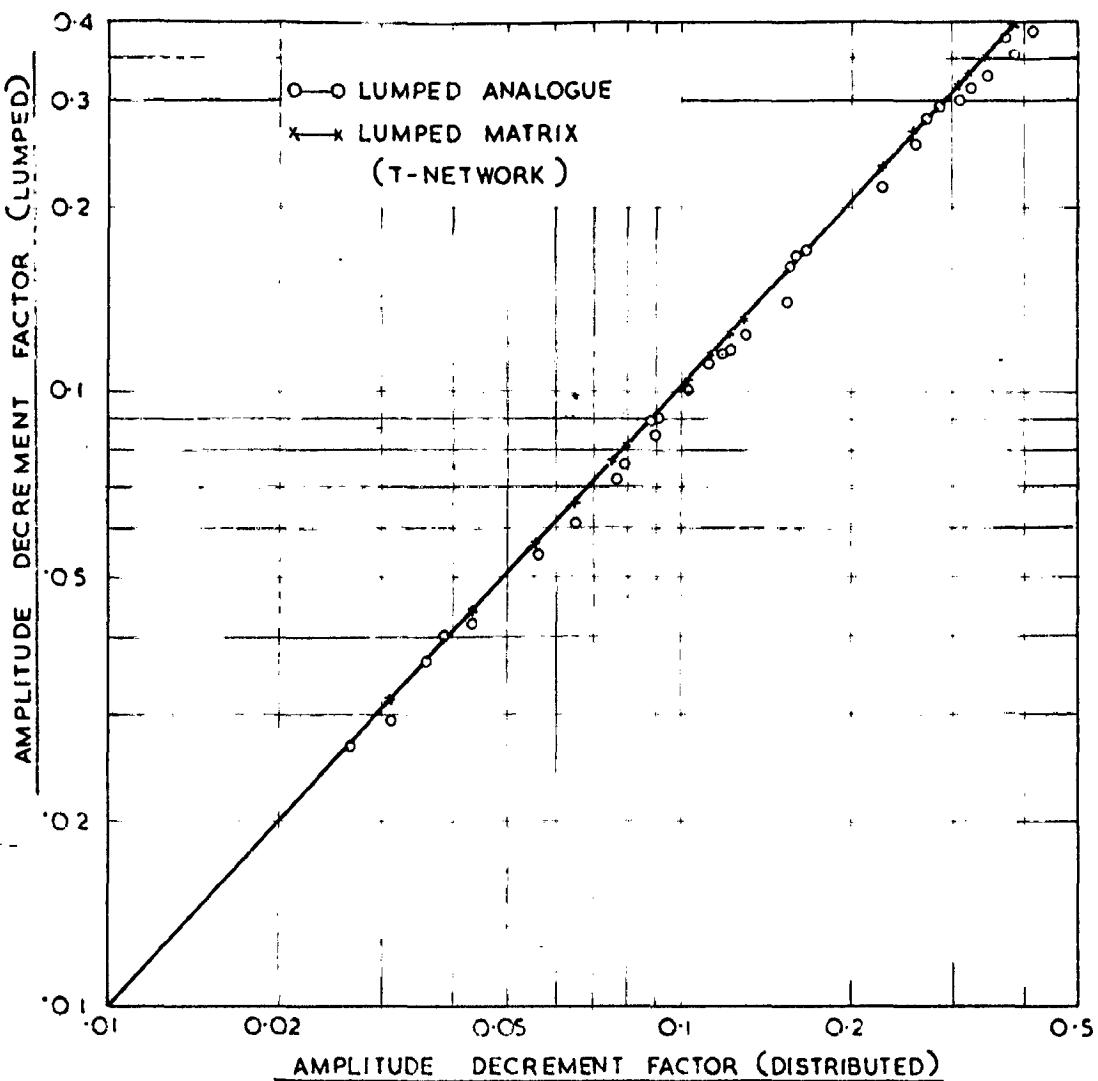


FIG.4.6 DISTRIBUTED VERSUS LUMPED (THEORETICAL & ANALOGUE)  
TRANSFER FUNCTIONS

given by

$$\begin{bmatrix} A' & B' \\ C' & D' \end{bmatrix} = \begin{bmatrix} 1 & 1/h_0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} 1 & 1/h_1 \\ 0 & 1 \end{bmatrix} \quad \dots(2)$$

For such a case the overall transfer function, ( $\lambda_L \angle \Phi_L$ ) is given by  $1/\Lambda'$

where

$$\Lambda' = A + B h_1 + C/h_0 + D h_1/h_0 \quad \dots(3)$$

The elements  $A$  and  $D$  of the matrix of the building component, depend only on the thermal time constant ( $\tau_C$ ) for a given frequency, while elements  $B$  and  $C$  will depend both on the thermal time constant and the thermal resistance ( $R$ ). Hence, the overall transfer function will not only depend on ' $R_C$ ' but also on the ratios of  $R/R_1$ ,  $R_2/R$  and  $R_0/R_1$  where  $R_0$  and  $R_1$  are the reciprocals of the surface heat transfer coefficients  $h_0$  and  $h_1$  respectively. This implies that building components having the same thermal time constant ( $\tau_C$ ) but different thermal resistances will give different transfer functions. Hence for evaluating the effect of film resistances or lumping errors, four specific values of  $R_C$  with fixed resistances, each representing a class of materials commonly used in practice, were

chosen. These are :-

- i)  $\text{RC} = 0.1$  and  $R = 0.1$  which represents highly conducting sections like Asbestos cement and Galvanised Iron sheets, glass panes etc.
- ii)  $\text{RC} = 1.0$   $R = 5.0$  which represents high resistive low capacity insulating materials like, slag wool, Thermocole etc.
- iii)  $\text{RC} = 10$ ;  $R = 1.0$  - Medium resistance and medium capacity type representing a range of light weight concretes.
- and iv)  $\text{RC} = 50$ ,  $R = 2.0$  - High capacity and high density type of heavy materials like Denso concrete, Brick, stone etc.

The number of lumps required for each case was fixed according to Table (4.2) and were 1, 2, 4 and 8 respectively for these four cases considered. These RC and R values chosen would provide a fairly representative idea of the lumping errors as influenced by the surface resistances for most of the commonly used building materials in their usual thicknesses.

For each case the transfer functions (amplitude decrement and phase lag) were computed for external surface heat transfer coefficient ( $h_o$ ) of values 1.0, 2.03, 3.5, 6.0, 8.0 and 10.0 with a fixed inside

film conductance ( $h_1$ ) of 1.6, and also for internal surface heat transfer coefficient ( $h_1$ ) of values 0.5, 1.0, 1.6, 2.0 and 5.0 with a fixed  $h_o$  of 3.5. Experimental and theoretical values of transfer functions for distributed and lumped systems, are compared in Tables (4.3 and 4.4).

The analogue and computed (matrix calculated) lumped circuit transfer functions were plotted against the distributed transfer functions, and shown in Fig. (4.6). The ideal correlation curve would be a straight line passing through the origin making  $45^\circ$  with the axes. The deviations in the amplitude decrement factor from the ideal line were only 2 to 3 percent, irrespective of the 'RC' value and the surface coefficient variation, when the above recommended number of lumps were used. The deviations of the phase lag angles from the ideal curve appear to be larger for low phase lag angles than for higher phase lags, in Fig. (4.6) but when plotted linearly this deviation is the same (within  $3^\circ$ ) all over. In all practical problems, such deviations are minute and are insignificant.

By a close examination of the equation (3) giving the overall transfer function in terms of the matrix coefficients for the building component and surface resistances, the following comments are

TABLE 4.3

Effect of Change in Heat Transfer Coefficient  $\alpha$  on Heat Transfer Coefficient  $\bar{\alpha}$  (For  $h_0 = 1.6$ )

Distributed Heat Transfer Coefficient $\bar{\alpha}$	Heat Transfer Coefficient $\alpha$											
1.	1.0	2.0	3.0	4.0	5.0	6.0	7.0	8.0	9.0	10.0	11.0	12.0
1.	1.0	2.0	3.0	4.0	5.0	6.0	7.0	8.0	9.0	10.0	11.0	12.0
2.	2.0	3.0	4.0	5.0	6.0	7.0	8.0	9.0	10.0	11.0	12.0	13.0
3.	3.0	4.0	5.0	6.0	7.0	8.0	9.0	10.0	11.0	12.0	13.0	14.0
4.	4.0	5.0	6.0	7.0	8.0	9.0	10.0	11.0	12.0	13.0	14.0	15.0
5.	5.0	6.0	7.0	8.0	9.0	10.0	11.0	12.0	13.0	14.0	15.0	16.0
6.	6.0	7.0	8.0	9.0	10.0	11.0	12.0	13.0	14.0	15.0	16.0	17.0

$J$  = Distributed,  $J_d$  = Local,  $J_c$  = Circuit Matrix,  $\Omega_d$  = Local,  $\Omega_c$  = Circuit analogue.

T-Blatt 44

OUTSIDE SURFACE COEFFICIENT (h<sub>1</sub>) OF LUMPING PLATE ON THE HOT SIDE

Inside		$\sigma = 0.1$		$\sigma = 1.0$		$\sigma = 10.0$		$\sigma = 50.0$	
Surface	coeff.	$R_1$	$T_A$	$R_1$	$T_A$	$R_1$	$T_A$	$R_1$	$T_A$
Trans. Function Trans. Function Trans. Function									
Modulus	Arg-phi	Modulus	Arg-phi	Modulus	Arg-phi	Modulus	Arg-phi	Modulus	Arg-phi
$R_1$	$\Delta_i$	$\lambda_{11}$	$\lambda_{12}$	$\lambda_{21}$	$\lambda_{22}$	$\lambda_{31}$	$\lambda_{32}$	$\lambda_{41}$	$\lambda_{42}$
1. 0.5	$\frac{2}{\pi}$	0.836	4.0	0.274	4.2	0.390	67.7	0.090	152.5
	$\frac{21}{\pi}$	0.840	4.1	0.274	4.3	0.392	67.6	0.039	152.0
	$\frac{22}{\pi}$	0.844	6.0	0.278	6.0	0.376	65.0	0.085	149.0
2. 1.0	D	0.721	3.6	0.169	3.6	0.315	60.7	0.076	146.0
	$\frac{21}{\pi}$	0.719	3.8	0.169	3.7	0.329	60.8	0.075	146.0
	$\frac{22}{\pi}$	0.722	4.0	0.168	5.0	0.303	58.3	0.071	142.0
3. 1.5	D	0.633	3.2	0.112	3.4	0.262	56.9	0.065	141.5
	$\frac{21}{\pi}$	0.633	3.5	0.112	3.4	0.253	56.3	0.064	140.6
	$\frac{22}{\pi}$	0.634	4.0	0.116	5.0	0.236	54.0	0.061	133.0
4. 2.0	D	0.564	3.9	0.033	3.2	0.224	53.9	0.056	137.9
	$\frac{21}{\pi}$	0.557	3.1	0.086	3.3	0.219	53.3	0.053	135.6
	$\frac{22}{\pi}$	0.567	4.0	0.039	4.0	0.216	51.0	0.034	133.0
5. 5.0	D	0.341	2.0	0.036	3.0	0.117	44.0	0.031	123.2
	$\frac{21}{\pi}$	0.361	3.2	0.036	3.0	0.114	44.0	0.030	123.0
	$\frac{22}{\pi}$	0.333	3.0	0.035	4.0	0.116	42.0	0.023	124.0

• Report of the Secretary of State [1873]

made :-

- i) The matrix element 'A' is not affected by the inclusion of surface resistances, whereas the coefficient 'B' is affected by inside surface film resistance by the ratio of  $R/R_1$ . Coefficient 'C' is affected by the external surface film resistance by the ratio  $R_o/R$  while the coefficient 'D' is affected by both the surface resistances by the ratio  $R_o/R_1$ .
- ii) If  $A_N$ ,  $B_N$ ,  $C_N$  and  $D_N$  be the elements of the transmission matrix of an 'N' lumped circuit, the errors due to lumping in individual matrix elements will be  $(A - A_N)$ ;  $(B - B_N) R/R_1$ ;  $(C - C_N) R_o/R$  and  $(D - D_N) R_o/R_1$  respectively. The total lumping error will then be cumulative.
- iii) For a given RC and surface resistances  $R_o$  and  $R_1$ , the error  $(B - B_N) R/R_1$  will increase with the increase of  $R$  while the error  $(C - C_N) R_o/R$  will decrease.
- iv) For a given RC and  $R$ , increase in  $R_o$  will increase the error in  $(C - C_N) R_o/R$  and  $(D - D_N) R_o/R_1$  value.

While an increase in  $R_1$  will decrease  $(S - D_{ij}) \frac{R}{R_1}$  and  $(S - D_{ij}) R_0/R_1$  the error in  $(A - A_{ij})$  will be unaltered. For a given number of lumps the error in all the four elements will increase with increasing RC and frequency.

As seen from the above, for any particular condition, the overall error in a lumped circuit will depend on the relative increase or decrease of the errors due to  $(S - D_{ij}) \frac{R}{R_1}$  and  $(S - S_{ij}) R_0/R_1$  and  $(S - D_{ij}) R_0/R_1$  as these errors oppose each other. Under natural conditions the external surface resistance will be usually less than that of the internal surface resistance. Hence, the net effect of surface resistances is to reduce the errors due to lumping. These conclusions can also be verified from the data given in Tables (4.3 and 4.4).

#### 4.5 Affect of Frequency on Lumping Errors

For a given thermal time constant (RC) and fixed number of lumps used, the error between the lumped and distributed transfer functions will increase with the increase of frequency. With the increase of frequency, the transfer constant  $\Theta$  which is given by  $\sqrt{\omega_{CR}}$  will also increase. For higher values of  $\Theta$  larger number of lumps are to be employed.

The effect of higher harmonics can be treated as if the 'RC' value is increased in harmonic steps and the frequency itself is unaltered. The advantage of this approach is that the curves already drawn between RC and transfer functions and errors for the fundamental frequency, may be utilised for the higher harmonics as well, to obtain the above mentioned quantities.

It follows that if the number of lumps to be used are selected on the basis of the fundamental frequency, increased errors are likely to occur for higher harmonic inputs. However, for practical case of a periodic soil-air temperature excitation, the overall errors will not be increased to any significant extent, as the higher harmonic components in the soil-air temperature will usually be of lower amplitudes than the fundamental and the attenuation in the amplitude for higher harmonics will be quite high.

#### 4.6 Composite Constructions

Evaluation of lumping error in composite constructions, theoretically is highly involved. In order to estimate the overall errors of the analogue method for composite constructions, the Matrix calculated distributed system transfer functions were

compared with Analogue results for some typical composite wall panels in Table (4.6). In setting up the Analogue model, each homogeneous section of the composite panel was represented by the number of lumps as recommended in Table (4.8). The comparisons indicate, that such a method of representation is permissible to obtain reasonably accurate results.

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COMPILATION OF THE STRENGTH AND STABILITY TABLES FOR COHERENT POINT DRIVING POLE

$h_0 = 3.5$   
 $h_1 = 1.5$

No.	Construction	External Driving point function	Transfer function	Internal driving point function
1.	$3'' p + 13''$ brick + $5''$ d	0.719 0.734	14.0 0.065 0.262	155.0 0.493 160.0
2.	$5'' c + 4''$ brick + $8''$ air space + $4_{1/2}''$ brick + $5''$ d	0.726 0.744	17.0 0.038 0.233	128.0 0.432 123.0
3.	$3''$ marble + $9''$ brick + $1_{1/2}''$ c	0.633 0.596	36.0 0.082 0.263	143.0 0.434 140.0
4.	$1'' c + 2''$ faced conc. + $4''$ dense conc. + $5''$ d	0.912 0.912	12.0 0.036 0.236	90.0 0.563 91.0
5.	$3'' c + 4_{1/2}''$ brick + $1''$ thermocouple + $1''$ d	0.719 0.734	26.0 0.052 0.255	85.0 0.800 87.0
6.	$3''$ lime conc. + $6''$ l.B. + $1_{1/2}''$ d	0.730 0.738	11.0 0.130 0.133	90.0 0.448 92.0
7.	$3''$ lime conc. + $4''$ mud phuska + $4_{1/2}''$ l.B. + $5''$ d	0.723 0.733	13.0 0.054 0.053	140.0 0.476 144.0
8.	$5''$ harbourd + $1''$ thor- mecole + $5''$ harbourd	0.900 0.913	4.0 0.034 0.080	23.0 0.370 30.0
				$0.875$

D = plaster, L.B. = "deforce" article

## **C H A P T E R   5**

**EXPERIMENTAL CHECK OF THE ANALOGUE BEHAVIOUR**

## CHAPTER 5

### EXPERIMENTAL CHECK OF THE ANALOGUE BEHAVIOUR

#### 5.1. Introduction

Before taking up a study of new problems on a large scale by the analogue method, it is necessary to check the analogue performance with experimental results, under different boundary conditions. It was therefore proposed to verify the R-C network analogue behaviour by comparing the analogue predicted diurnal temperature variations of incide surfaces with those actually measured ones on full scale experimental houses. For this purpose four cases, for which thermal data and the boundary conditions were available, have been chosen. These are given in Table (5.1).

TABLE (5.1)

Sl. No.	Building Element	Construction	External Boundary Condition	Internal Boundary Condition
1.	Homogeneous Roof slab	4.26 inch concrete	Surface input	Incide air temperature constant.
2.	Homogeneous Wall	9.25 inch Brick	"	"
3.	Homogeneous wall	"	Sol-air input	"
4.	Composite wall	1" plaster + 1" Thermocolo + 4" Brick + 3" plaster	Surface input	Indoor air temperature

## 6.2 Concrete Roof Slab with Surface Input

Raychoudhuri (65) has studied the heat flow variations through a concrete slab 4.26 inch thick which formed the roof of a wall insulated enclosure. The external surface was exposed to natural weather conditions, while the inside air temperature of the enclosure maintained nearly constant. The surface temperatures were measured continuously over 30 hours with thermocouples.

The physical properties of the concrete slab are given in Table (5.2).

Thickness                    (L) = 4.26 inches

Thermal resistance of the slab  $\lambda$  = 0.30

Thermal capacity of the slab     $C = 10.3$

This concrete slab (with an 'R.C' value of 4.0) can be adequately represented on the analogue by a lumped 'T' circuit network.

As the external surface temperature forms the external boundary condition, the outside surface coefficient ( $h_0$ ) need not be included in the analogue circuit. The inside surface coefficient ( $h_1$ ) is assumed to be constant, of value  $1.5 \text{ Btu/Pt}^2/\text{Fr}/{}^\circ\text{F}$ . The concrete slab was simulated on the analogue with the given boundary conditions and the transfer

TABLE 5.2

PHYSICAL PROPERTIES OF THE MATERIALS USED  
IN EXPERIMENTAL VALUE

No.	Material	K	$\rho$	s
1.	Concrete	10.8	132	0.22
2.	Brick	6.4	117	0.20
3.	Plaster	12.0	120	0.22
4.	Thermocole	0.2	1	0.32
5.	Brick	6.0	100	0.21

K = Thermal conductivity in Btu.in/Ft<sup>2</sup>.hr.<sup>0</sup>F

$\rho$  = Density in Lb/Ft<sup>3</sup>

s = Specific heat in Btu/Lb.<sup>0</sup>F

functions  $N_i \phi_i(t_{is}/t_{os})$   $t_{ia}$  constant for the fundamental and two higher harmonics were determined. These are given in Table (5.3).

The measured external surface temperatures were harmonically analysed and its Fourier representation is given by equation (1)

$$\begin{aligned} t_{os}(T) = & 65.30 + 18.67 \cos(15t - 54^\circ 43') \\ & + 6.7 \cos(30t - 54^\circ 54') \quad ..(1) \\ & + 1.63 \cos(45t - 47^\circ 30') \end{aligned}$$

Fundamental and two harmonic terms are found to be sufficient for a fairly accurate representation of the actual wave form.

The corresponding Fourier representation of the inside surface temperature variation, can be obtained with a knowledge of the 'U' value and transfer functions of the slab, by using Houghton et.al. (63) method. The Fourier equation of the inside surface temperature is obtained as

$$\begin{aligned} t_{is}(T) = & 65.30 + 13.6 \cos(15t - 82^\circ 43') \\ & + 3.3 \cos(30t - 106^\circ 54') \quad ..(2) \\ & + 0.6 \cos(45t - 121^\circ 30') \end{aligned}$$

The diurnal variation of inside surface temperatures were obtained by synthesising the equation (2) and are compared with the measured temperatures in

TABLE 5.3

OPTIMUM TRANSFER AND TRANSFER POINTS FOR THE HILF MELT TESTS

	Construction : $h_0$	$h_1$	Function : $M$	Fundamental	Second	Third	Fourth	Fifth	Sixth	Seventh	Eighth	Ninth	Tenth
1.	4.56" concrete	--	1.6 Transfer	0.726	26	0.482	49	0.368	70	...	...	...	...
	clab			3.734	23	0.483	52	0.386	74	...	...	...	...
2.	9.25" brick	..	1.4 -lo-	0.197	87	0.124	135	0.081	171	0.033	201	0.037	220
	wall			0.190	90	0.099	140	0.053	176	0.037	201	0.041	229
3.	9.25" brick	1.6	1.7 -lo-	0.091	126	0.039	150	0.021	193	0.012	229	0.013	250
	wall			0.039	103	0.039	160	0.020	198	0.013	229	0.014	254
4.	1" p + 1" thon-	..	1.5 Transfer	0.046	87	0.026	103	0.015	151	0.011	198	0.012	230
	p + " p + " p			0.043	90	0.024	117	0.014	154	0.013	198	0.014	234
	Int. crit.			0.531	21	0.390	34	0.339	35	0.290	36	0.296	39
	point			0.570	24	0.378	36	0.332	33	0.296	36	0.296	39

$P = \text{Plaster}.$        $M = \text{Matrix}$        $A = \text{Analogue}$

Fig. (5.1). From Fig. (5.1) it can be seen, that the agreement, between the measured and analogue predicted temperature, is good (within 3 per cent) throughout the daily cycle except at the initial hours.

### 5.3 Brick Wall with Surface Input

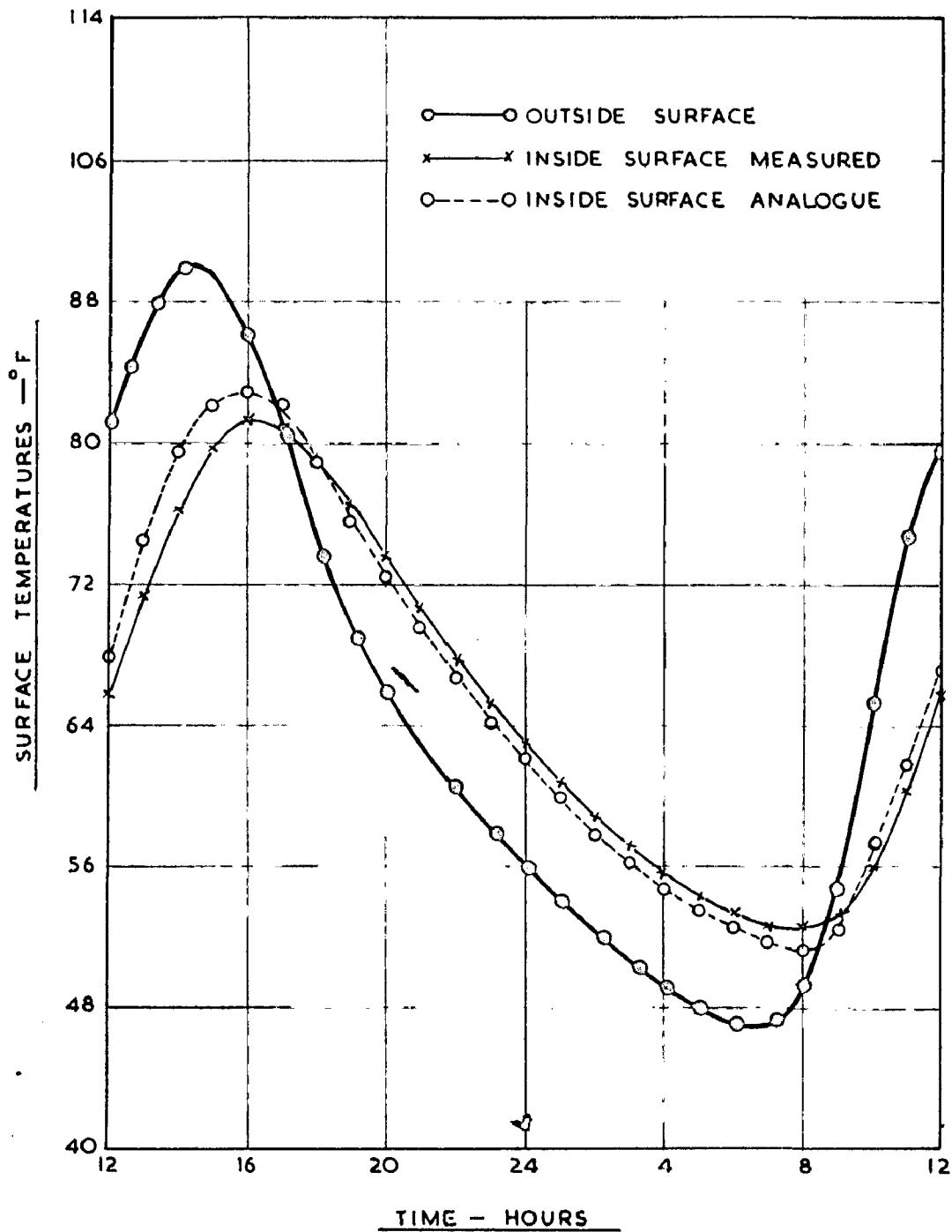
Roux (67) had investigated, how far the use of the analytical solutions developed, by Houghton et.al. (66) and Mackoy and Wright (29) for non-steady state heat flow through materials with all their simplifying assumptions are justified for practical building problems. In his investigations a room with all the four sides made up of walls of brick whose physical properties, are given in Table (5.2).

Thickness (L) = 0.25 inches

Thermal resistance of the wall  $A = 1.44$

Thermal capacity of the wall  $C = 18.0$

The inside air temperature was kept constant by an air-conditioning system and outside and inside surface temperatures were measured. The inside surface coefficient was taken  $1.4 \text{ Stu}/\text{sq.ft}/^{\circ}\text{F}$  and assumed constant. This brick wall of 0.25 inches thickness, having an 'RC' value of 35.9 can be represented on the analogue by 8 lumps of 'T' network with sufficient accuracy. The transfer functions (amplitude decrement



COMPARISON OF INSIDE SURFACE TEMPERATURES FOR A  
4.26" THICK CONCRETE SLAB

FIG. 5·1

factor and phase lag) of the brick wall were determined by the analogue method, for the fundamental and three higher harmonics. These are given in Table (5.3) alongwith the computed values by the analytical methods.

The Fourier analysis of the outside surface temperatures for a north wall on a summer test day is given by

$$t_{os} (T) = 70.73 + 15.03 \sin (15 t - 133^{\circ}19') \\ + 4.63 \sin (30 t - 303^{\circ}53') \\ + 1.47 \sin (45 t - 123^{\circ}57') \\ + 0.14 \sin (60 t - 61^{\circ}41') \quad ..(3)$$

Steady state component of the inside surface temperature is obtained from the equation

$$t_{is} (\text{mean}) = t_{ia} + U/h_i \left\{ t_{os} (\text{mean}) - t_{ia} \right\} \quad ..(4)$$

The harmonic components of the inside surface temperature are obtained with the help of transfer functions previously determined.

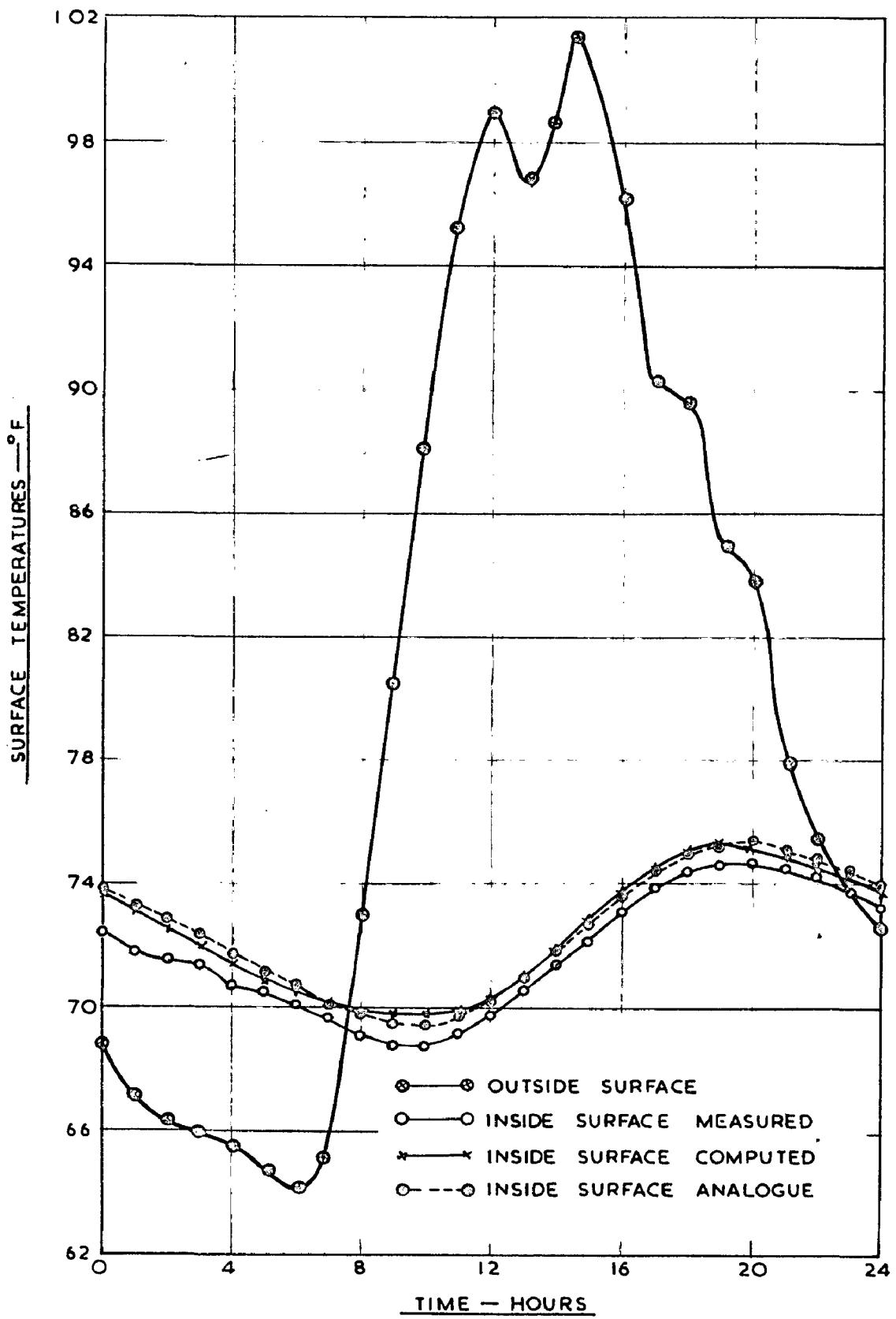
The Fourier equation for the inside surface temperatures corresponding to the input (equation - 3) is then obtained as

$$t_{is} (T) = 72.32 + 2.63 \sin (15 t - 233^{\circ}) \\ + 36.00 \sin (30 t - 92^{\circ}) \\ + 0.063 \sin (45 t - 309^{\circ}) \\ + 0.004 \sin (60 t - 276^{\circ}) \quad ..(5)$$

The hourly variation of inside surface temperatures obtained by synthesising the equation (3) are compared with the measured and the theoretically computed temperatures in Fig. (5.2). It can be seen from the Fig. (5.2) that the agreement between the analogue and theoretically computed results is excellent throughout daily cycle, while the deviations between the analogue and measured values, are more in the initial hours.

#### 5.4 Brick wall with sol-air temperature Input

The above two comparisons are for a particular external boundary condition, viz., outside surface temperature as the input. For this purpose, the outside surface temperature variations are to be known. These are not readily available. In order to make theoretical and Analogue methods, practicable the external boundary conditions should be computable from the available weather data. Mackey and Wright (1) have introduced sol-air temperature concept, which can be computed from the known weather data. They have also provided theoretical solutions taking sol-air temperature as the external boundary condition. Mour (63) had made use of the same room (9.25 in. brick walls) for the verification of the theoretical solutions, with this sol-air temperature input. In



COMPARISON OF INSIDE SURFACE TEMPERATURES FOR A  
9.25" THICK BRICK WALL

FIG.5·2

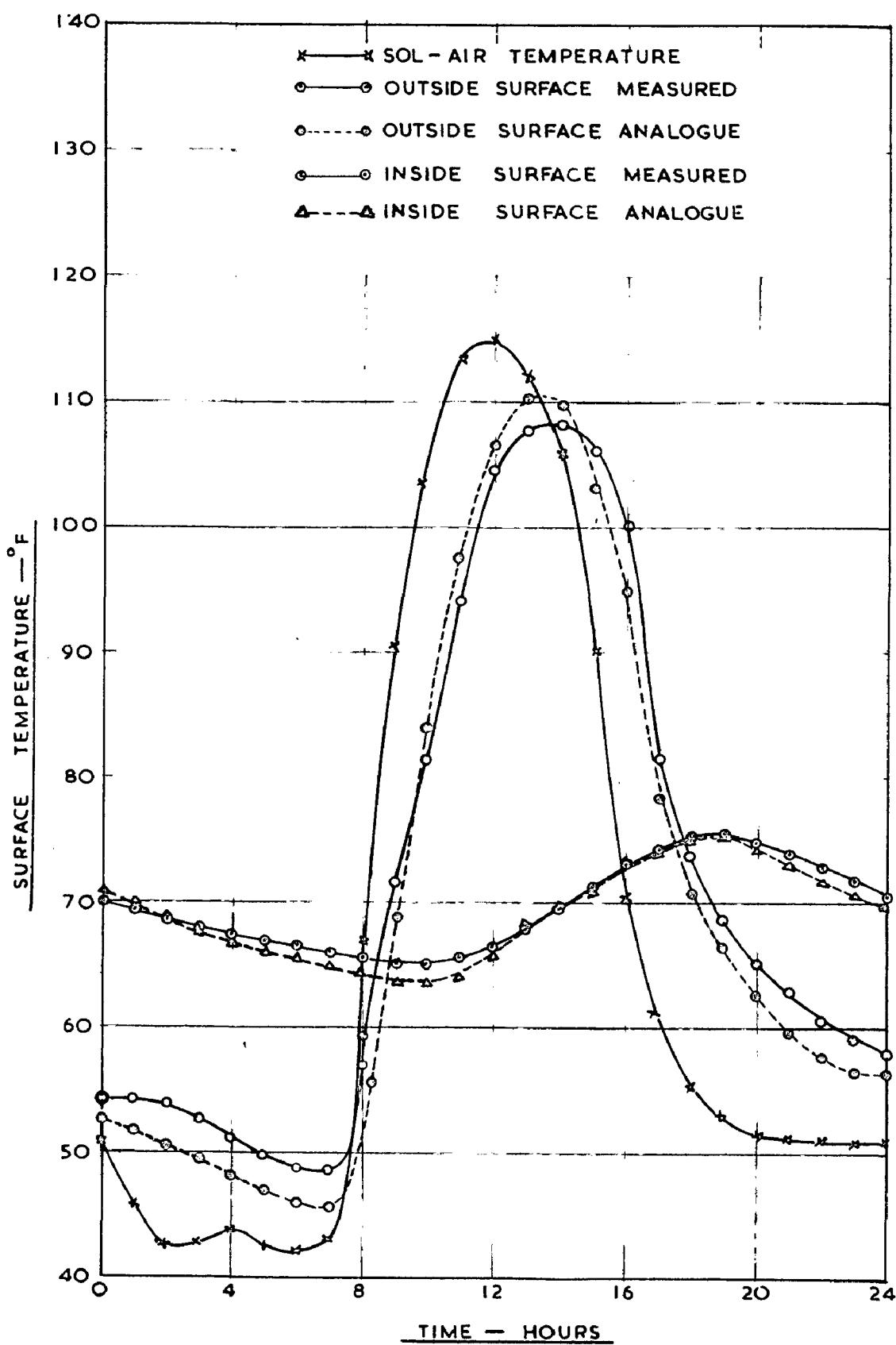
In this case the outside surface coefficient ( $h_o$ ) is to be included in the analogue circuit. The outside and inside surface coefficients were taken as 1.6 and 1.7  $\text{Btu}/\text{sq.ft.}/^{\circ}\text{F}$  respectively. The transfer functions  $N_i/\Phi_i(t_{is} / t_{sa})$  for this input boundary conditions will be different from the previous ones. These transfer functions determined by the analogue along with those theoretically computed are also given in Table (5.3). The Fourier equation of the sol-air temperature for the experimental conditions is given by

$$t_{sa} (T) = 68.27 + 61.48 \cos (16 t - 183^{\circ}31') \\ + 35.03 \cos (30 t - 1^{\circ}3') \\ + 5.75 \cos (46 t - 165^{\circ}80') \\ + 4.67 \cos (60 t - 199^{\circ}25') \quad ..(6)$$

The corresponding inside surface temperatures obtained, from the known 'U' value and the transfer functions, is given by equation (7)

$$t_{is} (T) = 68.70 + 4.9 \cos (16 t - 293^{\circ}) \\ + 1.12 \cos (30 t - 159^{\circ}) \\ + 0.098 \cos (46 t - 7^{\circ}) \\ + 0.047 \cos (60 t - 69^{\circ}) \quad ..(7)$$

The time temperature variations obtained by the analogue method and experimentally measured ones are compared in Fig. (5.3).



COMPARISON OF OUTSIDE AND INSIDE SURFACE TEMPERATURES FOR A 9.25" THICK BRICK WALL

WITH SOL-AIR INPUT

FIG. 5.3

It is also possible to determine analogically the outside surface temperatures of a building surface for a given sol-air temperature and outside surface coefficient ( $h_o$ ). For this purpose, external driving point transfer functions  $t_{os}/t_{sa}$  i.e.,  $\lambda_o/\phi_o$  are to be determined for the fundamental and higher harmonics. The mean outside surface temperature is obtained by the equation

$$t_{os} \text{ (mean)} = t_{sa} \text{ (mean)} - U/h_o \left\{ t_{sa} \text{ (mean)} - t_{in} \right\} \quad ..(8)$$

The external driving point transfer functions determined by the analogue are given in Table (5.5). The Fourier representation of the outside surface temperatures corresponding to the given sol-air temperature (equation - 6) for the brick wall is obtained as

$$\begin{aligned} t_{os} (T) = & 67.10 + 27.6 \cos (15 t - 207^\circ) \\ & + 13.3 \cos (30 t - 31^\circ) \\ & + 1.9 \cos (45 t - 197^\circ) \\ & + 1.46 \cos (60 t - 231^\circ) \end{aligned} \quad ..(9)$$

The comparison of the analogically determined and experimentally measured outside surface temperatures are also made in the above Fig. (5.3).

The deviations between the measured and analogue temperatures for the outside surface are greater than those for the inside surface. This may

be attributed to two factors, namely (i) as ' $h_o$ ' is directly dependent on the outside weather conditions, is liable to larger variations with time, than ' $h_i$ ', and (ii) the soil-air temperatures as suggested by Mackey and Wright need correction for low temperature radiation exchange taking place between the external surface and its surroundings. The computational methods employed (1) for low temperature radiation estimation are empirical and need improvement.

Inspite of those drawbacks the overall agreement between the measured and analogue predicted outside surface temperatures may be considered as quite satisfactory, in dealing with problems of day to day engineering practice.

#### 5.5 Composite Wall with External Surface Input and Indoor Air Temperature Variable

The above three instances are for homogeneous constructions. In practice many building elements are made of composite sections. The internal boundary condition, of constant indoor air temperature, is true for air conditioned enclosures only. But the main bulk of buildings are unconditioned and the indoor air temperatures fluctuate periodically. Hence a more general case of a composite construction with

variable indoor air temperature is also included for the verification of the Analogue.

Raychoudhri et al. (69) have reported experimental studies on the effect of walls and roofs (composite) on the indoor air temperatures for unconditioned rooms. The wall chosen for the present study has the following constructional details :-

1" plaster, 1" thermocole, 4.5" brick, ½" plaster.

The order of the layers being from outside to inside. The physical properties of the materials of each layer are given in Table (5.2). The equations of the external surface temperatures and the indoor air temperatures which form the boundary conditions, are given as

i) outside surface temperatures

$$t_{os}(T) = 105.6 + 21.33 \cos(15t - 52^{\circ}31') \\ + 8.41 \cos(30t - 100^{\circ}42') \quad \dots(10) \\ + 4.10 \cos(15t - 185^{\circ}24') \\ + 2.12 \cos(60t - 247^{\circ}24')$$

ii) indoor air temperatures

$$t_{ia}(T) = 23.08 + 3.67 \cos(15t - 125^{\circ}) \\ + 0.43 \cos(30t - 169^{\circ}12') \quad \dots(11) \\ + 0.003 \cos(45t - 247^{\circ}) \\ + 0.14 \cos(60t - 253^{\circ})$$

In representing the composite wall on the Analogue, each homogeneous layer, is represented by the number

of lumps (based on their NC values) as per earlier recommendations. Inside surface coefficient is taken as 1.5 Btu/Bq.ft./hr/ $^{\circ}$ F and assumed to be constant.

For a boundary condition where air temperatures on either side vary, the problem can be treated as consisting of two superposed parts, namely (i) outside air or surface temperature is variable, with a constant indoor air temperature, and (ii) indoor air temperature is variable with constant outside temperature. The actual inside surface temperature variation is then obtained by the super position of the solutions of the two parts. Hence not only the usual transfer functions  $\lambda_i / -\phi_i$  i.e.,  $(t_{is}/t_{os})$  but also the internal driving point transfer function  $\lambda'_i / -\phi'_i$  i.e.,  $(t_{is}/t_{ia})$  is required. These two sets of transfer functions are determined on the analogue and given in Table (5.3). The equation for the inside surface temperature in terms of the boundary variations and transfer functions can be expressed as

$$t_{is}(T) = t_{is}(\text{mean}) + \sum_{n=0}^{\infty} \lambda_{int} t_{os} n \cos(\omega_n t - \gamma_n - \phi_{in}) \\ + \lambda'_{in} t_{ia} n \cos(\omega_n t - \xi_n - \phi'_{in}) \quad \dots(12)$$

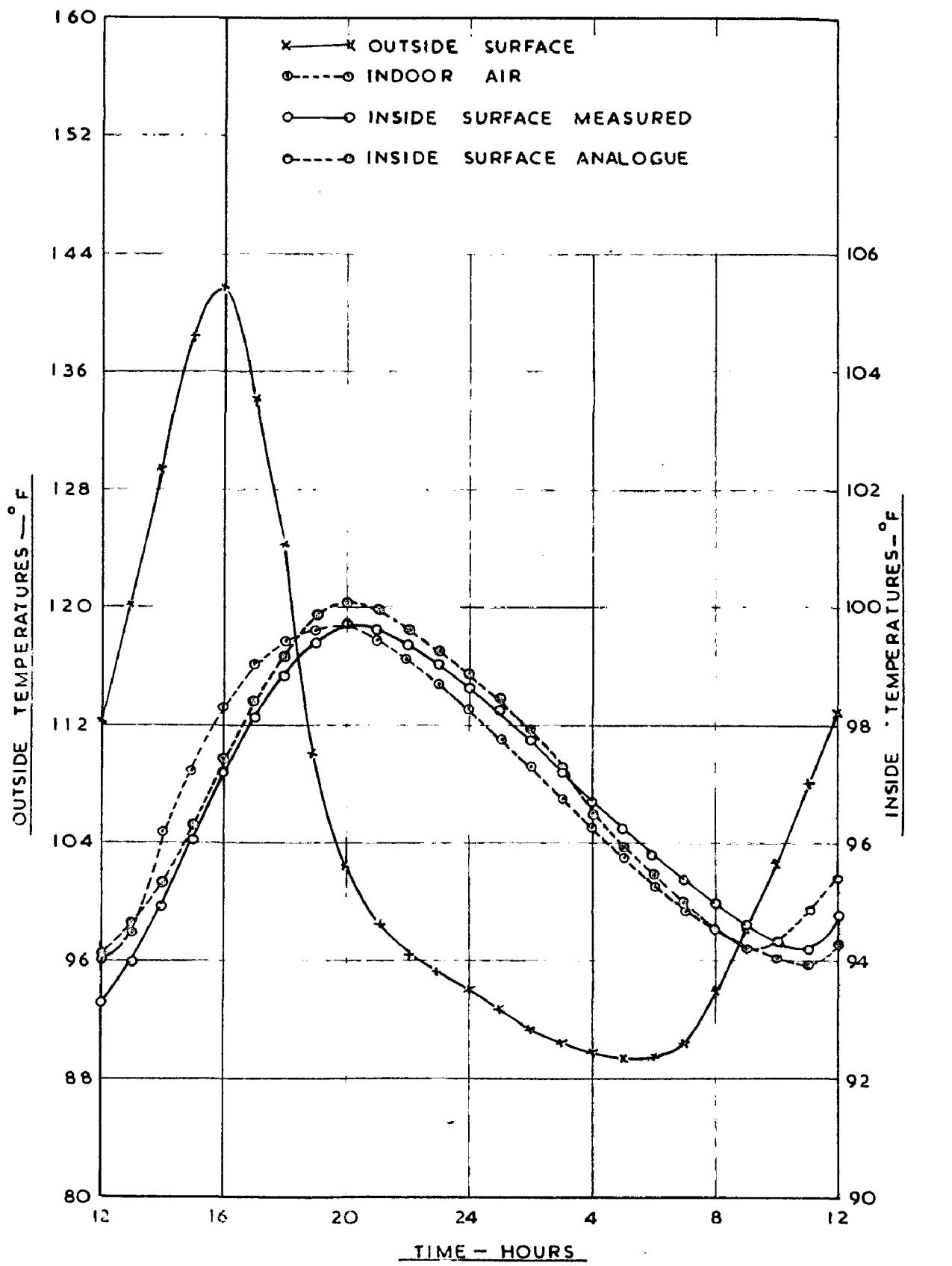
where  $t_{is}(\text{mean}) = t_{ia}(\text{mean}) + U/h_i \{t_{os}(\text{mean}) - t_{ia}(\text{mean})\}$

Substituting the numerical values for this present example, the inside surface temperature equation is

$$\begin{aligned}
 t_{18} (\tau) = 97.80 &+ 0.79 \cos(15\tau - 123^\circ) \\
 &+ 0.23 \cos(30\tau - 209^\circ) \\
 &+ 0.075 \cos(45\tau - 319^\circ) \\
 &+ 0.025 \cos(60\tau - 58^\circ) \\
 &+ 2.18 \cos(15\tau - 143^\circ) \\
 &+ 31 \cos(30\tau - 203^\circ) \\
 &+ 0.09 \cos(45\tau - 302^\circ) \\
 &+ 0.056 \cos(60\tau - 209^\circ)
 \end{aligned} \quad \dots (13)$$

The hourly variations of the inside surface temperatures as obtained by synthesising the equation (13) and compared with the measured values in Fig. (5.4). It can be seen from Fig. (5.4) that the analogue predictions agree closely with the measured temperatures as in the previous cases within 3 to 4 per cent. In the above computations it is expected that the deviations between the analogue predicted and theoretically computed values to be minimum as all the simplifying assumptions made in theoretical computations, apply to the analogue also. The deviations between the measured and the theoretical values may be attributed to the following assumptions made in the theoretical solution :-

- i) The surface coefficients are taken as of a constant value, though in practice they vary to a certain extent.
- ii) The radiation exchange at the inside surface has not been separately treated, but the inside surface coefficient is taken as the combined radiation and convection coefficient.



COMPARISION OF INSIDE SURFACE TEMPERATURES FOR A  
COMPOSITE WALL WITH VARIABLE INDOOR AIR  
TEMPERATURES

FIG. 5·4

In other words all the heat transfer from the inside surface to the indoor air is taking place directly through the film conductance, duly corrected for radiation transfer.

- iii) The soil-air or outside surface temperature wave forms, are identical on two successive days. This assumption is not strictly true. Non-conformity of this assumption is expected to introduce increased deviations in the first few hours corresponding to the time lag period.

However, all the above simplifying assumptions made in the theory and the analogue errors dealt in Chapter (4), combinedly do not exceed 3 to 4 per cent for all practical problems. The predictions of periodic temperatures within 5 per cent of the actual conditions can be considered as more than sufficient for engineering accuracies.

The above comparisons, which cover the commonly occurring boundary conditions and types of constructions, amply justify the validity of adopting the analogue method in the proposed form, for studying the periodic heat flow through building elements of different constructions and under different climatic conditions.

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## **CHAPTER 6**

**TERMINAL ELEMENT (CARRIER AND DRIVING POINT)  
FUNCTIONS OF HOMOGENEOUS BUILDING ELEMENTS**

## CHAPTER 6

### TERMAL SYSTEM (TRANSFER AND DRIVING POINT) FUNCTIONS OF HOMOGENEOUS BUILDING ELEMENTS

#### 6.2 Introduction

Comprehensive data on the thermal characteristics which completely define the thermal behaviour of building elements, under periodic heat flow and natural boundary conditions, are not as yet available. It has been shown (70) that a set of three system functions, i.e., one transfer and two driving point functions will completely characterise a four terminal network. Van Gorcum (23) and later Marmot (27) Kuncay (26) Pipoa (26) have shown that a homogeneous slab can be treated as a passive four terminal (two pair) two element (A and C) network and the same mathematical methods employed for solving such electrical network problems can also be applied to solve the problems of periodic heat flow through building elements. The network system functions, developed for electrical problems can profitably be used for the analogous thermal system functions. The computation of these functions for sinusoidal excitations, employing the well established mathematical methods used in electrical problems (Laplace transforms, symbolic Calculus,

Matrix) are hardly any simpler than analytical methods used for thermal problems. The most promising approach lies in the possibility of determining those system functions experimentally by the electrical analogue method which obviates complex calculations.

### 6.2 Thermal System Functions

The system functions depend not only on the network elements but also on its nature of termination. For a building element the surface heat transfer coefficients which are represented by pure resistances, form the termination on both ends. The boundary conditions, that are encountered in periodic heat flow problems of buildings and the corresponding system functions are derived in terms of general circuit parameters and given in Appendix (II). The three thermal functions determined by the analogue method are :-

- 1) Transfer temperature amplitude ratio

$$\lambda_i \angle \phi_i \text{ i.e. } \left( \frac{t_{ia}}{t_{sa}} \right) \quad t_{ia} = 0$$

- ii) External driving point temperature amplitude ratio

$$\lambda_o \angle \phi_o \text{ i.e. } \left( \frac{t_{oo}}{t_{sa}} \right) \quad t_{ia} = 0$$

- and iii) Internal driving point temperature amplitude ratio

$$\lambda'_i \angle \phi'_i \text{ i.e. } \left( \frac{t_{ia}}{t_{oo}} \right) \quad t_{sa} = 0$$

As these functions are vector quantities, for each function two quantities viz., the amplitude decrement and phase lag angle have to be measured. The procedure followed for the determination of those quantities have been described in Chapter 3.

Though the transfer and internal driving point admittance functions are of direct significance, in the heat flow calculations, the temperature amplitude ratios were chosen for the analogue determination, because these are easily measured. However, the admittance functions can be obtained from the following relationships between them.

i) Transfer admittance function

$$Y_1 \angle \psi_1 = \frac{q_{ia}}{t_{su}} \quad t_{ia} = 0$$

$$Y_1 \angle \psi_1 = h_1 \lambda_1 \angle \phi_1 \quad .. (1)$$

ii) Internal driving point admittance function

$$Y'_1 \angle \psi'_1 = \frac{q_{ia}}{t_{ia}} \quad t_{su} = 0$$

$$Y'_1 \angle \psi'_1 = h_1 (1 - \lambda_1 \angle \phi'_1) \quad .. (2)$$

The relation between these transfer and driving point thermal system functions and the generalised transfer matrix parameters have also been derived and presented in Appendix (II).

Once these thermal functions are determined

all other desired thermal quantities for a building element, such as the surface temperatures ( $\vec{t}_{os}$  and  $\vec{t}_{is}$ ) and the heat fluxes entering and leaving ( $\vec{q}_{os}$  and  $\vec{q}_{is}$ ) can be obtained. In practice the periodic driving temperatures are not simple sinusoidals, but can be expressed as a Fourier series. However, these can be adequately described by the fundamental and a few harmonics.

Thermal system functions of different types of materials commonly used in building practice, were determined for the fundamental and the first three harmonics on the electrical network analogue. As this data is of considerable practical utility, these along with their 'U' values are presented in Tabular form (Table 2) for reference purposes in Appendix (III). The physical properties, the characteristic thermal impedances and the propagation constants of these materials are also given therein (Table 1).

An inspection of Table (1) of Appendix (III) indicates that i) dense materials have lower characteristic impedance ( $Z_c$ ) and constants of propagation, and their variation from material to material is not significant; and ii) light weight materials have higher characteristic impedance and these vary considerably from material to material.

The following inferences are made from an analysis of the thermal system function data given in Table (2) of Appendix (III).

6.3

### $\lambda_o/\phi_o$ - External Driving Point Function

This function enables one to estimate the outside surface temperatures attained by different materials for a known soil-air temperature excitation. The magnitude of  $\lambda_o$  depends upon the coefficient of thermal absorption ' $p$ ' ( $\frac{\sqrt{\omega C}}{R}$ ) of the material. Materials with high ' $p$ ' value, will have lower  $\lambda_o$  and larger  $\phi_o$  for a given thickness and vice versa. For a given material and frequency, ' $p$ ' is independent of thickness, whereas  $\lambda_o$  and  $\phi_o$  depend on the thickness as well. With the increase of thickness both R and C will increase in the same proportion and the ratio C/R and hence ' $p$ ' is unaltered, for a given frequency.  $\lambda_o$  and  $\phi_o$  depend on R and C and also on the ratio of the material resistance ( $a$ ) to the surface resistances ( $R_s$  and  $R_t$ ). Hence  $\lambda_o$  and  $\phi_o$  depend on the thickness of the material. For higher harmonics the magnitude of  $\lambda_o$  decreases while  $\phi_o$  increases, slightly. Light weight materials having high characteristic impedance ( $\omega_c$ ) will have higher  $\lambda_o$  and lower  $\phi_o$  than for dense materials.

6.4

$\lambda_i/\phi_i$  - Transfer Function

- i) The magnitude of  $\lambda_i$  will depend not only upon 'RC' but also on the ratios of  $R/R_i$ ,  $R_o/R$  and  $R_o/R_i$ .
- ii) For a given thickness materials having large 'p' will have high phase lags and vice versa.
- iii) The decrease of  $\lambda_i$  and increase of  $\phi_i$  are more marked with the increase of thickness for dense materials than for light materials.
- iv) The effect of harmonics on  $\lambda_i$  and  $\phi_i$  is more pronounced than on  $\lambda_o$  and  $\phi_o$ . Higher harmonics have no significant effect on  $\lambda_i$  and  $\phi_i$  for very thin sections, such as glass panes, plasters, C.I. and A.C. sheets. (These have small C and R values).

