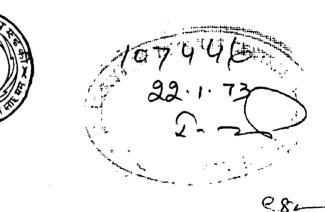
## **STUDIES ON REVERSION IN AGED ALLOYS**

A Dissertation submitted in partial fulfilment of the requirements for the degree of MASTER OF ENGINEERING in METALLURGICAL ENGINEERING (Physical Metallurgy)

> By S. K. SHARMA



DEPARTMENT OF METALLURGICAL ENGINEERING UNIVERSITY OF ROORKEE ROORKEE August, 1972

#### <u>**OERTIPICATE**</u>

Certified that the dissertation entitled "STUDIES ON REVERSION IN AGED ALLOYS", which is being submitted by Shri S. K. Sharma, in partial fulfilment for the award of the Degree of Unster of Engineering in Detallurgical Engineering (Physical Netallurgy) of the University of Roorkee, Roorkee, is a record of his own work, carried out by him under my supervision and guidance, from 3rd January, 1972 to 6th August, 1972.

not been subsitted for the sward of any other degree.

T. V. Rajan

[7. V. Rajen] Reader Department of Hotallurgical Engg., University of Roorkoo, Roorkoe.

Roorkoo. Dated: August /2,1972.

## ACKHOVLEDGEMENTS

I wish to express my deep sense of gratitude to Dr. T. V. Rajan, Reader in Department of Hetallurgical Engineering, University of Roorkee, Roorkee, for his valuable guidance, inspiration and useful discussions during the course of this work. Without his encouragement at every stage of the work, this thesis could not have been completed.

I am also very thankful to Dr. H. H. Samona, Professor and Head, Department of Hetallurgical Engineering, University of Reerkoe for providing the necessary laboratory facilities. By thanks are also due to my friends and the laboratory staff the assisted me in my work.

1. Te Shanna

[SUNIL KUMAR SHARMA]

#### ABSTRACT

'The experimental investigations presented in this thesis were undertaken to study the effects of homogenisation temperature, ageing temperature, pre-ageing time followed by step agoing and / or partial reversion of G.P. zones on reageing at temperatures in the vicinity of room temperature ( $0^{\circ}$ C to  $30^{\circ}$ C), in Aluminium-4.17 at. per cent Zine alloys by electrical resistivity studies. A comprehensive study was undertaken by choosing judicious combinations of the above factors, so that on the basis of findings reported in this thesis, there is possibility of a botter understanding of the mechanism of pre-precipitation stages in Aluminium-Zine alloys.'

## CONTENTS

.

• •

1.0	LITERATURE SURVEY			
1.1	Introduction	••	1.	
1.2	Historical Review	• •	2.	
1.3	Present understanding of precipitation phenome	enon	5.	
1.3.1	Introduction	• •	5.	
1.3.2	Clustering	••	6.	
	a) Nucleation of clusters	• •	7.	
	b) Growth and kinetics of clustering	• •	8.	
1.3.3	Formation of zones and their growth	• •	11.	
1.3.4	Intermediate precipitate	••	13.	
1.3.5	Equilibrium precipitate	• •	14.	
1.3.6	The sequence of precipitation	• •	15.	
1.3.7	Role of quench-in vacancies	• •	16.	
1.3.8	General conclusions on precipitation phenomen	on	20.	
1.4	Techniques for precipitation studies	• • .	22.	
1.5	Discussions on resistemetric studies on			
	precipitation		24.	
1.5.1	Theory	* *	24.	
1.5.2	Resistivity increase and the resistivity maximum 26			
1.5.3	Conclusions from resistivity studies	* *	28.	
1.6	Aluminium-zinc system	• •	30.	
1.7	Observations on resistivity studies on Al-Zn			
	precipitation system	* 4	35.	
1.8	Reversion and Al-Zn metastable phase diagram	••	37.	
1.8.1	Metastable phase diagram	• •	37.	
1.8.2	Reversion phenomenon	• •	39.	
2.0	EXPERIMENTAL PROCEDURE			
2.1	Preparation of alloy and specimen	* ÷	41.	
2.1.1	Welting and casting	••	41.	

Page.

2.1.2	Hot forging and wire drawing	• •	41.
2.1.3	Specimen preparation for resistivity		
	measurements	••	41.
2.1.4	Heat treatment	• •	41.
2.2	Resistivity measurements of specimen		
·	during ageing	• •	42.
2.2.1	Description of apparatus	• •	42.
2.2.2	Calculation of resistance and resistivity	••	43.
2.2.3	Mathematical exprapolation to find initial		
· .	resistivity ( Po)	* •	43.
2.2.4	Calculation of resistivity changes ( $\Delta \varrho$ )	••	
	and the thermal coefficient of electrical		
	resistivity (ä)	••	44.
2.3	Experimental