6.5

$\lambda'_i/\phi'_i$  - Internal Driving Point Function

- i) The variation of  $\lambda'_i$  and  $\phi'_i$  with ' $g$ ' and thickness are similar to that of  $\lambda_o$  and  $\phi_o$ .
- ii) For a given material and thickness  $\lambda'_i$  will be lower than the corresponding  $\lambda_o$ . This is because ' $h_i$ ' is less than ' $h_o$ '.

iii) When both the surface coefficients ( $h_0$  and  $h_1$ ) are equal,  $\lambda_o$  and  $\phi_o$  will also be equal to  $\lambda_i$  and  $\phi'_i$  respectively.

The over all decrement factor ( $\lambda_i$ ) is the product of decrement factors of the outside surface resistance ( $\lambda_o$ ) and the decrement factor of the material ( $\lambda_m$ ).

$$\text{i.e } \lambda_i \angle \phi_i = \lambda_o \angle \phi_o \times \lambda_m \angle \phi_m \quad .. (3)$$

Hence the decrement factor of the material or for the external surface input (i.e., without considering the outside surface coefficient) is obtained from the above equation as

$$\lambda_m \angle \phi_m = \frac{\lambda_i}{\lambda_o} \angle (\phi_i - \phi_o) \quad .. (4)$$

#### 6.6 Thermal Time Constant (T.C) Versus Transfer Function

Some investigators (71) (73) have attempted to correlate the ratio of thermal capacity (C) and the thermal conductance (K/L) i.e., (KC) of a building element with the amplitude decrement factor ( $\lambda_i$ ) and phase lag angle ( $\phi_i$ ) and suggest that the thermal time constant (T.C) can be taken as a single parameter to characterize the non-steady state behaviour of a homogeneous building element. This is true only when the material without surface resistances is considered,

as shown in Fig. (G.1). In order to find out how far this relationship holds good with the inclusion of surface resistances (which is necessary and essential when one considers the behaviour of a material under natural conditions) a graph is plotted between transfer function (fundamental)  $\lambda_i \angle \phi_i$  with surface resistance heat transfer coefficients  $h_0$  and  $h_1$  as 3.6 and 1.6 Dtu/Bq.ft./Hr/ $^{\circ}$ F and 'RC' in Fig. (6.2). This set of curves clearly indicate that such a correlation does not hold good, for these boundary conditions. There is a great scatter between points representing different materials and no single curve could be drawn. It can further be seen that each material forms a separate curve by itself and the curves of dense materials are far separated from those of the insulating materials. Thus two materials having same RC value but widely different it may very well have  $\lambda_i$  values differing from one another by a factor of even 3 or more. These large variations are due to the fact that  $\lambda_i$  depends not only on RC but also on the ratios of  $R/R_i$ ,  $R_0/R$  and  $R_0/R_1$ . This is evident from the equation 14 of Appendix II.

#### 6.7 Effect of Thickness on Transfer and Driving Point Functions

A graph of the fundamental transfer function ( $\lambda_i \angle \phi_i$ ) versus thickness; plotted for different

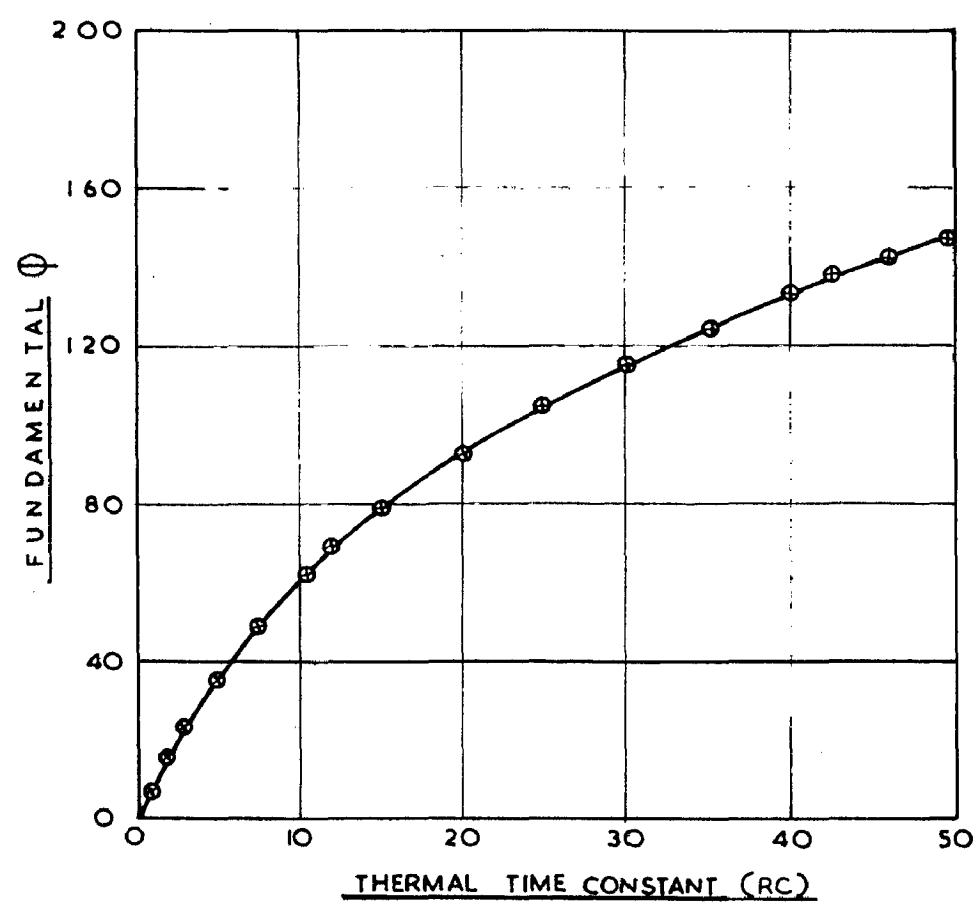
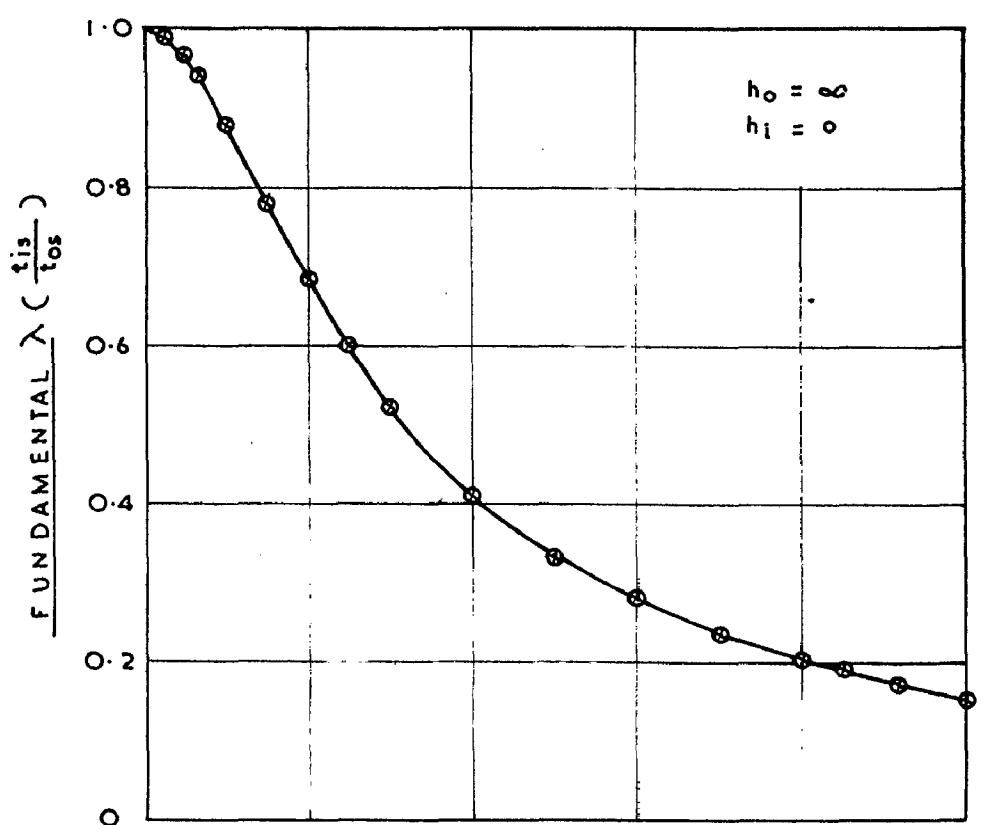


FIG 6.1. THERMAL TIME CONSTANT VERSUS  $\lambda$  AND  $\Phi$

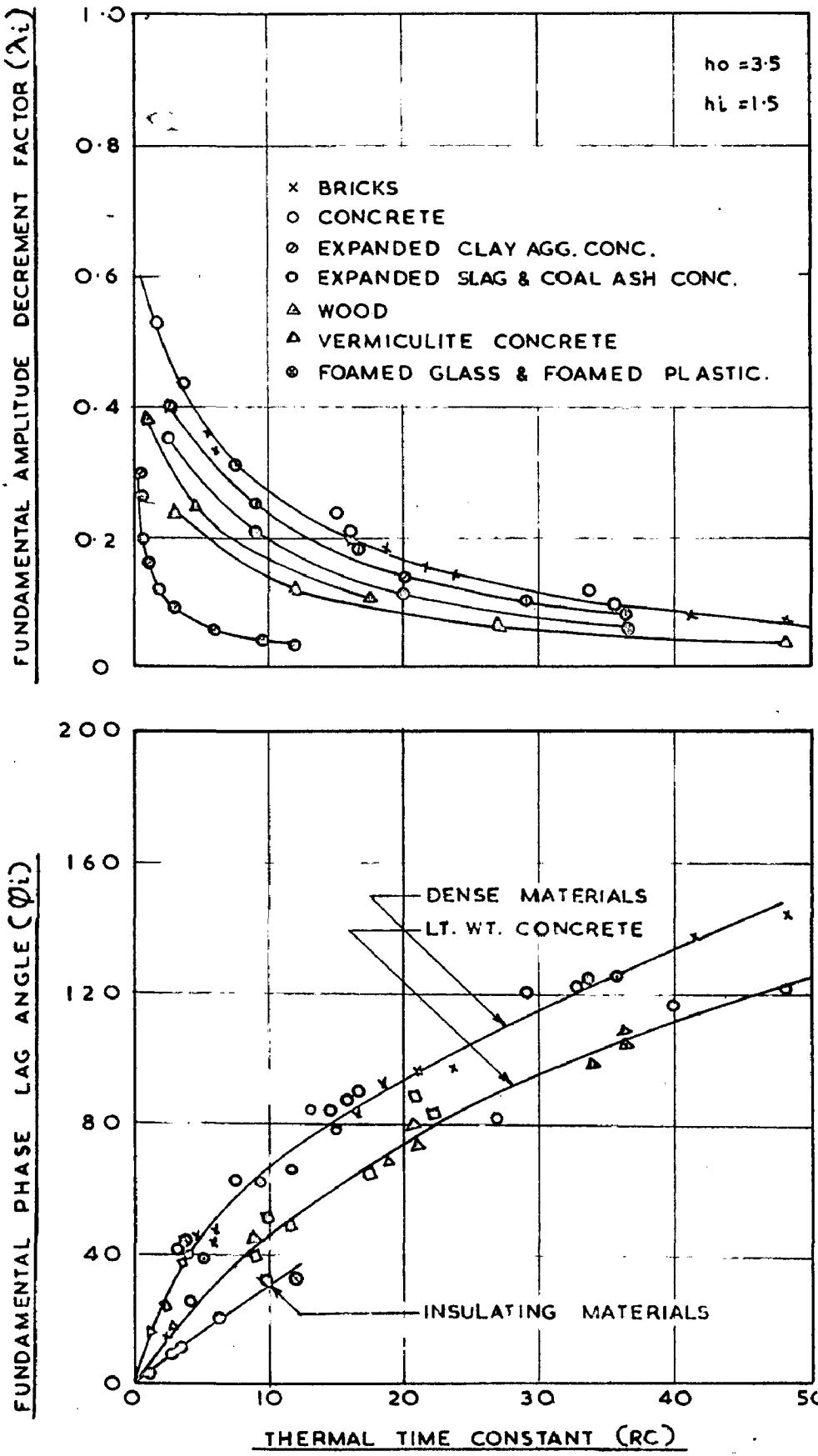


FIG. 6.2. THERMAL TIME CONSTANT VERSUS  $\lambda_i$  AND  $\phi_i$

materials is shown in Fig. (6.3). Here also each material behaves uniquely and no single curve will satisfy all the materials. For any one material the relation between thickness and phase lag is practically a straight line. The slope of these lines are different for different materials. For any material an increase of thickness will result in an increase of  $\lambda C$  and  $R/R_1$  and  $R/R_0$ . As this increase is different for different materials the variation in  $\lambda_i$  and  $\phi_i$  with thickness will also be widely different.

Bunter (73) has shown that the ratio of alternative heat transmittance (which is identical with the transfer admittance function  $Y_i$ ) and the steady state heat transmittance ( $U$ ) i.e.  $Y_i/U$ , versus thickness, provide a good single correlation curve. In his studies Bunter considered only five materials with densities ranging from 43 to 165 Lb/Cu.ft. and thermal conductivities ranging from 1.11 to 11.1 Stu in./sq.ft./hr/ $^{\circ}$ F. The validity of these conclusions were checked for a wider range of materials. For this purpose about twenty materials covering almost all possible values of density, thermal capacity ( $C$ ) and thermal resistance ( $R$ ) were included in this study. The transfer admittance functions  $Y_i/\phi_i$  were obtained by multiplying the analogically determined transfer function  $\lambda_i/\phi_i$  with  $h_i$ .

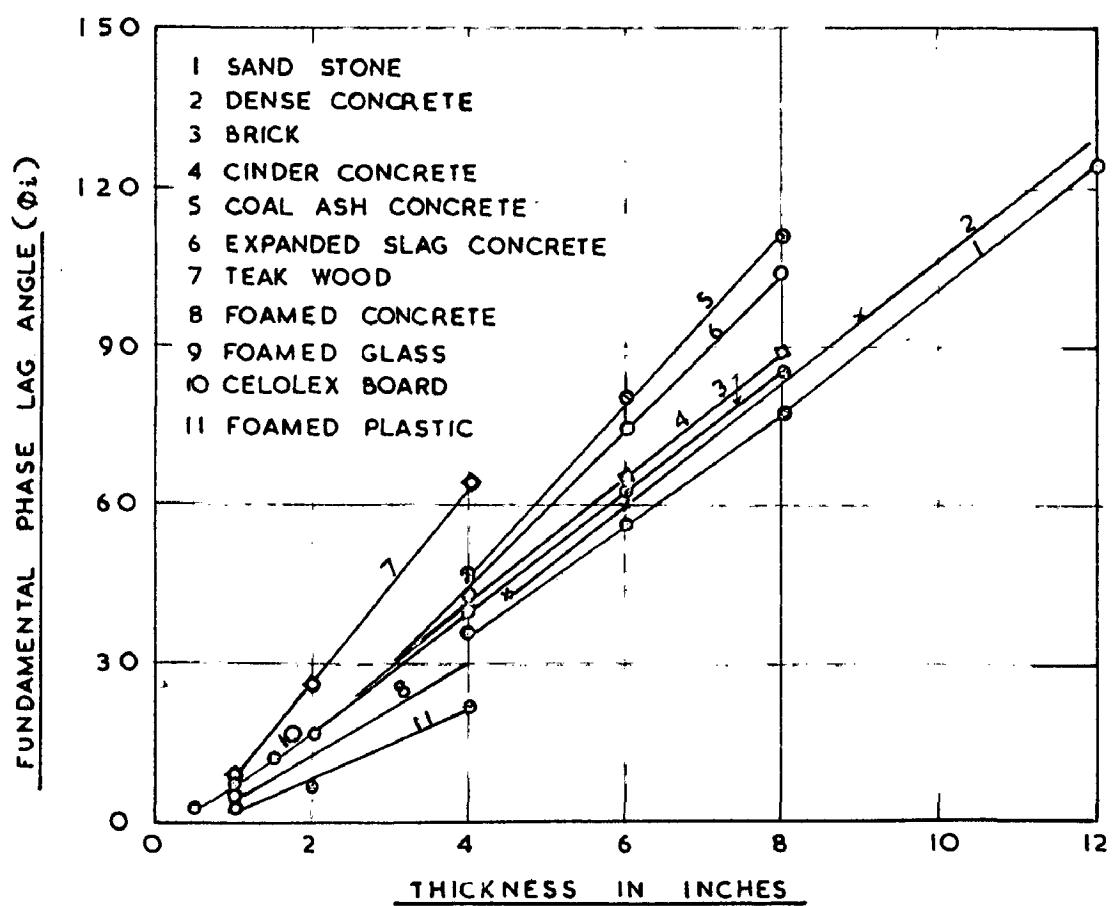
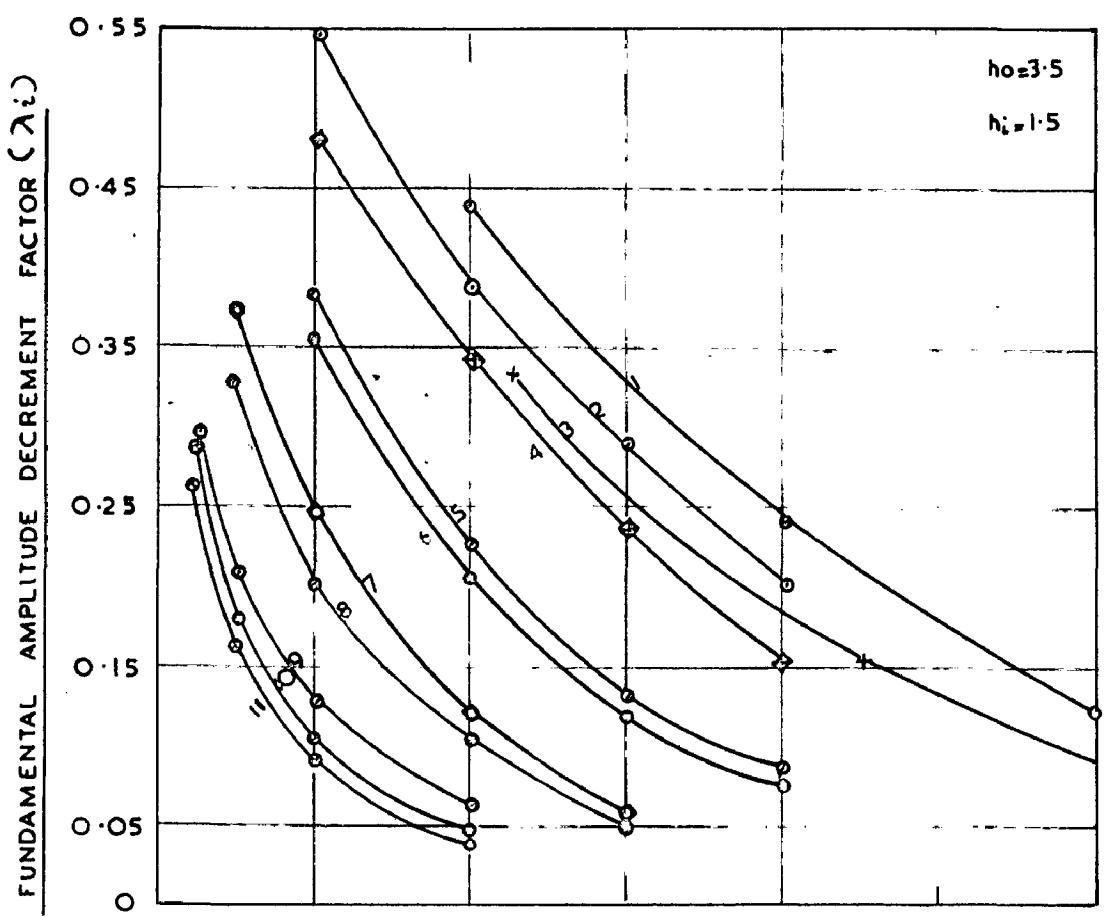


FIG.6.3 THICKNESS VERSUS  $\lambda_i$  AND  $\phi_i$

Taking the values of fundamental  $\frac{Y_1}{U}$  and  $\phi_i$  as ordinates and thicknesses as abscissa, graphs were plotted and shown in Fig. (6.4). All the materials irrespective of their physical properties fairly fit in one single curve, although the deviations of individual points are larger than those of Santor. The main advantage of this relationship is from a knowledge of steady state 'U' value and thickness of a material, the periodic transfer admittance functions are directly obtained. These correlations were checked for three higher harmonics and the shape of the curves are found to be similar, to those of the fundamental.

It will be of interest to find out whether similar relationships could be obtained for driving point admittance functions as well. The external and internal driving point admittance functions i.e.,  $\frac{Y_0}{U_0}$  and  $\frac{Y'_1}{U'_1}$  were obtained from the analogically determined driving point temperature amplitude ratio functions ( $\lambda_0 \angle \phi_0$  and  $\lambda'_1 \angle \phi'_1$ ) from the equations (5) and (6) respectively.

$$\frac{Y_0}{U_0} \angle \phi_0 = h_0 (1 - \lambda_0 \angle \phi_0) \quad \dots (5)$$

$$\text{and } \frac{Y'_1}{U'_1} \angle \phi'_1 = h_1 (1 - \lambda'_1 \angle \phi'_1) \quad \dots (6)$$

The relation between  $\frac{Y_0}{U}$  and  $\frac{Y'_1}{U}$  and thickness is shown in Figs. (6.5 and 6.6). It can be seen that

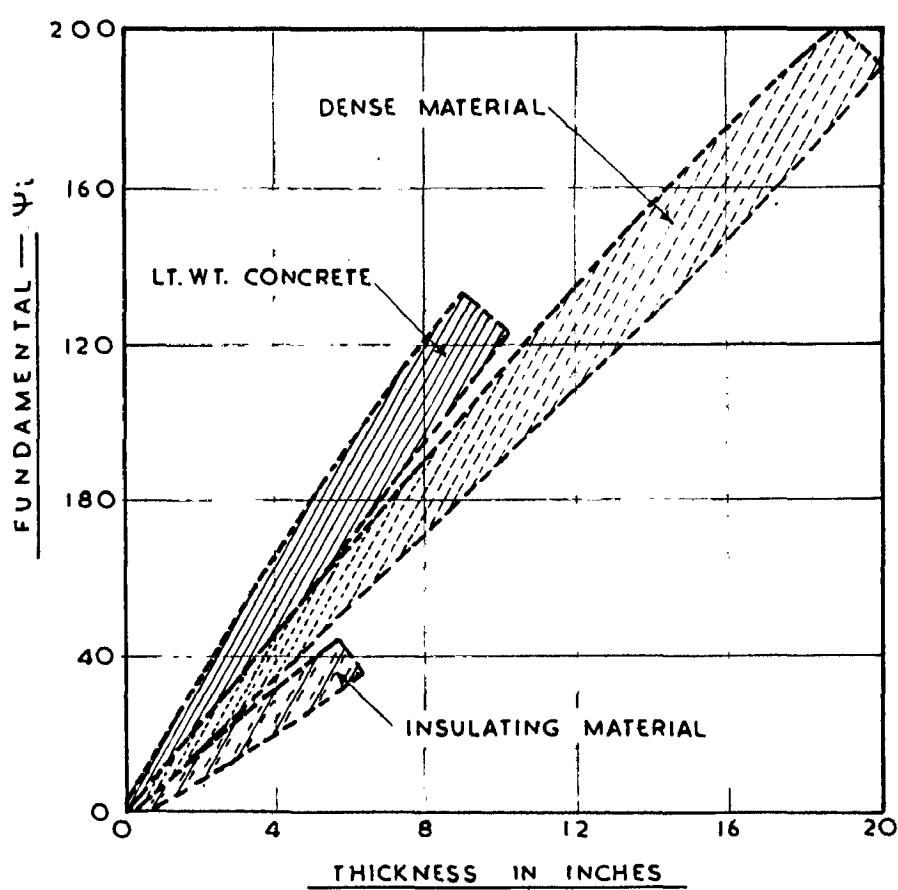
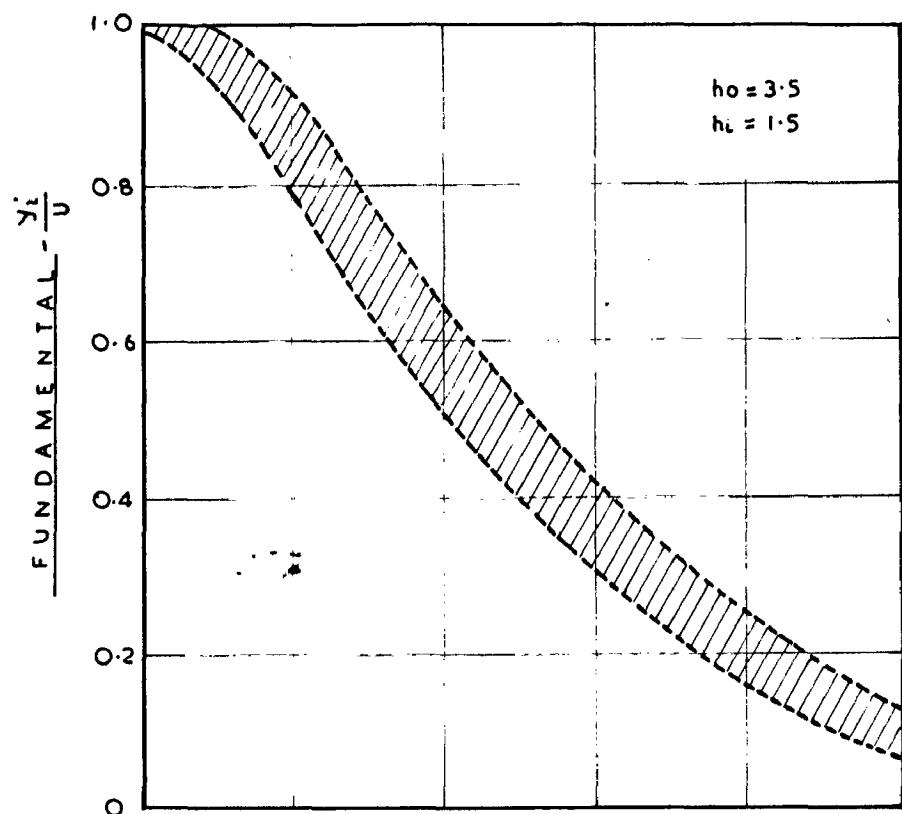


FIG.6.4 THICKNESS VERSUS  $\frac{y_i}{u}$  AND  $\Psi_i$

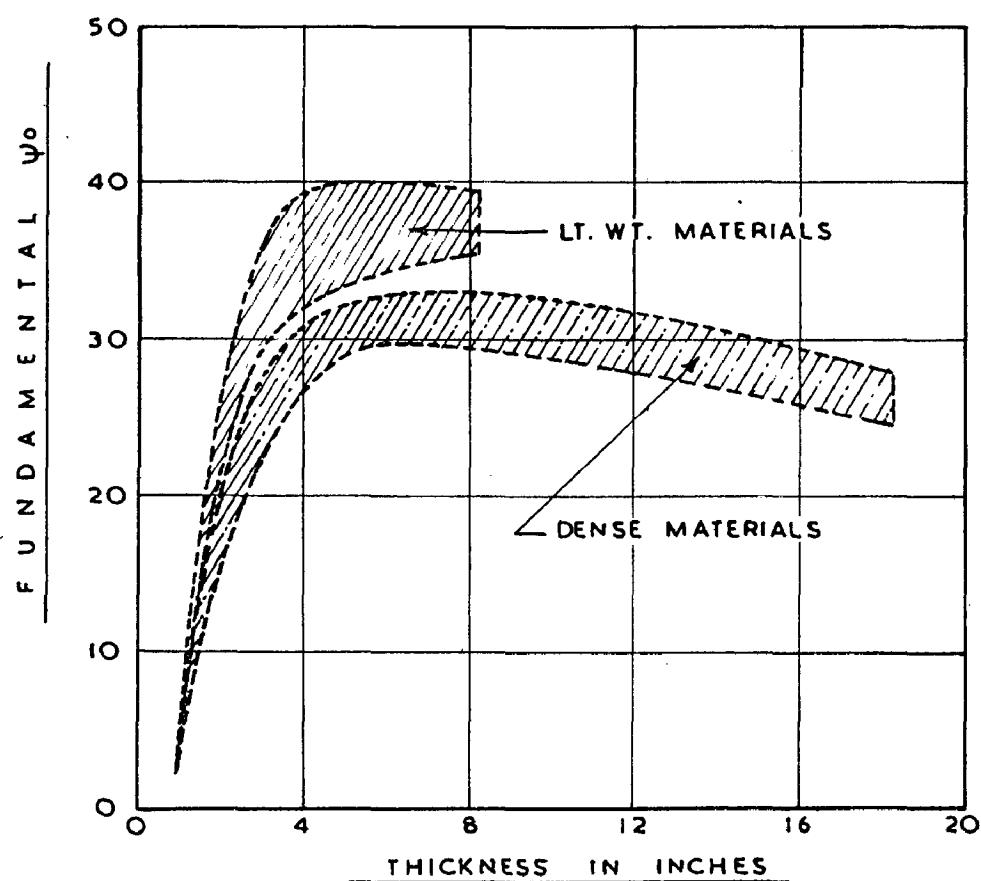
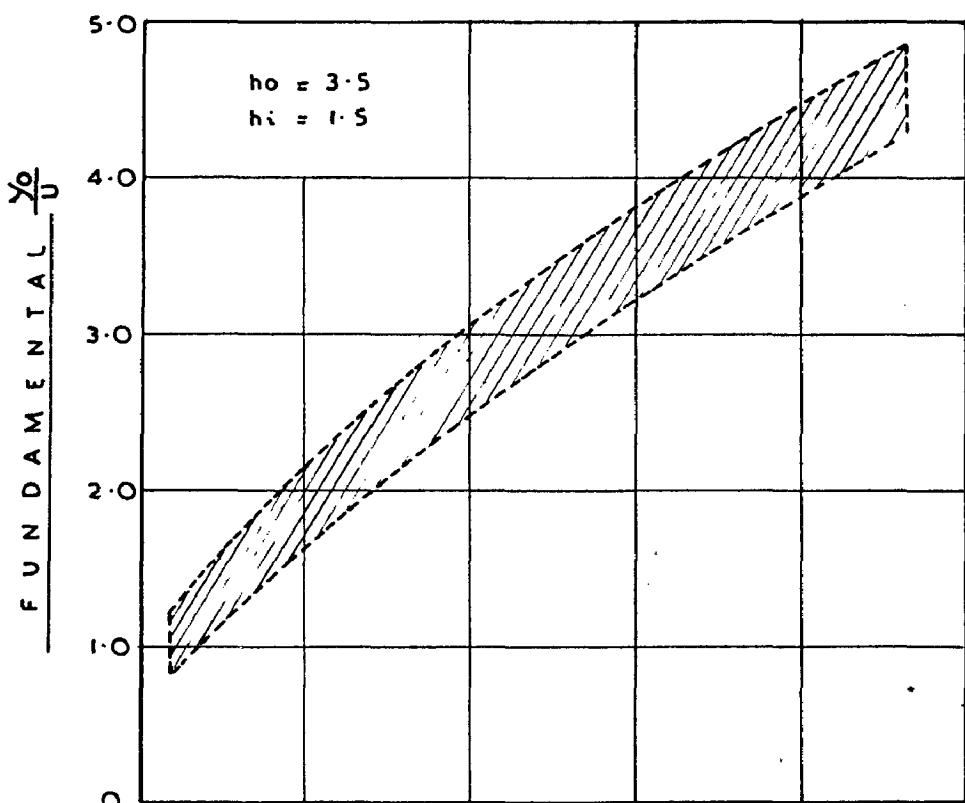


FIG. 65 THICKNESS VERSUS  $\frac{\gamma_o}{U}$  AND  $\psi_o$

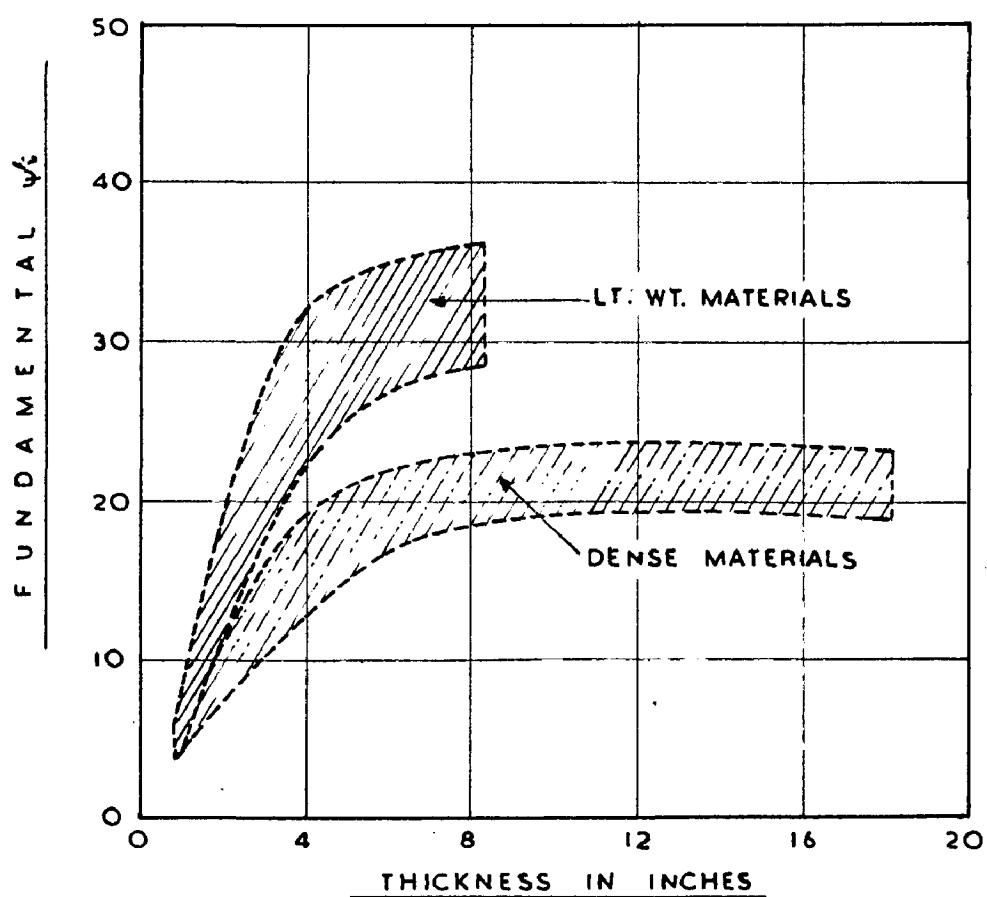
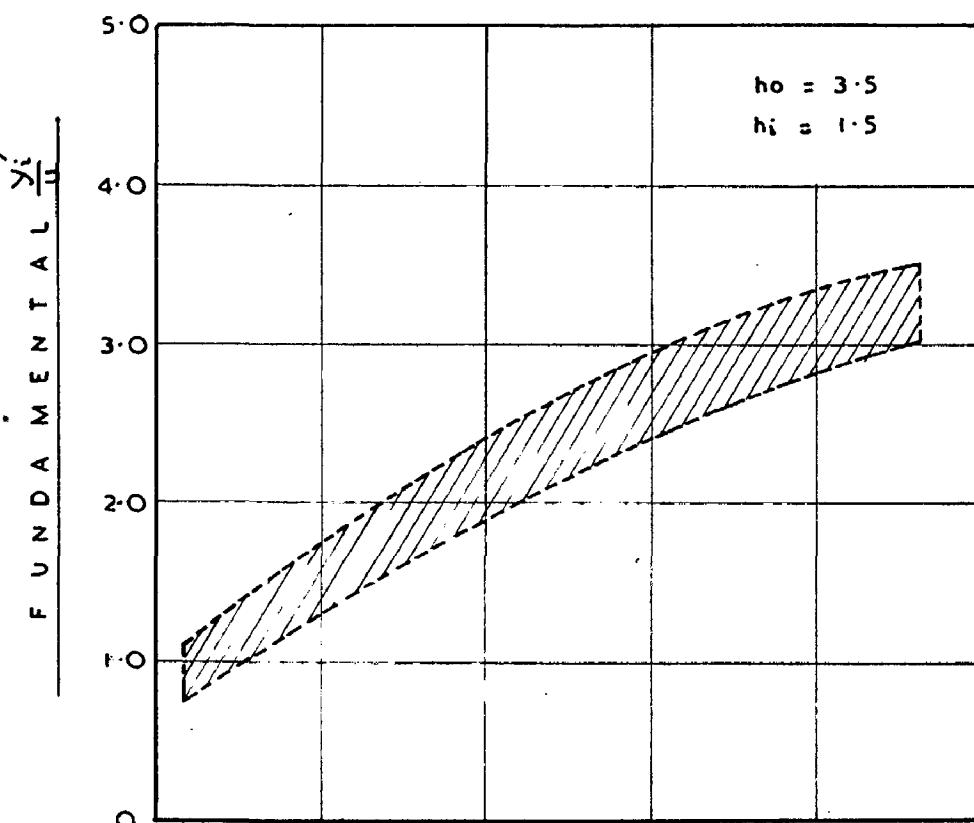


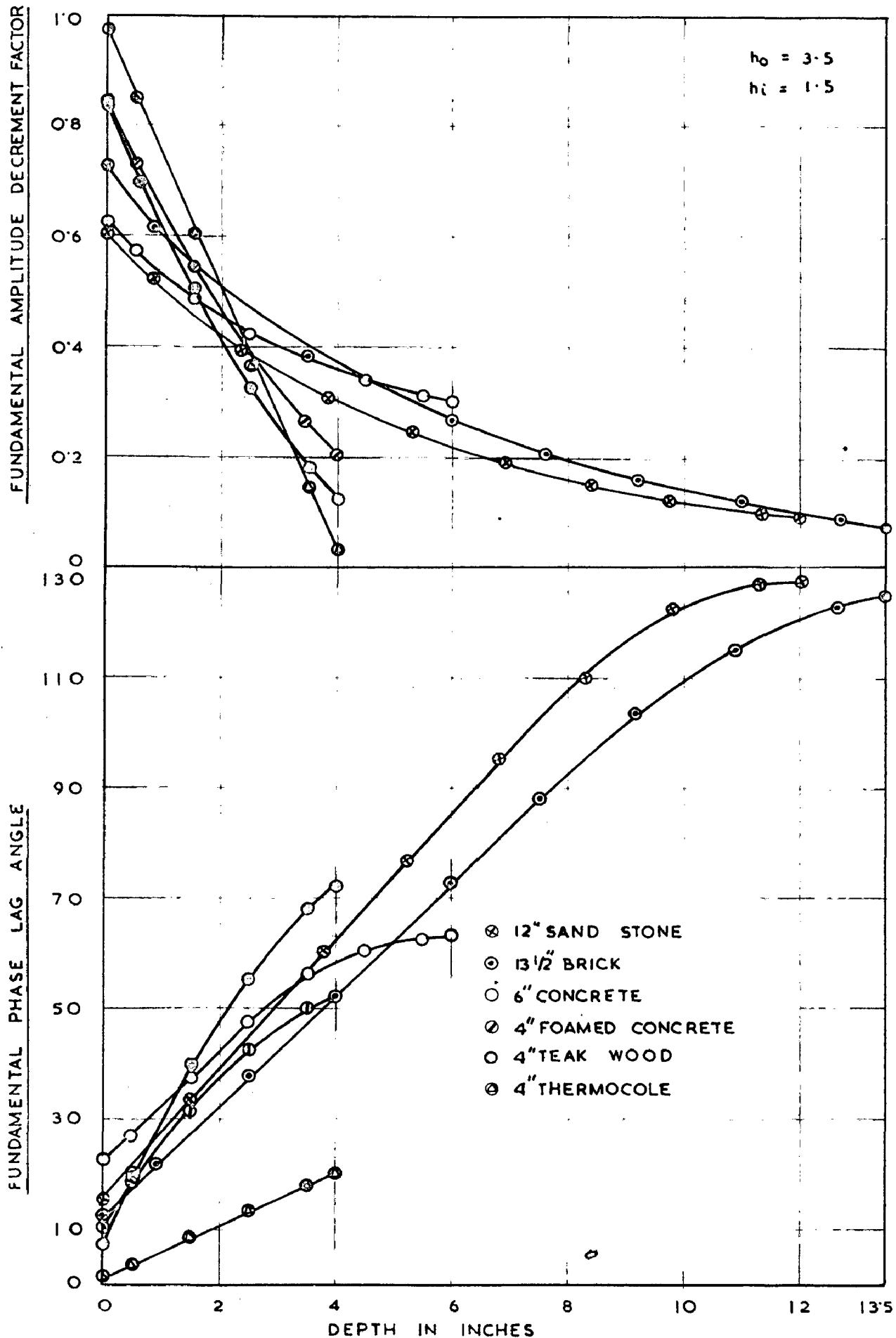
FIG. 6.6 THICKNESS VERSUS  $\frac{y_i}{U}$  AND  $\psi_i$

these also give a similar correlation (i.e., all the materials represented by a single mean curve) but deviations of the individual points from the mean curve are much larger. It can further be noted that the trend of these curves is ascending i.e.,  $\gamma_0/U$  and  $\gamma_1/U$  increase with the increase of thickness, in contrast to that of the curve in Fig. (G.4) where the curve is a descending one i.e.,  $\gamma_1/U$  decreases with the increase of thickness.

In all the three cases the phase angle versus thickness did not yield a single curve correction. Each class of materials give one line.

#### 6.9 Variation of Amplitude Decrement and Phase Inside the Material

The progressive damping and phase shift of a temperature wave, while it advances through a homogeneous material depends on the characteristics of the material. These can be obtained by determining the amplitude decrement factors and phase lag angles at different points inside the material. The analogue is convenient for this purpose as the lumping points are accessible for making these measurements. The amplitude decrement factors and phase lag angles were measured for different types of materials for the fundamental and higher harmonics. The fundamental values shown in Fig. (G.7) is a typical representation.



G. 6.7 VARIATION OF AMPLITUDE DECREMENT AND PHASE LAG WITH DEPTH

The general nature of the curves are similar for higher harmonics though the actual values are different. The decrease in  $\lambda_i$  for higher harmonics is not appreciable for insulating materials as compared to dense materials.

The drop in  $\lambda_i$  with depth is exponential for dense materials except at the boundaries whereas for insulating materials this tends to be a straight line with a steep slope. The increase in phase angle ( $\phi$ ) with depth is practically a straight line throughout except for the points approaching the other boundary.

#### 6.9 Generalized Charts for Transfer and Layering Point Functions

In the Table (3) of Appendix III the thermal function data are given for most of the commonly used homogeneous building materials with different thicknesses. For those studies, the thermal capacity ( $C$ ) and thermal resistance ( $R$ ) were obtained from the average value of the physical properties of the materials listed in Table (1) of the same appendix. As these depend upon many factors, wide variations of these properties (for the same type of material) are to be expected in practice. These will impose limitations on the use, the data presented. They may be taken only as indicative average values, but for any

specific material the precise transfer and driving point functions have to be determined every time experimentally. If generalised charts, relating those thermal functions and the physical properties of materials, are available, the above limitations will be overcome and the data will be of greater utility.

Mackoy and Wright had presented one such set of charts for  $\lambda_i$  and  $\phi_i$  as functions of thermal conductance ( $K/L$ ) and the specific thermal absorption coefficient ( $K\rho_s$ ). These were drawn for a fixed values of  $h_0$  and  $h_1$ , of 4.0 and 1.5 Btu/-q.ft./ $^{\circ}\text{F}$  respectively. These charts are useful only for the estimation of fabric cooling loads of air conditioned buildings. However, for a complete description of the thermal behaviour of a building element, three thermal system functions namely one transfer function and two driving point (external and internal) functions, are required. Further, the surface heat transfer coefficients for internal walls and intermediate floors will be different from those of the external building elements. As the overall thermal system functions are affected by the surface coefficients, the above charts of Mackoy and Wright will not be applicable for internal building elements like partitions.

Roux had found (63) that the most suitable

values for the outside and inside surface heat transfer coefficients for South African climatic conditions were 3.5 and 1.5 Btu/Eq.ft./hr/ $^{\circ}$ F respectively. These values are more appropriate for India, because South African conditions are close to Indian climatic conditions.

Five sets of such generalised thermal function charts are shown in Figs. (6.8 .. through 6.12). In each chart a family of curves one for each  $\alpha$  value ranging from 0.1 to 250, with thermal resistance ( $R$ ) as the abscissa (from 0.1 to 10) and the appropriate fundamental thermal function as the ordinate are drawn. Though ' $\alpha$ ' values greater than '50' are not likely to be met with in practice,  $\alpha$  values even upto '250' have been included as the same set of curves are intended to provide the thermal functions for higher harmonics as well. The transfer and driving point functions given in these charts are :

- 1)  $\lambda_i \angle \phi_i$
- 2)  $\lambda_o \angle \phi_o$  for  $h_o = 3.5$  and  $h_i = 1.5$
- 3)  $\lambda_i \angle \phi'_i$
- 4)  $\lambda_i \angle \phi_i$
- and 5)  $\lambda_i \angle \phi'_i$  for  $h_o = h_i = 1.5$

With these five sets of the charts, the thermal characteristics for periodic heat flow of any material

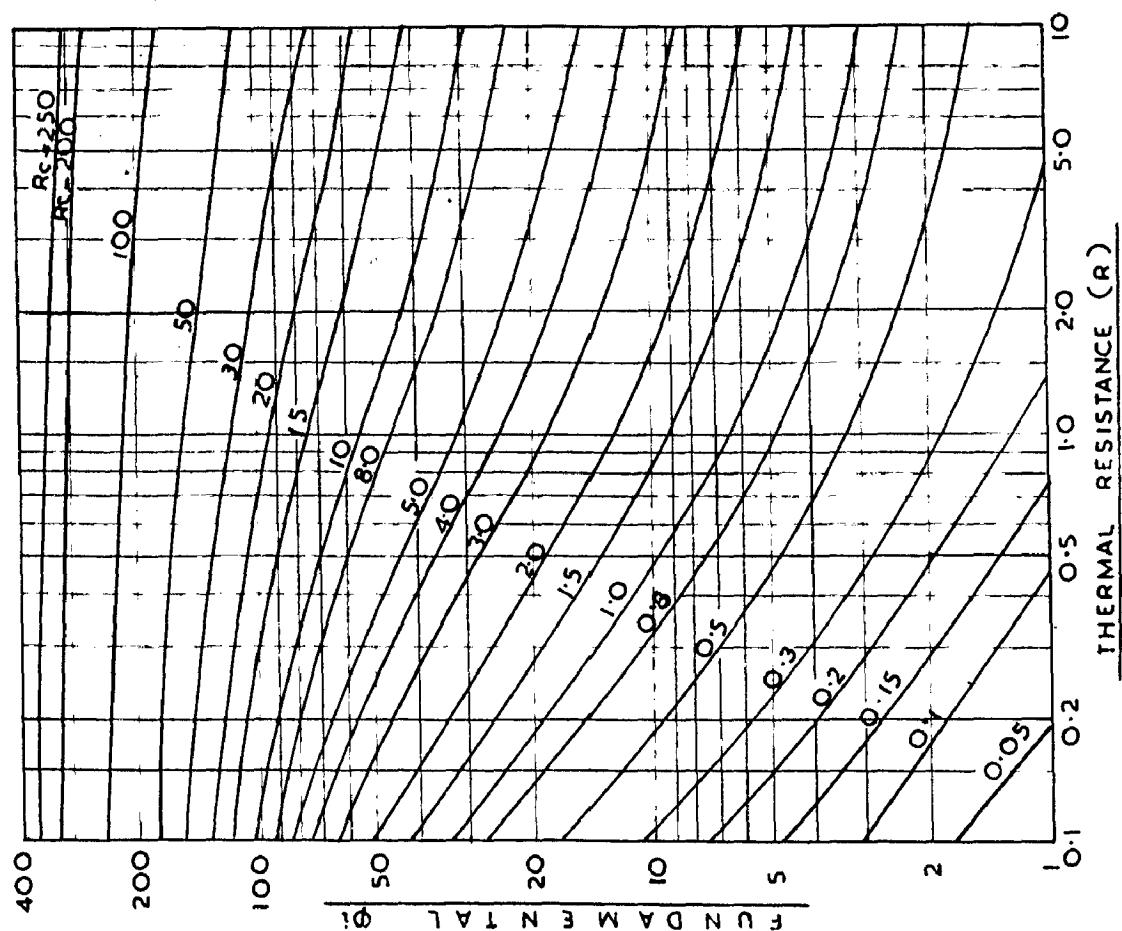
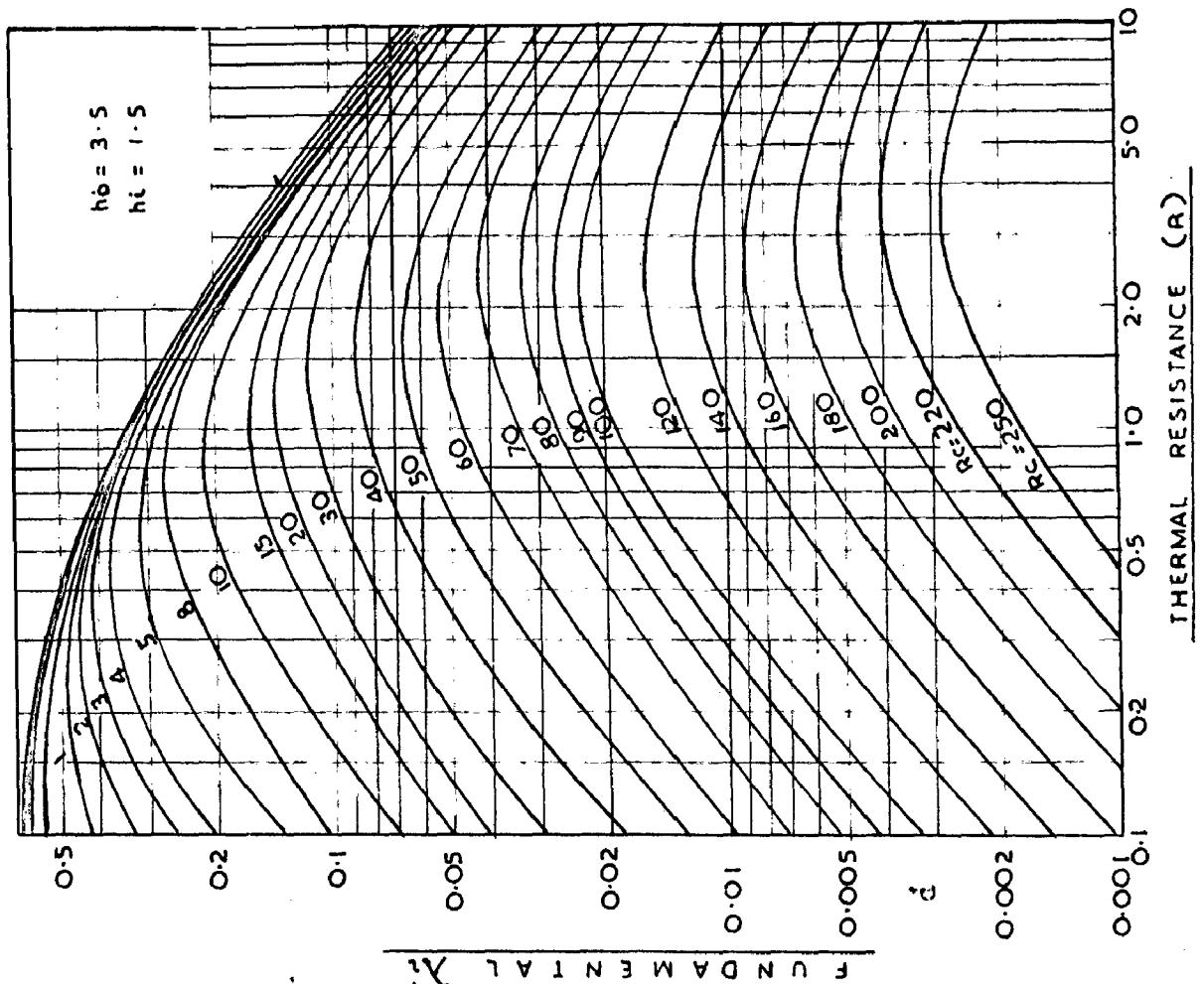


FIG.6.8. R Vs TRANSFER FUNCTION ( $\lambda_i \angle \Phi_i$ ) FOR RC VALUES OF 0.1 TO 250

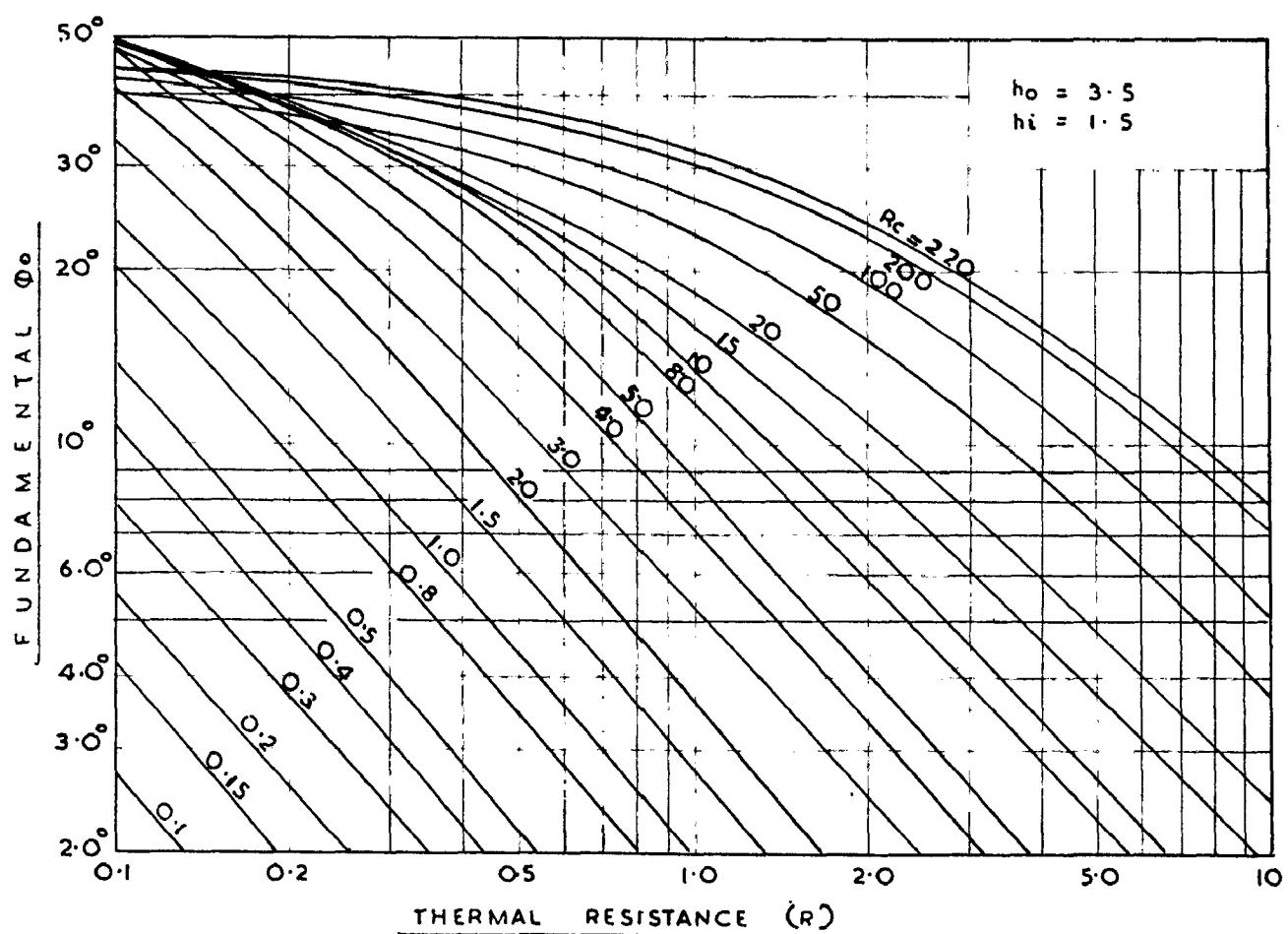
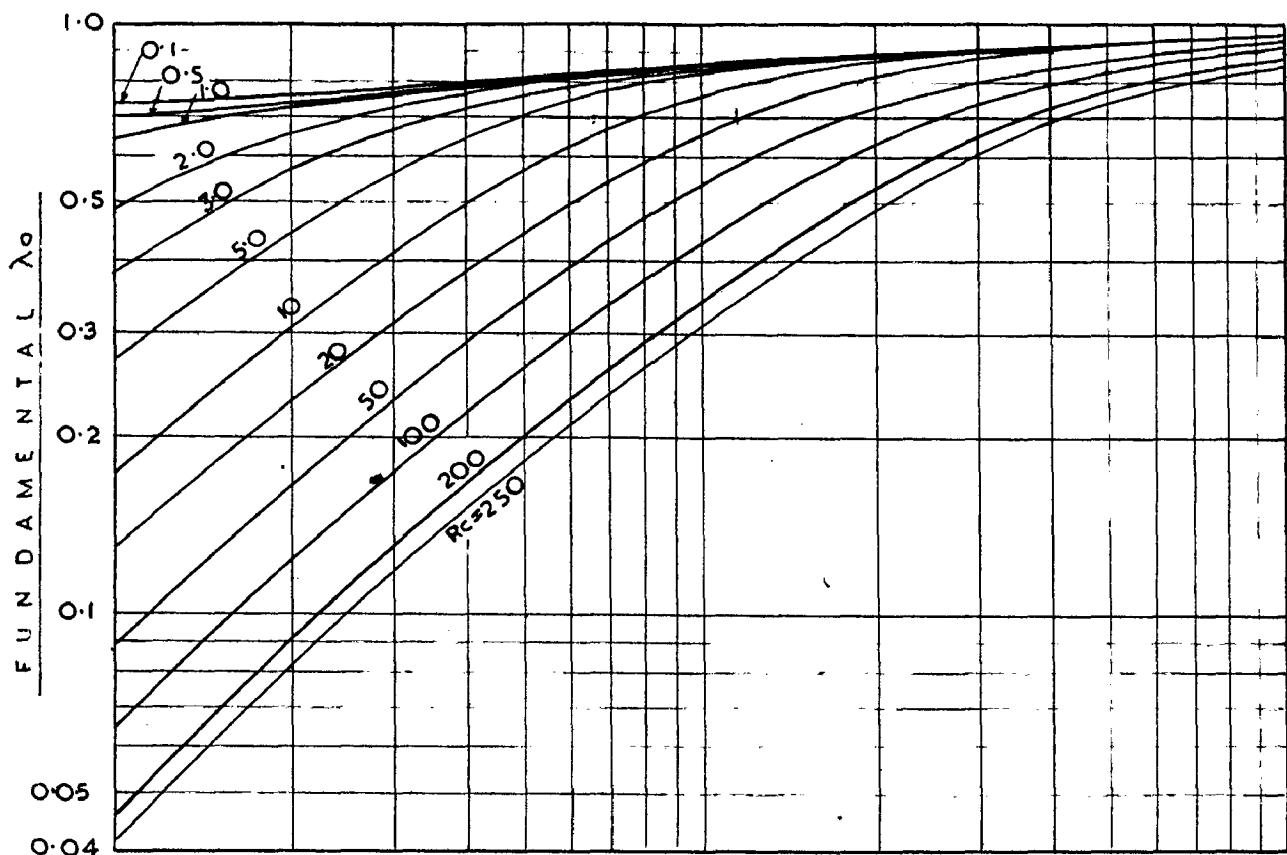


FIG. 6.9  $R$  Vs EXTERNAL DRIVING POINT- FUNCTION ( $\lambda_0 L - \phi_0$ )  
FOR  $RC$  VALUES 0.1 TO 250

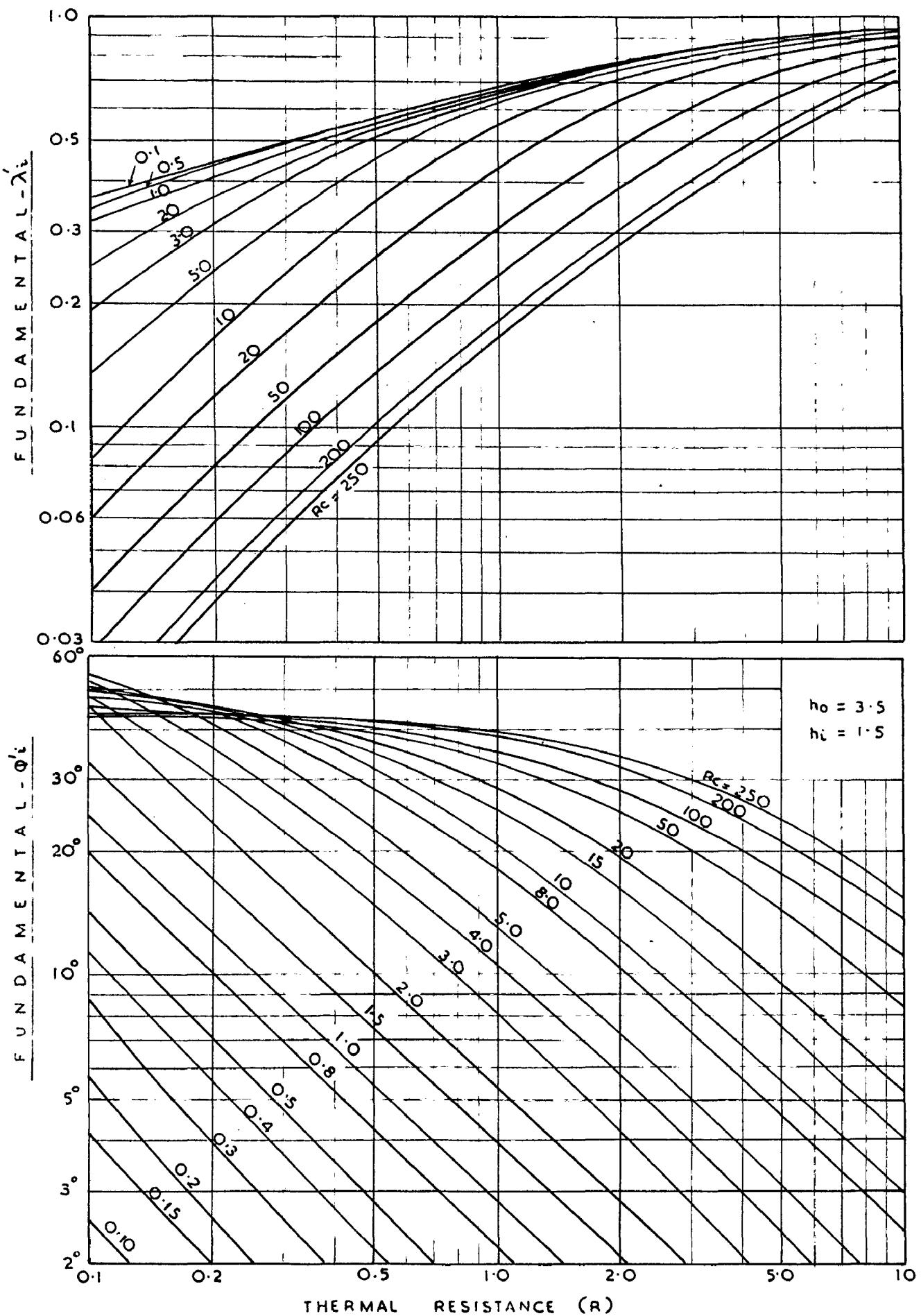


FIG. 6.10. R Vs INTERNAL DRIVING POINT FUNCTION ( $\lambda'_i L - \phi'_i$ )  
FOR RC VALUES 0.1 TO 250

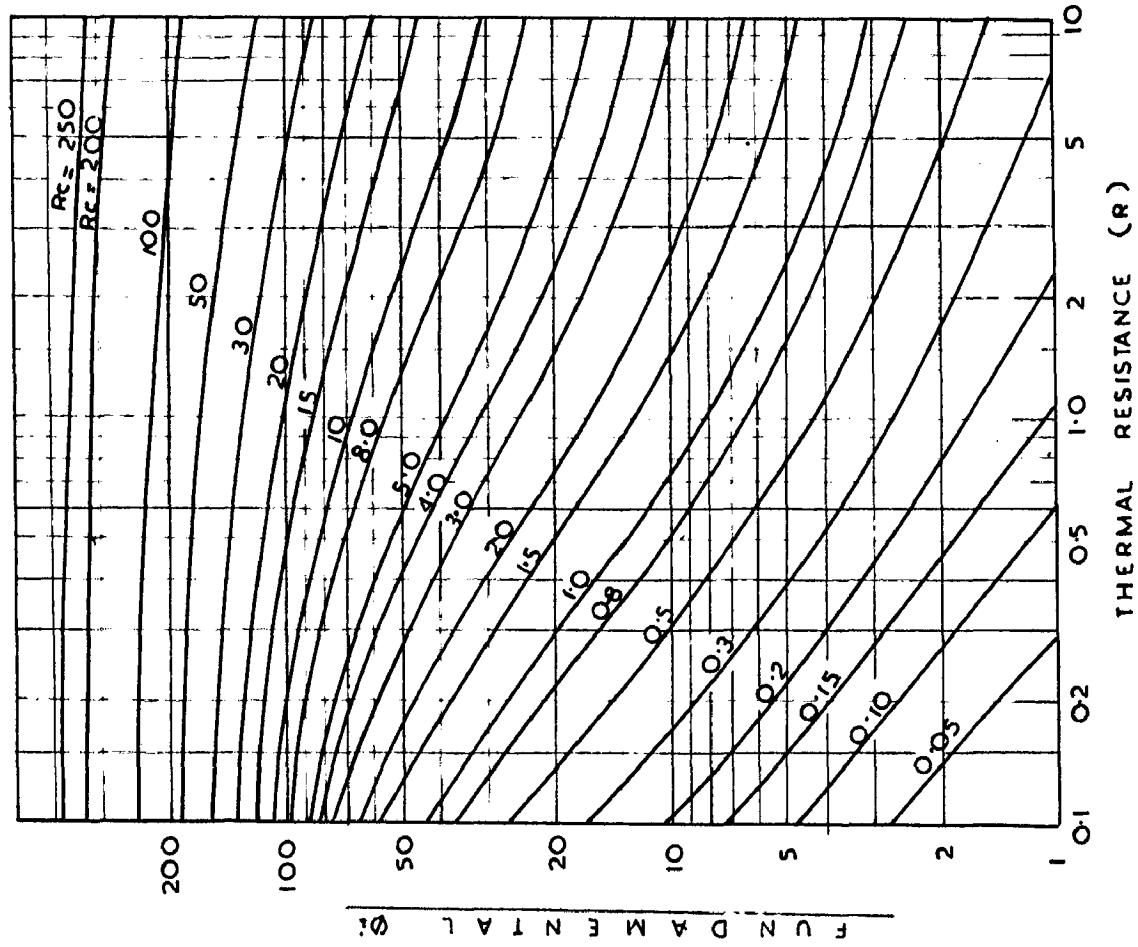
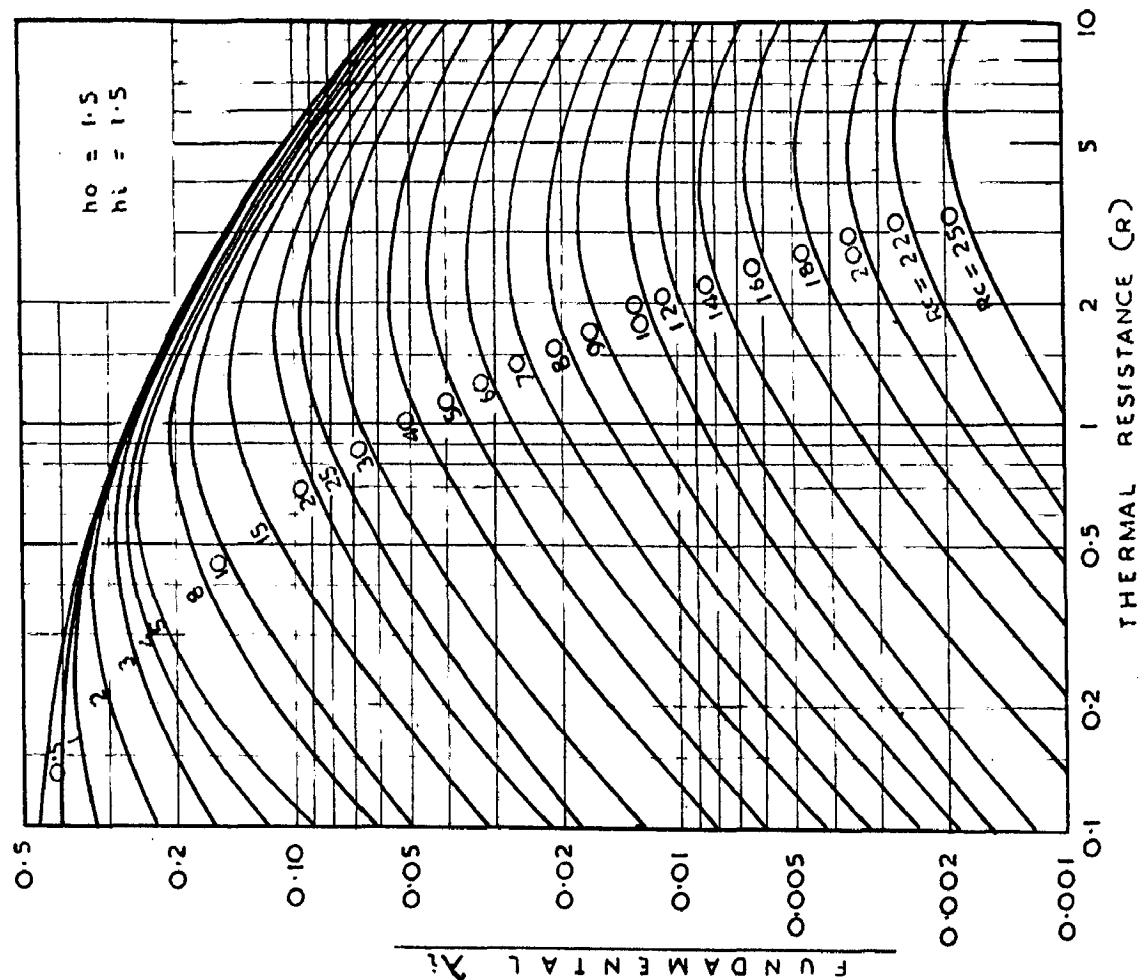


FIG. 6-II. R VS TRANSFER FUNCTION ( $\lambda_i \leq \theta_i$ ) FOR RC VALUES 0.1 TO 250

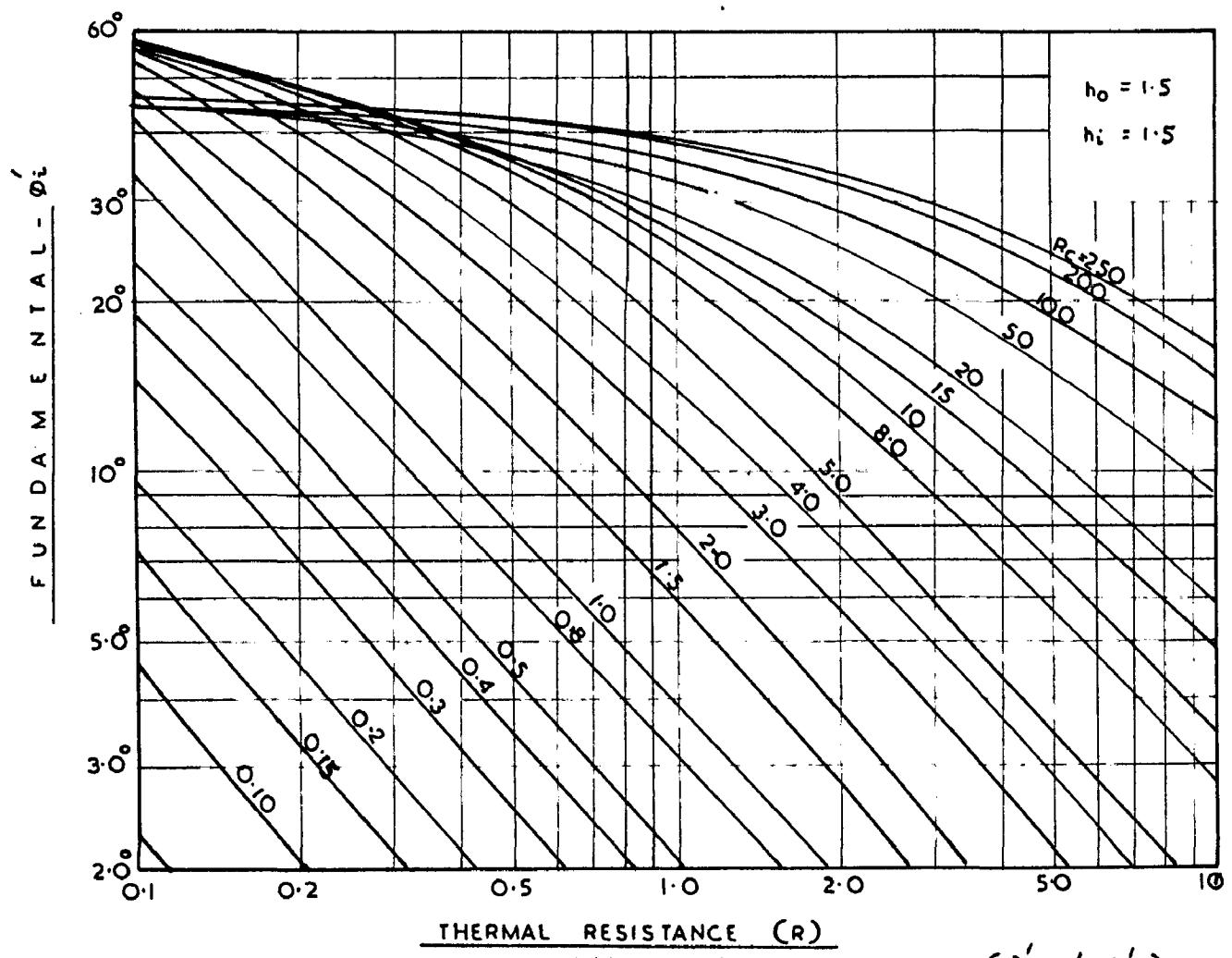
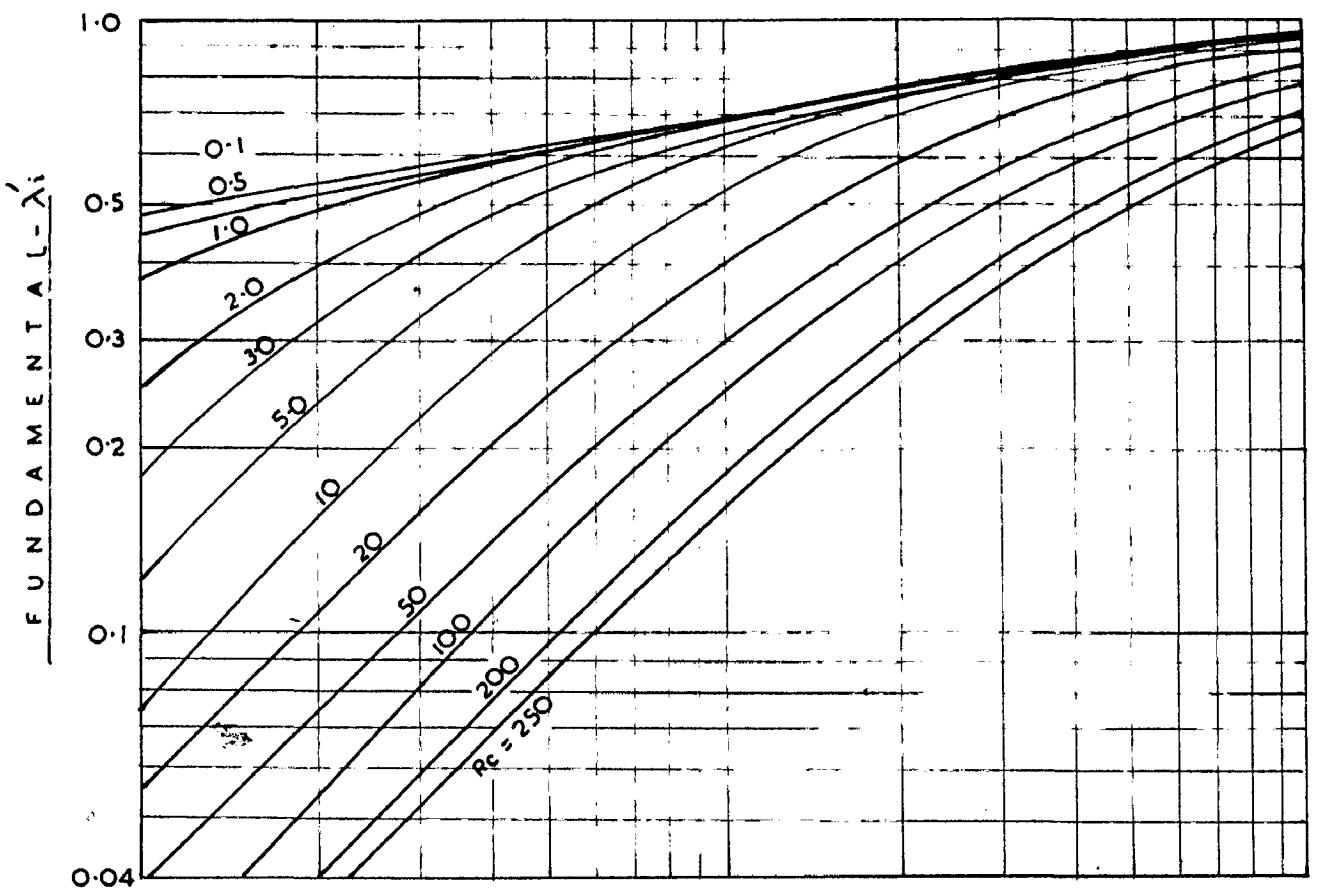


FIG. 6.12.  $R$  Vs INTERNAL DRIVING POINT FUNCTION ( $\chi_i \angle \phi_i$ )  
FOR  $R$  VALUES 0.1 TO 250

of known  $\kappa$ ,  $\rho$ ,  $s$  values and a given thickness ( $L$ ) whether used as an external or internal element can directly be obtained.

Though these charts are for the fundamental (thermal functions) the higher harmonics can also be obtained from them. The characteristics for the  $n$  th harmonic (period  $24/n$  hours) can be obtained from the charts by using a value of  $RC$  of  $n$  times the actual value for the same  $\kappa$ .

The use of these charts is explained below by an illustrative example. Suppose we are interested to find the thermal transfer and driving point functions for a material with  $\kappa = 8.0$  properties of

$$\rho = 120 \quad s = 0.20 \quad \text{and} \quad L = 10^{\prime\prime} \text{ inches.}$$

First of all the thermal resistance  $R$  and thermal capacity  $C$  are to be determined.

$$R = L/\kappa = 10.0/8.0 = 1.25$$

$$C = L\rho s = 10/12x120x0.2 = 20$$

$$\therefore RC = 1.25 \times 20 = 25$$

For those values of  $R$  and  $RC$ , the transfer and driving point functions read from the charts, for the fundamental and three higher harmonics, are given in Table (6.1).

These charts are mostly self explanatory. However, the following broad inferences can be made

TABLE 6.1

TERMAL SYSTEM FUNCTIONS OBTAINED FROM  
THE CHARTS FOR  $RC = 25$  AND  $R = 1.36$   
(Illustrated example)

Harmonic	External Driving point Function		Transfer Function		Internal Driving point Function	
	Modulus	Arg 'in deg'	Modulus	Arg 'in deg'	Modulus	Arg 'in deg.'
	$\lambda_0$	$-\phi_0$	$\lambda_i$	$-\phi_i$	$\lambda'_i$	$-\phi'_i$
Fundamental	0.67	17	0.14	100	0.46	26
Second	0.60	20	0.07	151	0.36	30
Third	0.55	22	0.03	194	0.32	32
Fourth	0.50	24	0.02	220	0.23	33

from them.

### 1. $\lambda_i / \phi_i$

- 1) For a given value of ' $\omega$ ',  $\lambda_i$  will decrease and  $\phi_i$  increases, with the increasing of  $\mu$ .
- 2) Materials with very low  $\mu$ , behave as if they are purely resistive.
- 3) For a given  $\mu$ ,  $\lambda_i$  is low for very low values of  $\omega$  and increases, with the increase of  $\omega$ , reaching a maximum and decreases with a further increase of  $\omega$ .
- 4) The occurrence of maximum is shifted towards higher values of  $\omega$  with the increase of ' $\mu$ '.
- 5) For a given  $\omega$ ,  $\phi_i$  will decrease with the increase of  $\mu$  and for a given  $\mu$  it increases with the increase of  $\omega$ .

### 2. $\lambda_o / \phi_o$

- 1) The magnitude of  $\lambda_o$  increases, with the increase of  $\omega$ , for a given  $\mu$ , while it decreases with the increase of  $\mu$  for a given  $\omega$ .
- 2) For large values of  $\omega$  all the curves for different values of  $\mu$ , tend to converge.
- 3)  $\phi_o$  decreases, with the increase of  $\omega$ , for a given  $\mu$ , while it increases with the increase of  $\mu$  for a given  $\omega$ .

3.  $\lambda'_1 / -\phi'_1$

The variation of  $\lambda'_1$  and  $\phi'_1$  with  $\alpha$  and  $\omega$  are similar to that of  $\lambda_0$  and  $\phi_0$ . The values of  $\lambda'_1$  is lower than  $\lambda_0$ , while  $\phi'_1$  are higher than  $\phi_0$ , for a given  $\alpha C$  and  $\beta$ .

6.10 Effect of Surface Heat Transfer Coefficients on Thermal System Functions

It is an accepted procedure to assume the surface heat transfer coefficients ( $h_o$  and  $h_i$ ) as constant in the computations of periodic heat flow. In the above studies the outside and inside surface heat transfer coefficients were assumed to be constant of values 3.5 and 1.5 Btu./sq.ft./hr/ $^{\circ}$ F, respectively. However, in practice these are not strictly constant and sometimes vary within wide limits, especially the outside surface coefficient ( $h_o$ ). Hence for more precise computations it is required to know quantitatively the effect of the variations of ' $h_o$ ' and ' $h_i$ ' on these thermal system functions. Johnson (74) has analysed this problem mathematically taking  $\lambda'_1 / -\phi'_1$ , as a function of three dimension less moduli. However, he did not study the effect of these surface heat transfer coefficients on the driving point functions. These have now been investigated for all the three sets of thermal system functions, as a function of

R (values ranging from 0.1 to 20) and RC (values ranging from 0.1 to 200) for a range of surface coefficients of practical importance (varying from 0.5 to 10). As these form a large number of family of curves only a few typical curves have been presented here. Those include two values of RC viz., 1.0, and 20 for values of R varying from 0.1 to 20. In one set ' $h_0$ ' was kept constant at 3.5 and ' $h_1$ ' varied (0.5, 1.5 and 5.0) and in another set ' $h_1$ ' was fixed as 1.5 and ' $h_0$ ' varied (1.0, 3.5 and 10). The results of those studies are presented graphically in Figs. (6.13 through 6.21). Though these sets of curves clearly bring out the effect of surface resistance and on thermal functions, the broad conclusions drawn are summarized below.

- 1) The influence of surface heat transfer coefficients, on the thermal system functions depends upon the RC and 'R' values of a given building element. Sections with small RC and 'R' are affected to a greater extent than those of with large RC and 'R'. For different building sections having the same RC, but of different 'R' the effect of  $h_0$  and  $h_1$  will be different. Their influence on thermal system functions, is not significant for sections with large ' $h_0$ '.

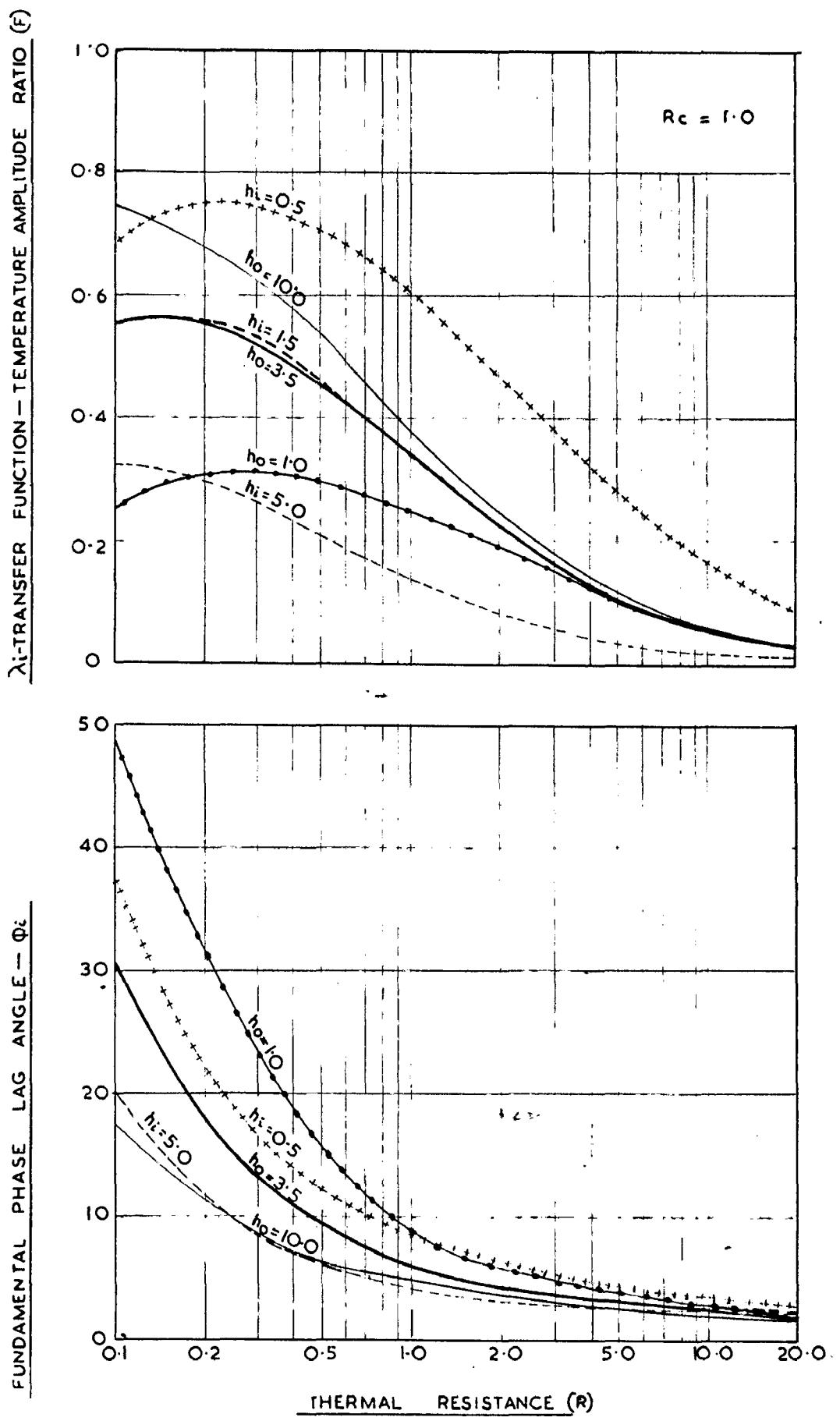


FIG.6.13 EFFECT OF  $h_o$  AND  $h_i$  ON  $\lambda_i$  AND  $\phi_i$  FOR HOMOGENEOUS SECTIONS WITH LOW TIME CONSTANT

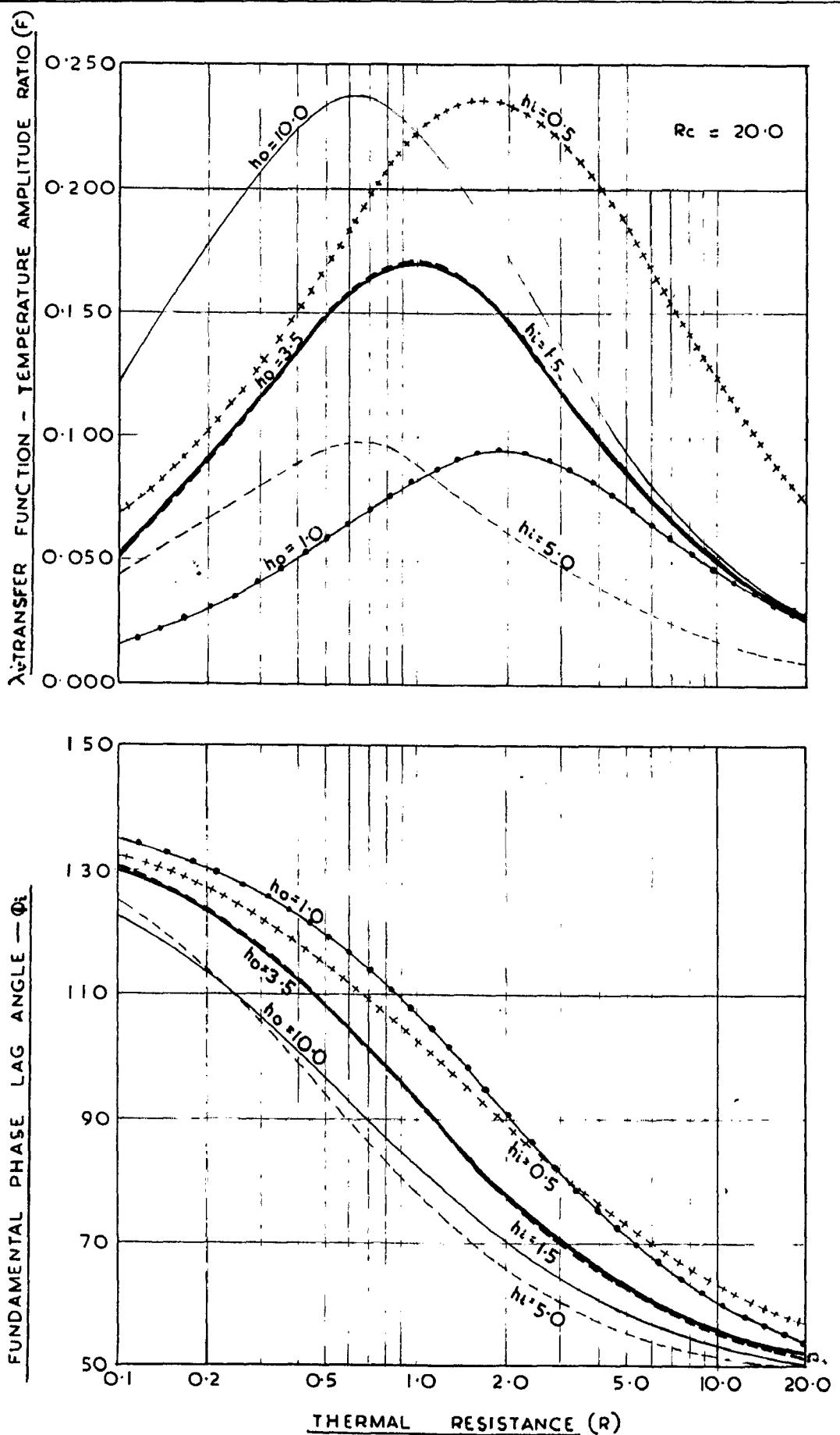


FIG.6.14 EFFECT OF  $h_o$  AND  $h_i$  ON  $\lambda_i$  AND  $\phi_i$  FOR HOMOGENEOUS  
SECTIONS WITH MEDIUM TIME  
CONSTANT

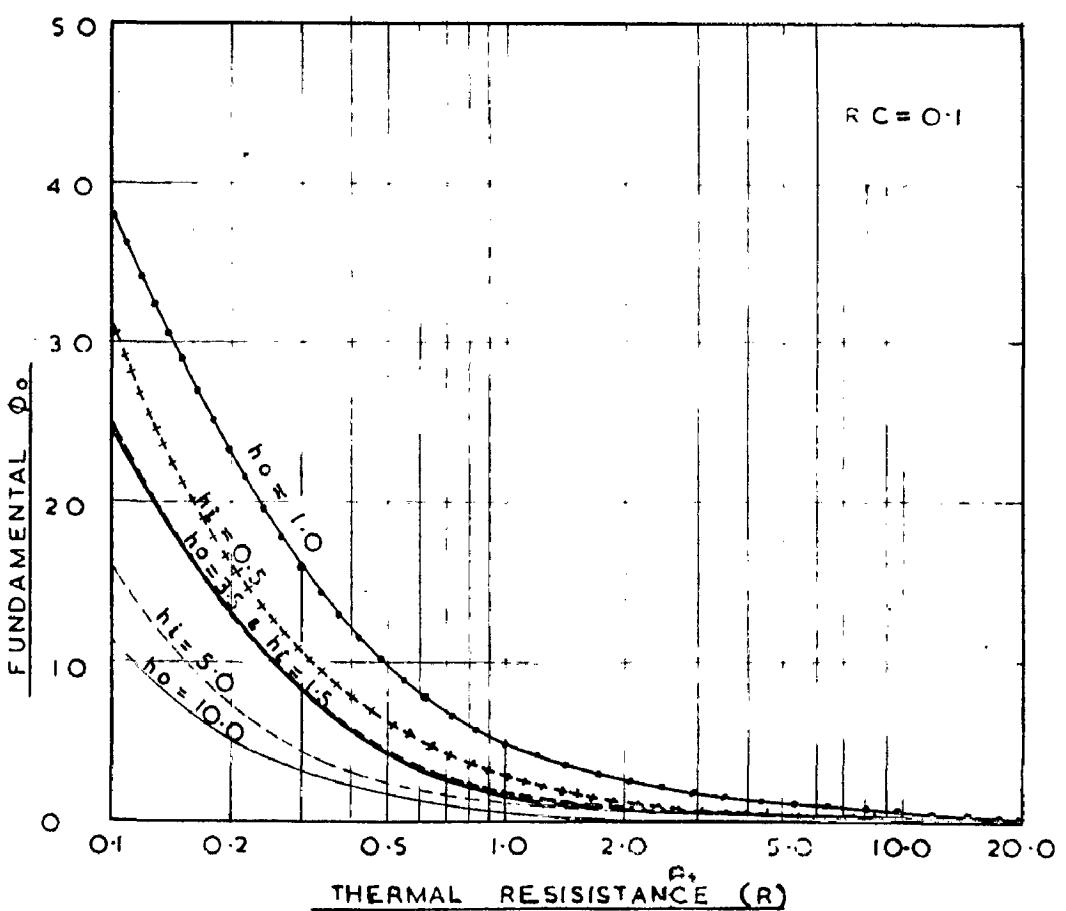
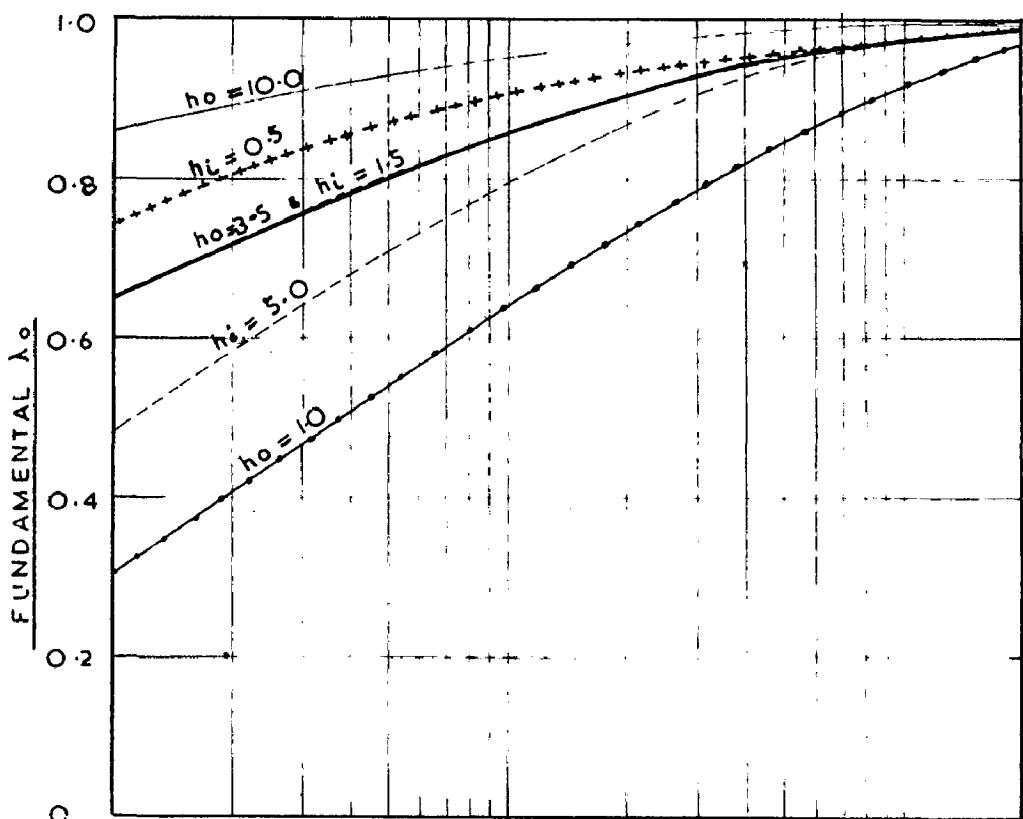


FIG.6.15 EFFECT OF  $h_o$  AND  $h_i$  ON  $\lambda_0$  AND  $\phi_0$  FOR HOMOGENEOUS SECTIONS WITH LOW TIME CONSTANT

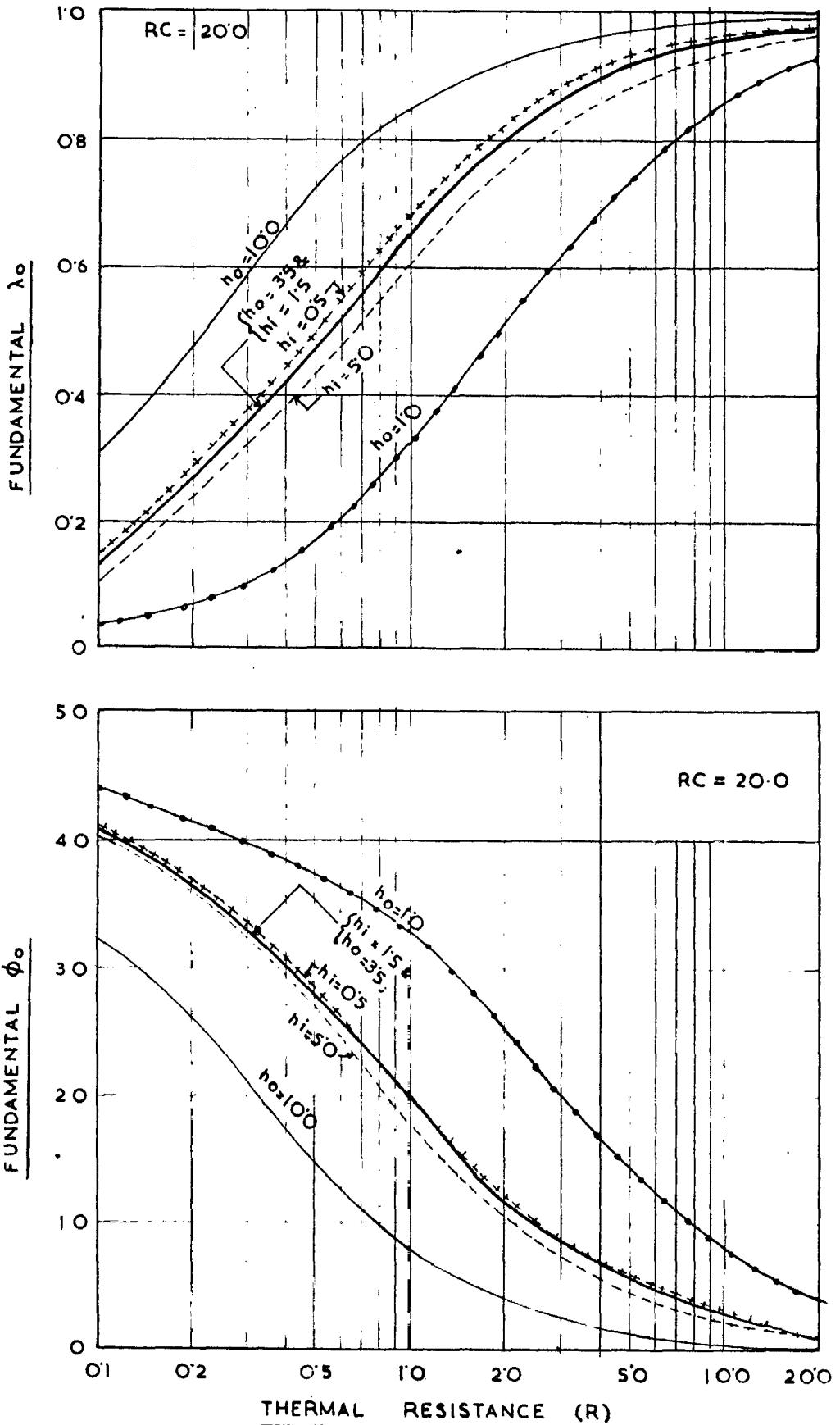


FIG. 6.16. EFFECT OF  $h_o$  AND  $h_i$  ON  $\lambda_0$  AND  $\phi_0$  FOR HOMOGENEOUS SECTIONS WITH MEDIUM TIME CONSTANT

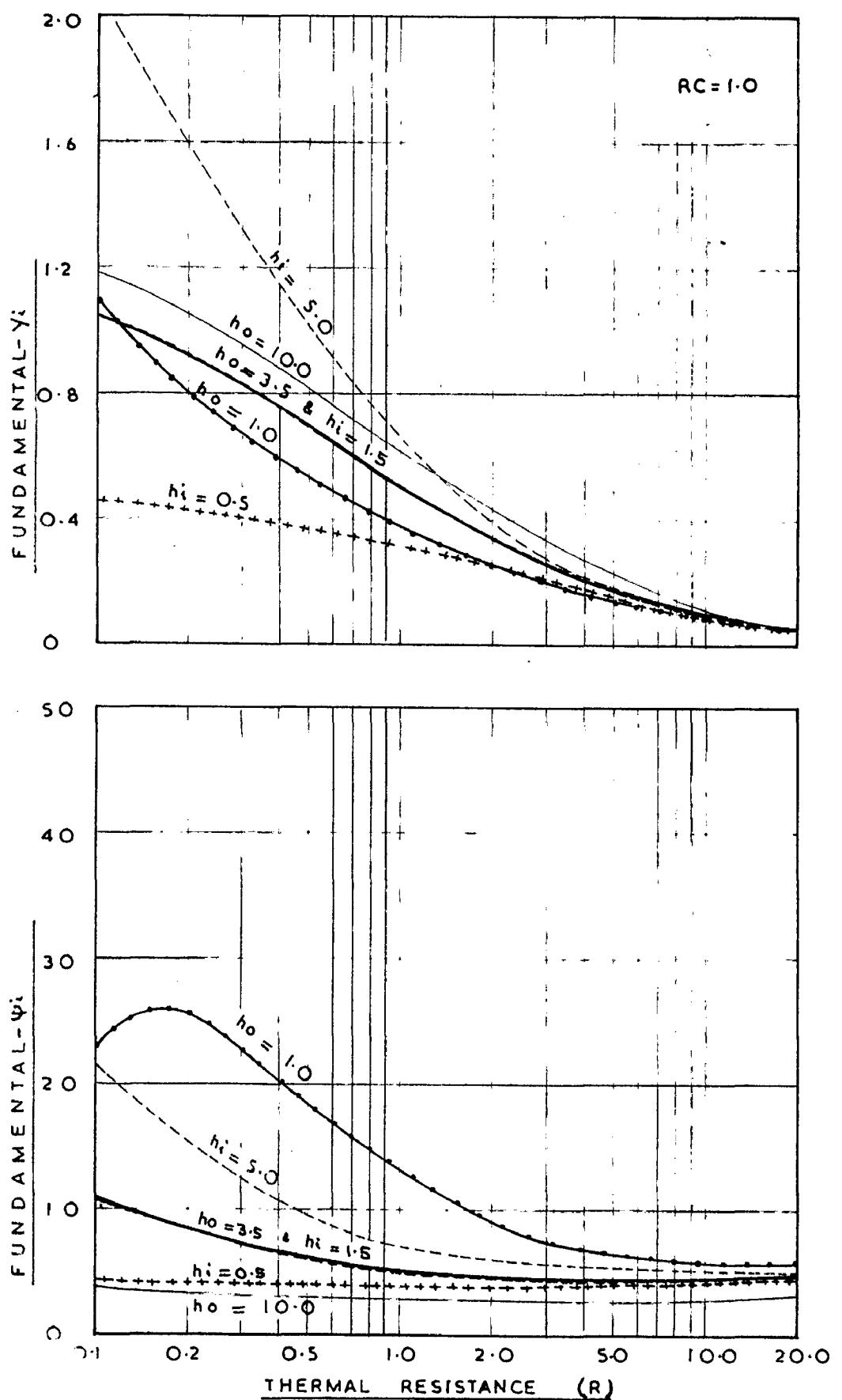


FIG.6.17 EFFECT OF  $h_o$  AND  $h_i$  ON  $Y'_i$  AND  $\Psi'_i$  FOR HOMOGENEOUS SECTIONS WITH LOW TIME CONSTANT

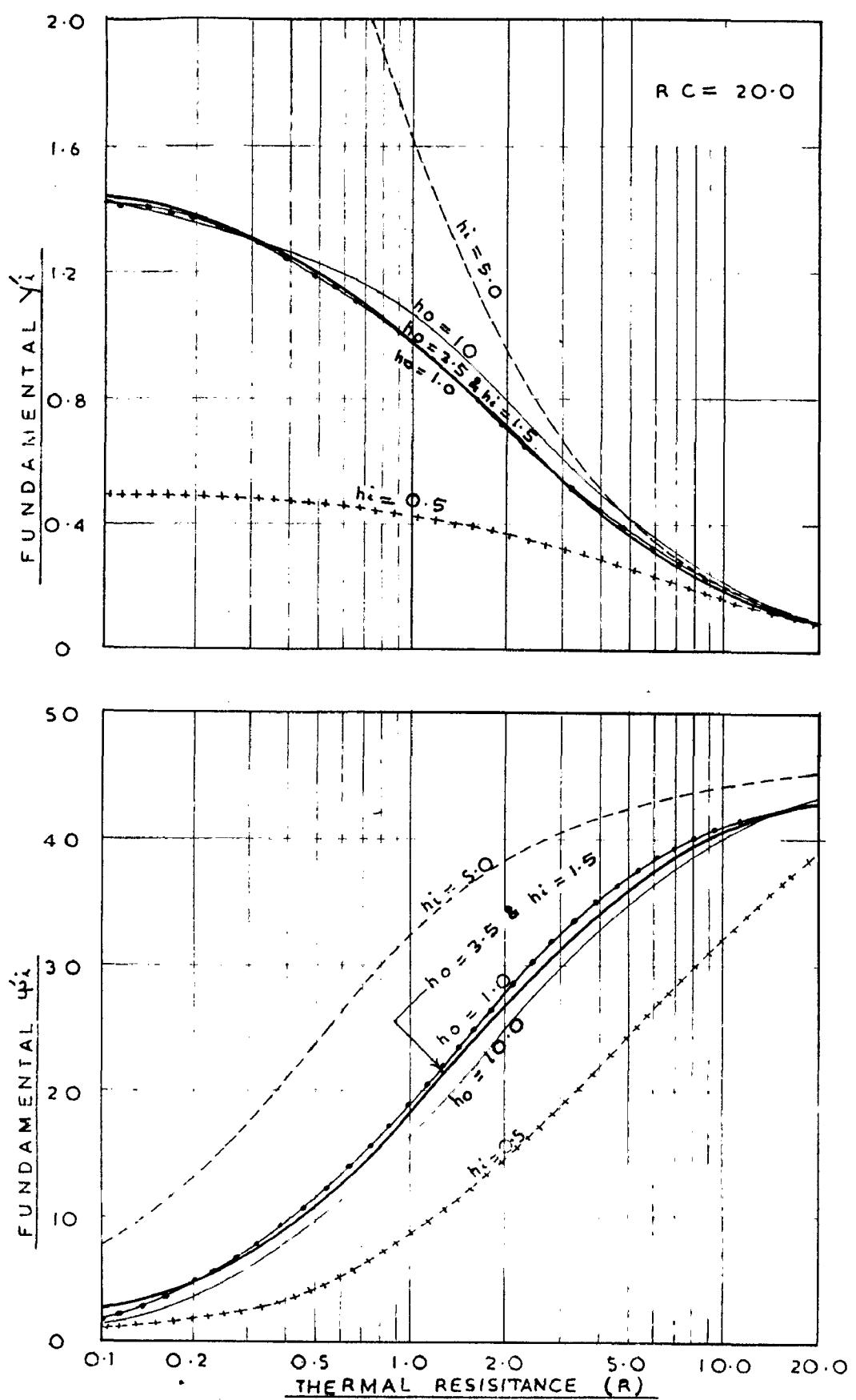


FIG.6.18 EFFECT OF  $h_o$  AND  $h_i$  ON  $Y'_i$  AND  $\Psi'_i$  FOR HOMOGENEOUS SECTIONS WITH MEDIUM TIME CONSTANT

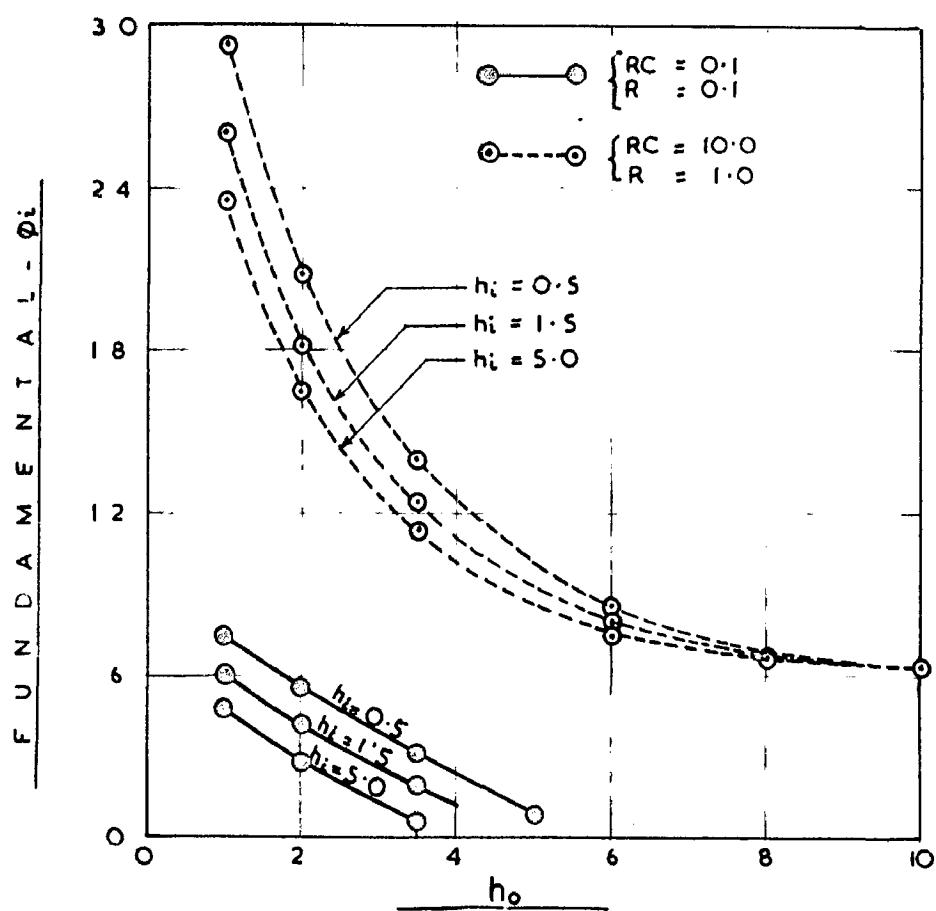
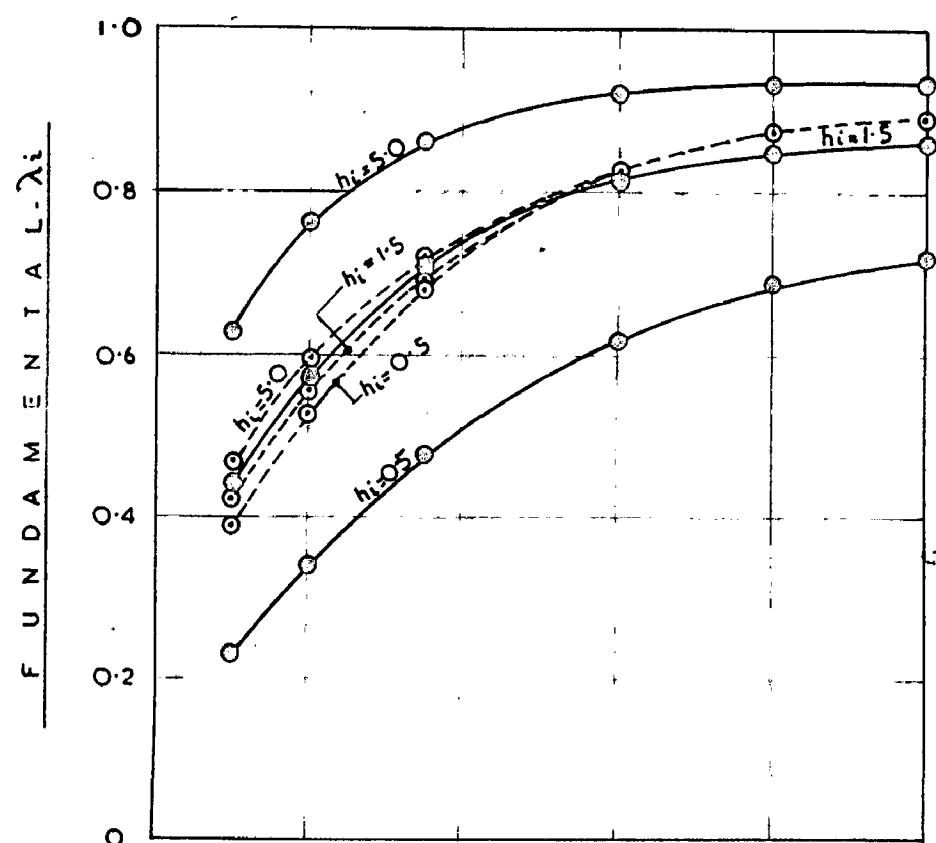


FIG.6.19 EFFECT OF  $h_o$  AND  $h_i$  ON  $\lambda_i$  AND  $\phi_i$  FOR THIN AND THICK SECTIONS OF DENSE MATERIALS

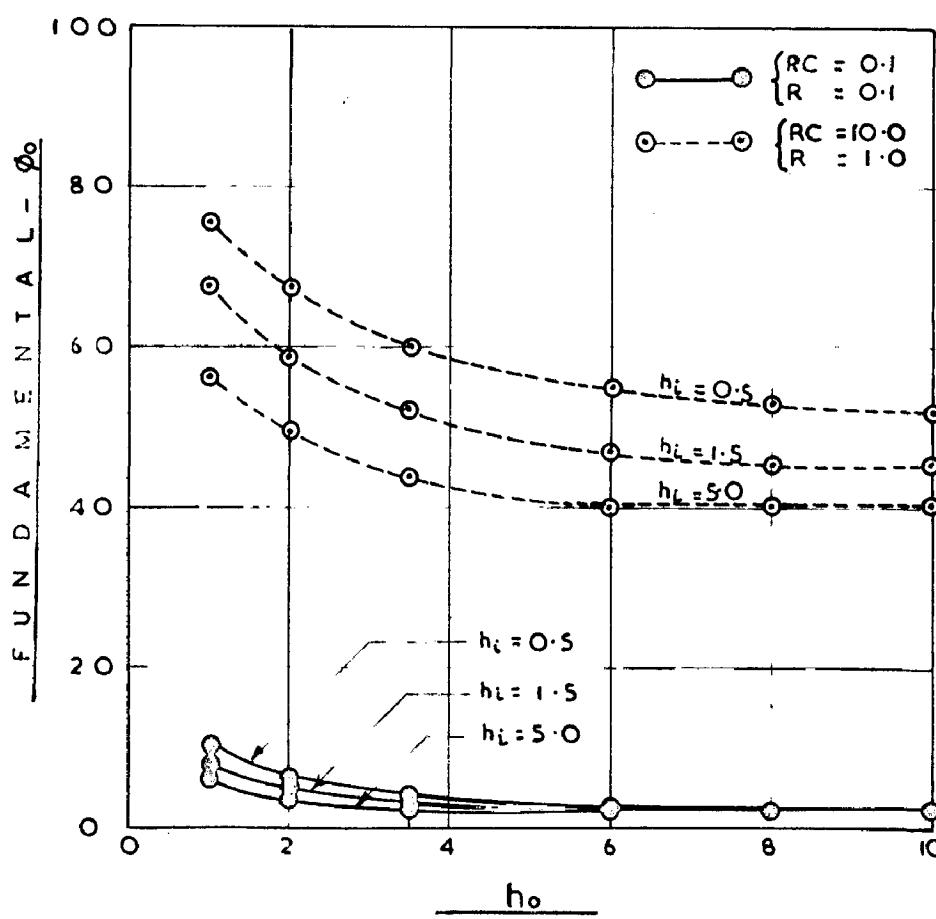
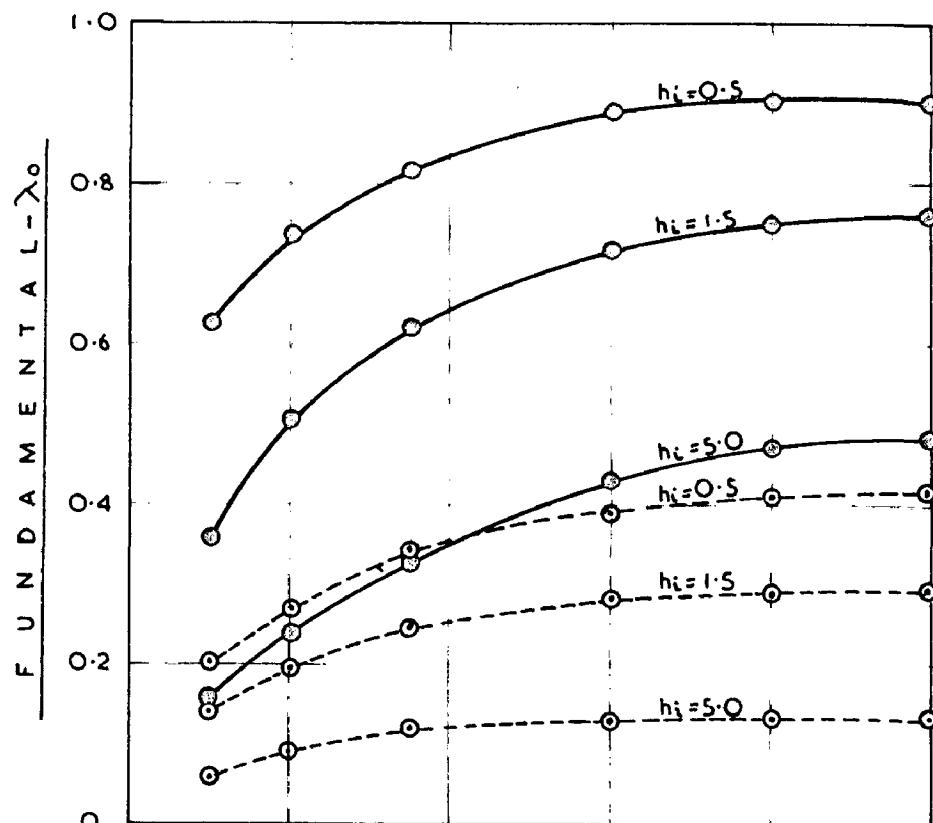


FIG.6.20 EFFECT OF  $h_o$  AND  $h_i$  ON  $\lambda_0$  AND  $\phi_0$  FOR THIN AND THICK SECTIONS OF DENSE MATERIALS

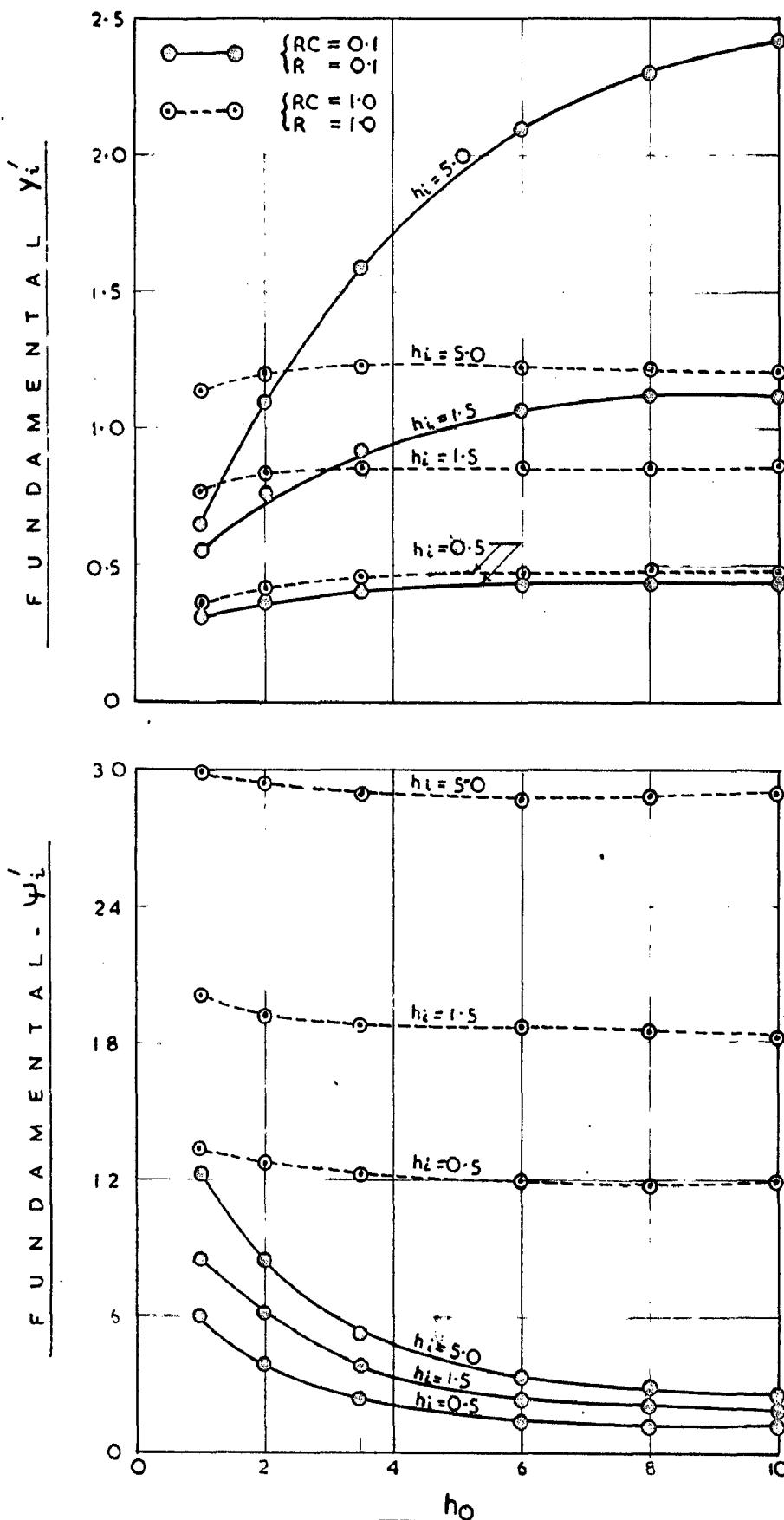


FIG. 6.21 EFFECT OF  $h_o$  AND  $h_i$  ON  $\gamma'_i$  AND  $\psi'_i$  FOR THIN AND THICK SECTIONS OF DENSE MATERIALS

- 2) Except for thin sections (which have small RC as well as R) the influence of the variation of ' $h_o$ ' on  $\lambda_i \angle \phi_i$  and of ' $h_i$ ' on  $\lambda_o \angle \phi_o$  is not significant.  $\lambda_o \angle \phi_o$  are influenced greatly by ' $h_o$ ' while  $\lambda_i \angle \phi_i$  are influenced ' $h_i$ ' considerably.
- 3) For thin sections increase of ' $h_o$ ' increases  $\lambda_o$  and  $\lambda_i'$  and decreases  $\phi_o$  and  $\phi_i'$ , whereas the increase of ' $h_i$ ' decreases  $\lambda_o$ ,  $\phi_o$  and  $\phi_i'$  and increases  $\lambda_i'$ .
- 4) For any given value of RC (small or large) both ' $h_o$ ' and ' $h_i$ ' have considerable effect on the transfer functions ( $\lambda_i \angle \phi_i$ ). Increase of ' $h_o$ ' increases  $\lambda_i$  and decreases  $\phi_i$ , whereas increase of ' $h_i$ ' decreases  $\lambda_i$  as well as  $\phi_i$ .
- 5) For building sections with large R and even with small RC, viz., insulating materials, ' $h_o$ ' and ' $h_i$ ' have no significant effect on thermal system functions.
- 6) For a building section with a given RC and a given ' $h_o$ ' and ' $h_i$ ',  $\lambda_i$  will be maximum for a particular resistance. For either an increase or decrease of the resistance from this optimum  $\lambda_i$  will decrease. This optimum resistance depends upon the RC value.

: 63 :

For higher values of RC, this R will also be higher. For a given C, the occurrence of maximum of  $\lambda_i$  shifts towards the lower resistance with the increase of ' $h_o$ ' as well as ' $b_i$ '.  
\_\_\_\_\_

## **C H A P T E R   7**

**THermal SYSTEM (TRANSFER AND DRIVING POINT)**

**FUNCTIONS OF COMPOSITE CONSTRUCTIONS**

## CHAPTER 7

### THERMAL SYSTEM (TRANSFER AND DRIVING POINT) FUNCTIONS OF COMPOSITE CONSTRUCTIONS

#### 7.1 Introduction

In the previous chapter (6) thermal functions for homogeneous constructions were dealt with. However, structural elements used in most of the present day buildings are composite (multi-layered). These composite constructions may consist of two or more layers of homogeneous materials in perfect thermal contact. There can also be a large number of combinations of layers made of different types of materials and thicknesses. In order to predict the thermal behaviour of such structural elements, the thermal functions of a wide range of composite constructions that are commonly met with in building practice are required.

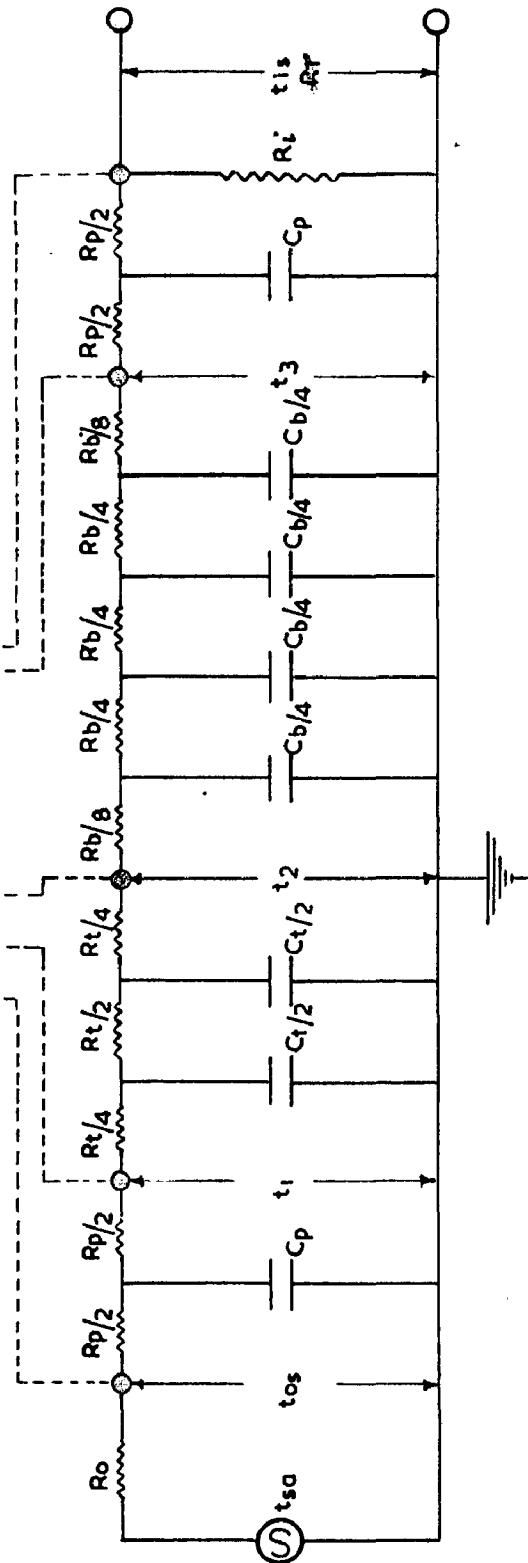
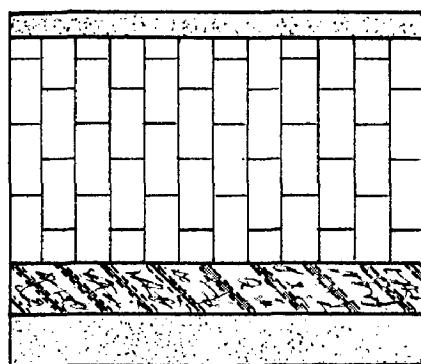
The type and thickness of each layer and the order of arrangement affect the thermal behaviour of composite elements under periodic heat flow conditions. Hence to make the most efficient use of the new as well as conventional materials, individually or in combination, the influence of the type and location

of the individual layers, on the overall thermal functions is to be known.

Analytical methods developed by Mackey and Wright (75) are complex even for homogeneous materials and become much more so for composite constructions. Other mathematical methods also developed by Skhloper (26) Marmet (37), no doubt, lessen the mathematical burden, but still are unwieldy. The matrix method (24) is more convenient, but when the number of layers is more than two, manual computation becomes impractical and a digital computer will be required, especially when the decrement factors of individual layers are also required to be determined. In contrast, the Electrical Analogue method requires only a few more measurements of the temperature amplitudes and phase lag angles at junction points of different layers. A typical composite construction and its equivalent network representation is shown in Fig. (7.1). The junction points at which temperature and phase measurements are made are indicated therein. For the representation of the composite element on the electrical network, the number of lumps employed for each layer were chosen according to the recommendations made in Chapter 4. The procedure, for the determination of the thermal transfer and driving point functions, is the same as that followed for homogeneous materials (described in Chapter 3).

COMPOSITE WALL

1 Inch Plaster + 1 Inch Thermocole  
+ 4  $\frac{1}{2}$  Inch Brick + 1/2 Inch Plaster



ANALOGUE REPRESENTATION OF A COMPOSITE WALL

FIG. 7.1

The thermal decrement factors of individual layers were obtained from the measurements of amplitude across them. For example for a four layer composite element (shown in Fig. 7.1) the decrement factor of the individual layers may be written as

$$\vec{\lambda}_0 = \frac{\vec{t}_{0s}}{\vec{t}_{0a}} \quad \vec{\lambda}_1 = \frac{\vec{t}_1}{\vec{t}_{0s}}$$

$$\vec{\lambda}_2 = \frac{\vec{t}_2}{\vec{t}_1} \quad \vec{\lambda}_3 = \frac{\vec{t}_3}{\vec{t}_2}$$

$$\vec{\lambda}_4 = \frac{\vec{t}_{is}}{\vec{t}_3}$$

The over all decrement factor (i.e. transfer function) of the whole building element is then obtained by the product of all the decrement factors of individual layers. Hence  $\lambda_1 \angle -\phi_1$  is given as

$$\lambda_1 \angle -\phi_1 = (\lambda_0 \angle -\phi_0)(\lambda_1 \angle -\phi_1)(\lambda_2 \angle -\phi_2)(\lambda_3 \angle -\phi_3)(\lambda_4 \angle -\phi_4) \quad ..(1)$$

From the principles of vector multiplication, we get

$$\lambda_1 = \lambda_0 \times \lambda_1 \times \lambda_2 \times \lambda_3 \times \lambda_4 \quad ..(2)$$

$$\text{and } \phi_1 = \phi_0 + \phi_1 + \phi_2 + \phi_3 + \phi_4 \quad ..(3)$$

The overall transfer and driving point functions for a large number of composite constructions were determined for fundamental and three higher harmonics, analogically. These alongwith 'S' values are given in Appendix (IV) Tables (1 to 6) with

illustrative diagrams. For the sake of convenience the above data is arranged in different groups in accordance with their functional use in a building.

It is well known that surface heat transfer coefficients vary with the inclination of the surface (i.e. vertical and horizontal surfaces will have different values). Since no precise data on these variations is available, it is the common practice to treat them as equal for most building calculations. With this simplifying assumption, the data (in Appendix IV), for a given construction used as a wall, may be used, without undue error either for roof or for walls.

The surface heat transfer coefficients used in the determination of the thermal functions are :

- 1) For external building elements

$$h_o = 3.5 \text{ Btu}/\text{sq.ft.}/\text{hr}/{}^{\circ}\text{F}$$

$$h_i = 1.5 \text{ " " " "$$

and      2) For internal structural elements

$$h_o = 1.5 \text{ Btu}/\text{sq.ft.}/\text{hr}/{}^{\circ}\text{F}$$

$$h_i = 1.5 \text{ " " " "$$

For the ground floor  $h_i$  is taken as 1.5 and the other surface (at a depth of 30 inches) is assumed to be at a constant temperature. This assumption

is based on experimental evidence (60).

For the calculation of thermal resistance ( $R$ ) and thermal capacity ( $C$ ) of individual layers, the average values of the physical properties of the materials listed in Table 1 of Appendix III were used. An examination of the thermal function data given in Appendix IV brings out the following points.

7.3

### $\lambda_o/\phi_o$ - External Driving Point Function

$\lambda_o$  and  $\phi_o$  mainly depend upon the type of the material and thickness of the upper most layer. If this layer is thick the preceding layers have no influence on them. When dense materials having large thermal absorption coefficient 'p' ( $\sqrt{\omega_e}$  for the top layer,  $\lambda_o$  will be low and  $\phi_o$  will be high while the reverse is true for light weight insulating materials (with small 'p').

With the increase of thickness of the top layer,  $\lambda_o$  decreases and  $\phi_o$  increases (slightly), for dense materials, while  $\lambda_o$  increases and  $\phi_o$  decreases for insulating materials. For higher frequencies,  $\lambda_o$  decreases and  $\phi_o$  increases, and this is more for dense materials than for light weight materials used as a top layer. The variations of  $\lambda_o$  and  $\phi_o$  are not large for different dense materials

as top layer (fundamental  $\lambda_o$  varied from 0.6 to 0.75 and  $\phi_o$  varied from 10 to 30°) while the variation between dense and insulating materials is quite large (fundamental  $\lambda_o$  varied from 0.6 to 0.95 and  $\phi_o$  varied from 1° to 30°). If the insulating layer forms the top layer, the phase lag angles will be very low (1° or less), whereas, if even a thin layer of a dense material like a plaster (which is usually the case) is placed above it,  $\phi_o$  increases upto 12 to 15°, while there is no appreciable change in  $\lambda_o$ . Intermediate floors and partition walls etc. where the surface heat transfer coefficients are equal (i.e.,  $h_o = h_1$ ) are of dense materials,  $\lambda_o$  and  $\phi_o$  vary within narrow limits (fundamental  $\lambda_o$  varied from 0.4° to 0.45° and  $\phi_o$  varied from 25° to 30°).

7.3

### $\lambda'_i / -\phi'_i$ - Internal Driving Point Function

The magnitude of  $\lambda'_i$  and  $\phi'_i$  is mainly determined by the inner most layer (as the heat flow direction is from inside to outside). These are influenced by the type of material and thickness of the inside layer in the same manner as the top layer affects  $\lambda_o$  and  $\phi_o$ .  $\lambda'_i$  is always lower than the corresponding  $\lambda_o$  and  $\phi'_i$  is greater than the corresponding  $\phi_o$ . For most of the conventional

types of structural elements (Bricks, stones, concrete etc.)  $\lambda_i$  varied within narrow limits, i.e., 0.42 to 0.52 for fundamental. With insulating layers on the inner side this varied from 0.7 to 0.94, depending upon the type of the insulating material and its thickness.

For partitions and intermediate floors  $\lambda_i / \phi_i$  will be equal to  $\lambda_o / \phi_o$  as both surfaces are interior ones and have equal surface coefficients (i.e.,  $h_o = h_i$ ).

#### 7.4 $\lambda_i / \phi_i$ - Transfer Function

As,  $\lambda_i$  and  $\phi_i$  give the overall decrement factor and phase lag of the entire composite element, these will depend upon the nature of all the layers. It is further noted that these quantities ( $\lambda_i$  and  $\phi_i$ ) are affected, not only by the type of the materials and thickness of the individual layers, but also on the arrangement of these layers (i.e., the order of the layers viewed in the direction of heat flow).

For the composite constructions studied, the fundamental  $\lambda_i$  varied from 0.3 to 0.5 for conventional types (i.e., combinations of brick, concrete, stone, etc. of different thicknesses) and  $\phi_i$  varied

from  $50^{\circ}$  to  $220^{\circ}$ . With the inclusion of insulating layers  $\lambda_i$  decrease considerably (as low as 0.009) even for smaller over all thickness.

A study of all the factors, that affect  $\lambda_i$  and  $\phi_i$  and the decrement factors of individual layers, will provide an insight of the mechanism of heat flow through composite constructions under periodic variations. These were studied and the results are given in Tables (7.1 through 7.6). A discussion, of the above results, follows in the next few pages.

2.5 Effect of the Location of the Insulating Layer on its Decrement Factor and on the Overall Thermal Functions

In this, a study of the effect of location of a highly insulating material (Thermocole) as outer, middle and inner layer and also on either side, in combination with a dense material (concrete) has been made. The decrement factors of the insulating layers and, the concrete layer are given in Table (7.1), along with the transfer and driving point functions.

It is seen from the table, that distributing the insulating layers on both sides of a dense material gives the lowest  $\lambda_i$  and ~~highest~~ phase lag ( $\phi_i$ ). But in this case  $\lambda_i$  will be quite high.

INFLUENCE OF THE LOCATION OF FLOOR-TRUSS LAYER ON INTERNAL DRIVING FORCES

1. Internal Driving Forces!		External Factors!		Internal Driving Forces!		External Factors!		Internal Driving Forces!		External Factors!	
Type of construction	Point Punction	Top	Bottom	Top	Bottom	Top	Bottom	Top	Bottom	Top	Bottom
1. 3" thoracoole + 4" denso concreto -φirg.	λ <sub>red.</sub>	0.839	0.059	...	0.649	0.031	0.541	0.041	0.009	0.041	0.033
2. 4" denso concreto + 2" thoracoole -φirg.	λ <sub>red.</sub>	0.735	0.913	...	0.086	0.044	0.055	0.035	0.009	0.041	0.033
3. 1" thoracoole + 4" denso concreto -φirg.	λ <sub>red.</sub>	0.848	0.099	0.346	0.117	0.225	0.041	0.009	0.041	0.033	0.033
4. 2" denso concreto + 2" thoracoole + 2" denso concreto -φirg.	λ <sub>red.</sub>	0.921	0.971	0.082	0.225	0.60	0.350	0.041	0.009	0.041	0.033
		10°	24°	29°	10°	10°	35°	6°	65°	32°	32°
		69°	69°	69°	69°	80°	65°	65°	65°	65°	65°
		39°	39°	39°	39°	39°	39°	39°	39°	39°	39°
		10°	10°	10°	10°	10°	10°	10°	10°	10°	10°
		64°	64°	64°	64°	64°	64°	64°	64°	64°	64°
		54°	54°	54°	54°	54°	54°	54°	54°	54°	54°
		34°	34°	34°	34°	34°	34°	34°	34°	34°	34°
		24°	24°	24°	24°	24°	24°	24°	24°	24°	24°
		14°	14°	14°	14°	14°	14°	14°	14°	14°	14°
		4°	4°	4°	4°	4°	4°	4°	4°	4°	4°

$$h_0 = 3.6 \\ h_1 = 1.5$$

The next best arrangement is obtained by placing all the insulation on the external side, where both  $\lambda_1$  and  $\lambda'_1$  are sufficiently low. The least advantageous way of utilising an insulating layer is to use it as an inside layer, if the main task is to prevent the entry of external heat, into the enclosure (for hot climates). On the other hand, if the main problem is to maintain higher indoor air temperatures and prevention of heat losses (for cold climates) placing the insulating layer on the inner side will be most advantageous.

If the insulating material is backed by a dense material the decrement factor ( $\lambda$ ) of the insulating layer, decreases slightly, whereas the phase lag angle increases considerably. When the dense material is backed by an insulating layer the decrement factor ( $\lambda$ ) of the dense material increases considerably, but the phase lag angle increases only slightly. For various possible arrangements, the greatest damping occurs when layers with small 'p' (insulating materials) alternate with layers of large 'p' (dense materials). In a two layer structure, when a layer with large 'p' is placed inside the transfer and internal driving point functions will be lower, though the external surface temperatures will be much higher, than for the reverse order.

### 7.6 Effect of the Back Layer on the Decrement Factor of an Insulating Upper Layer

In the above study it was found that placing insulation on the external side will provide maximum advantage. The decrement factor of the insulating layer depends upon the type of layer on which it is placed. The effect of the backing layer on the decrement factor of the insulating layer was studied. For this purpose a 3 inch thick Thermocole was taken as the outside layer and the following were used as backing layers :-

1. Thin asbestos cement sheet.
2. Dense materials - concrete and brick.
3. Insulating materials - foamed concrete and celotex board.

The results are shown in Table (7.2). These indicate that the decrement factor and phase lag across a top insulating layer are considerably affected by the type and thickness of the backing layer. With backing materials, are of dense type (large  $\rho$  value) the decrement factor of the insulating layer decreases and phase lag increases to a large extent. On the other hand with materials having small  $\rho$  value,  $\lambda$  increases considerably but not the phase lag ( $\phi$ ).

### 7.7 Effect of the Boundary layers on the Enclosed Air Space of Double wall Constructions

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which enclose air spaces are coming into use in building practice of this country. In order to determine the effect of the bounding layers on the damping across the air space bounded by different types of layers, the decrement factors and thermal functions were determined with different bounding layers. These are given in Table (7.3). It can be seen that the decrement factor and phase lag across an air space are affected to a considerable extent by the inner layer only but not by the outer one. If the inner layer is a thin sheet like, glass pane, galvanised iron sheet, and plywood sheet, practically no phase lag is introduced across the air space. If the backing layer is of a dense material like brick or concrato; sufficiently large phase lag angles (of the order of  $20^\circ$  to  $25^\circ$ ) are introduced. If the insulating materials form the backing layer, the decrement factor across the air space is increased ( $\lambda_{air}$  is increased) and only small phase lag angles ( $5^\circ$  or so) results.

The air space affects the decrement factor of the front bounding layer adversely (i.e.,  $\lambda$  increases) but increases its phase lag (by 10 to  $12^\circ$ ).

#### 7.8 Effect of the Order of Layers on the Thermal Functions

It was pointed out earlier that the

2.4DIL<sub>12</sub> 7.3

ASPECT OF BOUNDING LINES OF THE DUST COLLECTOR OR THE ENCLOSURE AND THE VACUUM SYSTEM (undecorated)

No.	Type of Construction	Decrement Factors of Overall Transfor- mation	External point function	Front Layer	Layer	space	Layer	Back	space + air + glass + 1/8"	0.478	0.321	0.740	0.556	0.852	0.493	0.220	0.595	0.490	0.043	0.40	0.934	0.683	0.99		
1.	Plywood	1/8" Glass + 2" air space + A mod.	0.896	0.10	0.10	0.478	0.321	0.974	0.767	0.167	0.167	0.740	0.556	0.852	0.493	0.220	0.595	0.490	0.043	0.40	0.934	0.683	0.99		
2.	Plywood	1/8" Glass + 2" air space + A mod.	0.890	0.10	0.10	0.795	0.581	0.636	0.30	0.10	0.10	0.740	0.556	0.852	0.493	0.220	0.595	0.490	0.043	0.40	0.934	0.683	0.99		
3.	Plywood	+ 1/8" colotek board + 1/8" concrete + 2" air space + A mod.	0.921	0.10	0.10	0.974	0.767	0.321	0.167	0.167	0.167	0.740	0.556	0.852	0.493	0.220	0.595	0.490	0.043	0.40	0.934	0.683	0.99		
4.	Plywood	space + 1/8" dense concrete - phi 16. phi 16. - 1/8" concrete + 2" air space + A mod.	0.910	0.70	0.70	0.776	0.710	0.556	0.40	0.220	0.220	0.556	0.40	0.539	0.40	0.20	0.250	0.280	0.490	0.043	0.40	0.934	0.683	0.99	
5.	Plywood	3" brick + 2" air space + A mod.	0.776	0.30	0.30	0.578	0.518	0.545	0.40	0.20	0.20	0.578	0.40	0.539	0.40	0.20	0.250	0.280	0.490	0.043	0.40	0.934	0.683	0.99	
6.	Plywood	4" brick + 2" air space + A mod.	0.776	0.70	0.70	0.578	0.518	0.490	0.40	0.20	0.20	0.578	0.40	0.539	0.40	0.20	0.250	0.280	0.490	0.043	0.40	0.934	0.683	0.99	
7.	Plywood	+ 2" wood board + 3" air space + A mod.	0.833	0.20	0.20	0.881	0.099	0.108	0.10	0.20	0.20	0.881	0.099	0.108	0.10	0.20	0.20	0.250	0.280	0.490	0.043	0.40	0.934	0.683	0.99
8.	Plywood	9" brick + 3" air space + A mod.	0.719	0.20	0.20	0.881	0.099	0.108	0.10	0.20	0.20	0.881	0.099	0.108	0.10	0.20	0.20	0.250	0.280	0.490	0.043	0.40	0.934	0.683	0.99

arrangement of the layers and their order (in the direction of heat flow) influence the overall thermal behaviour of the structural element. This effect was studied. The results, given in Table (7.4), clearly illustrate this effect on the transfer and driving point functions. If, insulating layers are also included the maximum damping occurs when the dense materials are used on the inner side. The decrement factors of the dense materials effect adversely when these are used as the outer layers.

#### 7.9 Effect of the Direction of Heat flow (Outside to Inside or Inside to Outside)

It is also interesting to know for a given composite construction, if the direction of the heat flow is reversed, how the decrement factors of individual layers and the overall thermal functions are affected. This has been studied for a half a dozen types of constructions and the results are given in Table (7.5). It is apparent that with the reversal of the heat flow directions (i.e., from inside to outside) for any given construction, the overall transfer function will be reduced by a ratio of  $h_0/h_1$  and the phase angle is not altered just as in the case of homogeneous elements.  $\lambda'_i$  is lower than the corresponding  $\lambda_o$  while  $\phi'_i$  is higher than  $\phi_o$ . The backing layers have considerable influence on the

TABLE 7.4

EFFECT OF THE ORDER OF LAYERS ON THE THERMAL SYSTEM FUNCTIONS (Fundamental)

$$h_0 = 3.5$$

$$h_1 = 1.5$$

S. No.	Type of Construction	External Driving point Function	Transfer Function	Internal Driving point Function
A. 1.	$\frac{1}{2}"$ plaster + $4\frac{1}{2}"$ brick $\lambda_{\text{Mod.}}$ 0.778 + 1" F.C. + $\frac{1}{2}"$ plaster - $\phi_{\text{A.R.G.}}$ 17°	0.159 75°	0.722 13°	
2.	$\frac{1}{2}"$ plaster + $4\frac{1}{2}"$ brick $\lambda_{\text{Mod.}}$ 0.756 + 2" F.C. $\frac{1}{2}"$ plaster - $\phi_{\text{A.R.G.}}$ 31°	0.099 94°	0.776 15°	
3.	$\frac{1}{2}"$ plaster + $4\frac{1}{2}"$ brick $\lambda_{\text{Mod.}}$ 0.766 + "4" F.C. + $\frac{1}{2}"$ plaster - $\phi_{\text{A.R.G.}}$ 22°	0.047 120°	0.822 22°	
4.	$\frac{1}{2}"$ plaster + $4\frac{1}{2}"$ brick $\lambda_{\text{Mod.}}$ 0.734 + 1" F.C. + 1" plaster - $\phi_{\text{A.R.G.}}$ 34°	0.055 87°	0.823 18°	
5.	$\frac{1}{2}"$ plaster + 4" dense Conc. + 2" F.C. $\frac{1}{2}"$ p - $\phi_{\text{A.R.G.}}$ 33°	0.113 33°	0.782 14°	
6.	$\frac{1}{2}"$ plaster + 2" dense Conc. + 2" F.C. + $\frac{1}{2}"$ p - $\phi_{\text{A.R.G.}}$ 30°	0.136 60°	0.718 14°	
B.	<u>Reverse Order</u>			
1.	$\frac{1}{2}"$ plaster + 1" F.C. + $\lambda_{\text{Mod.}}$ 0.900 $4\frac{1}{2}"$ brick + $\frac{1}{2}"$ plaster - $\phi_{\text{A.R.G.}}$ 16°	0.112 34°	0.576 30°	
2.	$\frac{1}{2}"$ plaster + 3" F.C. + $\lambda_{\text{Mod.}}$ 0.934 $4\frac{1}{2}"$ brick + $\frac{1}{2}"$ plaster - $\phi_{\text{A.R.G.}}$ 12°	0.073 100°	0.522 23°	
3.	$\frac{1}{2}"$ plaster + 4" F.C. + $\lambda_{\text{Mod.}}$ 0.950 $4\frac{1}{2}"$ brick + $\frac{1}{2}"$ plaster - $\phi_{\text{A.R.G.}}$ 7°	0.034 128°	0.533 27°	
4.	1" plaster + 1" F.C. + $\lambda_{\text{Mod.}}$ 0.934 $4\frac{1}{2}"$ brick + $\frac{1}{2}"$ plaster - $\phi_{\text{A.R.G.}}$ 10°	0.048 97°	0.570 24°	
5.	$\frac{1}{2}"$ plaster + 3" F.C. + $\lambda_{\text{Mod.}}$ 0.912 4" dense conc. + $\frac{1}{2}"$ p - $\phi_{\text{A.R.G.}}$ 14°	0.086 94°	0.573 36°	
6.	$\frac{1}{2}"$ plaster + 2" F.C. + $\lambda_{\text{Mod.}}$ 0.916 2" dense conc. + $\frac{1}{2}"$ p - $\phi_{\text{A.R.G.}}$ 7°	0.119 68°	0.573 32°	

F.C. = foamed concrete F.C. = thermocole p = plaster

TABLE 7.5

EFFECT OF THE DIRECTION OF HEAT FLOW ON THE  
THERMAL SYSTEM FUNCTIONS (Fundamental)

$$h_o = 3.5$$

$$h_i = 1.5$$

No.	Type of Construction	Heat flow from out- side to inside		Heat flow from inside to outside	
		Ext. Driv. transfer point function	Int. Driv. transfer point function	Reverse transfer function	Reverse transfer function
1.	1" p + 9" brick + 1" p - $\lambda_{\text{Mod.}} = 0.744$ $\phi_{\text{A.R.C.}} = 17^\circ$	0.123 $113^\circ$	0.500 $26^\circ$	0.053 $113^\circ$	
2.	1" p + 4½" brick + 1" p + 4" F.C. + 1" p - $\phi_{\text{A.R.C.}} = 23^\circ$	0.047 $120^\circ$	0.322 $24^\circ$	0.020 $120^\circ$	
3.	1" p + 4" D.C. + 1" p + 1" p - $\phi_{\text{A.R.C.}} = 36^\circ$	0.113 $36^\circ$	0.782 $14^\circ$	0.048 $86^\circ$	
4.	1" p + 1" T.C. + 1" p + 4½" brick + 1" p - $\phi_{\text{A.R.C.}} = 10^\circ$	0.048 $07^\circ$	0.570 $34^\circ$	0.031 $93^\circ$	
5.	1½" D.C. + 2" jax board + 4" brick + 1" p	$\lambda_{\text{Mod.}} = 0.934$ - $\phi_{\text{A.R.C.}} = 15^\circ$	0.039 $106^\circ$	0.511 $34^\circ$	0.018 $106^\circ$
6.	6" sandstone + 9" brick + 1" p - $\phi_{\text{A.R.C.}} = 18^\circ$	0.047 $170^\circ$	0.439 $23^\circ$	0.020 $170^\circ$	
7.	3" lime conc. + 4" R.C.C. + 1" p - $\phi_{\text{A.R.C.}} = 15^\circ$	0.200 $73^\circ$	0.423 $31^\circ$	0.036 $73^\circ$	
8.	2½" bricktilo + 3" lime conc. + 4½" brick + 1" p - $\phi_{\text{A.R.C.}} = 20^\circ$	0.037 $124^\circ$	0.439 $27^\circ$	0.037 $124^\circ$	

p = plaster F.C. = Formed concrete D.C. = Dense concrete  
 T.C. = Thermocolo R.C.C. = Reinforced cement concrete

decrement factors of the individual layers while front layers have no influence (viewed in the direction of heat flow). The decrement factor of the same layer is not same for either directions of heat flow. As a rule, the decrement factors for the same layer is lower, when the heat flows from inside to outside, than for the condition, of heat flow from outside to inside, if all the layers are of dense materials. If the insulating layers are also present the decrement factor of the layer located in front (viewed in the direction of heat flow) will have higher values.

#### 7.10 Damping Across Individual Layers as Affected by Their Location

It was noted earlier that a given layer will have different damping properties depending upon its position with respect to other layers. In order to obtain a better understanding of the behaviour of different types of layers, as affected by their location, the decrement factors were determined for a few types of materials, with different locations. These are given in Table (7.6). It is clear from the above data, that the inner most layer or a combination of two or more layers (in the same order) is not affected by the layers placed above them (for heat flow from outside to inside). On the other hand, the inner layers affect the decrement factors of the upper layer

TABLE 7-8

DECREMENTAL FACTOR OF INDIVIDUAL LAYERS AS AFFECTED BY THEIR LOCATION

 $h_0 = 3.5$  $h_1 = 1.5$ 

Sl.	Type of construction	Driving point function	Decrement Factor of Layers				Transfer function
			1st layer	2nd layer	3rd layer	4th layer	
A- BRICK							
1. 4" brick	Mod. - AFB - Φ	Mod. - 0.756 0.140	0.447 0.386	---	---	---	0.380 0.420
2. 4 1/2" brick + 3/4" plaster	Nod. - Φ FFB - Φ	0.756 0.120	0.445 0.350	0.342 0.10	---	---	0.316 0.475
3. 4 P. + 4 B. + 3/4" P. + 3/4" B. - Φ FFB - Φ	Nod. - 0.950 0.700	0.382 0.110	0.383 0.860	0.445 0.360	0.388 0.350	0.384 0.280	-
4. 1" P. + 2 1/2" C. + 3/4" B. + 3/4" P. - Φ FFB - Φ	Mod. - 0.934 0.100	0.383 0.20	0.125 0.420	0.445 0.350	0.340 0.310	0.348 0.300	-
5. 3" R. + 4 1/2" B. + 3/4" F. C. + 3/4" B. - Φ FFB - Φ	Nod. - 0.765 0.40	0.343 0.450	0.650 0.450	0.103 0.450	0.383 0.450	0.347 0.300	0.347 0.200
6. 3" P. + 4 1/2" B. + 3/4" F. C. + 3/4" B. - Φ FFB - Φ	Mod. - 0.734 0.30	0.315 0.380	0.739 0.350	0.125 0.360	0.384 0.380	0.355 0.370	-
B- CONCRETE							
1. 4" C-C - Φ	Mod. - 0.730 - Φ AFBB - Φ	0.552 0.230	0.552 0.300	0.545 0.300	0.343 0.10	---	0.403 0.400
2. 4" C-C - Φ	Mod. - 0.732 - Φ AFBB - Φ	0.545 0.300	0.343 0.10	---	---	---	0.376 0.430
3. 1" P. + 4" D-C - Φ Nod. - Φ FFB - Φ	Mod. - 0.732 0.40	0.362 0.320	0.342 0.10	---	---	---	0.360 0.540
4. 1" P. + 3" D-C - Φ Nod. - Φ FFB - Φ	Mod. - 0.612 0.10	0.382 0.10	0.186 0.530	0.345 0.300	0.340 0.260	0.386 0.340	0.386 0.340
5. 3" P. + 4" D-C - Φ Nod. - Φ FFB - Φ	Mod. - 0.772 0.30	0.682 0.30	0.793 0.290	0.205 0.230	0.383 0.350	0.313 0.260	0.313 0.260
6. 4" C-C - Φ	Mod. - 0.735 - Φ AFBB - Φ	0.618 0.230	0.618 0.10	0.368 0.10	0.344 0.10	0.344 0.10	0.344 0.10
C- STAINLESS CONCRETE							
1. 2" P-C. - Φ	Mod. - 0.700 - Φ AFBB - Φ	0.295 0.100	0.295 0.100	0.295 0.100	0.295 0.100	0.295 0.100	0.295 0.100
2. 2" P-C. + 2" P-C. - Φ	Mod. - 0.911 0.20	0.320 0.20	0.320 0.20	0.320 0.20	0.320 0.20	0.320 0.20	0.320 0.20
3. 3" P. + 2" P-C. - Φ	Mod. - 0.972 0.30	0.382 0.30	0.382 0.30	0.382 0.30	0.382 0.30	0.382 0.30	0.382 0.30
4. 1" 2" P. + 3" P-C. - Φ	Mod. - 0.912 0.20	0.312 0.20	0.312 0.20	0.312 0.20	0.312 0.20	0.312 0.20	0.312 0.20
5. 3" P. + 2" P-C. - Φ	Mod. - 0.780 0.20	0.322 0.20	0.322 0.20	0.322 0.20	0.322 0.20	0.322 0.20	0.322 0.20
D- PLASTIC							
1. 2" P-C. - Φ	Mod. - 0.780 - Φ AFBB - Φ	0.622 0.30	0.622 0.30	0.622 0.30	0.622 0.30	0.622 0.30	0.622 0.30
2. 1" I.D. Concrete Pipe - Φ	Mod. - 0.932 0.30	0.383 0.30	0.383 0.30	0.383 0.30	0.383 0.30	0.383 0.30	0.383 0.30
3. 1" I.D. Concrete Pipe - Φ	Mod. - 0.972 0.30	0.382 0.30	0.382 0.30	0.382 0.30	0.382 0.30	0.382 0.30	0.382 0.30
4. 1" 2" P. + 3" P-C. - Φ	Mod. - 0.912 0.20	0.312 0.20	0.312 0.20	0.312 0.20	0.312 0.20	0.312 0.20	0.312 0.20
5. 3" P. + 2" P-C. - Φ	Mod. - 0.780 0.20	0.322 0.20	0.322 0.20	0.322 0.20	0.322 0.20	0.322 0.20	0.322 0.20

C- 1. 2" P-C. -  
Φ

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0.384

to a considerable extent.

For example, a 4 $\frac{1}{2}$  in. brick wall alone has a decrement factor ( $\lambda$ ) of 0.467 and phase lag angle of 23°. When it is plastered inside ( $\frac{1}{2}$  inch),  $\lambda$  is decreased to 0.445 and  $\phi$  is increased to 35°. Those values did not change when additional plaster or other layers are placed on the external side. But when additional insulation layers are placed on the inner side the  $\lambda$  for the brick section is increased to 0.739 and  $\phi$  is increased to 41°. Other materials also behaved similarly.

#### 7.11 Equivalent Homogeneous Construction of Composite Structures

Though the transfer and driving point function data given in Appendix (IV) covers most of the commonly used composite constructions, there can be many more combinations possible. Any new specific case will have to be studied afresh. Moreover, the physical properties of the materials used in these studies are average values only. The variations in their properties should also be accounted for. This may be simplified by determining the equivalent homogeneous construction, which has the same decrement factor ( $\lambda_1$ ) and phase lag ( $\phi_1$ ) that of the composite construction. The two thermal properties

necessary to determine the equivalent homogeneous construction are

i) equivalent thermal resistance  $R_{eq}$  (eq)  
and ii) equivalent thermal capacitance  $C_{eq}$  (eq)

The equivalent resistance is obtained simply by adding all the resistances of the individual layers i.e.,  $\sum_{n=1}^{n=n} R_n$ . Mackey and Wright (78) have given an empirical relation for the determination of the equivalent  $(k\rho s)_{eq}$  for periodic heat flow conditions. Stewart (78) pointed out that it is difficult to establish limits for the empirical equations derived by Mackey and Wright. Bruckmayer (77) had also derived equations for equivalent capacity brick wall, based on Kirsch's equations (73) for free cooling of structures (transient response). Hofbauer (79) has derived equations for half value time for multilayered constructions and related it with Bruckmayer's equivalent capacity brick wall. The equivalent (RC) eq and hence the equivalent (C) eq can be obtained from these equations.

The greatest advantage of these equivalent homogeneous  $R_{eq}$  and  $C_{eq}$  is that the same set of charts of thermal functions prepared for homogeneous constructions, given in chapter 6, can be utilised for obtaining these functions for any composite constructions, whatever be the number, type and order of the

layers.

Skhloover (26) questioned the validity of those empirical relations and doubted the possibility of having an equivalent homogeneous construction which will under all circumstances, have the same damping and phase shift.

Although these methods yield approximate values, they are worth the effort as they enable a quick estimation of all the thermal functions and provide general solutions, provided the methods adopted for the calculation of the equivalent  $R_{eq}$  are reliable.

Mackay and Wright have checked the validity of their empirical relations by comparing the decrement factors for  $\lambda_i$  and  $\phi_i$  only with the analytical solutions. As the periodic heat flow characteristics of a building element are described by a set of three thermal functions, the validity of the equivalent homogeneous computational methods, should also be checked for the driving point functions. This check has been made for 10 types of composite constructions covering layers made of conventional and insulating materials. The equivalent  $R$  and  $IC$  for those constructions were calculated both by Mackay and Wright and Bruckmayers equations. The thermal functions were read from the respective charts (given in Chapter 6). Those were

compared with the corresponding functions determined by the analogue method. As a further check these thermal functions were computed by Matrix method also. In the matrix method the overall transmission matrix is obtained by the matrix multiplication of the transfer matrices of the individual layers, in the order of heat flow direction. The transfer functions and the driving point functions were determined with the equations given in Appendix (ii). The results obtained by the above four methods are compared in Table (7.7). From the table it can be seen that the analogue and Matrix values agree closely, as expected. However, results obtained by Hackney and Wright equations are only approximate and are useful only with layers of dense materials. The deviations become too large especially in phase lag angles when insulating layers are present and the method becomes unreliable. Druckmayers (77) equations also do not hold good for composite constructions which include insulating materials. From those studies it is apparent that as pointed out by Shiklover (23) it is not always possible to obtain an equivalent homogeneous construction, which gives the same thermal system functions, of a composite building element under all circumstances.

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TABLE 7.7

COMPARISON OF THE THERMAL SYSTEM FUNCTIONS OBTAINED BY EQUIVALENT  
HOMOGENEOUS, ANALOGUE AND MATRIX METHODS

(Fundamental)

$$h_0 = 3.5 \\ h_1 = 1.5$$

No.	Type of Construction	Thermal System Formula	Hockey and Bright	Bruckmeyer Formula	Analogous Method	Matrix Method		
		func-tion	Modu-lus	Arg. 'in deg'	Modu-lus	Arg. 'in deg'	Modu-lus	Arg. 'in deg'
1.	$\frac{1}{2}'' p + 9'' B + \frac{1}{2}'' p$	0.730 0.123 0.530	15 103 21	0.710 0.125 0.420	16 103 35	0.744 0.123 0.509	17 113 23	0.733 0.133 0.522
2.	$3'' D_o + 2'' air space + 3'' D$	0.805 0.170 0.625	11 63 19	0.810 0.105 0.600	11 73 19	0.773 0.175 0.635	10 73 22	0.781 0.179 0.634
3.	$6'' sand steno + 9'' brick + \frac{1}{2}'' plaster$	0.720 0.060 0.480	16 150 33	0.690 0.036 0.430	16 176 26	0.634 0.037 0.439	18 170 23	0.642 0.045 0.468
4.	$3'' L.C. + 4'' R.B. + \frac{1}{2}'' p$	0.740 0.175 0.635	15 24 24	0.740 0.170 0.520	15 24 24	0.733 0.162 0.489	13 73 22	0.719 0.184 0.500
5.	$\frac{1}{2}'' p + 8'' o.s.c. + \frac{1}{2}'' p$	0.850 0.061 0.670	8 118 16	0.820 0.059 0.700	8 125 16	0.839 0.035 0.639	11 115 23	0.865 0.064 0.670
6.	$\frac{1}{2}'' p + 4\frac{1}{2}'' B. + 4'' P.C. + \frac{1}{2}'' p$	0.970 0.027 0.740	7 155 15	0.910 0.055 0.800	6 93 11	0.766 0.047 0.822	22 127 22	0.736 0.040 0.812
7.	$1'' p + 1'' T.C. + 4\frac{1}{2}'' B. + \frac{1}{2}'' p$	0.830 0.037 0.763	7 137 13	0.840 0.018 0.690	8 100 10	0.924 0.043 0.670	10 90 34	0.933 0.046 0.533
8.	$2'' T.C. + 4'' D.C.$	0.930 0.013 0.795	5 150 11	0.890 0.005 0.760	7 235 13	0.269 0.031 0.544	1 69 39	0.959 0.034 0.540
9.	$2'' D.C. + 2'' T.C. + 2'' J.J.$	0.850 0.023 0.860	4 112 8	0.820 0.015 0.820	5 180 10	0.821 0.049 0.753	16 65 22	0.913 0.041 0.760
10.	$\frac{1}{2}'' C sheet + 1'' M.W. + \frac{1}{2}'' p.v.$	0.900 0.187 0.810	3 14 4	0.800 0.128 0.780	3 22 8	0.944 0.193 0.844	2 21 8	0.943 0.193 0.864

? = plaster p.v. = plywood D.C. = Dense concrete L.C. = Lime concrete R.B. = Reinforced brick  
o.s.c. = Expanded slag concrete F.C. = faced concrete B. = brick M.W. = mineral wool  
T.C. = Thoracolo a.c. = asbestos cement.

## **C H A P T E R   8**

### **A METHOD FOR THE PREDICTION OF INDOOR AIR TEMPERATURES OF ENCLOSURES**

## CHAPTER R

### A METHOD FOR THE PREDICTION OF INDOOR AIR TEMPERATURE OF ENCLOSURES

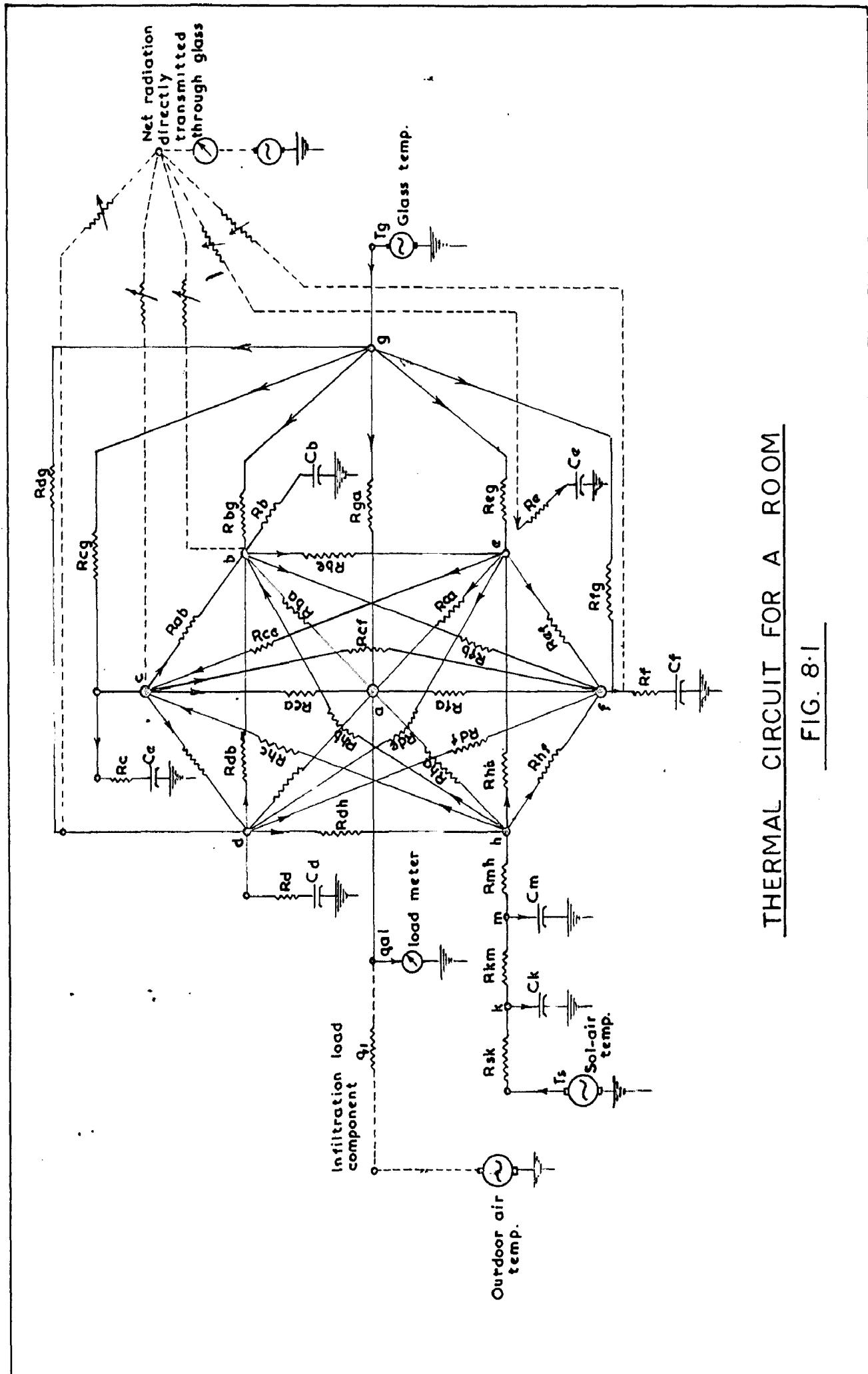
#### 8.1 Introduction

Chapters 6 and 7 deal with the transfer and driving point functions of a large number of common building elements (homogeneous and composite) which provide a quick means of comparing their thermal efficiencies. Criteria for the choice of a building component mostly consists of (i) Low 'U' values, (ii) Low transfer function ( $\lambda_1$ ) value, and (iii) Low internal driving point transfer function ( $\lambda'_1$ ). Such a choice will result in lower cooling loads for air conditioning and better thermal behaviour of unconditioned buildings. For a quantitative estimation, of the indoor air temperature, where due allowance has to be made for fluctuations of climate and variations in design a computational method using the above data is required.

Mathematical methods have been developed by Muney (25) Shklover (26) and Marmott (27). Muney has shown that the mathematical treatment of van Gorcum (28) could be extended to a complete building. Shklover and Marmott, by applying the symbolic calculus techniques

have developed methods for the estimation of indoor air temperatures of unconditioned buildings. But these computational procedures are quite involved and time consuming. Electrical analogue techniques have also been employed (48, 80) for the determination of indoor air temperatures. In these methods, the room with all its structural components, was simulated on the analogue model, coupled with proper radiative and convective resistances at the inner surfaces. The network with all the inputs provided by specially designed function generators at various boundaries becomes complex experimentally and economically. One such network arrangement is illustrated in Fig. (8.1). Another disadvantage of this method is that any specific variation in any one factor had to be represented on the analogue and studied separately. Since one comes across, quite a large number of individual variations of climate and design, such a method of analogue representation is not feasible to attain generalised solutions and some simpler and more flexible methods are to be aimed at.

A computational method utilising the (pre-tabulated) analogue transfer and driving point functions is given here. As buildings are subjected to periodic temperature and solar radiation, the indoor air temperature variations of enclosures are also



THERMAL CIRCUIT FOR A ROOM

FIG. 8.1

periodic. These can be represented by Fourier series.

$$t_{ia}(T) = t_{ia}(\text{mean}) + \sum_{n=0}^{\infty} t_{ia,n} \cos(\omega_n t - \xi_n) \quad (1)$$

Determination of  $t_{ia}(T)$  consists of two parts viz., (i) determination of the steady state temperature ( $t_{ia}(\text{mean})$ ) and (ii) the determination of the harmonic components.

### 8.2 Determination of the Mean Indoor Air Temperature

The steady state heat flow ( $\dot{Q}_{\text{mean}}$ ) through any building component 'k' due to the temperature differences between the mean soil-air temperature ( $t_{sa}(\text{mean})$ ) and the mean indoor air temperature ( $t_{ia}(\text{mean})$ ) is calculated from the equation (2)

$$\dot{Q}_k(\text{mean}) = U_k A_k \{ t_{sa,k}(\text{mean}) - t_{ia}(\text{mean}) \} \quad ..(2)$$

where  $U_k$  is the steady state over all thermal transmission coefficient of the building element in Btu/sq.ft./hr/ $^{\circ}$ F.  $A_k$  is the area of the exposed surface of the building element in sq.ft. The total steady state heat flow into the interior of a building, through all exterior bounding surfaces is obtained by the summation of all the individual quantities and is expressed by

$$Q(\text{mean})_{\text{Total}} = \sum A_k \text{ mean} = \sum U_k A_k \left\{ t_{sa,k}(\text{mean}) - t_{ia}(\text{mean}) \right\} \quad .. (3)$$

In order to satisfy the steady state heat balance conditions of the enclosure, the net heat gains of the enclosure should be zero i.e., the total heat gains of the enclosure should be equal to the total heat losses from it. If heat gains through ventilation and internal sources are also present, their steady state components are to be added to the total heat gains of the enclosure. Then the equation (3) becomes

$$Q(\text{mean}) = U_k A_k \left\{ t_{sa,k}(\text{mean}) - t_{ia}(\text{mean}) \right\} + q_g(\text{mean}) + q_v(\text{mean}) \quad .. (4)$$

According to the conditions for heat balance the  $Q(\text{mean})$  is to be equated to zero. Then we get

$$U_k A_k \left\{ t_{sa,k}(\text{mean}) - t_{ia}(\text{mean}) \right\} + q_g(\text{mean}) + q_v(\text{mean}) = 0 \quad .. (5)$$

This gives the mean indoor air temperature as

$$t_{ia}(\text{mean}) = \frac{U_k A_k \left\{ t_{sa,k}(\text{mean}) \right\} + q_g(\text{mean}) + q_v(\text{mean})}{U_k A_k} \quad .. (6)$$

### 8.3 Determination of the Harmonic Components of Indoor Air Temperature Variation

The next step is to determine the fundamental and higher harmonics of the indoor air temperature variations, which result from the harmonic heat flow through all the bounding elements. This can be solved in two steps, viz., (i) calculation of the heat flux entering the enclosure through each element due to the sol-air temperature variations of the exposed surfaces and (ii) calculation of the amount of heat flux absorbed by each internal surface due to indoor air temperature harmonic variation with unit amplitude, for the fundamental and higher harmonics.

i) The sinusoidal heat flux transmitted into the enclosed space through any exposed element  $k$  can be calculated, if the transfer admittance function  $\gamma_{nk}$  is known, by the equation

$$\vec{q}_{nk}(\tau) = \vec{t}_{sa nk} \cdot A_k \vec{\gamma}_{nk} \quad .. (7)$$

where  $t_{sa nk}$  is the  $n$  th harmonic component of the sol-air temperature of the element  $k$ .  $\gamma_{nk}$  is the transfer admittance of the  $n$  th harmonic for the same element  $k$ .

$A_k$  is the exposed surface area of the element  $k$ .

$\vec{t}_{\text{sank}}$  and  $\vec{Y}_{\text{ink}}$  are vectorial quantities which are expressed in polar co-ordinate form by

$t_{\text{sank}} \angle \theta_n$ ,  $Y_{\text{ink}} \angle -\psi_{\text{ink}}$  respectively.

The total amount of heat flux entered into the enclosed space through all the exposed elements, for any harmonic, is obtained by summation of all the individual quantities of heat flux i.e.,

$$\vec{q}_n(T) = \sum \vec{q}_{nk}(T) = \sum \vec{t}_{\text{sank}} \angle \theta_n \quad \vec{Y}_{\text{ink}} \quad ..(8)$$

Solar-air temperature variations will be different for different components of a building (i.e. roof, walls, door, windows etc.) even though the outside air temperature is the same. Hence their harmonic components will also have different amplitudes and phases. The summation of the heat flux quantities in equation (8) are to be carried out as complex additions. If any internal periodic heat source or heat flow due to ventilation are present, the harmonic components of their heat flux are to be added to the corresponding harmonic heat flux entering through all the exposed elements to obtain the overall heat flux entering the enclosure. This may be written as

$$\vec{q}_n(T) = \sum \vec{q}_{nk}(T) + \vec{q}_{ns}(T) + \vec{q}_{nv}(T) \quad ..(9)$$

- ii) This heat flux received in the room,

will produce a temperature variation of the indoor air. Let the corresponding harmonic variation of the indoor air temperature be  $\vec{t}_{ian}$ . This harmonic temperature variation of the indoor air, will in turn induce harmonic heat flux in all the interior surfaces, including the partitions, furniture etc. The quantity of heat flux absorbed, at the inside surface by any building element  $k$  can be calculated if the internal driving point admittance function  $\vec{Y}'_{ink}$  is known. The equation relating these quantities is given by

$$\vec{q}_k(T) = \vec{t}_{ian} A_k \vec{Y}'_{ik} \quad ..(10)$$

The total heat flux absorbed by all the interior surfaces of the enclosure is then obtained by summation of the individual quantities. This is expressed as

$$\vec{q}_n(T) = \sum \vec{t}_{ian} A_k \vec{Y}'_{ink} \quad ..(11)$$

Since all the interior surfaces are in contact with the same air, the driving temperature  $t_{ian}$  will be same for all surfaces, and the equation (11) can be written as

$$\vec{q}_n(T) = \vec{t}_{ian} \sum A_k \vec{Y}'_{ink} \quad ..(12)$$

As the condition for heat balance should also be satisfied for each harmonic, the total amount of

heat flux transmitted into the enclosure and that emitted by the internal periodic source, should be equal to the total heat flux absorbed by all the interior surfaces. This means

$$q_n(T) = q'_n(T)$$

i.e.,  $\sum t_{\text{sink}} A_k Y_{ik} + q_{ns} + q_{nv} = t_{ian} \sum A_k Y'_{ik}$  ..(13)

Then  $t_{ian}$  is obtained as

$$t_{ian} = \frac{\sum t_{\text{sink}} A_k Y_{ik} + q_{ns} + q_{nv}}{\sum A_k Y'_{ik}} \quad \text{..(14)}$$

Thus for obtaining  $t_{ian}$ , transfer admittance functions and internal driving point admittance functions are required. These can be obtained from the transfer and driving point functions (amplitude decrement and phase lag angles of temperatures) by the following relations

$$\vec{Y}_i = h_i \vec{\lambda}_i \quad \text{.. (15)}$$

$$\overleftarrow{Y}'_i = h_i (1 - \overleftarrow{\lambda}'_i) \quad \text{.. (16)}$$

By combining the steady state temperature and the harmonic components, the Fourier equation for the periodic indoor air temperatures of any enclosure can be obtained as

$$t_{ia}(T) = t_{ia}(\text{mean}) + \sum t_{ian} \cos(\omega_n t - \xi_n) \quad \text{..(17)}$$

This method gives the indoor air temperature variations in terms of the thermal system functions,

'U' values and the sol-air temperatures. Any changes in the building components can be incorporated in the calculations taking its corresponding 'U' value and thermal system function, while changes in climatic and other factors like orientation, surface treatments etc. can also be incorporated by corresponding changes in sol-air temperature. In this method the radiation exchanges between the interior surfaces are not separately considered. The surface coefficients were taken as the resultant of radiative and convective transfer coefficients and the heat is directly transferred to the room air, from the interior surfaces.

#### 3.4 Verification of the Method

To check the reliability of this method, the predicted temperatures were compared with the model measurements of Muncoy (25). Muncoy had used two models (i) a masonry one with heavy structural elements like brick, and (ii) a timber one with light structural parts. These models were subjected to the same external air temperature variation. The inside air temperatures were measured and computed by Matrix method. The physical properties of the materials of the model are given in Table (3.1). The 'U' values and the thermal transfer and driving point functions of the materials used determined by the analogue are

TABLE B.1

PHYSICAL PROPERTIES OF THE MATERIALS USED  
IN THE MODELS

Material	K	P	s
Brick	8.0	120	0.20
Concrete	10.0	150	0.20
Timber	1.0	50	0.30
Plywood	1.0	50	0.30
Mineral Wool	0.25	12	0.20
Cane Fibre Board	0.40	30	0.30
Glass	7.3	156	0.19

= Thermal conductivity in Btu.in./ft<sup>2</sup>.hr.<sup>o</sup>F

= Density in Lb/Ft<sup>3</sup>

= Specific heat in Btu/Lb.<sup>o</sup>F

given in Table (3.2 and 3.3) for the two models respectively. As all the six sides were exposed to the same air temperature cycle, the mean inside air temperature of the models will be equal to the mean outside air temperature.

$$\text{i.e., } t_{ia}(\text{mean}) = t_{oa}(\text{mean})$$

The heat flux entering the model through each exterior element due to external temperature variations, for the fundamental and three higher harmonics were calculated from the equation (6) and given in column 'A' of tables (3.4 and 3.5). The heat flux absorbed by each internal surface per unit amplitude variation of internal air temperature for the fundamental and three higher harmonics were also calculated from the equation (9) and given in columns 'B' of the same tables (3.4 and 3.5). The summation of all items, in column 'A' will give the total heat flux entering the model

$$\text{i.e., } \sum t_{oank} A_k Y_{ink} = Q_n$$

and the summation of all the items in column 'B' will give the total heat flux absorbed by all interior surfaces of the model for unit amplitude harmonic variation of inside air temperature

$$\text{i.e., } \sum A_k Y_{ink} = \overleftarrow{Q'_n}$$

If the resultant indoor air temperature variation had an amplitude  $t_{ian}$  then the total heat flux absorbed

TABLE 8.2

THE HARMONIC FUNCTIONS AND '0' VALUES OF THE  
COMPONENTS OF THE HARMONIC MODEL

$$h_0 = 2.5 \\ h_1 = 1.4$$

S. No.	'Component'	'value'	'Har-' no-	Transfer Function	'Internal Driving point Function'	
			$\lambda_i$	'Modulus ' $\Delta_i$	'Arg. - $\phi_i$ ' 'in deg.'	'Modulus ' $\lambda_i$
1.	2" brick wall	0.735	F	0.056	132	0.122
			H <sub>2</sub>	0.018	162	0.091
			H <sub>3</sub>	0.008	200	0.076
			H <sub>4</sub>	0.005	218	0.069
2.	2" concrete roof	0.760	F	0.046	130	0.100
			H <sub>2</sub>	0.016	164	0.075
			H <sub>3</sub>	0.007	208	0.062
			H <sub>4</sub>	0.004	223	0.053
3.	1" timber floor	0.473	F	0.038	124	0.333
			H <sub>2</sub>	0.034	134	0.230
			H <sub>3</sub>	0.014	204	0.240
			H <sub>4</sub>	0.007	226	0.213
4.	2" brick partition wall	0.593	F	..	..	0.123
			H <sub>2</sub>	..	..	0.100
			H <sub>3</sub>	..	..	0.030
			H <sub>4</sub>	..	..	0.072
5.	1/16" glass window	0.000	F	0.593	18	0.333
			H <sub>2</sub>	0.522	34	0.300
			H <sub>3</sub>	0.434	44	0.253
			H <sub>4</sub>	0.336	52	0.218

TABLE 8.2

TERMINAL FUNCTIONS AND 'U' VALUES OF THE  
COMPONENTS OF THE TIMBER MODEL

$$\begin{aligned} h_0 &= 2.5 \\ h_1 &= 1.4 \end{aligned}$$

S. No.	Component	'Hear-' value		Transfer Function	'Internal Driving point Function'	
		U mo- dulus	Arg- $\Phi$ $\lambda_i$		Modulus	Arg - $\Phi$ $\lambda_i$
1.	1/8" plywood + 1" mineral wool + 1/8" 0.186 plywood walls	P $H_2$ $H_3$ $H_4$	0.062 0.022 0.009 0.004	120 168 234 264	0.600 0.422 0.322 0.274	36 44 43 50
2.	2" Cane fibre board roof	P $H_2$ $H_3$ $H_4$	0.065 0.025 0.014 0.011	120 162 224 248	0.511 0.422 0.378 0.356	22 24 24 24
3.	1/8" plywood floor	P $H_2$ $H_3$ $H_4$	0.534 0.464 0.422 0.356	20 32 42 50	0.339 0.356 0.312 0.230	16 23 34 40
4.	1½" timber frame	P $H_2$ $H_3$ $H_4$	0.066 0.021 0.007 0.004	140 196 224 246	0.366 0.312 0.272 0.244	30 30 23 26
5.	1/16" Glass window	P $H_2$ $H_3$ $H_4$	0.800 0.532 0.433 0.367	18 34 44 52	0.334 0.300 0.266 0.218	20 36 46 54

TABLE 6.4

ESTIMATED HEAT FLUXES RECEIVED AND ABSORBED BY THE MALONEY MODEL AND THE BUILDING  
IN 10° MR WIND SPEEDS

No.	Construction	Fundamental		2nd Harmonic		3rd Harmonic		4th Harmonic	
		$F_{nh_1} \times \lambda_i$	$F_{nh_1}(1 - \lambda_i)$	$F_{nh_1} \times \lambda_i$	$F_{nh_1}(1 - \lambda_i)$	$F_{nh_1} \times \lambda_i$	$F_{nh_1}(1 - \lambda_i)$	$F_{nh_1} \times \lambda_i$	$F_{nh_1}(1 - \lambda_i)$
1. 2" brickwall 9.3 sq.ft.		-0.495- $j0.53$	11.77+ $j1.00$	-0.23- $j0.03$	12.10+ $j0.76$	-0.083+ $j0.083$	12.20+ $j0.69$	-0.037+ $j0.037$	12.24+ $j0.63$
2. 1" timber floor 7.1 sq.ft.		-0.400- $j0.720$	7.00+ $j1.06$	-0.38- $j0.09$	7.30+ $j1.50$	-0.126+ $j0.063$	7.83+ $j1.30$	-0.036+ $j0.037$	8.10+ $j1.0$
3. 2" concrete roof 7.1 sq.ft.		-0.230- $j0.310$	9.23+ $j0.70$	-0.15- $j0.05$	0.37+ $j0.49$	-0.033+ $j0.033$	0.43+ $j0.40$	-0.020+ $j0.020$	0.53+ $j0.33$
4. 1/10" glass window 1.7 sq.ft.		1.350- $j0.45$	1.65+ $j0.27$	1.02- $j0.70$	1.00+ $j0.49$	0.76+ $j0.70$	1.05+ $j0.46$	0.40- $j0.00$	2.00+ $j0.49$
5. 2" brick wall (partition) 4.0 sq.ft.		...	5.05+ $j0.47$	...	8.17+ $j0.24$	...	5.00+ $j0.20$	...	0.07+ $j0.23$
Total Heat Flux :		0.125- $j2.02$	24.78+ $j0.00$	0.33- $j0.03$	33.03+ $j3.51$	0.07- $j0.00$	33.77+ $j2.03$	0.37- $j0.02$	37.20+ $j2.40$
		3.02 $\angle -03^{\circ}30'$	30.0 $\angle 0^{\circ}$	0.075 $\angle -70^{\circ}00'$	33.02 $\angle 5^{\circ}30'$	0.76 $\angle 65^{\circ}0$	33.00 $\angle 0^{\circ}30'$	0.04 $\angle 33^{\circ}30'$	37.20 $\angle 30^{\circ}0$
Temperatura amplitudo and phase		0.057 $\angle -03^{\circ}30'$		0.027 $\angle -73^{\circ}18'$		0.021 $\angle 65^{\circ}33'$		0.021 $\angle 53^{\circ}18'$	

TBL 0.6

HARMONIC HEAT FLUXES RECEIVED AND ABSORBED BY THE TIMBER MODEL AND THE RESULTANT  
INCIDING AIR TEMPERATURES

$$h_0 = 2.5$$

$$h_1 = 1.4$$

No.	Construction	Fundamental		4th		2nd Harmonic		3rd Harmonic		5th Harmonic	
		A	B	A	B	A	B	A	B	A	B
1.	1/8" plywood + 1" mineral wool + 1/8" plywood walls 5.0 sq.ft.	$P_{th} \rightarrow \lambda_i$	$F_{th} h_1 (1 - \frac{1}{\lambda_i})$	$P_{th} \rightarrow \lambda_i$	$F_{th} h_1 (1 - \frac{1}{\lambda_i})$	$P_{th} \rightarrow \lambda_i$	$F_{th} h_1 (1 - \frac{1}{\lambda_i})$	$P_{th} \rightarrow \lambda_i$	$F_{th} h_1 (1 - \frac{1}{\lambda_i})$	$P_{th} \rightarrow \lambda_i$	$F_{th} h_1 (1 - \frac{1}{\lambda_i})$
2.	1/8" plywood floor 7.0 sq.ft.	6.20+j1.80	6.70+j1.20	6.30+j2.70	7.30+j1.00	3.30+j3.00	7.80+j1.85	0.40+j2.00	8.35+j1.00		
3.	1" cano fibro board roof 7.0 sq.ft.	-0.34+j0.60	5.60+j2.00	-0.24+j0.03	6.50+j1.00	-0.106+j0.103	7.00+j1.60	-0.003+j0.01	0.60+j0.60		
4.	1 1/2" timber frame 3.4 sq.ft.	-0.24+j0.20	3.80+j0.80	-0.003+j0.020	3.60+j0.75	-0.021+j0.021	3.66+j0.60	-0.003+j0.017	3.70+j0.00		
5.	1/16" glass window 1.7 sq.ft.	1.35-j0.46	1.67+j0.37	1.03-j0.70	1.82+j0.43	0.75-j0.70	1.97+j0.45	0.40-j0.60	2.10+j0.42		
Total Heat Flux :		6.72-j3.59	21.87+j7.22	4.61-j3.49	23.92+j7.10	3.88-j3.62	26.07+j0.55	2.86-j3.50	30.46+j3.12		
		6.63 $\angle 33^{\circ}12'$	22.73 $\angle 18^{\circ}30'$	5.70 $\angle 35^{\circ}33'$	26.02 $\angle 16^{\circ}12'$	5.86 $\angle 45^{\circ}10'$	27.8 $\angle 19^{\circ}42'$	4.61 $\angle 52^{\circ}56'$	30.9 $\angle 9^{\circ}30'$		
Temperatura amplitudo and phasos,		0.29 $\angle -50^{\circ}42'$		0.23 $\angle -62^{\circ}$		0.19 $\angle -53^{\circ}$		0.16 $\angle -60^{\circ}33'$			

will be

$$t_{ian} = \sum A_k Y_{ink}$$

To satisfy the heat balance equation the sums of column 'A' and the product of  $t_{ia}$  and the sum of the column 'B' should be equal for each harmonic.

$$\text{Then } \vec{\theta}_n = \vec{t}_{ian} \vec{\varphi}_n$$

$$\text{where } \vec{\theta}_n = \vec{\theta}_n \angle \theta_n' \quad \text{and} \quad \vec{\varphi}_n = \vec{\varphi}_n \angle \theta_n''$$

$$\text{This gives } t_{ian} = \frac{\vec{\theta}_n}{\vec{\varphi}_n}$$

The term  $\vec{\theta}_n / \vec{\varphi}_n$  will give the amplitude of the harmonic temperature variation and  $-(\theta_n' - \theta_n'')$  gives the phase lag. By dividing the total of column 'A' by the total of column 'B' for each harmonic, the amplitude and phase of the indoor air temperature variation are obtained. The amplitudes and phases for the fundamental and higher harmonics, obtained by the above procedure are also given in Tables (8.4 and 8.5). The Fourier equations of the inside air temperatures for the models are obtained as

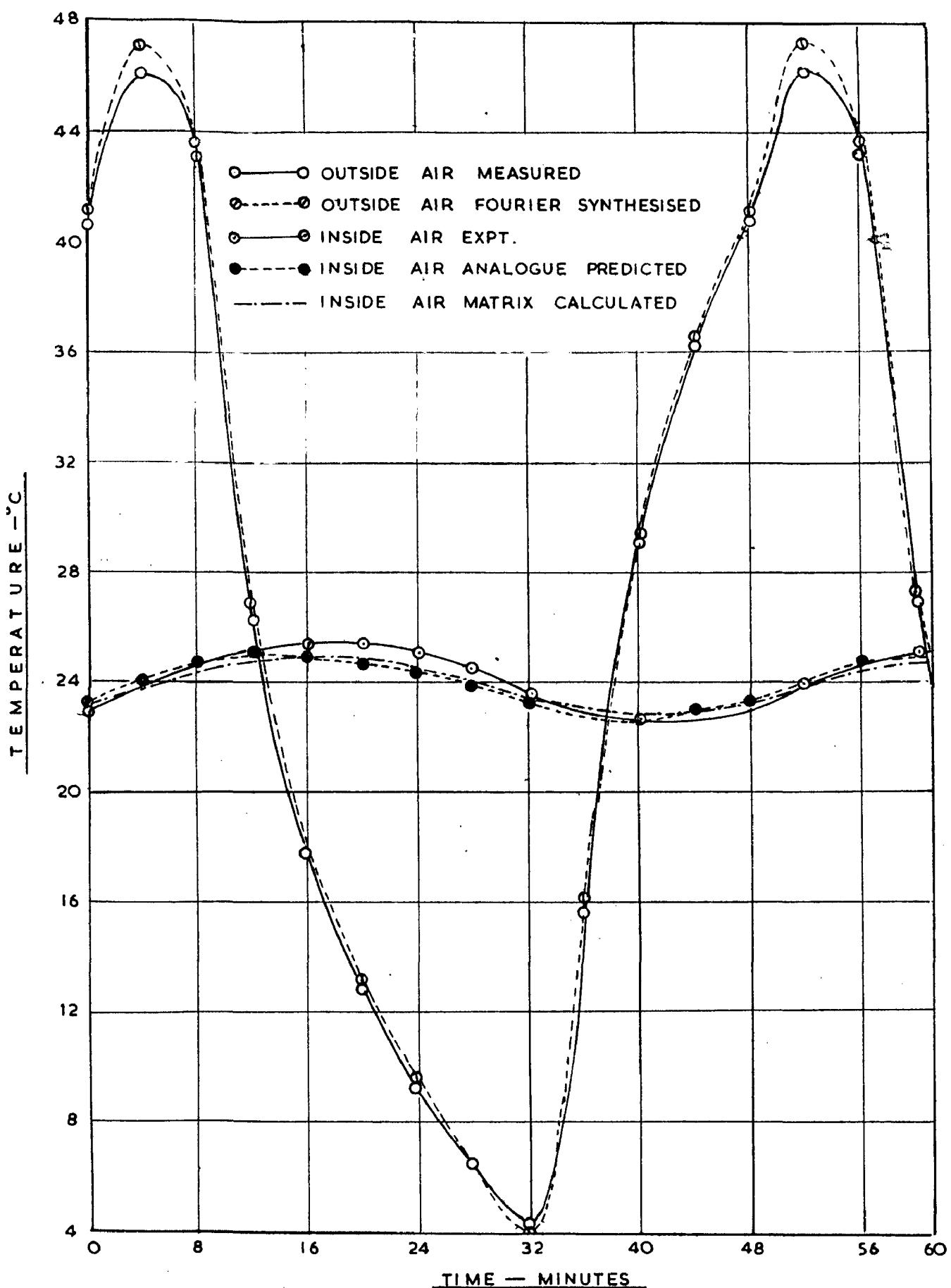
### i) Masonry Model

$$\begin{aligned} t_{iu}(T) = & 24.05 + 1.14 \cos(1ut - 115^\circ) \\ & + 0.06 \cos(2ut - 103^\circ) \\ & + 0.09 \cos(3ut - 219^\circ) \\ & + 0.02 \cos(4ut - 215^\circ) \end{aligned} \quad ..(13)$$

### ii) Timber Model

$$\begin{aligned} t_{ia}(T) = & 24.05 + 5.8 \cos(ut - 72^\circ) \\ & + 0.40 \cos(2ut - 85^\circ) \\ & + 0.68 \cos(3ut - 220^\circ) \\ & + 0.20 \cos(4ut - 218^\circ) \end{aligned} \quad ..(14)$$

By synthesising the above Fourier equations, the time temperature variations of the indoor temperatures were obtained and are compared with the measured and Matrix calculated values (Muncey) in Figs. (8.2 and 8.3) for Masonry and Timber Models respectively. The analogue predicted results found to be in close agreement with the Matrix calculated values than the measured ones, as expected. The deviations from the measured temperatures are due to the assumptions made in theory and are less significant, than what they appear to be, in view of the fact that the thermal properties of the materials could vary between wide limits. The main advantage of this method is the flexibility and ease of computation. The effect of almost any factor can be evaluated with a little additional effort. This is possible because it utilises a set of thermal system functions, which have been tabulated for a large number of commonly used structural elements and these transfer functions are not affected by the external climatic variations, as these are dependent only on the frequency of the wave form, but not on the amplitude. By adopting this procedure there is no need to represent the building, as a whole on the analogue model and hence the experimental part also is very much simplified and less expensive. This method is particularly suited for generalised design problems, where a quantitative knowledge of the improvement in the thermal efficiency versus economic factors, as affected by design variations, is of great practical importance.

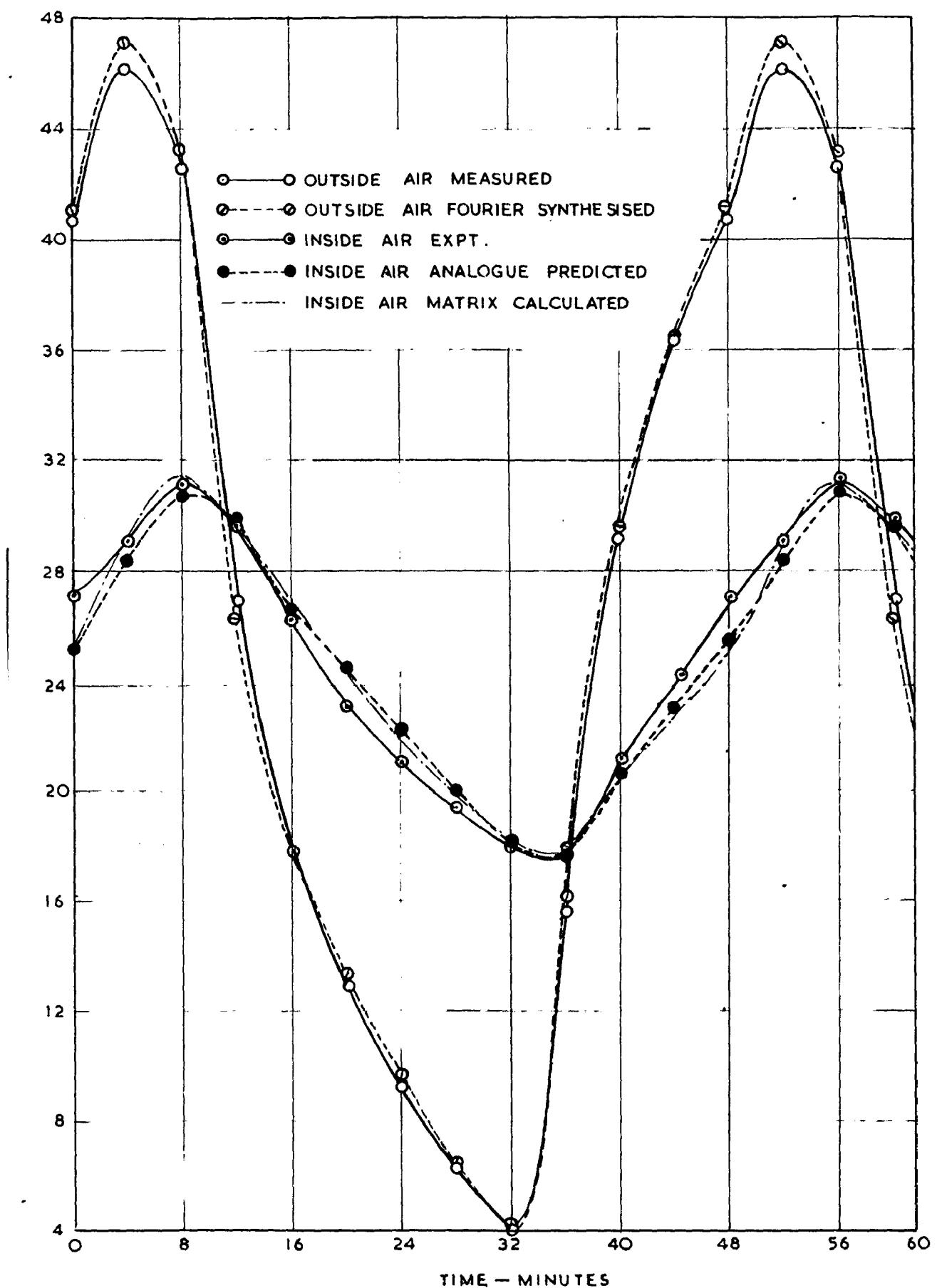


COMPARISION OF INSIDE AIR TEMPERATURES OF MASONRY MODEL

FIG.8.2

## **C H A P T E R   9**

### **APPLICATION OF THE THERMAL SYSTEM FUNCTION DATA TO DESIGN PROBLEMS**



COMPARISION OF INSIDE AIR TEMPERATURES OF TIMBER MODEL

FIG. 8.3

**C H A P T E R   9**

**APPLICATION OF THE THERMAL SYSTEM FUNCTION  
DATA TO DESIGN PROBLEMS**

## CHAPTER 9

### APPLICATIONS OF THE THERMAL SYSTEM FUNCTION

#### DATA TO DESIGN PROBLEMS

##### 9.1 Introduction

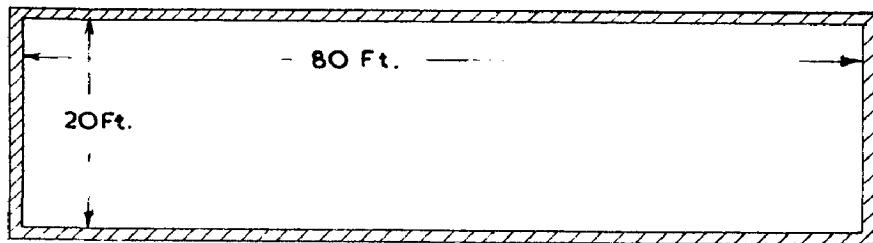
In the last Chapter (8) it has been shown that the analogically obtained, thermal system function data, for individual elements, can directly be applied for the prediction of indoor air temperatures with reasonable accuracy. The transfer and driving point function data also provide a quick and realistic estimates of cooling loads of the fabrics for air-conditioned buildings, taking the internal masses also into account.

The method described is highly flexible. The effects of many factors, such as, the type of materials and their combinations, design features, orientation surface colour treatments, regional and seasonal climatic variations, partitions, insulation and shading of roof, walls, windows and ventilation etc. on thermal efficiency, for both conditioned and unconditioned buildings can be evaluated with minimum effort.

A few typical applications of the transfer and driving point function data in solving the above mentioned problems are illustrated in this chapter.

#### 9.2 Effect of Orientation on Indoor Air Temperatures

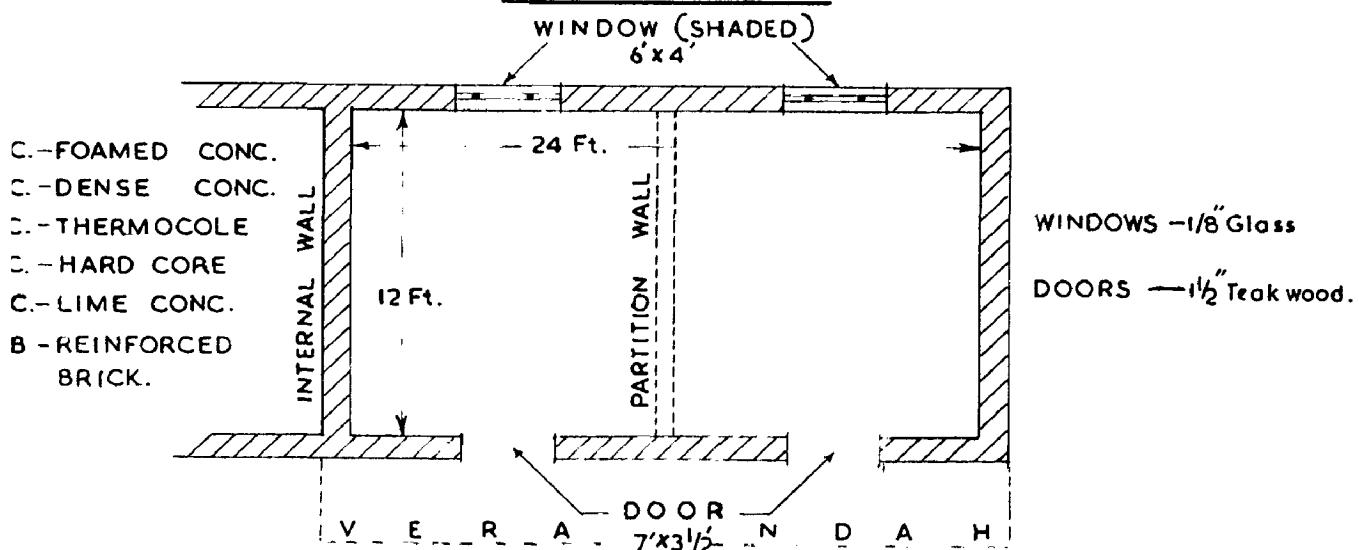
For evaluating the effect of orientation on the indoor air temperature variation a room of dimensions 30' x 20' x 12' has been considered (Fig. 9.1(a)). Physical properties and constants are given in Table 9.1. The doors and windows are also considered to be made of a construction which has the same U value and transfer and driving point functions as that of the walls. The surface absorption coefficient of wall exterior surfaces are taken as 0.5 while for the roof as 0.1. Roof being flat the sol-air temperatures of the roof will not be affected by orientation. Two orientations one with long walls facing East-West and another with long walls facing South-North have been considered. The example has been worked out for a place situated on a latitude of 29°N and for a typical summer day (May 16th). For these conditions, the sol-air temperatures were calculated. The sol-air temperature curves obtained for roof, and walls are shown in Fig. (9.2). These were harmonically analysed and their Fourier equations are in Table (9.2).



1. WALLS —  $1\frac{1}{2}$ " Plaster + 9" Brick +  $1\frac{1}{2}$ " Plaster
2. ROOF — 3" Lime Concrete + 4" R.C.C. +  $1\frac{1}{2}$ " Plaster
3. FLOOR —  $1\frac{1}{2}$ " D.C. + 3" L.C. + 6" H.C. + 20" Soil.
4. HEIGHT — 12 FT.

### ROOM CONSIDERED FOR THE ILLUSTRATIVE EXAMPLES FOR ORIENTATION

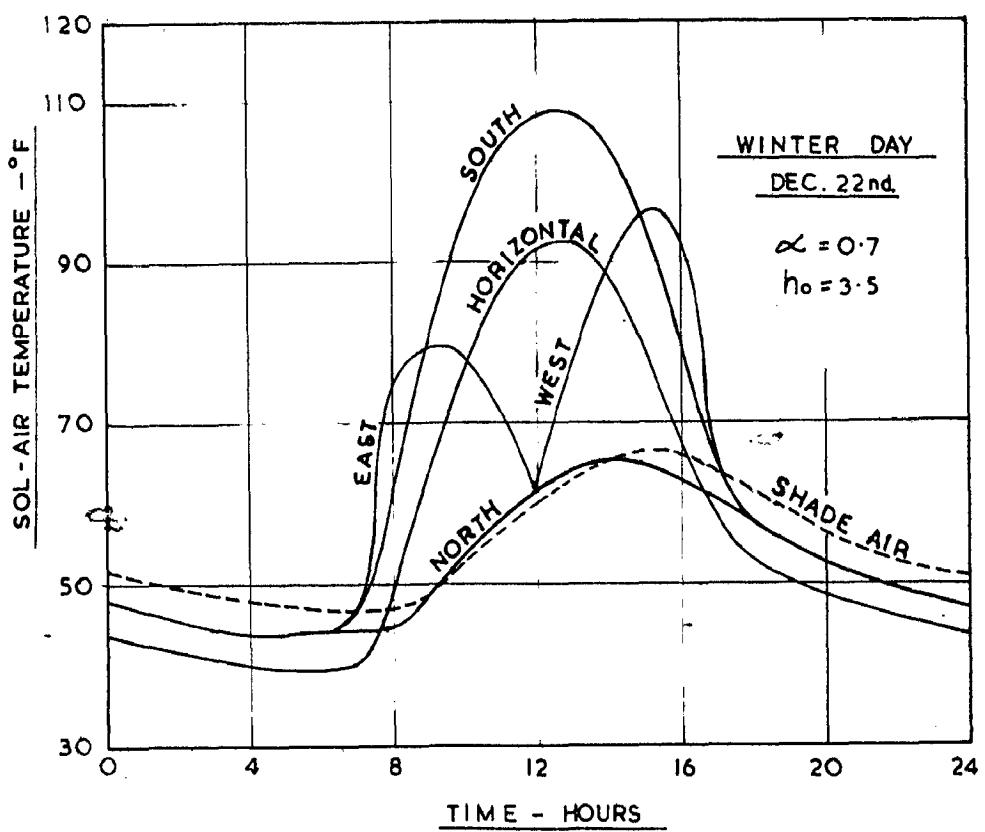
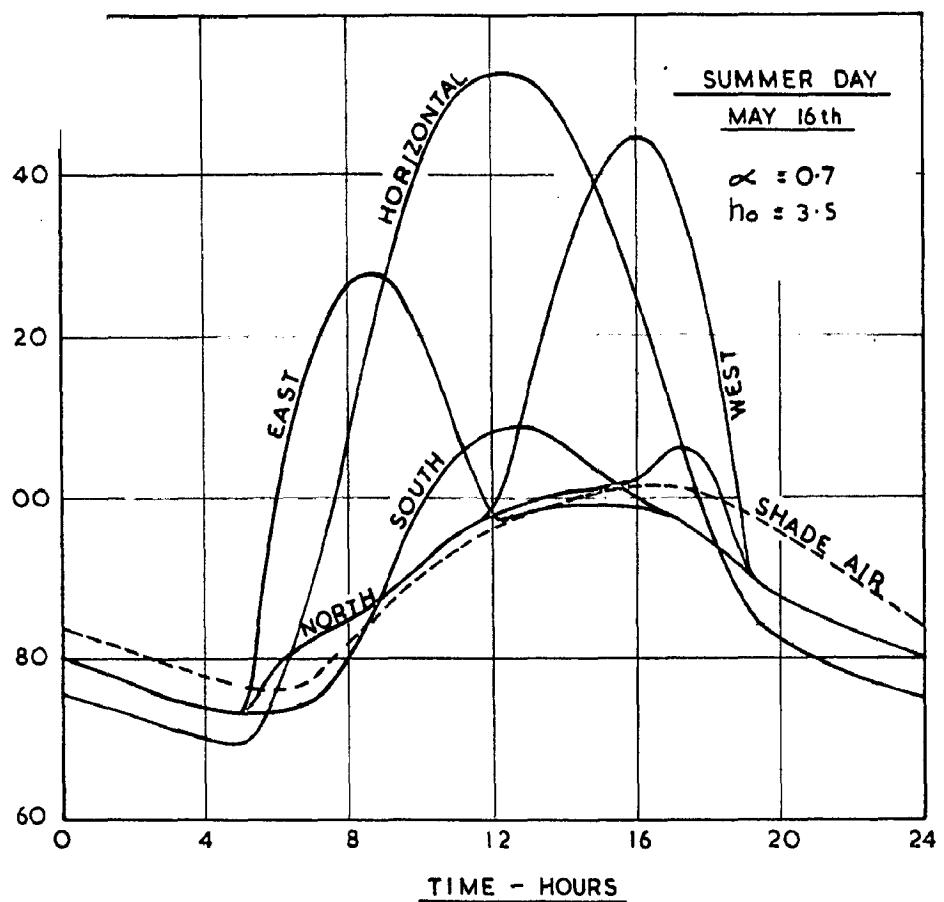
FIG. 9·1(a)



1. WALLS —  $1\frac{1}{2}$ " Plaster + 9" Brick +  $1\frac{1}{2}$ " Plaster
2. ROOF — (i) 3" Lime Concrete + 4" R.C.C. +  $1\frac{1}{2}$ " Plaster  
 (ii)  $1\frac{1}{2}$ " Plaster + 2" F.C. + 4" R.C.C. +  $1\frac{1}{2}$ " Plaster  
 (iii) 1" Plaster + 1" T.C. + 4 $\frac{1}{2}$ " R.B. +  $1\frac{1}{2}$ " Plaster
3. FLOOR —  $1\frac{1}{2}$ " D.C. + 3" L.C. + 6" H.C. + 20" Soil
4. HEIGHT — 10 Ft.

### ROOM CONSIDERED FOR OTHER ILLUSTRATIVE EXAMPLES

FIG. 9·1 (b)



SOL-AIR TEMPERATURES OF DIFFERENTLY ORIENTED SURFACES (LAT-29°N)

FIG. 9-2

TABLE 9.1

PHYSICAL PROPERTIES OF THE MATERIALS  
TAKEN IN THE EXAMPLE

No.	Material	K	$\rho$	s
1.	Brick	6.0	100	0.21
2.	Plaster	12.0	120	0.32
3.	Teakwood	1.2	40	0.30
4.	Thermocole	0.20	1	0.32
5.	Dense Concrete	8.0	120	0.21
6.	Foamed Concrete	0.75	30	0.25
7.	Lime Concrete	6.0	106	0.20
8.	Hard Coro	7.0	110	0.20
9.	Soil	8.0	110	0.20
10.	Glass	6.6	160	0.16

K = Thermal conductivity in Dtu.in/Ft<sup>2</sup>.Hr.<sup>0</sup>F

$\rho$  = Density in Lb/Ft<sup>3</sup>

s = Specific heat in Dtu/Lb.<sup>0</sup>F

TABLE 9.1

PHYSICAL PROPERTIES OF THE MATERIALS  
TAKEN IN THE EXAMPLE

No.	Material	K	$\rho$	s
1.	Brick	6.0	100	0.21
2.	Plaster	12.0	120	0.32
3.	Teakwood	1.2	40	0.39
4.	Thermocole	0.20	1	0.32
5.	Dense Concrete	8.0	120	0.21
6.	Foamed Concrete	0.75	30	0.25
7.	Lime Concrete	6.0	108	0.20
8.	Hard Core	7.0	110	0.20
9.	Soil	0.0	110	0.20
10.	Glass	6.6	160	0.16

K = Thermal conductivity in Btu.in/Ft<sup>2</sup>.Hr.<sup>0</sup>F

$\rho$  = Density in Lb/Ft<sup>3</sup>

s = Specific heat in Btu/Lb.<sup>0</sup>F

RADIANT 9.3OUTDOOR AIR AND SOL-AIR TEMPERATURE FOURIER EQUATIONS

$$\alpha = 0.7$$

$$b_0 = 3.6$$


---

A. SUMMER DAY (MAY 16TH)

1. Flat Roof	$104.7 + 40.1 \cos(15t - 180^\circ) + 13.9 \cos(30t - 30^\circ) + 1.7 \cos(45t - 57^\circ) + 2.2 \cos(60t - 365^\circ)$
2. East Wall	$95.1 + 17.6 \cos(15t - 170^\circ) + 10.6 \cos(30t - 286^\circ) + 0.0 \cos(45t - 14^\circ) + 3.3 \cos(60t - 56^\circ)$
3. West Wall	$86.1 + 20.0 \cos(15t - 234^\circ) + 12.2 \cos(30t - 103^\circ) + 10.8 \cos(45t - 42^\circ) + 10.6 \cos(60t - 189^\circ)$
4. South Wall	$83.8 + 14.3 \cos(15t - 203^\circ) + 5.3 \cos(30t - 90^\circ) + 2.0 \cos(45t - 135^\circ) + 0.6 \cos(60t - 273^\circ)$
5. North Wall	$88.4 + 14.3 \cos(15t - 229^\circ) + 0.0 \cos(30t - 115^\circ) + 2.8 \cos(45t - 39^\circ) + 1.6 \cos(60t - 337^\circ)$
6. Shaded Air Temperature	$89.0 + 10.1 \cos(15t - 249^\circ) + 1.6 \cos(30t - 1^\circ) + 1.7 \cos(45t - 60^\circ) + 0.0 \cos(60t - 235^\circ)$

B. WINTER DAY (DEC. 22ND)

1. Flat Roof	$57.3 + 22.3 \cos(15t - 100^\circ) + 11.3 \cos(30t - 13^\circ) + 3.3 \cos(45t - 170^\circ) + 0.3 \cos(60t - 170^\circ)$
2. East Wall	$83.8 + 12.6 \cos(15t - 180^\circ) + 4.7 \cos(30t - 305^\circ) + 6.0 \cos(45t - 83^\circ) + 6.1 \cos(60t - 189^\circ)$
3. West Wall	$83.8 + 17.7 \cos(15t - 230^\circ) + 10.4 \cos(30t - 77^\circ) + 6.3 \cos(45t - 303^\circ) + 6.4 \cos(60t - 123^\circ)$
4. South Wall	$86.6 + 20.4 \cos(15t - 100^\circ) + 13.9 \cos(30t - 13^\circ) + 2.3 \cos(45t - 60^\circ) + 1.8 \cos(60t - 365^\circ)$
5. Shaded Air Temperature	$86.4 + 11.4 \cos(15t - 232^\circ) + 1.7 \cos(30t - 46^\circ) + 0.3 \cos(45t - 100^\circ) + 0.5 \cos(60t - 337^\circ)$
6. North Wall	$82.7 + 9.9 \cos(15t - 229^\circ) + 3.6 \cos(30t - 53^\circ) + 0.7 \cos(45t - 183^\circ) + 0.6 \cos(60t - 277^\circ)$

C. FLAT ROOF - SUMMER DAY (MAY 16TH)  $b_0 = 3.6$ 

1. $\alpha = 0.1$	$83.2 + 13.6 \cos(15t - 210^\circ) + 6.8 \cos(30t - 50^\circ) + 1.3 \cos(45t - 310^\circ) + 0.8 \cos(60t - 223^\circ)$
2. $\alpha = 0.3$	$89.1 + 21.0 \cos(15t - 337^\circ) + 6.9 \cos(30t - 50^\circ) + 1.5 \cos(45t - 76^\circ) + 1.3 \cos(60t - 207^\circ)$
3. $\alpha = 0.5$	$95.8 + 30.3 \cos(15t - 103^\circ) + 7.2 \cos(30t - 00^\circ) + 1.0 \cos(45t - 66^\circ) + 4.0 \cos(60t - 167^\circ)$
4. $\alpha = 0.0$	$103.0 + 40.7 \cos(15t - 191^\circ) + 17.4 \cos(30t - 20^\circ) + 2.0 \cos(45t - 56^\circ) + 6.0 \cos(60t - 7^\circ)$

---

The outside and inside surface coefficients are taken as 0.5 and 1.6 Stu./sq.ft./hr./°F respectively. The 'U' values and the transfer and driving point admittance functions for the walls, roof and floor are given in Table (9.3). The steady state and harmonic heat flow rates, through each building element into the room and the heat absorption at the inside surfaces of all the bounding elements of the room, and the mean room air temperature and the amplitudes of the harmonic components, were determined by the procedure as outlined in Chapter (3). The Fourier representation of the indoor air temperatures for the two orientations are given below :

#### South Orientation

$$t_{ia}(T) = 63.6 + 2.45 \cos(18\pi t + 23^\circ) \\ + 0.25 \cos(30\pi t + 197^\circ) \\ + 0.07 \cos(48\pi t + 180^\circ) \\ + 0.03 \cos(60\pi t + 330^\circ)$$

#### West Orientation

$$t_{ia}(T) = 63.4 + 2.3 \cos(18\pi t + 27^\circ) \\ + 0.2 \cos(30\pi t + 219^\circ) \\ + 0.16 \cos(48\pi t + 118^\circ) \\ + 0.05 \cos(60\pi t + 333^\circ)$$

The hourly indoor air temperature variations obtained are shown in Fig. (9.3). This example clearly demonstrates the advantage of facing the long walls in South-North (orientation). The advantage of this orientation will be more marked, for buildings, with

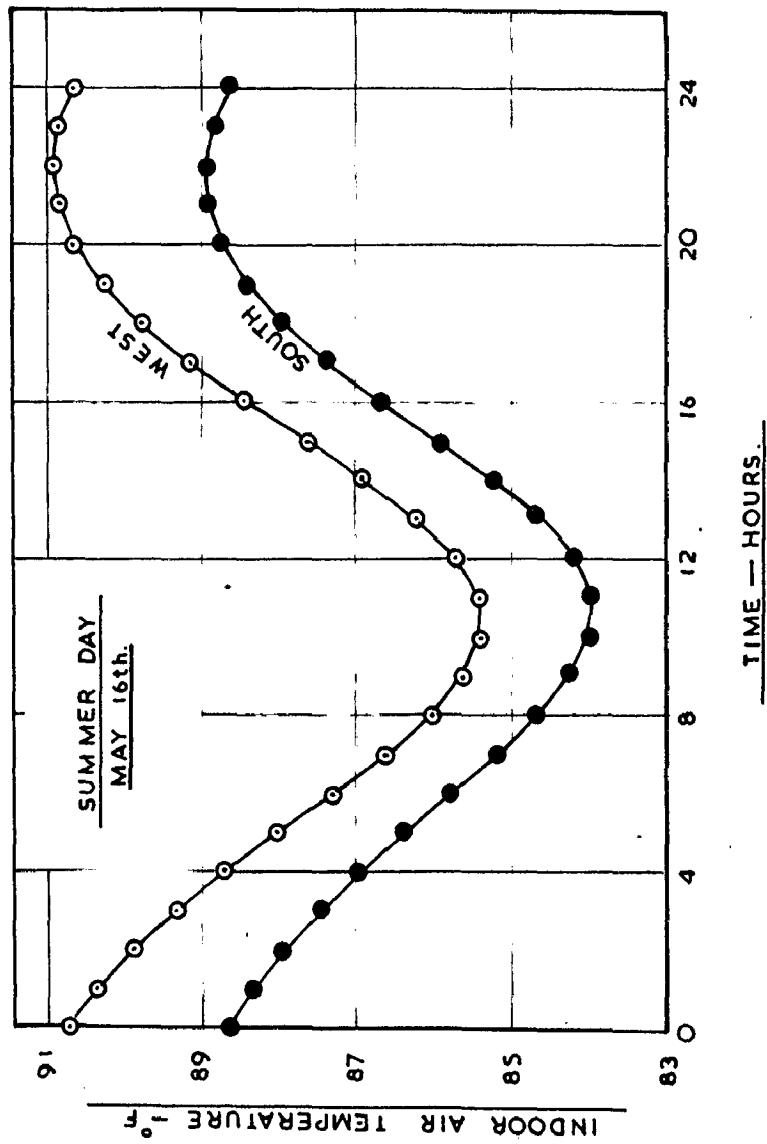


FIG. 9.3. EFFECT OF ORIENTATION ON INDOOR AIR TEMPERATURES

$$h_0 = 3.5$$

$$h_1 = 1.5$$

n-th Harmonic			
No.	Transfer admittance function	Internal driving point admittance function	
	Modu- lus 'in dB'	Arg. $-\psi_c$	
1.			
a)	0.243 0.063	160	$1.219+j0.233$
b)	0.240 0.030	160	$1.020+j0.301$
c)	0.235 0.016	202	$1.165+j0.279$
2.			
	0.294 0.026	223	$1.104+j0.370$
3.			
	0.244 ... ..		$1.169+j0.225$
4.			
	0.316 0.010	254	$1.104+j0.301$
5.			
	0.179 0.393	40	$0.010+j0.193$
6.			
	0.027 1.035	5	$1.034+j0.032$
7.			
	0.320 ... ..		$1.160+j0.303$

TABLE 9.3

TRANSFER AND DRIVING POINT ADMITTANCE FUNCTIONS OF THE CONCRETE  
USED IN THE SAMPLE  
(Fundamental)

 $b_0 = 0.6$  $b_1 = 1.6$ 

No.	Construction	U	Fundamental		2nd Harmonic		3rd Harmonic		4th Harmonic	
			Transfer admittance	Internal driving function						
			Modu.   $\psi_c$	Arg.   $\psi_c$						
			"Ins. in dog."	"Ins. in dog."						
			$y_c$	$1 - \psi_c$						

1. ROOF

a) 3" lime concrete + 4" R.C.C. + $\frac{1}{2}$ " plaster	0.602	0.300	70	$0.007 + j0.325$	0.100	116	$1.110 + j0.270$	0.100	140	$1.184 + j0.249$	0.030	160	$1.210 + j0.233$
b) $\frac{1}{2}$ " plaster + 3" foamed concrete + 4" R.C.C. + $\frac{1}{2}$ " plaster	0.333	0.117	00	$0.720 + j0.370$	0.033	103	$0.821 + j0.323$	0.036	123	$0.875 + j0.240$	0.030	100	$1.020 + j0.301$
c) $\frac{1}{2}$ " plaster + 1" thermo- cole + 4" R.B. + $\frac{1}{2}$ " plaster	0.037	0.003	03	$0.038 + j0.447$	0.003	110	$0.050 + j0.320$	0.021	154	$1.180 + j0.336$	0.014	202	$1.165 + j0.279$

2. WALL

$\frac{1}{2}$ " plaster + 3" brick + $\frac{1}{2}$ " plaster	0.323	0.104	103	$0.010 + j0.462$	0.037	152	$1.024 + j0.333$	0.046	200	$1.197 + j0.223$	0.023	220	$1.166 + j0.273$
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3. FLOOR

$\frac{1}{2}$ " concrete + 3" lime concrete + 0" hardcore + 20" soil	0.200	...	..	$0.803 + j0.818$	...	..	$1.032 + j0.373$	0..	..	$1.112 + j0.244$	...	..	$1.168 + j0.225$
--	-------	-----	----	------------------	-----	----	------------------	-----	----	------------------	-----	----	------------------

## 4.

$\frac{1}{2}$ " plaster + 3" brick + $\frac{1}{2}$ " plaster	0.345	0.117	123	$0.036 + j0.337$	0.034	176	$1.032 + j0.340$	0.020	234	$1.111 + j0.316$	0.010	284	$1.164 + j0.301$
---	-------	-------	-----	------------------	-------	-----	------------------	-------	-----	------------------	-------	-----	------------------

5. DOOR

$\frac{1}{2}$ " wood	0.302	0.000	13	$0.510 + j0.033$	0.003	24	$0.650 + j0.030$	0.004	33	$0.603 + j0.170$	0.003	40	$0.010 + j0.105$
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6. WINDOW

$\frac{1}{8}$ " glass	1.026	1.005	..	$1.034 + j0.010$	1.035	3	$1.034 + j0.021$	1.036	4	$1.030 + j0.037$	1.005	5	$1.024 + j0.002$
-----------------------	-------	-------	----	------------------	-------	---	------------------	-------	---	------------------	-------	---	------------------

7. REINFORCED WALL

$\frac{1}{2}$ " plaster + 3" brick + $\frac{1}{2}$ " plaster	0.345	...	..	$0.809 + j0.420$	...	..	$1.031 + j0.243$	...	..	$1.080 + j0.220$	...	..	$1.163 + j0.303$
---	-------	-----	----	------------------	-----	----	------------------	-----	----	------------------	-----	----	------------------

R.D. = Reinforced Brick    R.C.C. = Reinforced Cement Concrete.

sloping roofs than with flat ones especially for a large length to breadth ratio.

### 9.3 Effect of Seasonal Variations on Indoor Air Temperatures

For evaluating the thermal behaviour of any building in different seasons, at any particular geographical location, a single room structure, shown in Fig. (9.1(b)) was considered. The physical properties of the materials are given in Table (9.1). The 'U' values and the transfer functions and driving point admittance of the structural elements are given in Table (9.3). A typical summer day (May 16th) and a winter day (Dec. 23rd) have been taken. The orientation is fixed as South. The windows were considered to be completely shaded from the sun and no direct solar radiation penetrates into the room. The surface absorption coefficient ( $\lambda$ ) of the roof was taken as 0.7 and of the walls as 0.5. The sol-air temperatures for the roof, walls, doors and windows have been computed from the climatological data and their Fourier equations are given in Table (9.3).

The steady state and harmonic heat flow into the room through all external facing structural elements and the harmonic heat absorption at the internal surfaces for unit amplitude variation of indoor harmonic temperature variations, and the mean indoor air temperature and

harmonic variations have been determined for both the days. Their Fourier equations are given below.

Summer Day (May 16th)

$$t_{1u}(T) = 93.8 + 4.3 \cos(15 t + 57^\circ) \\ + 0.55 \cos(30 t + 233^\circ) \\ + 0.05 \cos(45 t + 190^\circ) \\ + 0.03 \cos(60 t + 146^\circ)$$

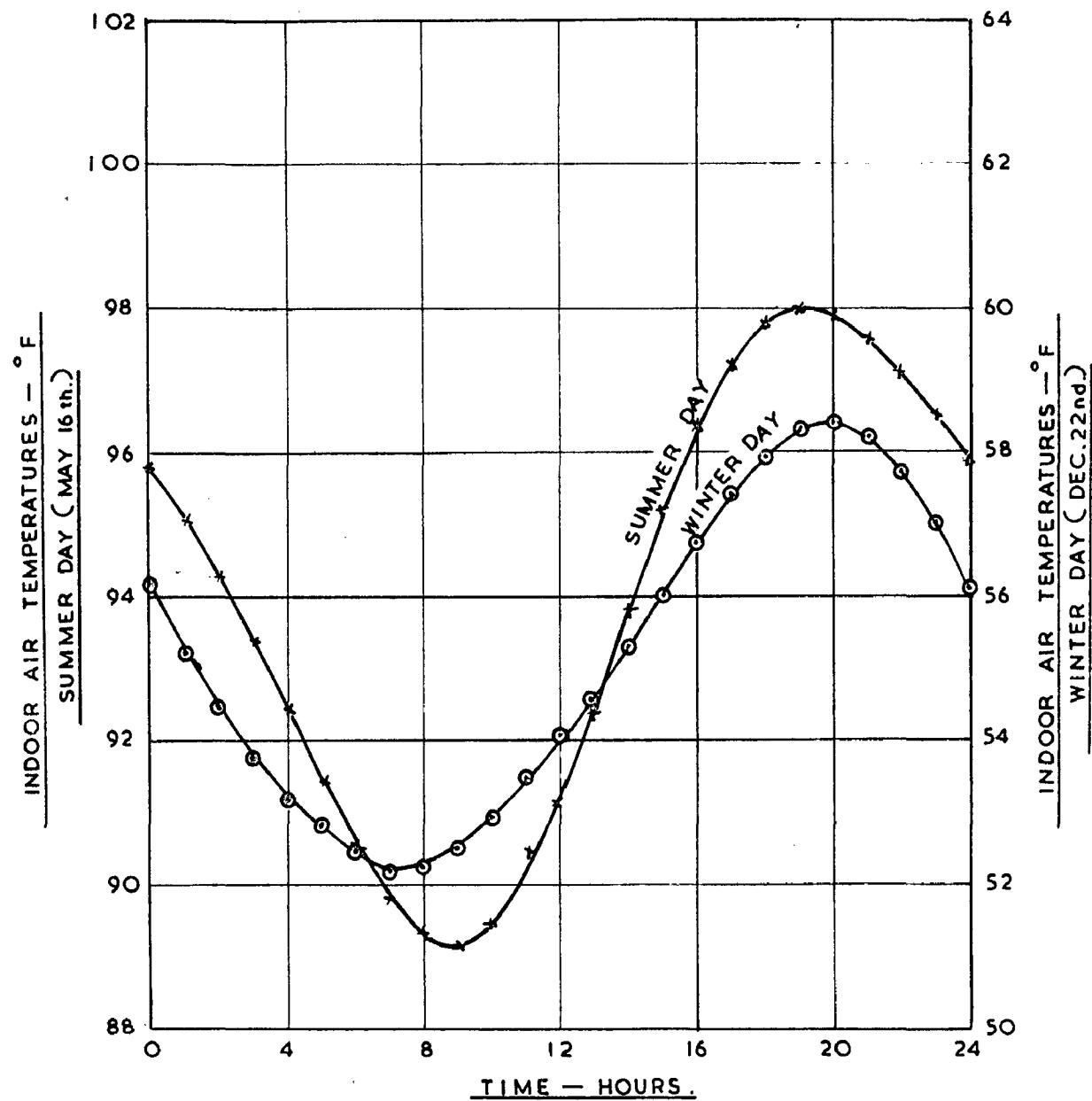
Winter Day (Dec. 23rd)

$$t_{1u}(T) = 55.2 + 3.9 \cos(15 t + 69^\circ) \\ + 0.40 \cos(30 t + 125^\circ) \\ + 0.03 \cos(45 t + 63^\circ) \\ + 0.00 \cos(60 t + 120^\circ)$$

The hourly indoor air temperatures as obtained from the equations are shown in Fig. (3.4).

9.4 Effect of surface treatments on Indoor Air Temperatures

It is interesting to know the influence of surface colour of a building element exposed to the solar radiation on indoor air temperatures. Such a study will provide an indication of the effectiveness of different surface treatments in improving the indoor thermal conditions. For this purpose the same single room structure Fig. (3.1 (b)) was considered for a typical summer day and south orientation. Four surface absorption coefficients viz., 0.1 0.5, 0.7 and 0.9, for roof only were considered. The solair temperature wave



SEASONAL VARIATIONS OF INDOOR AIR TEMPERATURES

FIG. 9-4

forms of the roof with these four surface absorption coefficients are shown in Fig. (9.5). Here heat flow through all other building elements except the roof are, the same as that of the summer day, in the previous example (2). The fourier equations of the indoor air temperatures obtained for these four variations of  $\alpha$ , are given below.

For  $\alpha = 0.1$ ;

$$t_{ia}(T) = 87.2 + 2.6 \cos(15 t + 40^\circ) \\ + 0.32 \cos(30 t + 242^\circ) \\ + 0.04 \cos(45 t + 188^\circ) \\ + 0.02 \cos(60 t + 78^\circ)$$

For  $\alpha = 0.5$

$$t_{ia}(T) = 91.1 + 3.7 \cos(15 t + 51^\circ) \\ + 0.23 \cos(30 t + 200^\circ) \\ + 0.05 \cos(45 t + 188^\circ) \\ + 0.07 \cos(60 t + 53^\circ)$$

For  $\alpha = 0.7$

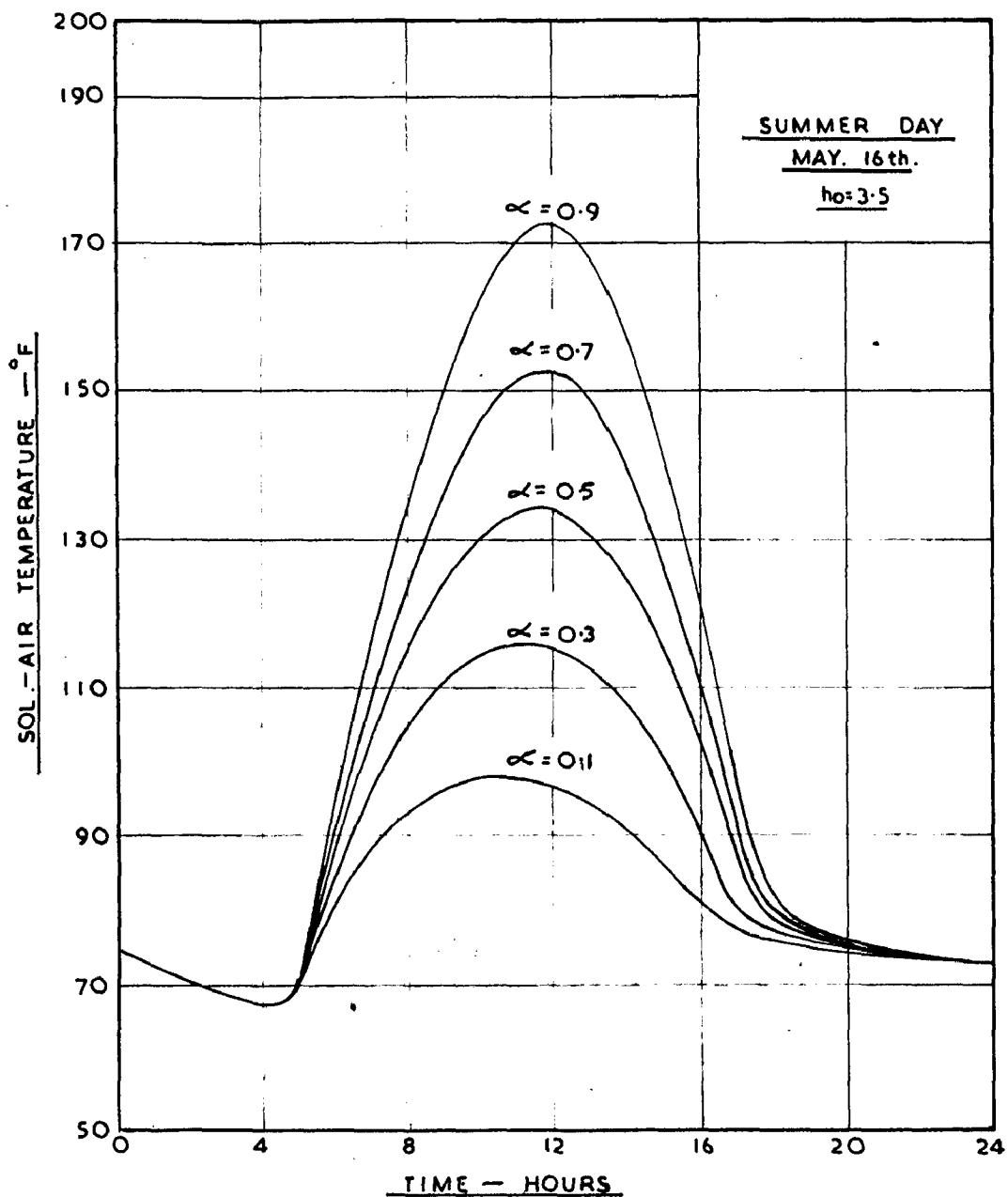
$$t_{ia}(T) = 93.8 + 4.3 \cos(15 t + 57^\circ) \\ + 0.56 \cos(30 t + 235^\circ) \\ + 0.06 \cos(45 t + 190^\circ) \\ + 0.03 \cos(60 t + 146^\circ)$$

For  $\alpha = 0.9$

$$t_{ia}(T) = 94.9 + 5.0 \cos(15 t + 59^\circ) \\ + 0.67 \cos(30 t + 235^\circ) \\ + 0.06 \cos(45 t + 187^\circ) \\ + 0.07 \cos(60 t + 155^\circ)$$

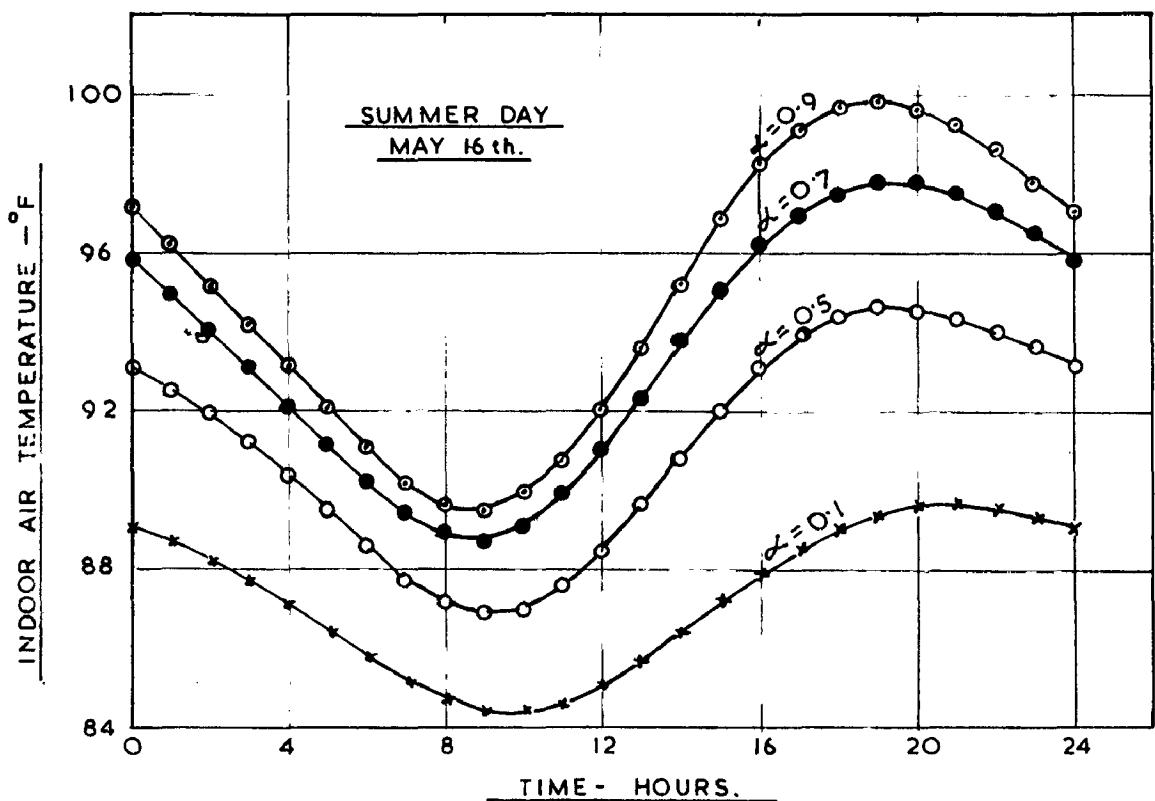
The hourly indoor air temperature variations are calculated and compared in Fig. (9.6).

It is clear from this study that highly reflective surface finishes have a marked influence in



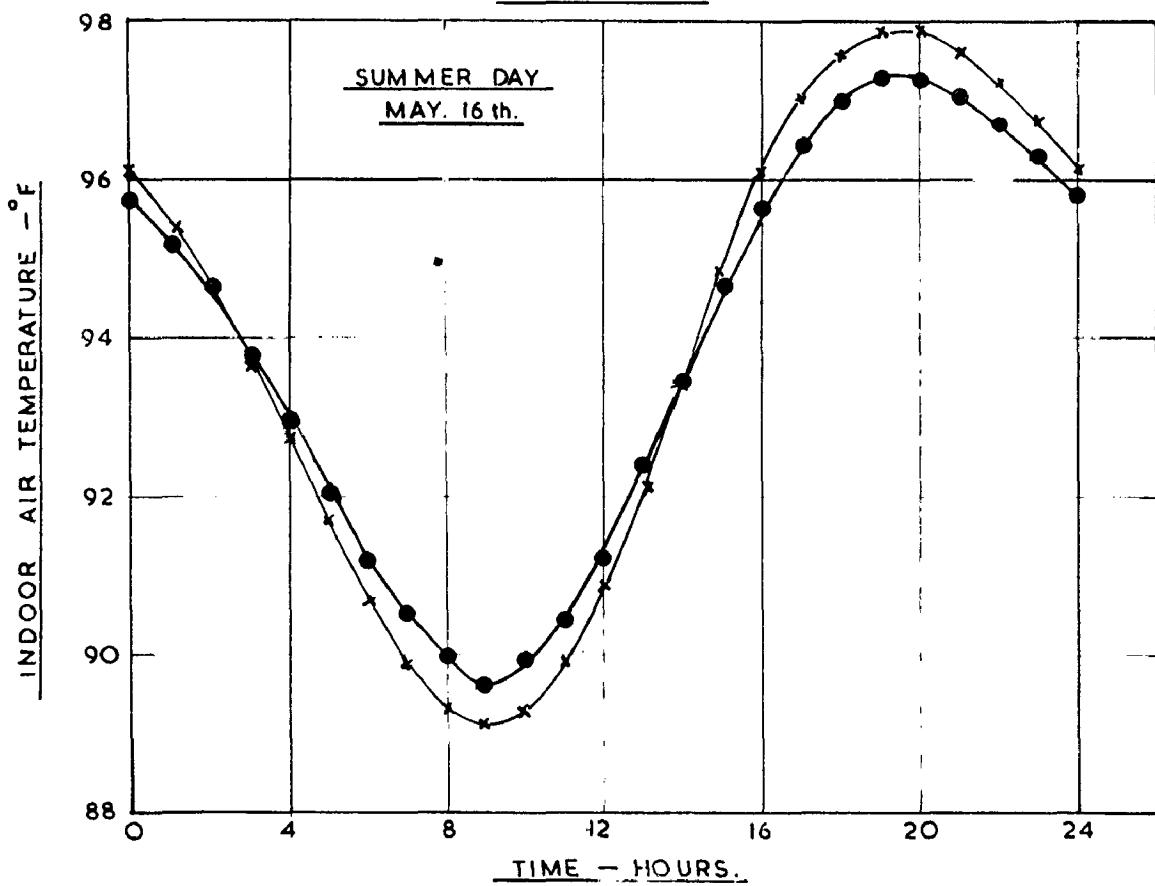
EFFECT OF SURFACE ABSORPTION COEFFICIENT ON  
SOL-AIR TEMPERATURE OF A FLAT ROOF

FIG. 9·5



EFFECT OF SURFACE TREATMENT ON INDOOR AIR TEMPERATURES

FIG. 9·6



EFFECT OF INTERNAL MASS ON INDOOR AIR TEMPERATURES

FIG. 9·7

reducing the indoor air temperatures. The differences between two extreme cases  $\alpha = 0.9$  and 0.1 was nearly  $10^{\circ}\text{F}$ , which is a very significant reduction. The major part of the reduction is affected in the steady state level ( $7.7^{\circ}\text{F}$ ).

#### 9.5 Effect of Internal Mass on Indoor Air Temperature

Internal masses like partitions and furniture, have thermal capacity and hence a part of the heat that enters into the room is absorbed by them to modify the indoor air temperature. The magnitude of the reduction in indoor air temperature will depend upon the internal driving point admittance function and area of those internal masses. As an illustrative example, the same single room structure taken in earlier examples (Fig. 9.1(b)) has been considered with the inclusion of a partition wall of 9" brick with 1/2" plaster on either side as shown by dotted lines in the same figure. This being an internal wall the surface coefficients on either side can be taken to be the same. The driving point and transfer admittance functions of the partition wall with these surface coefficients will be different from that of the external bounding wall though of same construction. There will be no change in the mean indoor air temperature with and without the partition, while the amplitudes of harmonic components will be reduced. These

have been determined and the corresponding Fourier equation of the indoor air temperature is given below.

$$\begin{aligned} t_{iu} (\tau) = & 93.8 + 3.62 \cos(15\tau + 56^\circ) \\ & + 0.46 \cos(30\tau + 236^\circ) \\ & + 0.04 \cos(45\tau + 190^\circ) \\ & + 0.03 \cos(60\tau + 146^\circ) \end{aligned}$$

The hourly indoor air temperatures obtained with and without the partition are compared in Fig. (9.7). These indicate that the reduction in temperatures due to the addition of a partition (internal mass) is not very significant in small buildings. This may reach a significant value for large buildings with many internal walls and larger surface areas.

#### 9.6 Effect of the Type of Roof on Indoor Air Temperature

In order to find the effect of the type of roof, on the indoor air temperatures of unconditioned buildings, two types of insulated roofs were considered and compared with the normal roof taken in the previous examples, for the same room under similar conditions. The two roofs considered are (i)  $\frac{1}{2}$ " plaster + 2" foamed concrete +  $4\frac{1}{2}$ " reinforced brick work +  $\frac{1}{2}$ " plaster, and (ii) 1" plaster + 1" thermocole +  $4\frac{1}{2}$ " reinforced brick work +  $\frac{1}{2}$ " plaster.

The physical properties of the materials and the 'U' values and transfer driving point functions

are given in Tables (9.1. and 9.4) respectively. The heat flow through the roof entering the room and the rate of heat flux absorption at the internal surfaces of these roof will be different while the rest will be same as in the case of normal roof. The resultant indoor air temperatures are obtained, and their Fourier equations are given below.

Type A

$$\begin{aligned} t_{ia}(T) = & 01.7 + 0.8 \cos(15t + 43^\circ) \\ & + 0.3 \cos(30t + 253^\circ) \\ & + 0.04 \cos(45t + 200^\circ) \\ & + 0.03 \cos(60t + 120^\circ) \end{aligned}$$

Type B

$$\begin{aligned} t_{ia}(T) = & 80.7 + 3.2 \cos(15t + 52^\circ) \\ & + 0.83 \cos(30t + 253^\circ) \\ & + 0.03 \cos(45t + 216^\circ) \\ & + 0.03 \cos(60t + 113^\circ) \end{aligned}$$

The hourly indoor air temperatures for these three roof types are obtained and compared in Fig. (9.8). Insulating the building components, especially the roof will improve the indoor thermal conditions, to considerable extent, but not as effective as surface treatments.

### 9.7 Effect of Ventilation on Indoor Air Temperatures

It is known, that the indoor air temperature of a room, will not only depend upon the heat flux received through the building fabrics, but also on ventilation rate. Because of the temperature differences

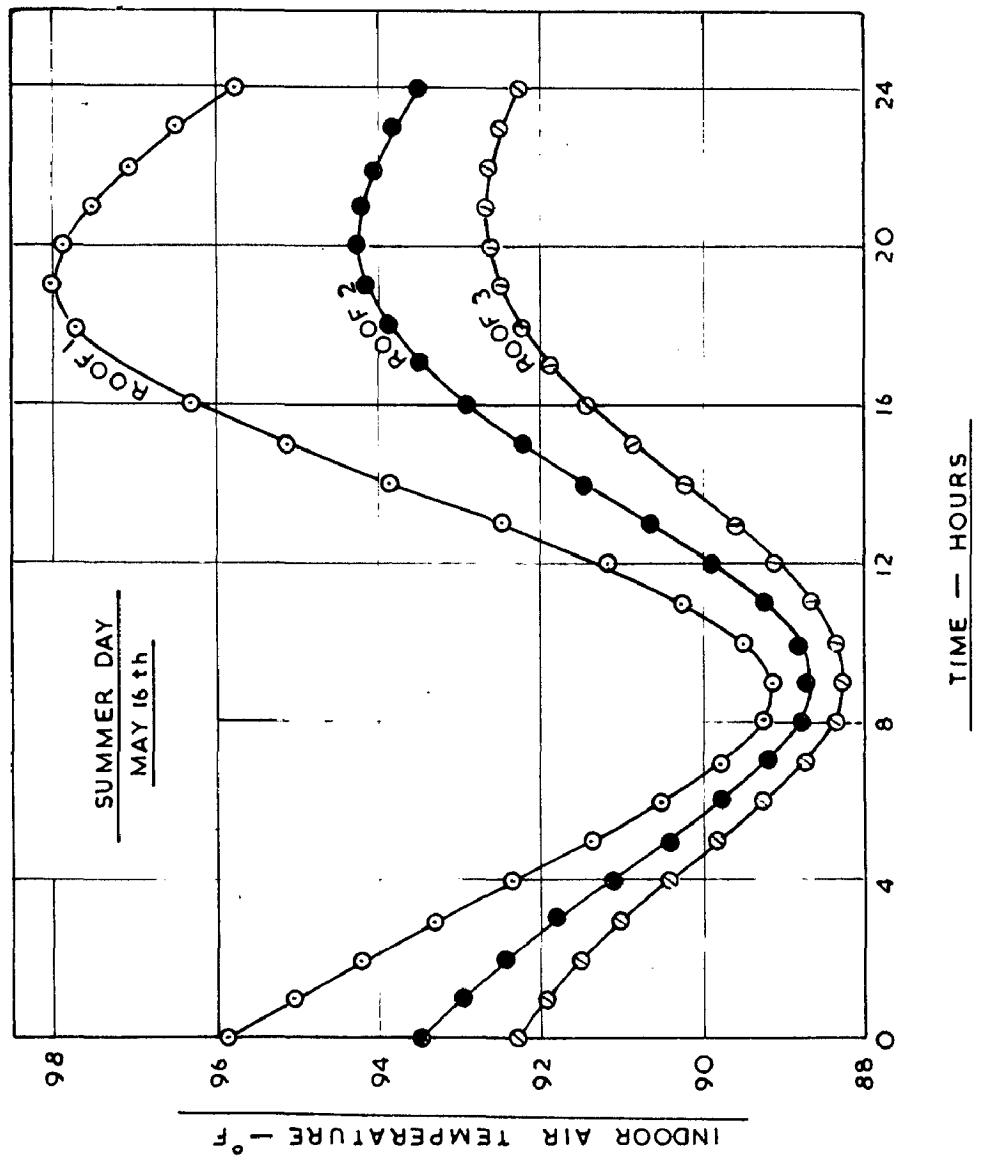


FIG. 9.8. EFFECT OF ROOF ON INDOOR AIR TEMPERATURES

TABLE 9.4

INTERNAL FUNCTIONS OF THE 9007- USED IN THE INPUTS

No.	Construction value	Module	Arg	Modulus	Arg in deg	Arg in rad	Modulus	Arg in deg	Arg in rad	Modulus	Arg in deg	Arg in rad	Modulus	Arg in deg	Arg in rad
1.	3" 14no concrete + 4" A.C.C. + 1" plaster	P H2 H3 H4	0.733 0.649 0.600 0.567	12 17 19 21	0.200 0.111 0.067 0.044	78 116 140 160	0.422 0.316 0.268 0.244	31 36 38 40	0.086 0.047 0.029 0.019	73 113 140 160	- 150	-	-	-	-
2.	{" planter + 3" 10noad concrete + 4" A.C.C. + 0.238 1" plaster	P H2 H3 H4	0.812 0.828 0.866 0.818	14 17 21 26	0.078 0.041 0.030 0.020	82 104 125 160	0.573 0.433 0.378 0.351	38 41 43 44	0.034 0.018 0.013 0.009	80 104 123 160	-	-	-	-	-
3.	1" planter + 1" chamoco10 + 4" 2" brick + 1" 0.147 plaster	P H2 H3 H4	0.834 0.844 0.782 0.734	11 20 28 35	0.037 0.024 0.014 0.009	88 110 154 202	0.507 0.378 0.322 0.286	34 36 38 39	0.016 0.010 0.003 0.004	88 110 164 202	-	-	-	-	-

that exist between the indoor and outdoor air heat is either supplied to or extracted from, the room, by the process of ventilation. The exact quantity of heat transfer by the ventilation process depends upon the volume of the room and the rate of ventilation (number of air changes per hour).

In order to illustrate the effect of ventilation on the indoor air temperature, the same single room structure taken in earlier examples (Fig. 9.1(b)) has been considered. The steady state and the harmonic components of the heat flux received in the room by ventilation, have been calculated, for an assumed ventilation rate of 4 volume changes per hour. The indoor air temperature harmonic variations are determined as outlined in Chapter (8). The Fourier equations of the indoor air temperatures obtained for the conditions of, with and without ventilation are given below.

#### 1. With Ventilation

$$\begin{aligned} t_{ia}(T) = & 03.1 + 7.13 \cos(15t + 70^\circ) \\ & + 0.52 \cos(30t + 260^\circ) \\ & + 0.25 \cos(45t + 276^\circ) \\ & + 0.12 \cos(60t + 132^\circ) \end{aligned}$$

#### 2. Without Ventilation

$$\begin{aligned} t_{ia}(T) = & 03.0 + 4.0 \cos(15t + 70^\circ) \\ & + 0.65 \cos(30t + 233^\circ) \\ & + 0.03 \cos(45t + 193^\circ) \\ & + 0.03 \cos(60t + 146^\circ) \end{aligned}$$

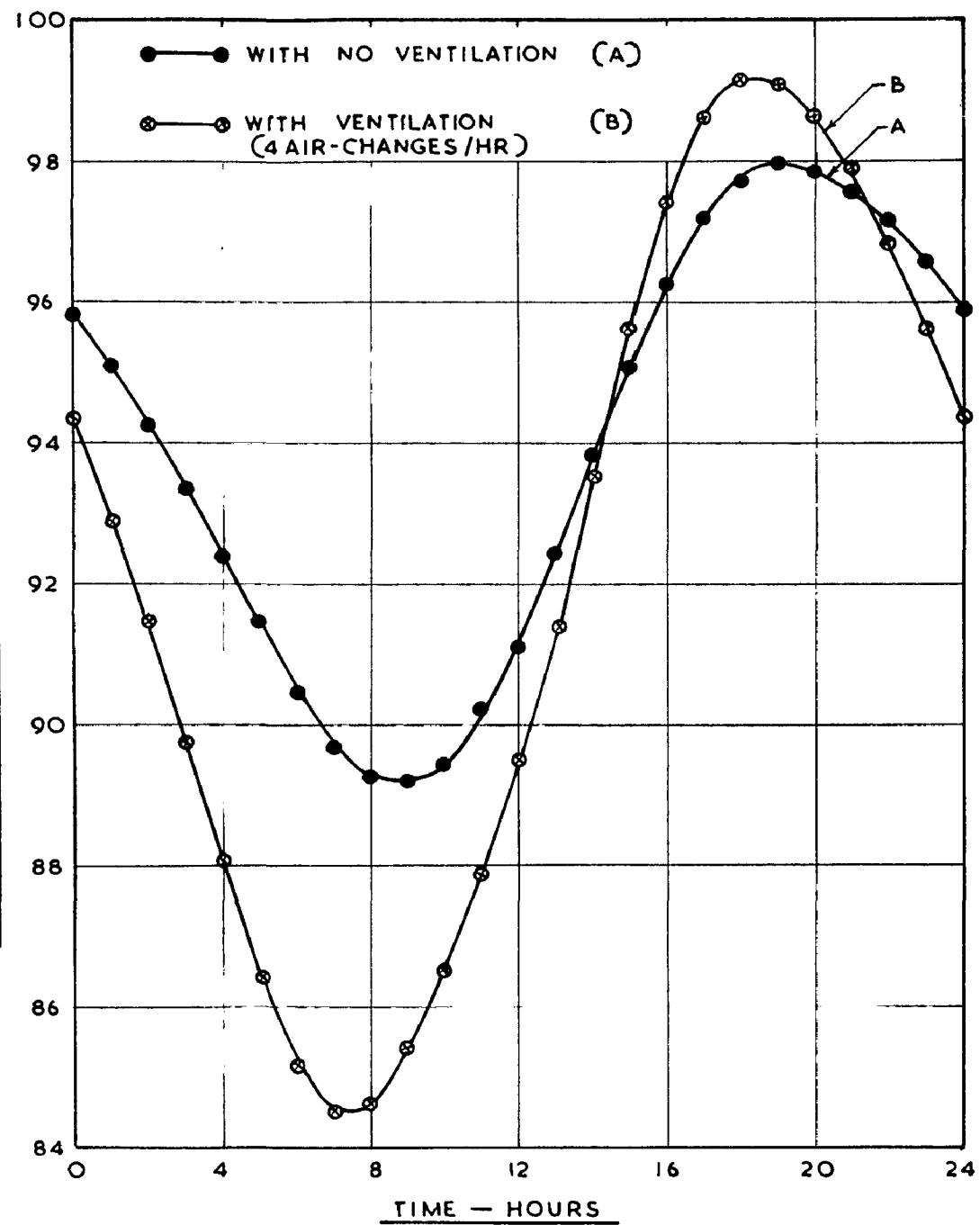
The hourly indoor air temperature variations are also

calculated and compared in Fig. (9.9). It can be seen from this, that the ventilation, will tend to reduce the temperature differences between the indoor and outdoor air, and also the time lag in the occurrence of their maxima.

### 9.8 Estimation of Fabric Cooling Loads for Air-Conditioning

Most of the methods employed for the estimation of fabric cooling loads do not take the internal masses into account. These are approximate methods and sizeable errors are likely to occur. The thermal circuit techniques (81) which do take these internal masses into account are complex and involve considerable labour. The transfer function approach will simplify the computations to a minimum. The application of transfer functions for these cooling load calculations, is explained below. For a given building, under any known climatic conditions, the amount of heat flux entering the building through all the external bounding surfaces can be obtained as illustrated in the previous examples. Let us suppose that the room air is to be maintained constant at a temperature  $75^{\circ}\text{F}$ ; we have to calculate the amount of heat that is to be extracted from the room. The steady state part of the cooling load is obtained by

$$\sum u_k A_k \{t_{sak(\text{mean})} - 75^{\circ}\}$$



EFFECT OF VENTILATION ON INDOOR AIR TEMPERATURES

FIG. 9.9

The total amount of the harmonic heat flux entering the room through all exposed building elements can be computed with transfer admittance function  $\vec{Y}_i$  in the same way, as done for non air conditioned buildings. All the amount of heat flux entering will not appear as the cooling load to the air-conditioning plant. The floor and the interior masses like partitions and furniture, will absorb a certain amount of the heat flux. Thus the cooling load transferred to the plant is the heat entering the enclosure less that absorbed by the internal masses. The quantity of heat flux absorbed by these internal elements can be obtained from the internal driving point admittance function  $\overleftarrow{Y}_i'$ .

As an illustration, the same room taken in the previous examples for unconditioned buildings, (Fig. 9.1(b)) with internal partition wall is considered here. The quantities of heat flow into the room, (steady state as well as harmonic components) and the quantity of heat flux absorbed by the floor and partition wall and the balance of the heat quantity to be extracted by the plant to maintain the indoor air temperature constant at  $75^{\circ}$  are calculated. The Fourier equations for the cooling load with and without considering the internal masses are given below.

1. Without taking the internal masses into account

$$Q(T) = 3836 + 5390 \cos(15t + 79^{\circ}) \\ + 788 \cos(30t + 252^{\circ}) \\ + 76 \cos(45t + 203^{\circ}) \\ + 50 \cos(60t + 153^{\circ})$$

### 2. Internal masses taken into account

$$\begin{aligned} Q(T) = & 8836 + 3245 \cos(15t + 79^\circ) \\ & + 455 \cos(30t + 252^\circ) \\ & + 46 \cos(45t + 203^\circ) \\ & + 29 \cos(60t + 161^\circ) \end{aligned}$$

The hourly variations of cooling load with and without for a condition, considering the internal masses are shown in Fig. (9.10).

Increase of internal masses will make this difference proportionately more. This method of approach gives more realistic picture of fabric cooling loads of buildings for air conditioning estimations.

### 0.9. Estimation of Outside Surface Temperatures

The above examples deal with the applications of only two thermal system functions namely, i) the transfer admittance function, and ii) internal driving point admittance function. The use of third function i.e., external driving point function ( $t_{os}/t_{sa}$ ) is explained here. It was mentioned earlier that outside surface temperatures of building elements are not usually available. These have to be determined by actual measurement. Different building elements attain different (outside) surface temperatures when exposed to same weather conditions. A knowledge of these external surface temperatures, is a prerequisite for computing the thermal stresses induced in the structural elements and

the expansions and contractions that are likely to take place during diurnal and seasonal variations. The external driving point functions will provide the external surface temperatures attained by different constructions under different climatic exposures. As an illustrative example, the case of (S.G) has been considered here. The  $\vec{\lambda}_o$  for the three types of roofs for the fundamental and three higher harmonics are given in Table (C.4). The Fourier equations of the external surface temperatures obtained for the three types of roofs for the corresponding sol-air temperature are given below.

#### Uninsulated roof

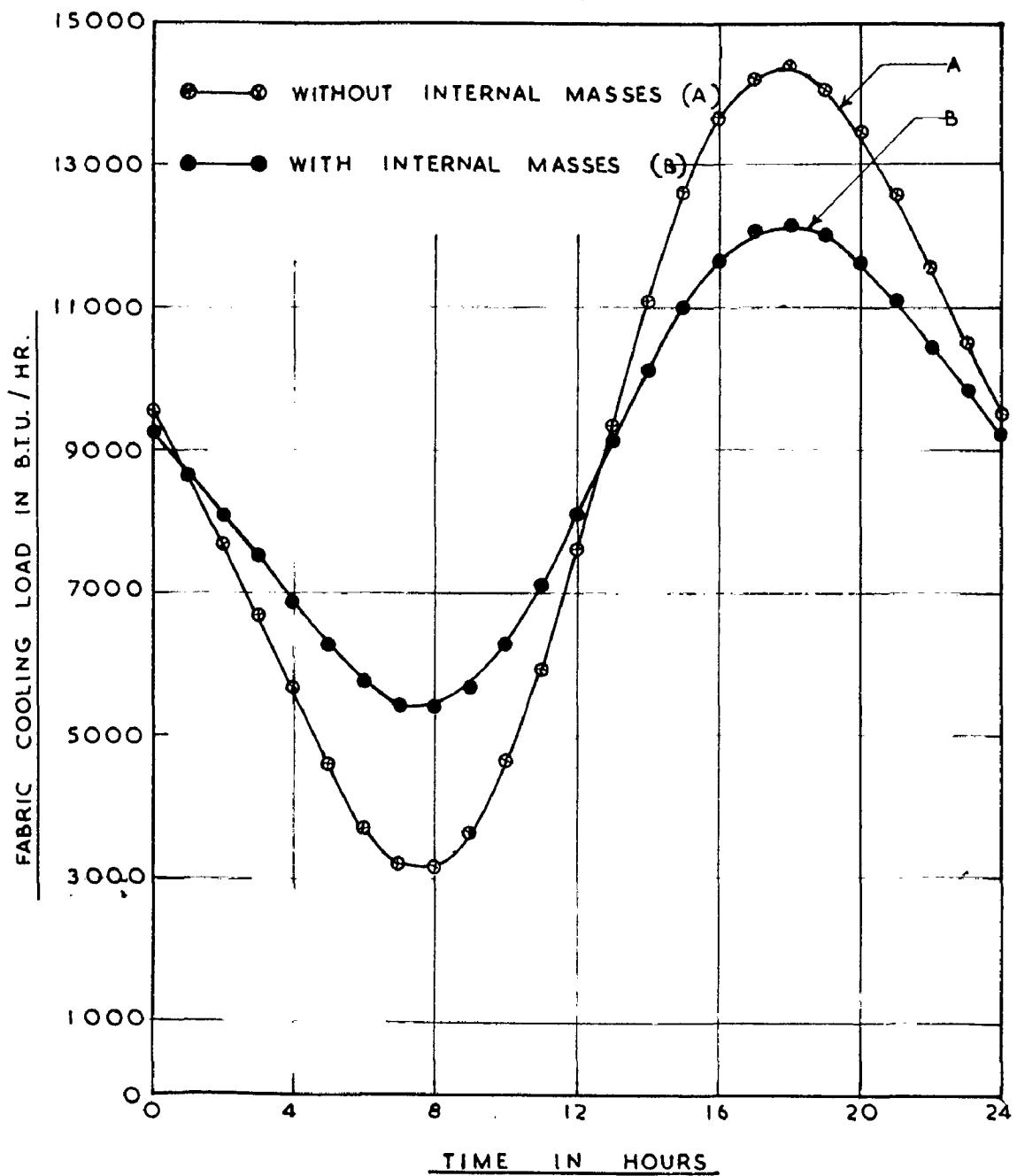
$$t_{os}(T) = 103.3 + 30.4 \cos(15t - 207^\circ) \\ + 9.1 \cos(30t - 200^\circ) \\ + 1.0 \cos(45t - 74^\circ) \\ + 1.3 \cos(30t - 6^\circ) \\ + 0.37 \cos(15t - 21^\circ)$$

#### Insulated type roof A

$$t_{os}(T) = 103.0 + 33.5 \cos(15t - 203^\circ) \\ + 12.4 \cos(30t - 20^\circ) \\ + 1.4 \cos(45t - 76^\circ) \\ + 1.8 \cos(60t - 11^\circ) \\ + 0.00 \cos(15t - 31^\circ)$$

#### Insulated type roof B

$$t_{os}(T) = 104.1 + 37.4 \cos(15t - 210^\circ) \\ + 11.7 \cos(30t - 23^\circ) \\ + 1.3 \cos(45t - 33^\circ) \\ + 1.6 \cos(60t - 23^\circ) \\ + 0.03 \cos(15t - 45^\circ)$$



EFFECT OF INTERNAL MASSES ON FABRIC COOLING LOAD

FIG. 9·10

the expansions and contractions that are likely to take place during diurnal and seasonal variations. The external driving point functions will provide the external surface temperatures attained by different constructions under different climatic exposures. As an illustrative example, the same case of (0.6) has been considered here. The  $\vec{\lambda}_o$  for the three types of roofs for the fundamental and three higher harmonics are given in Table (C.4). The Fourier equations of the external surface temperatures obtained for the three types of roofs for the corresponding sol-air temperature are given below.

#### Uninsulated roof

$$t_{os}(T) = 103.3 + 20.4 \cos(15t - 207^\circ) \\ + 0.1 \cos(30t - 230^\circ) \\ + 1.0 \cos(45t - 74^\circ) \\ + 1.2 \cos(60t - 6^\circ) \\ + 0.37 \cos(15t - 21^\circ)$$

#### Insulated type roof A

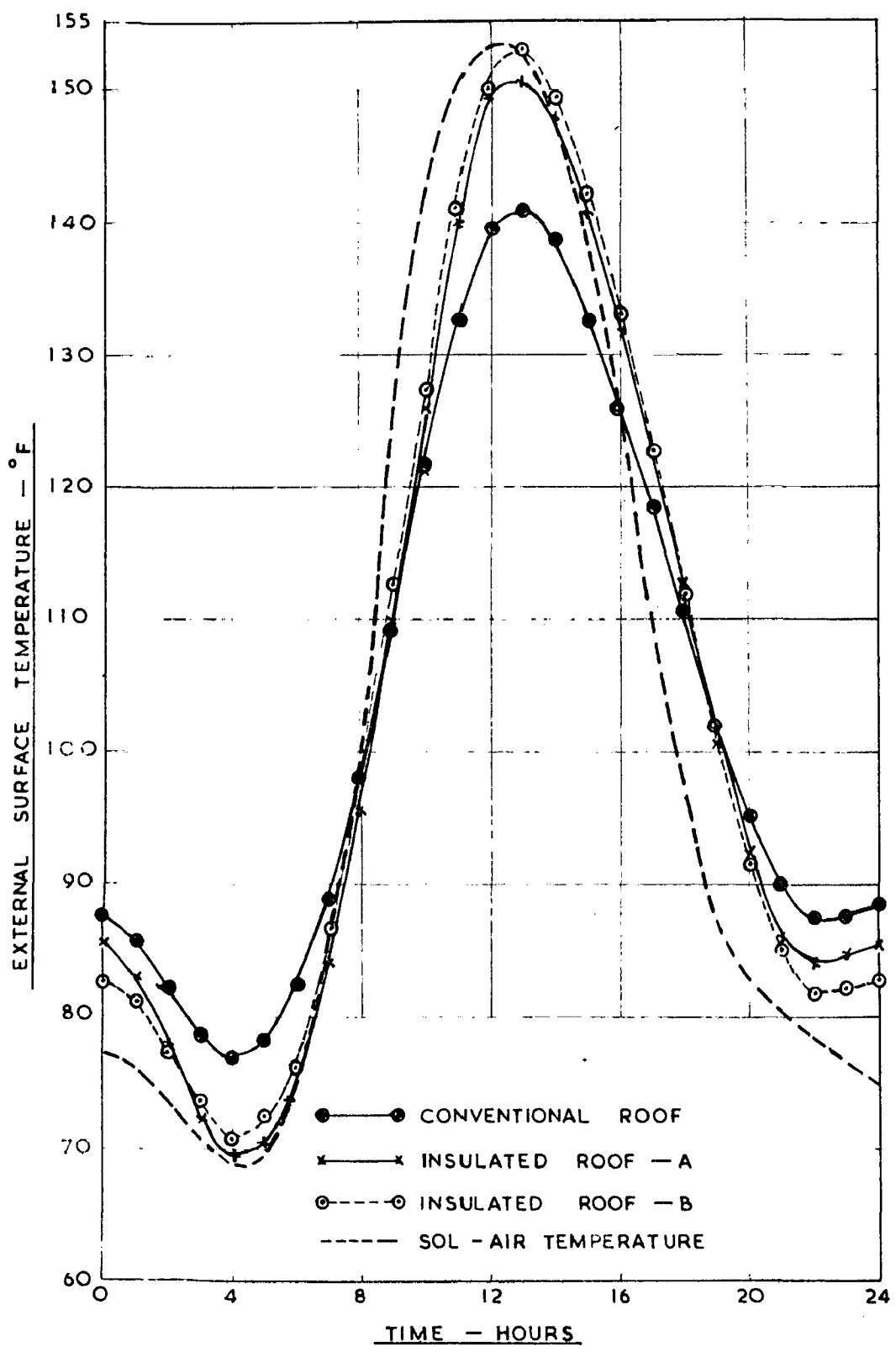
$$t_{os}(T) = 103.9 + 33.5 \cos(15t - 203^\circ) \\ + 13.4 \cos(30t - 20^\circ) \\ + 1.4 \cos(45t - 73^\circ) \\ + 1.8 \cos(60t - 11^\circ) \\ + 0.00 \cos(15t - 31^\circ)$$

#### Insulated type roof B

$$t_{os}(T) = 104.1 + 37.4 \cos(15t - 210^\circ) \\ + 11.7 \cos(30t - 23^\circ) \\ + 1.3 \cos(45t - 63^\circ) \\ + 1.6 \cos(60t - 23^\circ) \\ + 0.03 \cos(15t - 45^\circ)$$

The hourly temperature variations were calculated from those equations and compared in Fig. (9.11). The insulated type roofs will attain higher temperatures than the uninsulated ones, when the insulating layer is on the top.

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EXTERNAL SURFACE TEMPERATURES OF DIFFERENT TYPES  
OF ROOFS

FIG. 9·11

## C O N C L U S I O N S

CONCLUSIONS

In this thesis an adequate quantitative data, on the transient thermal characteristics of building sections, and a simple and flexible computational procedure, which will enable a designer to calculate the heat transfer and temperatures of enclosures, under given conditions, are established.

An analysis of the electrical analogue studies made on the transient thermal behaviour of building elements, with special reference to a) the effect of surface heat transfer coefficients, b) the effect of different combinations of layers of various materials and c) lumped network representation of the building section, lead to the following conclusions :-

1. The steady state sinusoidal approach adopted here separates the effects of the climatic factors from the thermal characteristics of building elements, so that each could be studied independently and then combined to get the thermal behaviour of the fabric under any climatic conditions.
2. A knowledge of transfer function only, is not sufficient to specify the transient thermal behaviour of any building element completely and two more thermal system

functions viz., i) external driving point function and ii) internal driving point function are also required. With these three system functions the surface temperatures and heat fluxes of the outer and inner surfaces can easily be obtained for any boundary condition.

3. a) The influence surface heat transfer coefficients on the thermal system functions depends upon both  $\lambda C$  and  $R$  of a given building section.
- b) The external and internal surface heat transfer coefficients influence their respective driving point functions for sections with large  $\lambda C$  and  $R$ . For thin sections (with small  $\lambda C$  and  $R$ ) however, both surface coefficients affect both driving point functions.
- c) The transfer functions of any building section are affected by both the surface coefficients considerably.
4. a) In hot climates, placing the insulation on the exposed side is more advantageous than placing it inside.
- b) For various possible arrangements, the greatest damping occurs when insulating layers alternate with layers of dense materials.

- c) Backing layers have considerable influence on the decrement factors of the front layers while the converse is not true.
  - d) If the heat flow direction is reversed (i.e. from inside to outside) the overall decrement factor will be reduced by the ratio of the outside to inside surface heat transfer coefficients ( $h_o/h_i$ ) while the phase lag is not altered.
  - e) The decrement factor of the individual layers in a composite section will not be the same when the direction of heat flow is reversed.
  - f) It is not always possible to obtain an equivalent homogeneous, construction, which will have the same thermal system functions of heat of a composite construction.
3. The number of lumps of T network required for an adequate representation of a building section, is proportional to the square root of the thermal time constant (RC) of the section.
6. This thesis has several practical uses. For instance :
- a) Reference data on transient thermal characteristics of building components commonly used in India,

given in Appendices III and IV, enables to make the right choice of the materials for a given situation.

- b) The data given on external driving point functions provides a basis for the easy computation of outside surface temperatures attained by building components under different climatic exposures.
7. With the aid of the data on the matrix coefficients of the distributed and lumped systems (Appendix I) it is possible to evaluate the errors due to lumping on any of the thermal quantities of interest (surface temperatures, heat fluxes, or the thermal system functions) for any homogeneous layer.
8. Moreover as the method, put forth for the computation of indoor air temperatures of enclosures, utilises pretabulated data on thermal system functions and 'U' values; almost all types of conditions that may be encountered in practice can be analysed by means of the graphs and tables given here, without constructing an analogue for each case. Thus the method suggested here offers a number of advantages over the analytical or field measurement techniques.

A P P E N D I C E S

I TO IV

## APPENDIX I

### TRANSIENT INVERSES OF DISTRIBUTED AND LIMITED NETWORKS

(FOR SINUSOIDAL TEMPERATURE VARIATIONS)

#### Distributed System

$$\begin{bmatrix} A_D & B_D \\ C_D & D_D \end{bmatrix} = \begin{bmatrix} \cosh(1+\gamma)\theta & R \sinh(1+\gamma)\theta \\ (1+\gamma)\theta \sinh(1+\gamma)\theta & \cosh(1+\gamma)\theta \end{bmatrix}$$

where  $\omega = \sqrt{\frac{\omega CR}{2}}$

$$= \frac{2\pi}{P}$$

P = Period of the thermal cycle.

C = Total thermal capacity ( $\text{L}^2\text{F}$ )

R = Total thermal resistance ( $\text{L}/\text{K}$ )

L = Thickness in Ft.

K = Thermal conductivity of the material  
in Btu/Ft./hr./°F

P = Density in Lb/Cu.Ft.

s = Specific heat in Btu/Lb/°F

$$\cosh(1+\gamma)\theta = \left\{ \left( 1 - \frac{\theta^4}{0} + \frac{\theta^8}{3620} \dots \right) + \gamma \left( \theta^2 - \frac{\theta^6}{90} + \dots \right) \right\}$$

$$\frac{\sinh(1+\gamma)\theta}{(1+\gamma)\theta} = \left\{ \left( 1 - \frac{\theta^4}{30} + \frac{\theta^8}{32330} \dots \right) + \gamma \left( \frac{\theta^2}{3} - \frac{\theta^6}{620} \dots \right) \right\}$$

$$\text{and } (1+\gamma)\theta \sinh(1+\gamma)\theta = 2 \left\{ \left( -\frac{\theta^4}{3} + \frac{\theta^8}{630} \dots \right) + \gamma \left( \theta^2 - \frac{\theta^6}{30} \dots \right) \right\}$$

### T - Network

The transfer matrices for one lump and N - lumped circuits (Figs. 1 & 2 respectively) are given by :-

#### One Lump Circuit

$$\begin{bmatrix} A_T & B_T \\ C_T & D_T \end{bmatrix} = \begin{bmatrix} \left(1 + \frac{j\omega CR}{2}\right) & R \left(1 + \frac{j\omega CR}{4}\right) \\ j\omega C & \left(1 + \frac{j\omega CR}{2}\right) \end{bmatrix}$$

#### N Lumped Circuit

$$\begin{bmatrix} A_T & B_T \\ C_T & D_T \end{bmatrix} = \begin{bmatrix} \left(1 + \frac{j\omega CR}{2N^2}\right) & \frac{R}{N} \left(1 + \frac{j\omega CR}{4N^2}\right) \\ -\frac{j\omega C}{N} & \left(1 + \frac{j\omega CR}{2N^2}\right) \end{bmatrix}$$

where C = Total thermal capacitance of the building element.

R = Total thermal resistance of the building element.

T - NET WORK

I - LUMP CIRCUIT

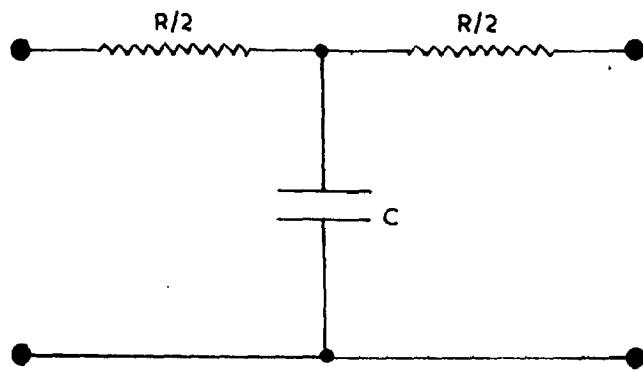


FIG. 1 .

N - LUMPED CIRCUIT

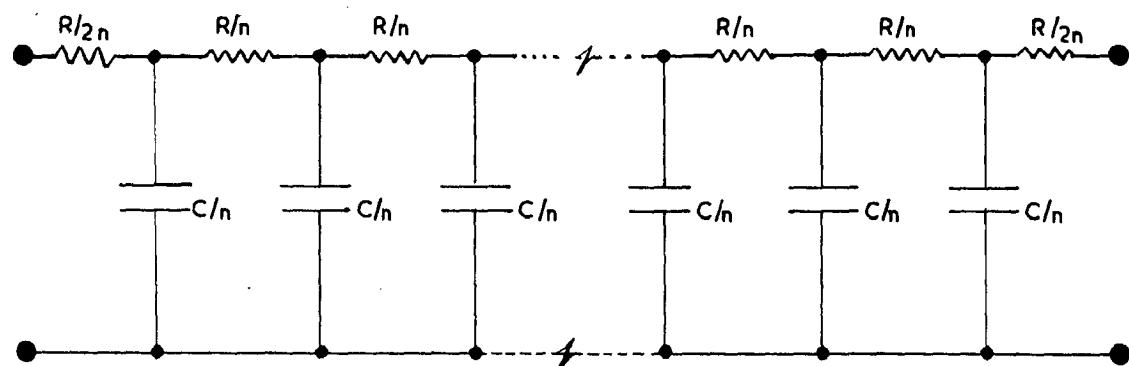


FIG. 2 .

### π - Network

The transfer matrices for one lump and N lumped circuits (Figs. 3 & 4 respectively) are given by :-

#### One Lump Circuit

$$\begin{bmatrix} A_{\pi} & B_{\pi} \\ C_{\pi} & D_{\pi} \end{bmatrix} = \begin{bmatrix} \left(1 + \frac{j\omega_{CR}}{2}\right) & R \\ j\omega_C \left(1 + \frac{j\omega_{CR}}{4}\right) & \left(1 + \frac{j\omega_{CI}}{2}\right) \end{bmatrix}$$

#### N Lumped Circuit

$$\begin{bmatrix} A_{\pi} & B_{\pi} \\ C_{\pi} & D_{\pi} \end{bmatrix} = \begin{bmatrix} \left(1 + \frac{j\omega_{CR}}{\pi N^2}\right) & \frac{R}{N} \\ \frac{j\omega_C}{\pi} \left(1 + \frac{j\omega_{CI}}{4N^2}\right) & \left(1 + \frac{j\omega_{CR}}{\pi N^2}\right) \end{bmatrix}$$

where  $C$  = Total thermal capacitance of the building element.

$R$  = Total thermal resistance of the building element.

## TT - NET WORK

### I - LUMP CIRCUIT

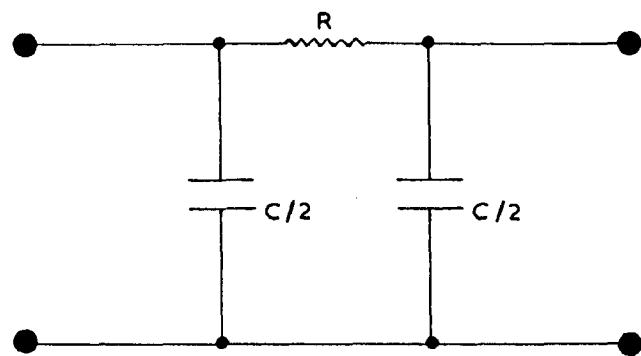


FIG. 3.

### N - LUMPED CIRCUIT

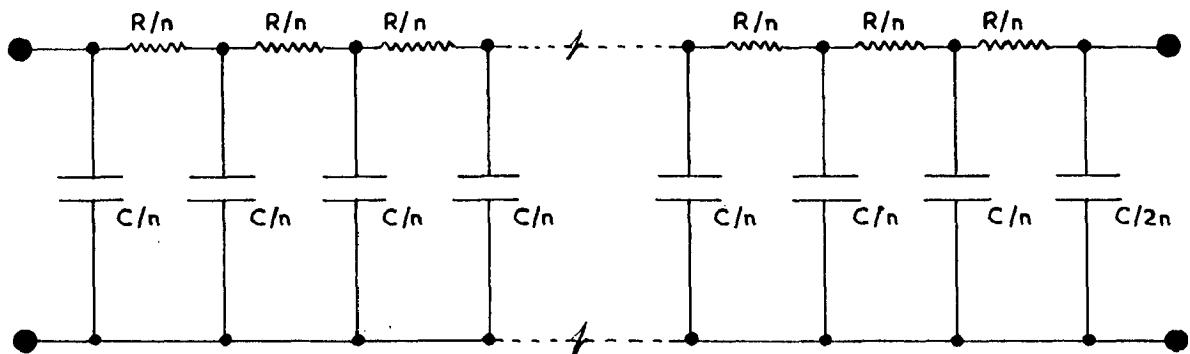


FIG.4

L - Network

The transfer matrices for one lump and N lumped circuits (Figs. 5 & 6 respectively) are given by :-

One Lump Circuit

$$\begin{bmatrix} A_L & B_L \\ C_L & D_L \end{bmatrix} = \begin{bmatrix} (1 + j\omega CR) & R \\ j\omega C & 1 \end{bmatrix}$$

N Lumped Circuit

$$\begin{bmatrix} A_L & B_L \\ C_L & D_L \end{bmatrix} = \begin{bmatrix} \left(1 + \frac{j\omega CR}{N^2}\right) & -\frac{R}{N} \\ \frac{j\omega C}{N} & 1 \end{bmatrix}$$

where  $C$  = Total thermal capacitance of the building element.

$R$  = Total thermal resistance of the building element.

L - NET WORK

I - LUMP CIRCUIT

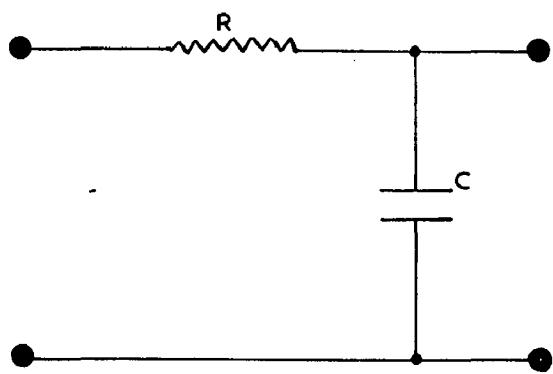


FIG. 5.

N - LUMPED CIRCUIT

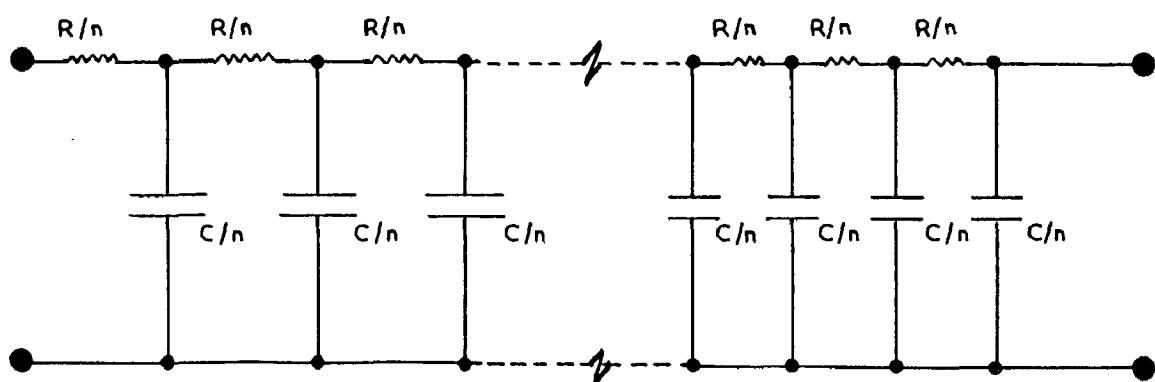


FIG. 6.

CELESTE LEE OF BOSTON, MASS., TO DIRECT THE  
SCHOOL OF MUSIC, BOSTON UNIVERSITY.

TRANSPORTS, MIGRATION, AND CONNECTIVITIES OF DIADURID FISHES IN ECUADOR

ESTATE PLANNING - A



## (e) Matrix Conditioning - G.R.A.

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Physical size - connection RC	Diameter - rotating in degrees	No. of Lancs	Lanc Technique	Lanc Mechanics	Lanc Capacitance	Lanc in degrees	Lanc in degrees
0.1	0.00242	50.23					
0.5	0.1252	23.30					
1.0	0.2502	23.47					
5.0	1.3247	103.00					
10.0	2.6063	115.00					
20.0	5.0140	137.01					
40.0	10.0240	191.00					
80.0	20.0240	331.00					
160.0	40.0240	662.00					
320.0	80.0240	1324.00					
640.0	160.0240	2648.00					
1280.0	320.0240	5296.00					
2560.0	640.0240	10592.00					
5120.0	1280.0240	21184.00					
10240.0	2560.0240	42368.00					
20480.0	5120.0240	84736.00					
40960.0	10240.0240	169472.00					
81920.0	20480.0240	338944.00					
163840.0	40960.0240	677888.00					
327680.0	81920.0240	1355760.00					
655360.0	163840.0240	2711520.00					
1310720.0	327680.0240	5423040.00					
2621440.0	655360.0240	10846080.00					
5242880.0	1310720.0240	21692160.00					
10485760.0	2621440.0240	43384320.00					
20971520.0	51485760.0240	86768640.00					
41943040.0	102971520.0240	173537280.00					
83886080.0	205943040.0240	347074560.00					
167772160.0	411886080.0240	694149120.00					
335544320.0	823772160.0240	1388298240.00					
671088640.0	1647544320.0240	2776596480.00					
1342177280.0	3295088640.0240	5553192960.00					
2684354560.0	6580177280.0240	11106385920.00					
5368709120.0	13160354560.0240	22212771840.00					
10737418240.0	26320709120.0240	44425543680.00					
21474836480.0	52641418240.0240	88851087360.00					
42949672960.0	105282836480.0240	177702174720.00					
85899345920.0	210565672960.0240	355404349440.00					
171788691840.0	421131345920.0240	710808698880.00					
343577383680.0	842262691840.0240	142161739760.00					
687154767360.0	1684545783680.0240	284323479520.00					
1374309534720.0	3369091567360.0240	568646959040.00					
2748619069440.0	673818291567360.0240	1137293858080.00					
5497238138880.0	13476365891567360.0240	2274587716160.00					
10994476277760.0	2695273178291567360.0240	4549175432320.00					
2198895255520.0	539054635658291567360.0240	9098350864640.00					
4397785511040.0	1078109271178291567360.0240	1819670216160.00					
8795571022080.0	215621854235658291567360.0240	3639340432320.00					
17591142044160.0	43124370847135658291567360.0240	7278680864640.00					
35182284088320.0	8624874169427135658291567360.0240	1455736172920.00					
70364568176640.0	1724974833885427135658291567360.0240	2911472345840.00					
140729136353280.0	344994966777085427135658291567360.0240	5822944691680.00					
281458272706560.0	689989933554175427135658291567360.0240	1164588938320.00					
562916545413120.0	1379979867108355427135658291567360.0240	2329176876640.00					
1125833090826240.0	27599597342167108355427135658291567360.0240	4658353753280.00					
2251666181652480.0	551991946843342167108355427135658291567360.0240	9316707506560.00					
4503332363304960.0	11039839348866843342167108355427135658291567360.0240	1863341501280.00					
9006664726609920.0	220796786977338866843342167108355427135658291567360.0240	3726683002560.00					
18013329453219840.0	4415935739546754738866843342167108355427135658291567360.0240	7453366005120.00					
36026658906439680.0	88318714790935094738866843342167108355427135658291567360.0240	14906732000240.00					
72053317812879360.0	17663742981987018738866843342167108355427135658291567360.0240	29813464000480.00					
14410663562559680.0	353274859639540374738866843342167108355427135658291567360.0240	59626928000960.00					
28821327125119360.0	70654971927858074738866843342167108355427135658291567360.0240	119253856001920.00					
57642654250238720.0	14130994385571614738866843342167108355427135658291567360.0240	238507712003840.00					
115285388504577440.0	28261828771143229738866843342167108355427135658291567360.0240	477015424007680.00					
230570777009154880.0	56523657542286459738866843342167108355427135658291567360.0240	954030848015360.00					
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922283108036619520.0	226094630169065839738866843342167108355427135658291567360.0240	381612336061440.00					
184456621607323840.0	452189260338131679738866843342167108355427135658291567360.0240	763224672032880.00					
368913243214647680.0	904378520676263359738866843342167108355427135658291567360.0240	1526449344065760.00					
737826486429295360.0	1808757041352526719738866843342167108355427135658291567360.0240	3052898688131520.00					
1475652972858590720.0	3617514082705053439738866843342167108355427135658291567360.0240	6105797376263040.00					
2951305945717181440.0	7235028165410106879738866843342167108355427135658291567360.0240	12215957524526080.00					
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23610447565737471520.0	5788225327528087039738866843342167108355427135658291567360.0240	97727660196208640.00					
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3022136968414396434560.0	741292845123593861119738866843342167108355427135658291567360.0240	125091405146025280.00					
6044273936828792869120.0	148258569224718772239738866843342167108355427135658291567360.0240	250182810292050560.00					
12088547736457965738240.0	296517138549437544479738866843342167108355427135658291567360.0240	500365620584101120.00					
24177095472915931476480.0	593034277098875088959738866843342167108355427135658291567360.0240	100073124168202240.00					
48354190945831862952960.0	118606855419775177819738866843342167108355427135658291567360.0240	200146248336404480.00					
96708381891663725905920.0	237213710839550355639738866843342167108355427135658291567360.0240	400292496672808960.00					
19341676378332745811840.0	474427421679100711279738866843342167108355427135658291567360.0240	800584993345617920.00					
38683348756665491763680.0	948854843358201422559738866843342167108355427135658291567360.0240	160116998691235840.00					
77366697513330983527360.0	189770968671640284519738866843342167108355427135658291567360.0240	320233997382471680.00					
154733395026661967054720.0	379541937343280569039738866843342167108355427135658291567360.0240	640467994764943360.00					
309466790053323934109440.0	758883874686561138079738866843342167108355427135658291567360.0240	128093598953986720.00					
61893358010664786821880.0	1517767549773122270159738866843342167108355427135658291567360.0240	256187197907973440.00					
12378671611329415643760.0	3035535099546244540319738866843342167108355427135658291567360.0240	512374395815946880.00					
24757343222678831127520.0	6071070199092489080639738866843342167108355427135658291567360.0240	102474791163189360.00					
49514686445357662255040.0	1214214039818497816129738866843342167108355427135658291567360.0240	204949582326378720.00					
99029372895313124500080.0	2428428079636995632259738866843342167108355427135658291567360.0240	409898165532757440.00					
19805874578662624900160.0	4856856159273991264519738866843342167108355427135658291567360.0240	819796331065514880.00					
39611749157325249800320.0	9713709318827982529039738866843342167108355427135658291567360.0240	163959266531029760.00					</

(a) REVIEW COMPLEX EK - D

APPENDIX II

DERIVATION OF THE THERMAL SYSTEM FUNCTIONS OF BUILDING

ELEMENTS BY MATRIX ANALYSIS

Von Gorchum (23) and Pipes (24) have shown that a homogeneous building element is analogous to a two element ( $R \& C$ ) four terminal passive network (quadrupole) and it is possible to use all the mathematical techniques employed in electrical circuit analyses in the study of linear heat conduction problems.

Analogous to the passive network system functions we have a set of thermal system functions, which will specify completely the transient behaviour of a building element. The system function is defined as a function representing the ratio of the Laplace transformations of response variable to excitation variable. In the formation of this ratio it is assumed that all the initial conditions within the system are zero.

There are two types of system functions of interest, namely (i) the transfer functions and (ii) the driving point functions. Each of these may have the dimensions of impedance, admittance, in addition to the dimensionless quantities. For a driving point function the input and output (excitation and response) are measured at the same pair of terminals, while for a transfer function the input and output are measured at two different pairs of terminals.

The transfer functions are given as :-

- i)  $\frac{t_1}{t_0}$  = temperature ratio.
- ii)  $\frac{q_1}{q_0}$  = heat flux ratio.
- iii)  $\frac{q_1}{t_0}$  = transfer admittance.
- iv)  $\frac{t_1}{q_0}$  = transfer impedance.

The driving point functions are given as :-

- i)  $\frac{t_0}{q_0}$  = external driving point impedance.
- ii)  $\frac{q_0}{t_0}$  = external driving point admittance.
- iii)  $\frac{t_1}{q_1}$  = internal driving point impedance.
- iv)  $\frac{q_1}{t_1}$  = internal driving point admittance.

The ratios of the input to output are called the reciprocal system functions. All these functions are not independent and only three functions (one transfer and two driving point) are sufficient to characterise a building element. The rest of the functions can be obtained from these three as they are interrelated.

The two linear equations that relate the input and output temperatures and heat fluxes (analogous to the quadri-pole) are given by :-

$$t_o = A t_i + B q_i \quad \dots (1)$$

$$q_o = C t_i + D q_i \quad \dots (2)$$

where  $A$ ,  $B$ ,  $C$  and  $D$  are the general circuit parameters.  $A$  and  $D$  are dimensionless numbers,  $B$  and  $C$  have dimensions of impedance and admittance respectively. These general circuit parameters are the reciprocal system functions and are given by :-

$$A = \frac{t_o}{t_i} \quad q_i = 0 \quad \text{i.e., open circuit output.}$$

$$B = \frac{t_o}{q_i} \quad t_i = 0 \quad \text{i.e., short circuit output.}$$

$$C = \frac{q_o}{t_i} \quad q_i = 0 \quad \text{i.e., open circuit output.}$$

$$D = \frac{q_o}{q_i} \quad t_i = 0 \quad \text{i.e., short circuit output.}$$

For a sinusoidal excitation all these parameters will be complex quantities.

The above equations (1) and (2) may also be expressed in the matrix form as :-

$$\begin{bmatrix} t_o \\ q_o \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} t_i \\ q_i \end{bmatrix} \quad \dots (3)$$

where

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix}$$

is the transfer matrix of the building element.

The determinant of the above matrix is equal to unity as the building element is considered as a passive network. i.e.,  $(AD - BC) = 1$

For a symmetrical network  $A = D$

In addition to the above matrix equation (3) there are two other useful forms of matrices, relating the input and output quantities of a network namely the impedance matrix and the admittance matrix. The corresponding matrix equations are given by :-

$$\begin{bmatrix} t_0 \\ t_1 \end{bmatrix} = \begin{bmatrix} z_{11} & z_{12} \\ z_{21} & z_{22} \end{bmatrix} \begin{bmatrix} q_0 \\ q_1 \end{bmatrix} \dots (4)$$

and

$$\begin{bmatrix} q_0 \\ q_1 \end{bmatrix} = \begin{bmatrix} y_{11} & y_{12} \\ y_{21} & y_{22} \end{bmatrix} \begin{bmatrix} t_0 \\ t_1 \end{bmatrix} \dots (5)$$

The relationship between the coefficients of the transfer matrix and the Z's and Y's of the impedance and admittance matrices may be simply derived as :-

$$z_{11} = \frac{A}{C} \quad z_{12} = -\frac{1}{C}$$

$$z_{21} = \frac{1}{C} \quad z_{22} = -\frac{D}{C}$$

$$\text{and } Y_{11} = \frac{D}{B} \quad Y_{12} = -\frac{1}{B}$$

$$Y_{21} = \frac{1}{B} \quad Y_{22} = -\frac{A}{B}$$

The overall transfer matrices and the system functions of a network will not only depend upon the general circuit parameters but also on the nature of the termination. The heat transfer between the ambient air and the building element, takes place through the surface film conductances ( $h_0$  and  $h_1$ ). The reciprocals of those surface conductances ( $1/h_0$  and  $1/h_1$ ) may be considered as analogous to the source impedance and load respectively in electrical circuits.

The overall thermal system functions, for building elements in terms of general circuit parameters and the surface heat transfer coefficients, for boundary conditions of practical importance, are derived below :-

- 1) External surface subjected to temperature variations (sinusoidal) while the heat flow from the inside surface is prevented. This condition is illustrated by Fig. 1. The corresponding matrix equation is given by :-

$$\begin{bmatrix} t_{os} \\ q_{os} \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} t_{is} \\ 0 \end{bmatrix} \dots (6)$$

$$\text{This gives } t_{os} = A t_{is}$$

$$q_{os} = C t_{is}$$

### Transfer Function

$$\vec{\lambda}_i = \lambda_i / \varphi_i = \frac{t_{is}}{t_{os}} = \frac{1}{A} \quad \dots(7)$$

### Driving Point Admittance Function

$$\vec{Y}_o = Y_o / \varphi_o = \frac{q_{os}}{t_{os}} = \frac{C t_{is}}{t_{os}} = \frac{C}{A} \quad \dots(8)$$

ii) Internal Surface Temperature Variation with a Constant Indoor Air Temperature. (Fig. 2).

This condition is expressed by the Matrix equation

$$\begin{bmatrix} t_{os} \\ q_{os} \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} t_{is} \\ q_{is} \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} t_{is} \\ q_{is} \end{bmatrix} \quad \dots(9)$$

where  $t_{is} = q_{is}/h_1$

$$q_{is} = q_{ia} = t_{is} \times h_1$$

Then we get

$$t_{os} = t_{is} (A + B h_1)$$

$$q_{os} = t_{is} (C + D h_1)$$

Transfer Function is given by :-

$$\vec{\lambda}_i = \lambda_i / \varphi_i = \frac{t_{is}}{t_{os}} = \frac{1}{(A + B h_1)} \quad \dots(10)$$

Transfer admittance function is given by

$$\vec{Y}_i = Y_i / \varphi_i = \frac{q_{is}}{t_{os}} = \frac{t_{is} \times h_1}{t_{os}} = \frac{h_1}{(A + B h_1)} = \vec{\lambda}_i h_1 \quad \dots(11)$$

Driving point admittance function is given by :-

$$\vec{Y}_o = Y_o / \psi_o = \frac{q_{os}}{t_{os}} = \frac{t_{is} (C + D h_1)}{t_{os}} = \frac{(C + D h_1)}{(A + B h_1)} \dots (12)$$

iii) Sol-Air temperature input with a constant indoor air temperature. (Fig. 3).

For this condition the matrix equation takes the form

$$\begin{bmatrix} t_{sa} \\ q_{sa} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} t_{os} \\ q_{os} \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} t_{is} \\ q_{is} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_1 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ q_{ia} \end{bmatrix} \dots (13)$$

$$= \begin{bmatrix} 1 & 1/h_0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} t_{os} \\ q_{os} \end{bmatrix}$$

$$\text{where } t_{os} = t_{is} (A + B h_1)$$

$$q_{os} = t_{is} (C + D h_1)$$

$$\text{Therefore } t_{sa} = t_{is} (A + B h_1 + C/h_0 + D h_1/h_0)$$

$$q_{sa} = t_{is} (C + D h_1)$$

The transfer function is obtained as :-

$$\vec{\lambda}_i = \lambda_i / \psi_i = \frac{t_{is}}{t_{sa}} = \frac{1}{(A + B h_1 + C/h_0 + D h_1/h_0)} \dots (14)$$

Transfer admittance function is obtained as :-

$$\vec{Y}_i = Y_i / \psi_i = \frac{q_{is}}{t_{sa}} = \frac{t_{is} \times h_1}{t_{sa}} = \frac{h_1}{(A + B h_1 + C/h_0 + D h_1/h_0)}$$

Driving point admittance function is given by :-

$$\vec{Y}_o = \frac{\vec{I}_o / \psi_o}{t_{sa}} = \frac{q_{os}}{t_{sa}} = \frac{t_{is} (C + D h_i)}{t_{sa}} = \vec{\lambda}_i (C + D h_i)$$

Driving point function is given by :-

$$\vec{\lambda}_o = \lambda_o \left[ \frac{t_{os}}{t_{sa}} \right] = \frac{t_{is} (\Delta + B h_i)}{t_{sa}} = \vec{\lambda}_i (\Delta + B h_i)$$

- iv) Indoor air temperature is variable and external air temperature is constant (heat flow direction is reversed) (Fig. 4).

The matrix equation for this condition is given by:-

$$\begin{bmatrix} t_{ia} \\ -q_{ia} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_i \\ 0 & 1 \end{bmatrix} \begin{bmatrix} t_{is} \\ q_{is} \end{bmatrix} = \begin{bmatrix} D & D \\ C & A \end{bmatrix} \begin{bmatrix} t_{is} \\ q_{is} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_o \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 0 \\ -q_{ea} \end{bmatrix} \quad \dots(15)$$

Then we get

$$t_{os} = -q_{ea}/h_o$$

$$-q_{os} = -q_{ea} = t_{os} \times h_o$$

and

$$\begin{bmatrix} t_{ia} \\ -q_{ia} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_i \\ 0 & 1 \end{bmatrix} \begin{bmatrix} t_{is} \\ -q_{is} \end{bmatrix}$$

where

$$t_{is} = t_{os} (D + B h_o)$$

$$-q_{is} = t_{os} (C + \Delta h_o) \quad \dots(16)$$

$$\text{Therefore } t_{ia} = t_{os} (D + B h_o + C/h_i + \Delta h_o/h_i)$$

$$-q_{ia} = t_{os} (C + \Delta h_o)$$

The transfer function is obtained as :-

$$\overleftarrow{\lambda}_o = \lambda_o / -\phi'_o = \frac{t_{os}}{t_{ia}} = \frac{1}{h_0/h_1 (A + B h_1 + C/h_0 + D h_1/h_0)}$$

Transfer admittance function will be

$$\overleftarrow{Y'_o} = Y'_o / -\psi'_o = \frac{-q_{os}}{t_{ia}} = \frac{t_{os} \times h_0}{t_{ia}} = \frac{h_1}{(A + B h_1 + C/h_0 + D h_1/h_0)}$$

Driving point admittance function is obtained as

$$\overleftarrow{Y'_1} = Y'_1 / -\psi'_c = \frac{-q_{1o}}{t_{ia}} = \frac{t_{os} (C + A h_0)}{t_{ia}} = \lambda'_o (C + A h_0)$$

and the driving point function will be

$$\overleftarrow{\lambda'_c} = \lambda'_c / -\phi'_c = \frac{t_{1o}}{t_{ia}} = \frac{t_{os} (D + B h_0)}{t_{ia}} = \lambda'_o (D + B h_0)$$

v) Variable air temperature both sides (external and internal) (Fig. 5).

The matrix equation for this condition is given by

$$\begin{bmatrix} t_{os} \\ q_{os} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} t_{1o} \\ q_{1o} \end{bmatrix} = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} t_{1o} \\ q_{1o} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_1 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} t_{ia} \\ q_{ia} \end{bmatrix} \quad \dots (17)$$

Then

$$t_{1o} = t_{ia} + q_{ia}/h_1$$

$$q_{1o} = q_{ia}$$

... (18)

$$\text{and} \quad t_{os} = t_{ia} A + q_{ia}/h_1 (A + B h_1)$$

$$q_{os} = t_{ia} C + q_{ia}/h_1 (C + D h_1) \quad \dots (19)$$

In the above equations  $q_{1a}$  is an unknown quantity and has to be expressed in terms of known parameters. This can be obtained from the matrix equation given below :-

$$\begin{bmatrix} t_{0a} \\ q_{0a} \end{bmatrix} = \begin{bmatrix} \frac{1}{h_0} (C + A h_0) & \frac{1}{h_1} (A + B h_1 + \frac{C}{h_0} + D \frac{h_1}{h_0}) \\ C & \frac{1}{h_1} (C + D h_1) \end{bmatrix} \begin{bmatrix} t_{1a} \\ q_{1a} \end{bmatrix}$$

This gives  $t_{0a}$  and  $q_{0a}$  as

$$t_{0a} = \frac{t_{1a}}{h_0} (C + A h_0) + \frac{q_{1a}}{h_1} (A + B h_1 + \frac{C}{h_0} + D \frac{h_1}{h_0})$$

$$q_{0a} = t_{1a} C + \frac{q_{1a}}{h_1} (C + D h_1)$$

From the above equation we got

$$q_{1a} = \frac{t_{0a} - \frac{t_{1a}}{h_0} (C + A h_0)}{\frac{1}{h_1} (A + B h_1 + \frac{C}{h_0} + D \frac{h_1}{h_0})}$$

Substituting this value of  $q_{1a}$  in the equation (18) and (19) by rearranging the terms we got

$$t_{1a} = t_{0a} \xrightarrow{\rightarrow} \lambda_c + t_{1a} \xleftarrow{\leftarrow} \lambda'_c$$

$$q_{1a} = t_{0a} \xrightarrow{\rightarrow} \lambda'_1 - t_{1a} \xleftarrow{\leftarrow} \lambda'_1$$

$$t_{0a} = t_{0a} \xrightarrow{\rightarrow} \lambda_o + t_{1a} \xleftarrow{\leftarrow} \lambda'_o$$

$$q_{0a} = t_{0a} \xrightarrow{\rightarrow} \lambda'_o - t_{1a} \xleftarrow{\leftarrow} \lambda'_o$$

This condition of both sides subjected to variable air temperature can be considered as the super position of cases (iii) and (iv) (i.e., outside air temperature is variable with inside constant and vice-versa). By applying the principle of super position (which is permissible for linear circuits) we obtain the above results for the surface temperatures and heat fluxes in a much simpler way.

- vi.) The air temperature on both sides is variable with the same amplitude and phase (this applies for partition walls) (Fig. 6).

For this special case the matrix equation becomes

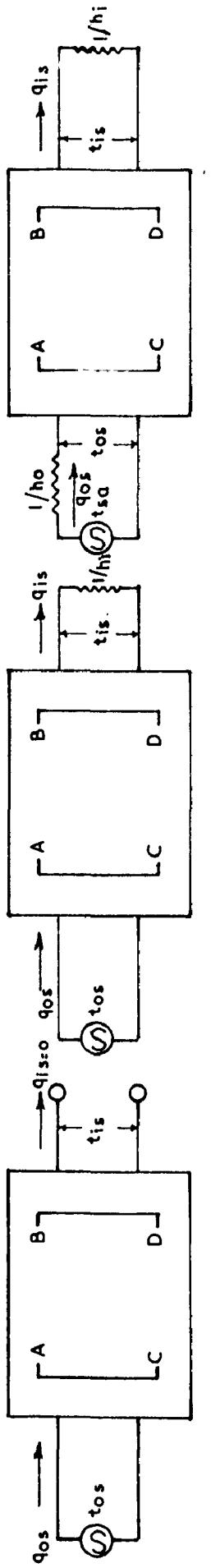
$$\begin{bmatrix} t_{ia} \\ q_{ia} \end{bmatrix} = \begin{bmatrix} 1 & 1/h_i \\ 0 & 1 \end{bmatrix} \begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} 1 & 1/h_i \\ 0 & 1 \end{bmatrix} \begin{bmatrix} t_{ia} \\ -q_{ia} \end{bmatrix}$$

In this case we are interested in finding out the  $q_{ia}$  i.e., the heat flux entering into the partition from either side. This is obtained as :-

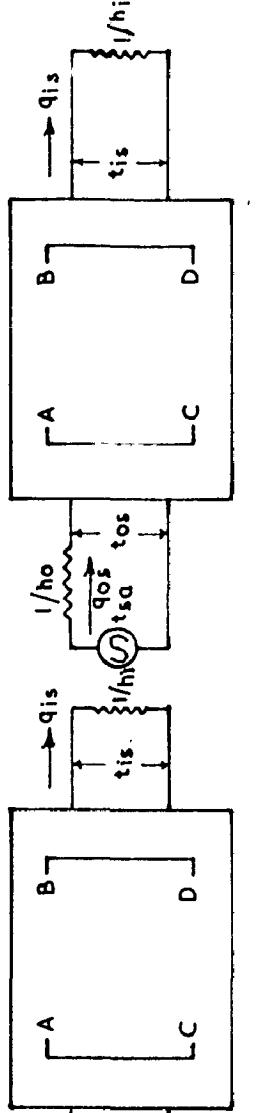
$$q_{ia} = t_{ia} (\overleftarrow{Y'_i} - \overrightarrow{Y'_i})$$

The total quantity of heat flux entering the partition i.e., from both sides will be equal to  $2 q_{ia}$ . If the building section is symmetrical with respect to the central plane, the heat flow at that plane of symmetry will be zero.

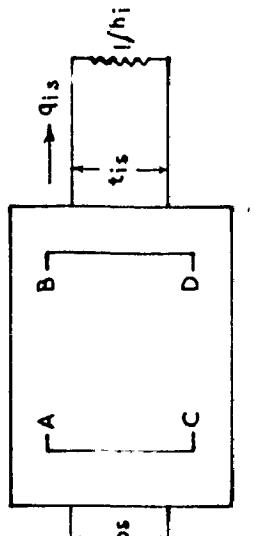
The inter relationship between the thermal system functions are easily established from the above derivations as given below :-



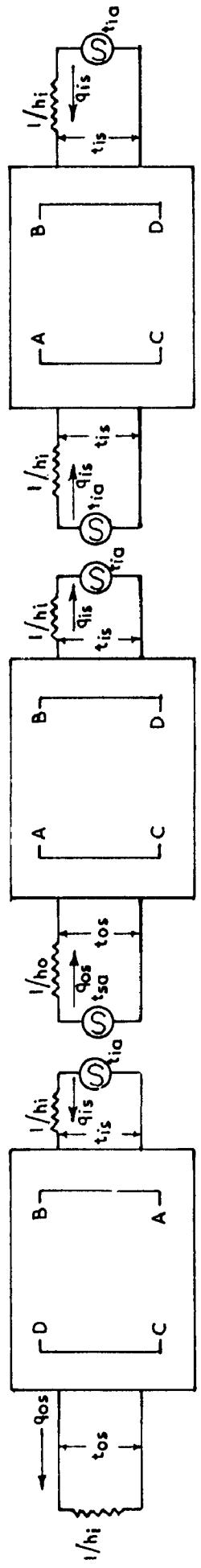
①



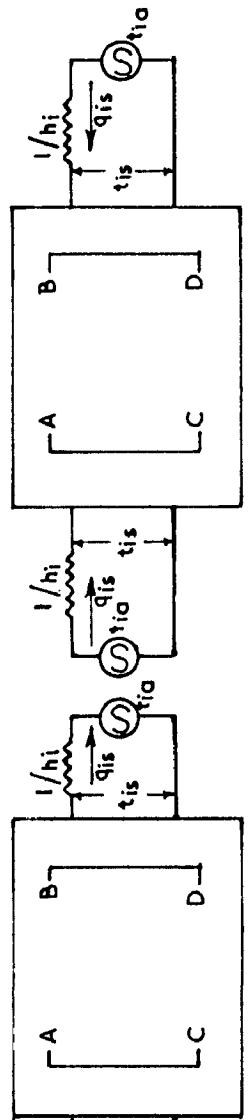
②



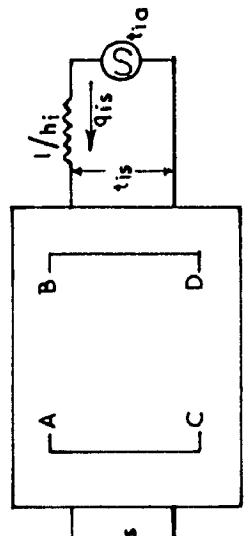
③



④



⑤



⑥

FIGS. I TO 6. DIAGRAMATIC REPRESENTATION OF DIFFERENT BOUNDARY CONDITIONS

: 180 :

### Transfer Functions

$$1. \lambda_c \underline{-\varphi_c} = \left( \frac{t_{ia}}{t_{sa}} \right) \quad t_{ia} = 0$$

$$2. \lambda_o \underline{-\varphi_o} = \left( \frac{t_{os}}{t_{ia}} \right) \quad t_{sa} = 0$$

$$\lambda'_o \underline{-\varphi'_o} = \frac{1}{h_o} \lambda_c \underline{-\varphi_c}$$

### Transfer Admittance Functions

$$1. Y_1 \underline{-\psi_c} = \left( \frac{q_{ic}}{t_{sa}} \right) \quad t_{ia} = 0$$

$$Y_c \underline{-\psi_c} = h_1 \times \lambda_c \underline{-\varphi_c}$$

$$2. Y'_o \underline{-\varphi'_o} = \left( \frac{q_{os}}{t_{ia}} \right) \quad t_{sa} = 0$$

$$Y'_o \underline{-\varphi'_o} = Y_1 \underline{-\psi_c}$$

### Driving Point Functions

$$1. \lambda_o \underline{-\varphi_o} = \left( \frac{t_{os}}{t_{sa}} \right) \quad t_{ia} = 0$$

$$2. \lambda'_c \underline{-\varphi'_c} = \left( \frac{t_{ia}}{t_{sa}} \right) \quad t_{sa} = 0$$

### Driving Point Admittance Functions

$$1. Z_o \underline{\psi_o} = \left( \frac{q_{os}}{t_{sa}} \right) \quad t_{ia} = 0$$

$$Y_o \underline{\psi_o} = h_o (1 - \lambda_o \underline{-\varphi_o})$$

$$2. Y'_1 \underline{\psi'_c} = \left( \frac{q_{ic}}{t_{ia}} \right) \quad t_{sa} = 0$$

The above relationships between the surface temperatures and heat fluxes and the thermal system functions (in terms of general circuit parameters and surface heat transfer coefficients) will hold good not only for homogeneous building elements but also for composite (multi layered) building sections. In the case of multi layered sections the transfer matrix coefficients (general circuit parameters A, B, C and D) are obtained by the matrix multiplication of the transfer matrices of the individual layers. The overall transfer matrix of a composite building element is then given by

$$\begin{bmatrix} A & B \\ C & D \end{bmatrix} = \begin{bmatrix} A_1 & B_1 \\ C_1 & D_1 \end{bmatrix} \begin{bmatrix} A_2 & B_2 \\ C_2 & D_2 \end{bmatrix} \dots \begin{bmatrix} A_n & B_n \\ C_n & D_n \end{bmatrix}$$

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## APPENDIX III

TABLE I  
PHYSICAL AND PHYSICAL-CHIMICAL PROPERTIES OF VARIOUS BUILDING MATERIALS

No.	Material	Specific Gravity	Specific Heat	Conductivity	Strength per square inch	lb./sq.in.	Strength per cu. in.	lb./cu.in.	Modulus of elasticity	lb./sq.in.	Modulus of rigidity	lb./sq.in.	Modulus of resilience	lb./sq.in.	Modulus of toughness	lb./sq.in.
<b>1.0.1</b> <b>Brick</b>																
1.0.2	Common Concrete	2.0	1.20	0.21	2.097	0.477	0.023	3.147	2.823							
1.0.3	"	1.0	1.30	0.21	2.430	0.479	0.031	2.831	3.033							
1.0.4	"	1.2	1.10	0.21	2.774	0.381	0.034	2.775	3.201							
1.0.5	"	1.2	1.60	0.21	2.939	0.335	0.034	2.760	3.218							
1.0.6	"	1.0.8	1.32	0.22	3.030	0.325	0.030	2.972	2.566							
<b>2.</b> <b>Light weight concrete</b>																
2.0.1	Concrete	2.0	0.21	0.21	0.721	1.297	0.012	4.633	1.020							
2.0.2	Cinder	5.0	0.21	0.21	1.475	0.673	0.021	3.541	2.608							
2.0.3	Lime-cement cinder	2.8	0.21	1.074	0.835	0.012	4.814	1.929								
2.0.4	Expansive cement concrete	1.1	0.20	0.490	2.041	0.009	5.350	1.631								
2.0.5	Cow dung	2.6	0.23	1.055	0.943	0.110	4.879	1.821								
2.0.6	Coal ash (1)	2.0	0.21	0.701	1.349	0.013	4.454	1.907								
2.0.7	Coal ash (11)	2.5	0.21	0.937	1.045	0.012	4.663	1.033								
<b>3.</b> <b>Plastered Concrete</b>																
3.0.1	Common	0.75	3.0	0.25	0.350	2.854	0.008	6.613	1.636							
3.0.2	"	1.45	50	0.25	0.820	1.590	0.010	6.200	1.706							
3.0.3	"	2.30	70	0.25	0.917	1.030	0.011	4.924	1.775							
3.0.4	"	3.30	80	0.25	1.310	0.763	0.013	4.463	1.977							
<b>4.</b> <b>Bricks</b>																
4.0.1	Common	6.0	100	0.21	1.653	0.603	0.023	3.317	2.677							
4.0.2	6.0	117	0.20	1.669	0.635	0.034	2.757	3.245								
4.0.3	6.0	120	0.21	2.097	0.377	0.025	3.147	2.821								
4.0.4	"	130	0.21	2.340	0.410	0.031	2.831	3.033								
4.0.5	"	170	0.20	4.303	0.282	0.031	3.033	4.300								
4.0.6	"	100	0.10	3.177	0.459	0.033	2.033	3.035								
4.0.7	"	150	0.22	2.914	0.355	0.023	3.071	2.033								
4.0.8	"	170	0.20	3.045	0.190	0.031	2.031	1.760								
<b>5.</b> <b>Common</b>																
5.0.1	Wood	1.2	40	0.39	0.339	1.635	0.006	6.503	1.360							
5.0.2	Teak	1.0	35	0.39	0.545	1.835	0.003	6.530	1.353							
5.0.3	Teakwood	1.26	42	0.57	0.809	1.233	0.004	7.733	1.124							
5.0.4	Oak	2.0	65	0.20	0.753	1.321	0.003	5.470	1.033							
5.0.5	A.C. Shoot	2.0	470	0.12	22.740	0.003	0.650	13.375								
5.0.6	G.I. Shoot	2.0	160	0.16	1.920	0.521	0.021	3.494	2.541							
5.0.7	Globe Shoot	2.0	124	0.20	2.163	0.454	0.030	2.835	2.835							
5.0.8	Un Dried Bricks	2.0														

TABLES I (Cont'd.)

Sl. No.	Material	K	P	S	C	E <sub>c</sub>	E <sub>a</sub>	E <sub>r</sub>	E <sub>t</sub>	S
11.	Soil (compacted)	3.0	110	0.20	1.853	0.513	0.031	2.930	3.032	
12.	Lime Concrete	6.1	103	0.20	1.674	0.593	0.024	3.320	2.630	
13.	Concrete	12.0	140	0.20	2.759	0.382	0.036	2.707	2.216	
14.	Monolithic tile	5.8	120	0.17	1.607	0.621	0.024	3.327	2.663	
15.	Plasterboard	12.0	120	0.22	2.629	0.390	0.038	2.630	3.372	
16.	Lime Concrete	3.0	110	0.22	2.055	0.487	0.027	3.084	2.850	
17.	Gypsum	3.5	70	0.23	1.103	0.801	0.318	3.302	2.336	
18.	Stucco	12.0	110	0.23	2.679	0.388	0.043	2.457	3.616	
19.	Painted Plaster (Plastered)	4.84	90	0.25	1.921	0.521	0.011	4.733	1.833	
20.	Found Glass	0.40	10	0.13	0.135	0.003	0.010	3.760	2.368	
21.	Former Plastic	0.23	3	0.23	0.125	0.300	0.008	5.357	1.363	
22.	Thermocol	0.20	1	0.22	0.083	27.830	0.053	2.215	4.134	
23.	Insulation Boards	0.34	24	0.25	0.173	5.339	0.007	6.012	1.477	
24.	Color Coated Metal	0.34	30	0.25	0.080	1.459	0.008	5.394	1.595	
25.	Aluminum Foil	0.23	36	0.34	0.207	4.332	0.006	6.963	1.287	
26.	Polycarbonate	0.24	17	0.24	0.248	6.737	0.307	6.140	1.846	
27.	Double Pole	0.25	21	0.33	0.103	7.143	0.006	6.852	1.355	
28.	Double Pole (44)	0.24	30	0.25	0.173	5.339	0.007	6.012	1.477	
29.	Double Pole (44)	0.24	30	0.25	0.173	5.339	0.007	6.012	1.477	
30.	Double Pole	0.25	21	0.33	0.103	7.143	0.006	6.852	1.355	
31.	Double Pole	0.25	21	0.33	0.103	7.143	0.006	6.852	1.355	
32.	Double Pole	0.25	21	0.33	0.103	7.143	0.006	6.852	1.355	
33.	Double Pole	0.25	21	0.33	0.103	7.143	0.006	6.852	1.355	
34.	Smoothed Wood	0.81	10	0.45	0.384	2.604	0.008	5.730	1.654	
35.	Bricklaying	1.61	31	0.60	0.811	1.233	0.007	6.043	1.433	
36.	Bricklaying	1.21	37	0.30	0.70	1.282	0.005	7.440	1.195	
37.	Bricklaying	1.21	37	0.30	0.70	1.282	0.005	7.440	1.195	

TABLE 2

EFFECTS OF EXTERNAL CONSTRUCTION AND DRIVING POINTS ON INSIDE AND OUTSIDE SURFACE HEAT TRANSFER COEFFICIENTS									
No.	Building	Thickness in inches	Value	EXTERNAL DRIVING POINTS				INTERNAL DRIVING POINTS	
				H3	H4	H5	H6	Function	Phi
2.	Brick	2	0.633	10	17	22	23	0.464	0.460
		4	0.633	10	17	23	23	0.464	0.464
2.	Concrete	4	0.633	10	17	23	23	0.464	0.464
		6	0.637	10	17	23	23	0.464	0.464
2.	Cinder Concrete	4	0.633	10	17	23	23	0.464	0.464
		6	0.637	10	17	23	23	0.464	0.464
2.	Steel	4	0.570	10	17	23	23	0.464	0.464
		6	0.634	10	17	23	23	0.464	0.464
2.	Concrete	4	0.352	10	17	23	23	0.464	0.464
		6	0.612	10	17	23	23	0.464	0.464
2.	Steel	4	0.330	10	17	23	23	0.464	0.464
		6	0.603	10	17	23	23	0.464	0.464

**TABLE 2** (Cont'd.)

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		Structural System			Material Properties			Geometric Properties			Boundary Conditions			Transfer Function			Information Processing		
		Concrete	Steel	Wood	Concrete	Steel	Wood	Concrete	Steel	Wood	Concrete	Steel	Wood	Concrete	Steel	Wood	Concrete	Steel	Wood
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TABLE 2 (cont'd.)

$b_0 = 3.5$   
 $b_1 = 1.5$

S.No.	Balances	Internal Driving Function		External Function		Transition Function		Internal Driving Function		External Function		Transition Function	
		Thickness	Conc.	Conc.	Conc.	Conc.	Conc.	Thickness	Conc.	Conc.	Conc.	Conc.	Conc.
10.	Hand Recons.	3	0.543	12	0.433	15	0.375	16	0.333	13	0.333	10	0.333
11.	Can Drilled brick	13	0.405	14	0.355	15	0.333	16	0.333	13	0.333	10	0.333
12.	L40 Concrete	11	0.610	12	0.570	13	0.530	14	0.500	11	0.500	8	0.500
13.	Cell Concrete	10	0.630	11	0.590	12	0.550	13	0.520	10	0.520	7	0.520
14.	20x20x20	12	0.660	13	0.620	14	0.580	15	0.550	12	0.550	9	0.550
15.	4x6x12	13	0.680	14	0.640	15	0.600	16	0.570	13	0.570	10	0.570

FIGURE 3 (Cont'd.)

$H_0 = 3.6$

$B_1 = 1.6$

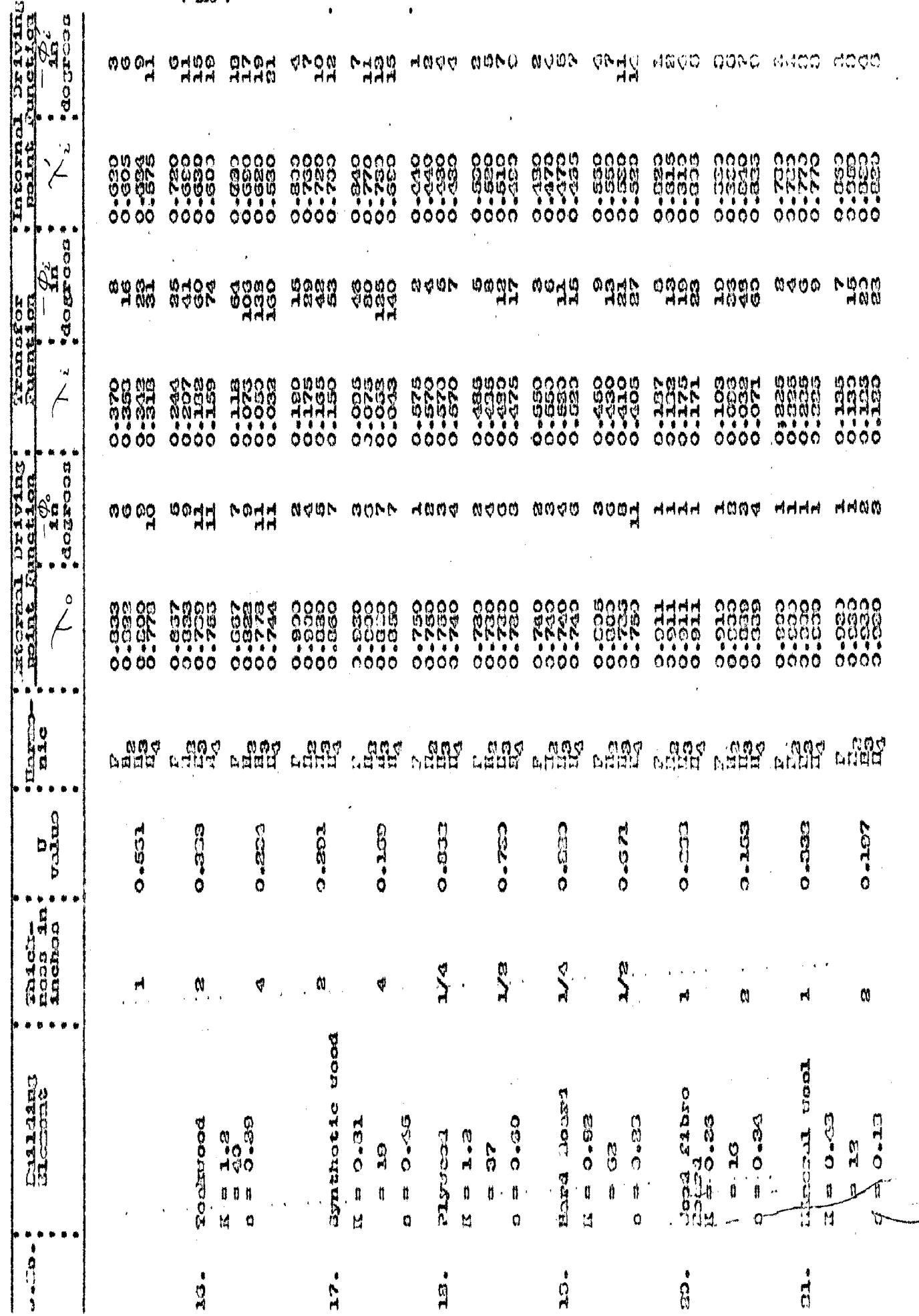


TABLE 2 (CONT'D.)

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APPENDIX IV

TELLING ALGORITHM (THERMAL AND DRIVING POINT) FUNCTION DATA  
FOR COMPOSITE BUILDING ELEMENTS  
(FOR SINUSOIDAL TEMPERATURE VARIATION)

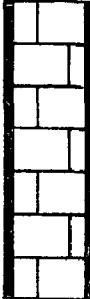
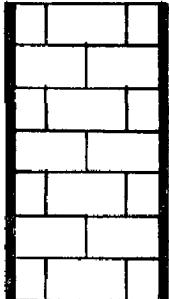
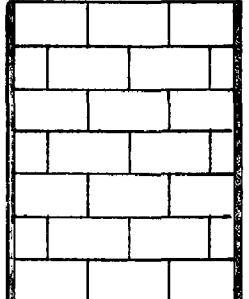
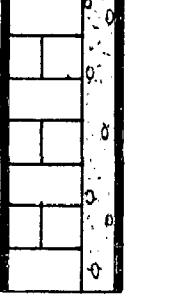
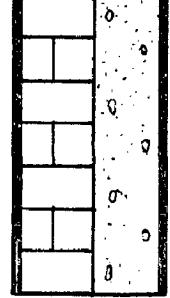
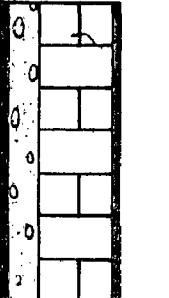
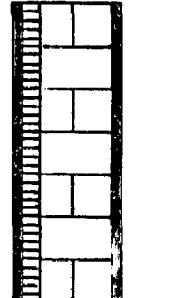
In building practice a large number of composite constructions (made up of two or more layers of homogeneous materials of different thicknesses) are met with. The traditional building materials like brick, stone, timber, concrete etc. still form the basic core in these constructions. Of late, a large variety of lightweight concretes, insulating boards, sandwich, panel constructions etc. have come into use and their use in combination with the conventional materials in a variety of ways, are becoming popular.

In this Appendix the thermal system functions for a number of composite constructions (both traditional and modern) that are employed in the present day building practice are given along with the steady state thermal transmission coefficients ( $U$  - values).

For the sake of convenience these data are divided into different sections, according to their functional use viz., walls, roofs, floors, doors and windows.



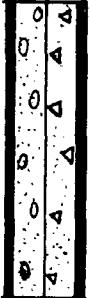
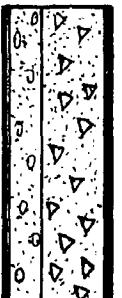
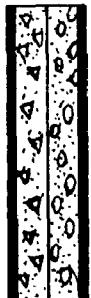
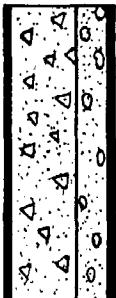
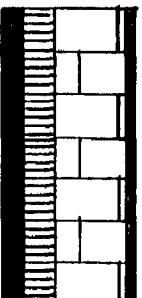
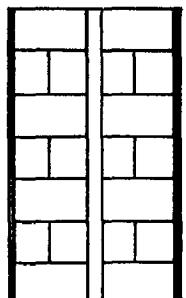
## WALL SECTIONS

$\frac{1}{2}'' p + 4\frac{1}{2}'' \text{brick}$ $+ \frac{1}{2}'' p$ 	$\frac{1}{2}'' p + 9'' \text{brick}$ $+ \frac{1}{2}'' p$ 	$\frac{1}{2}'' p + 13\frac{1}{2}'' \text{brick}$ $+ \frac{1}{2}'' p$ 
$\frac{1}{2}'' p + 4\frac{1}{2}'' \text{brick}$ $+ 1'' \text{F.C.} + \frac{1}{2}'' p$ 	$\frac{1}{2}'' p + 4\frac{1}{2}'' \text{brick}$ $+ 2'' \text{F.C.} + \frac{1}{2}'' p$ 	$\frac{1}{2}'' p + 4\frac{1}{2}'' \text{brick}$ $+ 4'' \text{F.C.} + \frac{1}{2}'' p$ 
$\frac{1}{2}'' p + 1'' \text{F.C.} +$ $4\frac{1}{2}'' \text{brick} + \frac{1}{2}'' p$ 	$\frac{1}{2}'' p + 2'' \text{F.C.} +$ $4\frac{1}{2}'' \text{brick} + \frac{1}{2}'' p$ 	$\frac{1}{2}'' p + 1'' \text{r.b.} +$ $4\frac{1}{2}'' \text{brick} + \frac{1}{2}'' p$ 

p = plaster   F.C. = Foamed concrete  
 r.b. = reed board



## WALL SECTIONS

$\frac{1}{2}" p + 3" F.C. + \frac{1}{2}" p$ 	$\frac{3}{4}" p + 2" r.b. + \frac{3}{4}" p$ 	$\frac{3}{4}" p + 4" r.b. + \frac{3}{4}" p$ 
$\frac{1}{2}" p + 2" F.C. + \frac{1}{2}" D.C. + \frac{1}{2}" p$ 	$\frac{1}{2}" p + 2" F.C. + 4" D.C. + \frac{1}{2}" p$ 	$\frac{1}{2}" p + 2" D.C. + 2" F.C. + \frac{1}{2}" p$ 
$\frac{1}{2}" p + 4" D.C. + 2" F.C. + \frac{1}{2}" p$ 	$1\frac{1}{2}" conc. + 2" jax board + 4\frac{1}{2}" brick + \frac{1}{2}" p$ 	$\frac{1}{2}" p + 4\frac{1}{2}" brick + 2" airspace + 4\frac{1}{2}" brick + \frac{1}{2}" p$ 

F.C. = Foamed concrete    r.b. = reed board  
D.C. = Dense concrete    p = plaster

D<sub>2</sub> = 0.5D<sub>3</sub> = 1.5

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No.	Unit	Location	Material	Description	Quantity	Dimensions	Notes
18.	1"	Plasterer + 1/2" Plaster	0.147	Plasterer	0.373	3/4" Plaster + 1/2" Plaster	18.
19.	1"	Plasterer + 1/2" Plaster	0.147	Plasterer	0.373	3/4" Plaster + 1/2" Plaster	19.
20.	1/2"	Plasterer + 1/2" Plaster	0.147	Plasterer	0.373	3/4" Plaster + 1/2" Plaster	20.
21.	1/2"	Plasterer + 1/2" Plaster	0.147	Plasterer	0.373	3/4" Plaster + 1/2" Plaster	21.
22.	3/4"	Plasterer + 1/2" Plaster	0.203	Plasterer	0.606	3/4" Plaster + 1/2" Plaster	22.
23.	3/4"	Plasterer + 1/2" Plaster	0.203	Plasterer	0.606	3/4" Plaster + 1/2" Plaster	23.
24.	1/2"	Plasterer + 1/2" Plaster	0.203	Plasterer	0.606	1/2" Plaster + 1/2" Plaster	24.
25.	1/2"	Plasterer + 1/2" Plaster	0.203	Plasterer	0.606	1/2" Plaster + 1/2" Plaster	25.
26.	3/4"	Plasterer + 1/2" Plaster	0.203	Plasterer	0.606	3/4" Plaster + 1/2" Plaster	26.
27.	3"	3/2" Stone + 3" Brick	0.317	Brick	0.317	3" Stone + 3/2" Brick	27.

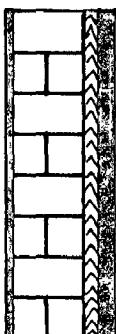
## WALL SECTIONS

1" p + 1" T.C. +  
 $\frac{1}{2}$ " p +  $4\frac{1}{2}$ " brick +  
 $4\frac{1}{2}$ " brick +  $\frac{1}{2}$ " p



(19)

$\frac{1}{2}$ " p +  $4\frac{1}{2}$ " brick +  
 1" T.C. +  $\frac{1}{2}$ " p



(20)

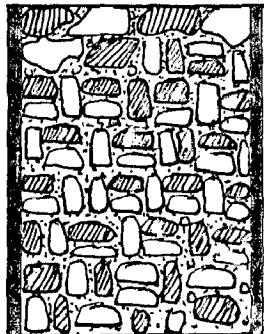
15" R.M. +  $\frac{3}{4}$ " p



(21)

$\frac{3}{4}$ " p + 15" R.M. +  
 $\frac{1}{2}$ " p

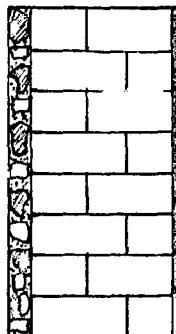
$1\frac{1}{2}$ " stone + 9"  
 brick +  $\frac{1}{2}$ " p



(22)



(23)

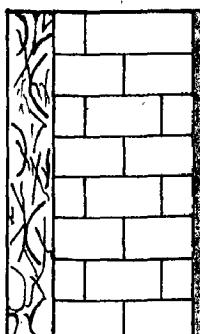


(24)

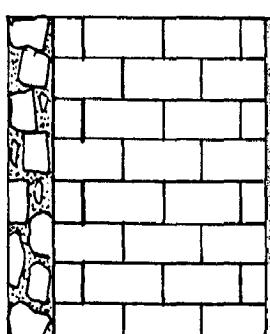
3" stone + 9" brick  
 $+\frac{1}{2}$ " p

3" stone +  $13\frac{1}{2}$ "  
 brick +  $\frac{1}{2}$ " p

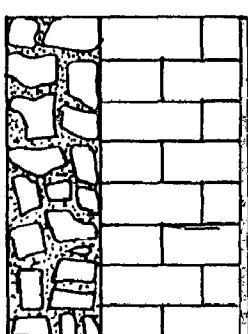
6" stone + 9" brick  
 $+\frac{1}{2}$ " p



(25)



(26)



(27)

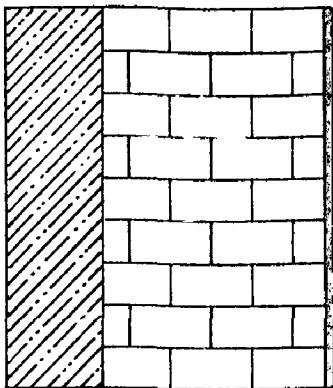
p = plaster    T.C. = Thermocole    R.M. = Rubble Masonry

## TABLE 3 (continued)

No.	Roll Section	Weld Value	Horiz. Open Cuts	Vertical Function	Function $\phi_1$	Function $\phi_2$	$b_1 = 1.0$	$b_2 = 3.0$
23.	6" beam + 1/2" plate + 12" track	0.283	0.013	220	0.437	0.273	194	194
24.	1" thick + 1/2" plate + 12" track	0.300	0.013	273	0.437	0.289	204	204
25.	2" thick + 1/2" plate + 12" track	0.316	0.013	310	0.437	0.289	213	213
26.	1" thick + 1/2" plate + 12" track	0.333	0.013	350	0.437	0.295	223	223
27.	2" thick + 1/2" plate + 12" track	0.350	0.013	393	0.437	0.295	232	232
28.	1" thick + 1/2" plate + 12" track	0.367	0.013	433	0.437	0.295	241	241
29.	2" thick + 1/2" plate + 12" track	0.383	0.013	473	0.437	0.295	250	250
30.	1" thick + 1/2" plate + 12" track	0.400	0.013	513	0.437	0.295	259	259
31.	2" thick + 1/2" plate + 12" track	0.416	0.013	553	0.437	0.295	268	268
32.	1" thick + 1/2" plate + 12" track	0.433	0.013	593	0.437	0.295	277	277
33.	2" thick + 1/2" plate + 12" track	0.450	0.013	633	0.437	0.295	286	286
34.	1" thick + 1/2" plate + 12" track	0.467	0.013	673	0.437	0.295	295	295
35.	2" thick + 1/2" plate + 12" track	0.483	0.013	713	0.437	0.295	304	304
36.	1" thick + 1/2" plate + 12" track	0.500	0.013	753	0.437	0.295	313	313
37.	2" thick + 1/2" plate + 12" track	0.516	0.013	793	0.437	0.295	322	322
38.	1" thick + 1/2" plate + 12" track	0.533	0.013	833	0.437	0.295	331	331
39.	2" thick + 1/2" plate + 12" track	0.550	0.013	873	0.437	0.295	340	340
40.	1" thick + 1/2" plate + 12" track	0.567	0.013	913	0.437	0.295	349	349
41.	2" thick + 1/2" plate + 12" track	0.583	0.013	953	0.437	0.295	358	358
42.	1" thick + 1/2" plate + 12" track	0.600	0.013	993	0.437	0.295	367	367
43.	2" thick + 1/2" plate + 12" track	0.616	0.013	1033	0.437	0.295	376	376
44.	1" thick + 1/2" plate + 12" track	0.633	0.013	1073	0.437	0.295	385	385
45.	2" thick + 1/2" plate + 12" track	0.650	0.013	1113	0.437	0.295	394	394
46.	1" thick + 1/2" plate + 12" track	0.667	0.013	1153	0.437	0.295	403	403
47.	2" thick + 1/2" plate + 12" track	0.683	0.013	1193	0.437	0.295	412	412
48.	1" thick + 1/2" plate + 12" track	0.700	0.013	1233	0.437	0.295	421	421
49.	2" thick + 1/2" plate + 12" track	0.716	0.013	1273	0.437	0.295	430	430
50.	1" thick + 1/2" plate + 12" track	0.733	0.013	1313	0.437	0.295	439	439
51.	2" thick + 1/2" plate + 12" track	0.750	0.013	1353	0.437	0.295	448	448
52.	1" thick + 1/2" plate + 12" track	0.767	0.013	1393	0.437	0.295	457	457
53.	2" thick + 1/2" plate + 12" track	0.783	0.013	1433	0.437	0.295	466	466
54.	1" thick + 1/2" plate + 12" track	0.800	0.013	1473	0.437	0.295	475	475
55.	2" thick + 1/2" plate + 12" track	0.816	0.013	1513	0.437	0.295	484	484
56.	1" thick + 1/2" plate + 12" track	0.833	0.013	1553	0.437	0.295	493	493
57.	2" thick + 1/2" plate + 12" track	0.850	0.013	1593	0.437	0.295	502	502
58.	1" thick + 1/2" plate + 12" track	0.867	0.013	1633	0.437	0.295	511	511
59.	2" thick + 1/2" plate + 12" track	0.883	0.013	1673	0.437	0.295	520	520
60.	1" thick + 1/2" plate + 12" track	0.900	0.013	1713	0.437	0.295	529	529
61.	2" thick + 1/2" plate + 12" track	0.916	0.013	1753	0.437	0.295	538	538
62.	1" thick + 1/2" plate + 12" track	0.933	0.013	1793	0.437	0.295	547	547
63.	2" thick + 1/2" plate + 12" track	0.950	0.013	1833	0.437	0.295	556	556
64.	1" thick + 1/2" plate + 12" track	0.967	0.013	1873	0.437	0.295	565	565
65.	2" thick + 1/2" plate + 12" track	0.983	0.013	1913	0.437	0.295	574	574
66.	1" thick + 1/2" plate + 12" track	1.000	0.013	1953	0.437	0.295	583	583

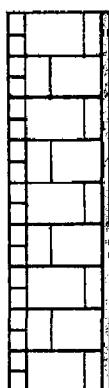
## WALL SECTIONS

6" stone + 13 $\frac{1}{2}$ " brick +  $\frac{1}{2}$ " p



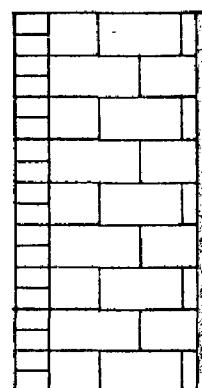
(28)

1" B.T. + 9" brick +  $\frac{1}{2}$ " p



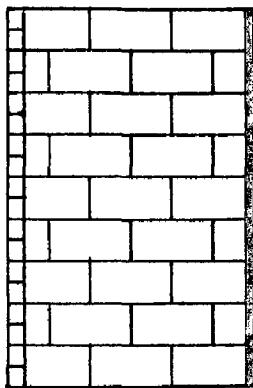
(29)

2" B.T. + 9" brick +  $\frac{1}{2}$ " p



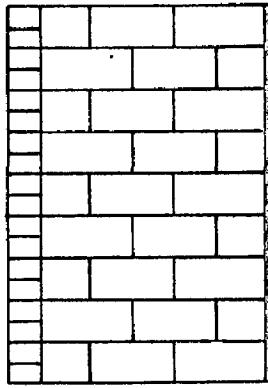
(30)

1" B.T. + 13 $\frac{1}{2}$ " brick +  $\frac{1}{2}$ " p



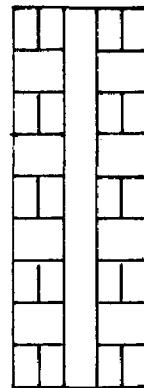
(31)

2" B.T. + 13 $\frac{1}{2}$ " brick +  $\frac{1}{2}$ " p



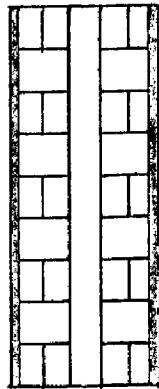
(32)

3" brick + 2" air space + 3" brick



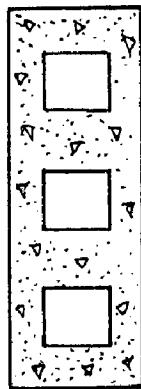
(33)

$\frac{1}{2}$ " p + 3" brick + 8" H.Block  
2" airspace + 3" br.  
 $\frac{1}{2}$ " p

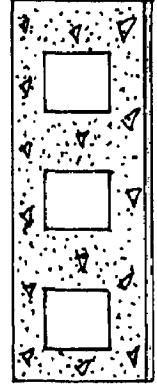


(34)

8" H. Block +  $\frac{1}{2}$ " p



(35)



(36)

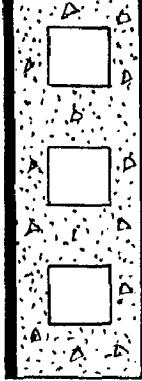
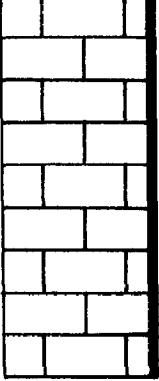
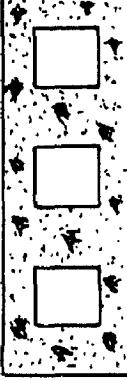
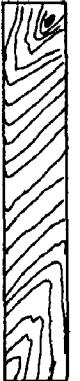
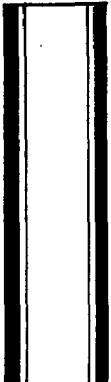
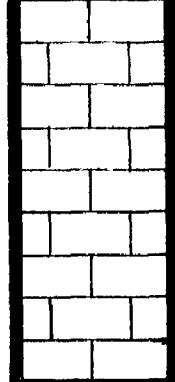
p = plaster    B.T. = Bricktile    br. = brick  
H. Block = Hollow block

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Sect. 1 (cont'd.)

## WALL SECTIONS

$\frac{1}{2}'' p + 8'' H.\text{block}$ $+ \frac{1}{2}'' p$  <p>(37)</p>	$9'' \text{brick} + \frac{1}{2}'' p$  <p>(38)</p>	$8'' H.G. \text{block}$  <p>(39)</p>
$1/8'' W.B. + 4\frac{1}{2}'' H.G.\text{block} + 1/8'' W.B.$  <p>(40)</p>	$4'' \text{Syntheticwood panel}$  <p>(41)</p>	$1\frac{1}{2}'' \text{conc.} + 2'' \text{air space} + 1\frac{1}{2}'' \text{conc.}$  <p>(42)</p>
$3/4'' p + \frac{1}{2}'' r.b. +$ $4'' \text{airspace} + \frac{1}{2}'' r.b. +$ $\frac{3}{4}'' p$  <p>(43)</p>	$3'' m.p. + 1'' b.b.$ $+ \frac{3}{4}'' m.p.$  <p>(44)</p>	$\frac{3}{4}'' m.p. + 9'' S.B.$ $+ \frac{3}{4}'' m.p.$  <p>(45)</p>

p = plaster    H.block = Hollow block    H.G. = Hollow gypsum  
 r.b. = reed board    m.p. = mud plaster    b.b. = bamboo board

EXHIBIT 2 (Continued)

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Transfer of Motor vehicle  
Registration from one state to another state.

46. 3" and Plaster of Paris  
applied thick & 3/4" and  
Plaster

### 43. *Vera* Plaster • "A" + Vise, Pudding concrete

49.  $\sqrt{2}$ " planter + "B" orange  
diag concrete +  $\sqrt{2}$ "  
planter

33. User-defined classes & objects

51. *Maya*, *Parasurama* → *Uma*, *Shiva* (see *Diagram 1*)

32.  $\Delta S = \Delta H - T\Delta S_{\text{cond}}$  +  $T\Delta S_{\text{diss}}$  +  $T\Delta S_{\text{electr}}$ .

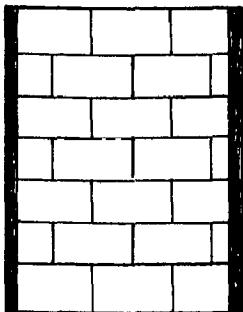
### 63. Caja de leche + 27 aler spaco

FACON TOUZÉ & CIE  
STRASBOURG

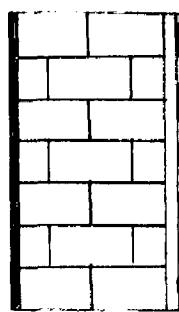
## WALL SECTIONS

3" m.p. + 13 $\frac{1}{2}$ " S.B.  
+  $\frac{3}{4}$ " m.p.

$\frac{1}{2}$ " p + 9" brick +  $\frac{1}{2}$ " p + 4" e.s.c. +  
 $\frac{1}{4}$ " airspace +  $\frac{1}{2}$ " p  
p.w.



(46)



(47)

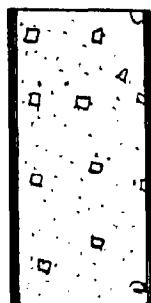


(48)

$\frac{1}{2}$ " p + 8" e.s.c. +  
 $\frac{1}{2}$ " p

$\frac{1}{2}$ " p + 4" c.c. +  
 $\frac{1}{2}$ " p

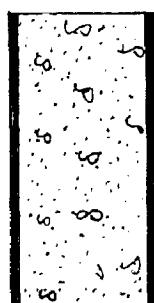
$\frac{1}{2}$ " p + 3" c.c. +  
 $\frac{1}{2}$ " p



(49)



(50)

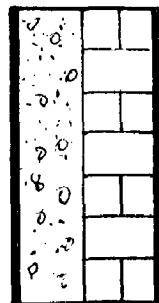


(51)

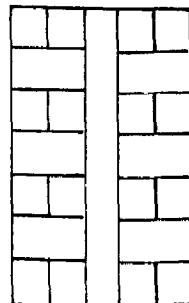
$\frac{1}{2}$ " p + 4" F.C. +  
4 $\frac{1}{2}$ " brick +  $\frac{1}{2}$ " p

4 $\frac{1}{2}$ " brick + 2" air space + 4 $\frac{1}{2}$ " brick

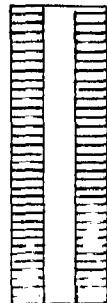
2" r.b. + 2" air space + 2" r.b.



(52)



(53)



(54)

p = plaster, r.b. = reed board m.p. = mud plaster

S.B. = sun dried brick p.w. = plywood c.c. = cinder conc.

e.s.c. = expanded slag concrete. F.C. = Foamed concrete



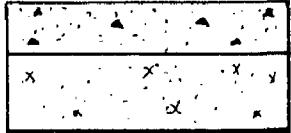
## ROOF SECTIONS

3" L.C. + 4" RCC  
+  $\frac{1}{2}$ " p



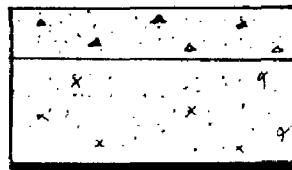
(1)

3" L.C. + 4 $\frac{1}{2}$ " R.B.  
+  $\frac{1}{2}$ " p



(2)

3" L.C. + 6" R.B.  
+  $\frac{1}{2}$ " p



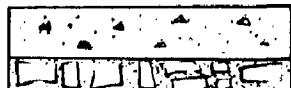
(3)

3" L.C. + 7 $\frac{1}{2}$ " R.B.  
+  $\frac{1}{2}$ " p



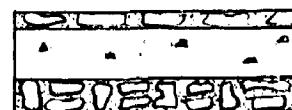
(4)

3" L.C. + 2" s.s.



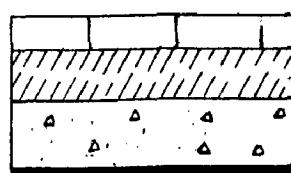
(5)

1" s.s. + 3" L.C.  
+ 2" s.s.



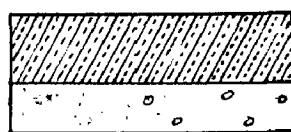
(6)

2" b.f. + 3" M.P.  
+ 4" RCC +  $\frac{1}{2}$ " p

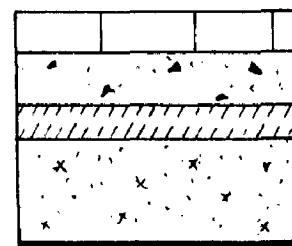


(7)

4" F.C. + 3" RCC +  $\frac{1}{2}$ " p  
+ 3 $\frac{1}{2}$ " b.t. + 3" L.C.  
2" M.P. + 6" R.B.  
+  $\frac{1}{2}$ " p



(8)



(9)

L.C. = Lime concrete    RCC = Reinforced cement concrete

R.B. = Reinforced brick    s.s. = sand stone    M.P. = mudphusk

F.C. = Foamed concrete    b.t. = bricktile    p = plaster

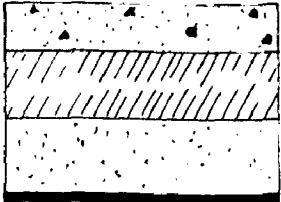
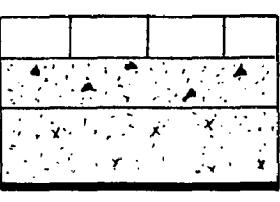
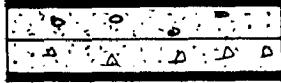
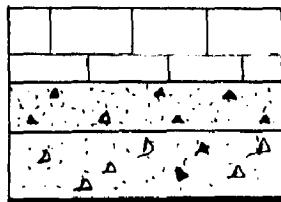
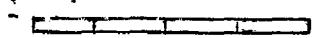
b.f. = brickflat

TABLE 2 (Cont'd.)

 $\nu_0 = 3.5$  $\nu_1 = 1.5$ 

10.	$3\pi/2$ , 140°, 0.300	$17$	$0.730$	$17$	$0.051$	$144$	$0.404$	$24$
11.	$2\pi/3$ , 180°, 0.300	$21$	$0.804$	$20$	$0.037$	$124$	$0.499$	$27$
12.	$2\pi/3$ , 180°, 0.325	$22$	$0.837$	$21$	$0.118$	$64$	$0.553$	$30$
13.	$4\pi/5$ , 180°, 0.325	$23$	$0.839$	$22$	$0.073$	$124$	$0.423$	$24$
14.	$4\pi/5$ , 180°, 0.350	$24$	$0.837$	$23$	$0.023$	$200$	$0.339$	$25$
15.	$2\pi/5$ , 180°, 0.350	$25$	$0.839$	$24$	$0.023$	$200$	$0.300$	$26$
16.	$2\pi/5$ , 180°, 0.375	$26$	$0.839$	$25$	$0.023$	$200$	$0.271$	$20$
17.	$2\pi/5$ , 180°, 0.400	$27$	$0.839$	$26$	$0.023$	$200$	$0.242$	$19$
18.	$2\pi/5$ , 180°, 0.425	$28$	$0.839$	$27$	$0.023$	$200$	$0.213$	$18$
19.	$2\pi/5$ , 180°, 0.450	$29$	$0.839$	$28$	$0.023$	$200$	$0.184$	$17$
20.	$2\pi/5$ , 180°, 0.475	$30$	$0.839$	$29$	$0.023$	$200$	$0.155$	$16$
21.	$2\pi/5$ , 180°, 0.500	$31$	$0.839$	$30$	$0.023$	$200$	$0.126$	$15$
22.	$2\pi/5$ , 180°, 0.525	$32$	$0.839$	$31$	$0.023$	$200$	$0.097$	$14$
23.	$2\pi/5$ , 180°, 0.550	$33$	$0.839$	$32$	$0.023$	$200$	$0.068$	$13$
24.	$2\pi/5$ , 180°, 0.575	$34$	$0.839$	$33$	$0.023$	$200$	$0.040$	$12$
25.	$2\pi/5$ , 180°, 0.600	$35$	$0.839$	$34$	$0.023$	$200$	$0.011$	$11$
26.	$2\pi/5$ , 180°, 0.625	$36$	$0.839$	$35$	$0.023$	$200$	$0.000$	$10$
27.	$2\pi/5$ , 180°, 0.650	$37$	$0.839$	$36$	$0.023$	$200$	$-0.011$	$9$
28.	$2\pi/5$ , 180°, 0.675	$38$	$0.839$	$37$	$0.023$	$200$	$-0.040$	$8$
29.	$2\pi/5$ , 180°, 0.700	$39$	$0.839$	$38$	$0.023$	$200$	$-0.071$	$7$
30.	$2\pi/5$ , 180°, 0.725	$40$	$0.839$	$39$	$0.023$	$200$	$-0.102$	$6$
31.	$2\pi/5$ , 180°, 0.750	$41$	$0.839$	$40$	$0.023$	$200$	$-0.133$	$5$
32.	$2\pi/5$ , 180°, 0.775	$42$	$0.839$	$41$	$0.023$	$200$	$-0.164$	$4$
33.	$2\pi/5$ , 180°, 0.800	$43$	$0.839$	$42$	$0.023$	$200$	$-0.195$	$3$
34.	$2\pi/5$ , 180°, 0.825	$44$	$0.839$	$43$	$0.023$	$200$	$-0.226$	$2$
35.	$2\pi/5$ , 180°, 0.850	$45$	$0.839$	$44$	$0.023$	$200$	$-0.257$	$1$
36.	$2\pi/5$ , 180°, 0.875	$46$	$0.839$	$45$	$0.023$	$200$	$-0.288$	$0$
37.	$2\pi/5$ , 180°, 0.900	$47$	$0.839$	$46$	$0.023$	$200$	$-0.319$	$-1$
38.	$2\pi/5$ , 180°, 0.925	$48$	$0.839$	$47$	$0.023$	$200$	$-0.350$	$-2$
39.	$2\pi/5$ , 180°, 0.950	$49$	$0.839$	$48$	$0.023$	$200$	$-0.381$	$-3$
40.	$2\pi/5$ , 180°, 0.975	$50$	$0.839$	$49$	$0.023$	$200$	$-0.412$	$-4$
41.	$2\pi/5$ , 180°, 1.000	$51$	$0.839$	$50$	$0.023$	$200$	$-0.443$	$-5$
42.	$2\pi/5$ , 180°, 1.025	$52$	$0.839$	$51$	$0.023$	$200$	$-0.474$	$-6$
43.	$2\pi/5$ , 180°, 1.050	$53$	$0.839$	$52$	$0.023$	$200$	$-0.505$	$-7$
44.	$2\pi/5$ , 180°, 1.075	$54$	$0.839$	$53$	$0.023$	$200$	$-0.536$	$-8$
45.	$2\pi/5$ , 180°, 1.100	$55$	$0.839$	$54$	$0.023$	$200$	$-0.567$	$-9$
46.	$2\pi/5$ , 180°, 1.125	$56$	$0.839$	$55$	$0.023$	$200$	$-0.600$	$-10$
47.	$2\pi/5$ , 180°, 1.150	$57$	$0.839$	$56$	$0.023$	$200$	$-0.631$	$-11$
48.	$2\pi/5$ , 180°, 1.175	$58$	$0.839$	$57$	$0.023$	$200$	$-0.662$	$-12$
49.	$2\pi/5$ , 180°, 1.200	$59$	$0.839$	$58$	$0.023$	$200$	$-0.693$	$-13$
50.	$2\pi/5$ , 180°, 1.225	$60$	$0.839$	$59$	$0.023$	$200$	$-0.724$	$-14$
51.	$2\pi/5$ , 180°, 1.250	$61$	$0.839$	$60$	$0.023$	$200$	$-0.755$	$-15$
52.	$2\pi/5$ , 180°, 1.275	$62$	$0.839$	$61$	$0.023$	$200$	$-0.786$	$-16$
53.	$2\pi/5$ , 180°, 1.300	$63$	$0.839$	$62$	$0.023$	$200$	$-0.817$	$-17$
54.	$2\pi/5$ , 180°, 1.325	$64$	$0.839$	$63$	$0.023$	$200$	$-0.848$	$-18$
55.	$2\pi/5$ , 180°, 1.350	$65$	$0.839$	$64$	$0.023$	$200$	$-0.879$	$-19$
56.	$2\pi/5$ , 180°, 1.375	$66$	$0.839$	$65$	$0.023$	$200$	$-0.910$	$-20$
57.	$2\pi/5$ , 180°, 1.400	$67$	$0.839$	$66$	$0.023$	$200$	$-0.941$	$-21$
58.	$2\pi/5$ , 180°, 1.425	$68$	$0.839$	$67$	$0.023$	$200$	$-0.972$	$-22$
59.	$2\pi/5$ , 180°, 1.450	$69$	$0.839$	$68$	$0.023$	$200$	$-1.003$	$-23$
60.	$2\pi/5$ , 180°, 1.475	$70$	$0.839$	$69$	$0.023$	$200$	$-1.034$	$-24$

## ROOF SECTIONS

<p>3" L.C. + 4" M.P. + 4½" R.B. + ½" p</p>  <p>(10)</p>	<p>2½" b.t. + 3" L.C. + 4½" R.B. + ½" p</p>  <p>(11)</p>	<p>½" p + 2" F.C. + 2" R.C.C. + ½" p</p>  <p>(12)</p>
<p>4½" brick + 3" L.C. + 4" R.C.C. + ½" p</p>  <p>(13)</p>	<p>½" slate + 4" air space + ½" p.w.</p>  <p>(14)</p>	<p>½" w.s. + 4" air space + ½" p.w.</p>  <p>(15)</p>
<p>AC sheet + 4" air space + ½" p.w.</p>  <p>(16)</p>	<p>½" Mangalore tile</p>  <p>(17)</p>	<p>½" m.t. + 4" air space + ½" p.w.</p>  <p>(18)</p>

L.C. = Lime concrete    M.P. = Mud phuska    b.t.=bricktile  
 R.B. = Reinforced brick    R.C.C.= Reinforced cement conc.  
 w.s. = wood shingle    p.w. = plywood    m.t. = Mangalore tile

G.I. = Gold standard mean.  
D.D. = Duplicate mean.

A-C. = Incomplete Concourse.

B-C. = Incomplete Concave.

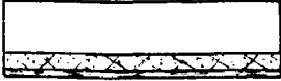
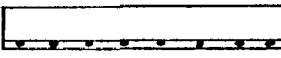
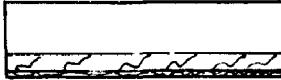
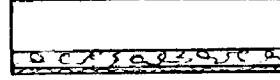
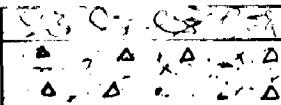
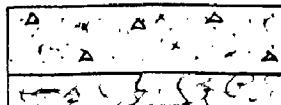
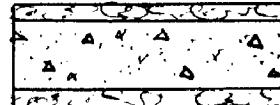
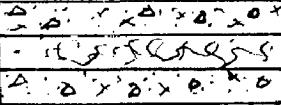
C-C. = Incomplete Convex.

27. 21. 22. 23. 24. 25. 26. 27. 28. 29. 30. 31. 32. 33. 34. 35. 36. 37. 38. 39. 40. 41. 42. 43. 44. 45. 46. 47. 48. 49. 50. 51. 52. 53. 54. 55. 56. 57. 58. 59. 60. 61. 62. 63. 64. 65. 66. 67. 68. 69. 70. 71. 72. 73. 74. 75. 76. 77. 78. 79. 80. 81. 82. 83. 84. 85. 86. 87. 88. 89. 90. 91. 92. 93. 94. 95. 96. 97. 98. 99. 100.

|                |                |                |                |                |                |                |                |                |                |  |
|----------------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|----------------|--|
| 10. G.I. - 3.6 | 11. G.I. - 1.6 | 12. D.D. - 3.6 | 13. D.D. - 1.6 | 14. A-C. - 3.6 | 15. A-C. - 1.6 | 16. B-C. - 3.6 | 17. B-C. - 1.6 | 18. C-C. - 3.6 | 19. C-C. - 1.6 | 20. 21. 22. 23. 24. 25. 26. 27. 28. 29. 30. 31. 32. 33. 34. 35. 36. 37. 38. 39. 40. 41. 42. 43. 44. 45. 46. 47. 48. 49. 50. 51. 52. 53. 54. 55. 56. 57. 58. 59. 60. 61. 62. 63. 64. 65. 66. 67. 68. 69. 70. 71. 72. 73. 74. 75. 76. 77. 78. 79. 80. 81. 82. 83. 84. 85. 86. 87. 88. 89. 90. 91. 92. 93. 94. 95. 96. 97. 98. 99. 100. |
| 10. G.I. - 3.6 | 11. G.I. - 1.6 | 12. D.D. - 3.6 | 13. D.D. - 1.6 | 14. A-C. - 3.6 | 15. A-C. - 1.6 | 16. B-C. - 3.6 | 17. B-C. - 1.6 | 18. C-C. - 3.6 | 19. C-C. - 1.6 | 20. 21. 22. 23. 24. 25. 26. 27. 28. 29. 30. 31. 32. 33. 34. 35. 36. 37. 38. 39. 40. 41. 42. 43. 44. 45. 46. 47. 48. 49. 50. 51. 52. 53. 54. 55. 56. 57. 58. 59. 60. 61. 62. 63. 64. 65. 66. 67. 68. 69. 70. 71. 72. 73. 74. 75. 76. 77. 78. 79. 80. 81. 82. 83. 84. 85. 86. 87. 88. 89. 90. 91. 92. 93. 94. 95. 96. 97. 98. 99. 100. |
| 10. G.I. - 3.6 | 11. G.I. - 1.6 | 12. D.D. - 3.6 | 13. D.D. - 1.6 | 14. A-C. - 3.6 | 15. A-C. - 1.6 | 16. B-C. - 3.6 | 17. B-C. - 1.6 | 18. C-C. - 3.6 | 19. C-C. - 1.6 | 20. 21. 22. 23. 24. 25. 26. 27. 28. 29. 30. 31. 32. 33. 34. 35. 36. 37. 38. 39. 40. 41. 42. 43. 44. 45. 46. 47. 48. 49. 50. 51. 52. 53. 54. 55. 56. 57. 58. 59. 60. 61. 62. 63. 64. 65. 66. 67. 68. 69. 70. 71. 72. 73. 74. 75. 76. 77. 78. 79. 80. 81. 82. 83. 84. 85. 86. 87. 88. 89. 90. 91. 92. 93. 94. 95. 96. 97. 98. 99. 100. |
| 10. G.I. - 3.6 | 11. G.I. - 1.6 | 12. D.D. - 3.6 | 13. D.D. - 1.6 | 14. A-C. - 3.6 | 15. A-C. - 1.6 | 16. B-C. - 3.6 | 17. B-C. - 1.6 | 18. C-C. - 3.6 | 19. C-C. - 1.6 | 20. 21. 22. 23. 24. 25. 26. 27. 28. 29. 30. 31. 32. 33. 34. 35. 36. 37. 38. 39. 40. 41. 42. 43. 44. 45. 46. 47. 48. 49. 50. 51. 52. 53. 54. 55. 56. 57. 58. 59. 60. 61. 62. 63. 64. 65. 66. 67. 68. 69. 70. 71. 72. 73. 74. 75. 76. 77. 78. 79. 80. 81. 82. 83. 84. 85. 86. 87. 88. 89. 90. 91. 92. 93. 94. 95. 96. 97. 98. 99. 100. |

TABLE II (cont'd.)

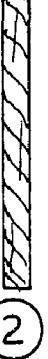
## ROOF SECTIONS

|  |  |   |
|--|--|---|
| <p>G.I.sheet + 3" air space + 1" m.w. + <math>\frac{1}{4}</math>" p.w.</p>  <p>(19)</p>   | <p>G.I.sheet + 3" air space + 2" m.w. + <math>\frac{1}{4}</math>" p.w.</p>  <p>(20)</p> | <p>A.C.sheet + 2" air space + <math>\frac{1}{2}</math>" c.b.</p>  <p>(21)</p> |
| <p>A.C.sheet + 3" air space + 1" m.w. + <math>\frac{1}{4}</math>" p.w.</p>  <p>(22)</p> | <p>A.C.sheet + 3" air space + 1" T.C. + <math>\frac{1}{4}</math>" p.w.</p>  <p>(23)</p>    | <p>2" T.C. + 4" D.C.</p>  <p>(24)</p>                               |
| <p>4" D.C. + 2" T.C.</p>  <p>(25)</p>  | <p>1" T.C. + 4" D.C. + 1" T.C.</p>  <p>(26)</p>  | <p>2" D.C. + 2" T.C. + 2" D.J.</p>  <p>(27)</p>                           |

m.w. = mineral wool p.w. = plywood c.b. = celotex board  
 D.C. = Dense concrete T.C. = Thermocole

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## DOOR & WINDOW SECTIONS

|  |  |   |
|--|--|---|
| 1 $\frac{1}{4}$ " Teakwood<br><br>①                                   | 1 $\frac{1}{2}$ " Teakwood<br><br>②   | 1 $\frac{3}{4}$ " Teakwood<br><br>③                              |
| $\frac{1}{2}$ " p.w. + 1" air space + $\frac{1}{2}$ " p.w.<br><br>④ | $\frac{1}{2}$ " h.b. + 1" air space + $\frac{1}{2}$ " h.b.<br><br>⑤           | $\frac{1}{2}$ " p.w. + 1" T.C. + $\frac{1}{2}$ " p.w.<br><br>⑥ |
| $\frac{1}{2}$ " h.b. + 1" T.C. + $\frac{1}{2}$ " h.b.<br><br>⑦      | $\frac{1}{2}$ " p.w. + 1 $\frac{7}{8}$ " m.w. + $\frac{1}{2}$ " p.w.<br><br>⑧ | $\frac{1}{2}$ " h.b. + 1" m.w. + $\frac{1}{2}$ " h.b.<br><br>⑨ |

p.w. = plywood    h.b. = hardboard    T.C. = Thermocole  
 m.w. = mineral wool

TABLE I. SURFACE FRICTION COEFFICIENTS FOR  
INDUSTRIAL DRIVING

$D_0 = 1.6$   
 $D_4 = 1.6$

| No. | Internal Friction<br>Coefficient                               | Floor<br>Location | Volume<br>Capacity | Point Contact | Friction<br>Coefficient | Internal<br>Friction<br>Coefficient | Point Contact | Friction<br>Coefficient | Volume<br>Capacity | Point Contact | Friction<br>Coefficient | Internal<br>Friction<br>Coefficient | Point Contact | Friction<br>Coefficient | Volume<br>Capacity | Point Contact | Friction<br>Coefficient | Internal<br>Friction<br>Coefficient | Point Contact | Friction<br>Coefficient |
|-----|--|-------------------|--------------------|---------------|-------------------------|-------------------------------------|---------------|-------------------------|--------------------|---------------|-------------------------|-------------------------------------|---------------|-------------------------|--------------------|---------------|-------------------------|-------------------------------------|---------------|-------------------------|
| 1.  | 1½" Concrete<br>+ 4" R.C.C.<br>+ 1/2" Plaster                  | 0.480             |                    |               | 0.432                   | 0.432                               |               |                         |                    |               |                         | 0.422                               |               |                         |                    |               |                         | 0.333                               |               |                         |
| 2.  | 3½" concrete<br>+ 4" R.C.C.<br>+ 1/2" Plaster.                 | 0.337             | F 2<br>H 2<br>D 2  |               | 0.443                   | 0.443                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 3.  | 3½" concrete<br>+ 1¾" concrete<br>+ 4" R.B. + 1/2"<br>Plaster. | 0.365             | F 2<br>H 2<br>D 2  |               | 0.444                   | 0.444                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.320                               |               |                         |
| 4.  | 1½" concrete<br>+ 1/2"<br>Plaster.                             | 0.325             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 5.  | 1½" concrete<br>+ 1/2"<br>Plaster.                             | 0.323             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 6.  | 1½" concrete + G<br>+ D.G. + 1/2"<br>Plaster.                  | 0.491             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 7.  | 2½" concrete<br>+ 1/2"<br>Plaster.                             | 0.422             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 8.  | 2½" terrazzo<br>+ 1/2"<br>Plaster.                             | 0.490             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 9.  | 2½" concrete<br>+ 1/2"<br>Plaster.                             | 0.489             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 10. | ½" block stone<br>+ 1/2"<br>Plaster.                           | 0.475             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |
| 11. | ½" block<br>+ 1/2"<br>Plaster.                                 | 0.424             | F 2<br>H 2<br>D 2  |               | 0.445                   | 0.445                               |               |                         |                    |               |                         | 0.416                               |               |                         |                    |               |                         | 0.300                               |               |                         |

TABLE 5

## THEORETICAL DESIGN FUNCTIONS FOR PARTITION WALLS

 $n_0 = 1.05$   
 $n_2 = 1.05$ 

| No.     | Post-tension<br>wall<br>section  | G<br>value | Normal<br>size | In conform<br>connection |                                  | Internal driving<br>moment, $F_m$ |                                  | Angle, $\phi_c$<br>in<br>degrees | Angle, $\phi_c$<br>in<br>degrees |
|---------|--|------------|----------------|--------------------------|----------------------------------|-----------------------------------|----------------------------------|----------------------------------|----------------------------------|
|         |  |            |                | $\lambda_c$              | $\lambda_c$<br>in<br>degrees     | $\lambda_c$                       | $\lambda_c$<br>in<br>degrees     |                                  |                                  |
| • 202 • |  |            |                |                          |                                  |                                   |                                  |                                  |                                  |
| 1.      | $1/2''$ Plaster + $3/8''$ brick<br>$\Rightarrow 1/2''$ plaster + $3/8''$ brick | 0.200      |                | P<br>H2<br>H3<br>H4      | 0.204<br>0.122<br>0.076<br>0.054 | 53<br>61<br>105<br>130            | 0.511<br>0.337<br>0.309<br>0.300 | 30<br>32<br>33<br>33             |                                  |
| 2.      | $1/2''$ Plaster + $3/8''$ brick<br>$\Rightarrow 1/2''$ plaster                 | 0.239      |                | P<br>H2<br>H3<br>H4      | 0.073<br>0.029<br>0.013<br>0.009 | 116<br>162<br>214<br>238          | 0.499<br>0.382<br>0.339<br>0.303 | 30<br>31<br>35<br>40             |                                  |
| 3.      | $1/2''$ plaster + $1/8''$ plaster<br>coated + $3/8''$ plaster                  | 0.185      |                | P<br>H2<br>H3<br>H4      | 0.107<br>0.073<br>0.051<br>0.033 | 32<br>80<br>118<br>140            | 0.322<br>0.281<br>0.256<br>0.209 | 21<br>23<br>24<br>30             |                                  |
| 4.      | $2/4''$ plaster + $3/8''$ wood<br>beam + $3/8''$ plaster                       | 0.155      |                | P<br>H2<br>H3<br>H4      | 0.038<br>0.034<br>0.034<br>0.034 | 35<br>63<br>88<br>104             | 0.733<br>0.634<br>0.570<br>0.500 | 19<br>31<br>35<br>35             |                                  |
| 5.      | $1/2''$ plaster + $3/8''$ cylinder<br>block + $1/2''$ plaster                  | 0.331      |                | P<br>H2<br>H3<br>H4      | 0.114<br>0.085<br>0.069<br>0.046 | 100<br>139<br>173<br>210          | 0.532<br>0.503<br>0.373<br>0.339 | 21<br>27<br>32<br>33             |                                  |
| 6.      | $4/4''$ synthetic wood panel<br>and plaster                                    | 0.150      |                | P<br>H2<br>H3<br>H4      | 0.090<br>0.035<br>0.016          | 53<br>82<br>103<br>125            | 0.856<br>0.323<br>0.277<br>0.235 | 13<br>16<br>17<br>18             |                                  |
| 7.      | $3/4''$ hard plaster + $3/4''$<br>base battening + $3/4''$<br>and plaster.     | 0.453      |                | P<br>H2<br>H3<br>H4      | 0.273<br>0.213<br>0.159          | 30<br>52<br>75<br>93              | 0.622<br>0.523<br>0.400          | 17<br>22<br>33<br>30             |                                  |
| 8.      | $1/2''$ hard plaster + $3/8''$ base board<br>and plaster.                      | 0.316      |                | P<br>H2<br>H3<br>H4      | 0.191<br>0.162<br>0.124          | 30<br>52<br>72<br>93              | 0.729<br>0.677<br>0.603<br>0.584 | 19<br>25<br>31<br>36             |                                  |

**THE TIME ESTIMATE FOR EXECUTING THE PROJECT IS 12 MONTHS.**

卷之五

R E F E R E N C E S

LITERATURE

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## NOMENCLATURE

|          |  |
|----------|--|
| K        | Thermal conductivity of the material<br>in Btu/Ft./Hr. $^{\circ}$ F  |
| $\rho$   | Density of the material in Lb/Ft <sup>3</sup>  |
| s        | Specific heat of the material in Btu/Lb. $^{\circ}$ F  |
| a        | Temperature conductivity (thermal diffusivity)<br>of the material in Ft <sup>2</sup> /Hr ( a = K/ $\rho s$ )   |
| p        | Coefficient of thermal absorption of the<br>material in Btu/Hr.Ft <sup>2</sup> . $^{\circ}$ F ( p = $\sqrt{\omega K \rho s}$ )                               |
| b        | Specific coefficient of thermal absorption of<br>the material in Btu/Hr.Ft <sup>2</sup> . $^{\circ}$ F ( b = $\sqrt{\kappa \rho s}$ )                        |
| U        | Coefficient of thermal transmission of the wall<br>in Btu/Hr.Ft <sup>2</sup> . $^{\circ}$ F  |
| L        | Thickness of a layer in Ft.  |
| $R_s$    | Thermal Resistance of a building section<br>in Hr.Ft <sup>2</sup> . $^{\circ}$ F/Btu. (L/K)  |
| C        | Thermal Capacity of a building section<br>in Btu/Ft <sup>2</sup> . $^{\circ}$ F (L $\rho s$ )  |
| h        | Coefficient of heat transfer from the surrounding<br>medium (air) to the surface or from the surface<br>to the air in Btu/Hr.Ft <sup>2</sup> . $^{\circ}$ F. |
| t        | Temperature variation (sinusoidal) in $^{\circ}$ F   |
| P        | Fundamental period (24 hours)  |
| T        | Time in hours.   |
| q        | Variable thermal current (heat flux)<br>in Btu/Hr.Ft <sup>2</sup>  |
| B        | Thermal absorption capacity of a building section<br>in Btu/Hr. $^{\circ}$ F   |
| $\gamma$ | Propagation const. $\sqrt{j\omega \rho s / K}$   |
| $\Theta$ | Dimensionless parameter characteristic of the<br>building section ( $j\omega \beta R_s$ )  |
| $j$      | $\sqrt{-1}$  |

- $\delta$  Wave length of the temperature wave in a material in Ft.  
 $\sigma$  Stephen Boltzman constant ( $0.1712 \times 10^{-3}$  Btu/Ft<sup>2</sup>.hr. $^{\circ}$ F<sup>4</sup>)  
 $T$  Absolute temperature ( $^{\circ}$ F + 459.69)  
 $\epsilon$  Emissivity of the surface.  
 $\alpha$  Absorbtivity of the surface to solar radiation.  
 $I_p$  Intensity of incident solar radiation in Btu/Ft<sup>2</sup>.hr.  
 $I_{LT}$  Difference between low temperature radiation emitted and received by the surface in Btu/Ft<sup>2</sup>.hr.  
 $F$  Area of the surface in Ft<sup>2</sup>.  
 $p_v$  Vapour pressure in inches of mercury.  
 $\omega$  Thermal impedance in hr.Ft<sup>2</sup>. $^{\circ}$ F/Btu.  
 $\beta_c$  Characteristic thermal impedance in hr.Ft<sup>2</sup>. $^{\circ}$ F/Btu.  
 $\chi$  Thermal admittance in Btu/hr.Ft<sup>2</sup>. $^{\circ}$ F  
 $N$  Number of lumps  
 $P$  Fundamental frequency.  
 $E_2$  Second harmonic  
 $E_3$  Third harmonic  
 $E_4$  Fourth harmonic  
 $R_o$  Outside surface resistance ( $1/h_o$ ) in hr.Ft<sup>2</sup>. $^{\circ}$ F/Btu.  
 $R_i$  Inside surface resistance ( $1/h_i$ ) in hr.Ft<sup>2</sup>. $^{\circ}$ F/Btu.  
 $\left. \begin{matrix} A \\ B \\ C \end{matrix} \right\}$  Coefficients of transfer matrix for a building section.  
 $\lambda_o$  Modulus of the external driving point function  $\left\{ \frac{t_{os}}{t_{ss}} \right\}$   
 $\phi_o$  Argument of the external driving point function.  
 $\lambda_i$  Modulus of the transfer function ( $t_{is}/t_{ss}$ )  
 $\phi_i$  Argument of the transfer function (

- $\lambda'_i$  Modulus of the Internal Driving point Function  $\{t_{i_m}/\bar{t}_{i_m}\}$
- $\phi'_i$  Argument of the internal driving point function.
- $t_0$  Modulus of the external driving point admittance function ( $q_{os}/t_{sa}$ ).
- $\psi_0$  Argument of the external driving point admittance function.
- $\epsilon_i$  Modulus of the transfer admittance function ( $q_{is}/t_{sa}$ )
- $\psi_i$  Argument of the transfer admittance function.
- $\lambda'_i$  Modulus of the internal driving point admittance function ( $q_{ig}/t_{ia}$ ).
- $\phi'_i$  Argument of the internal driving point admittance function.

Subscripts

- in internal air
- ca soil-air
- os external air
- es external surface
- is internal surface
- n harmonic of the frequency

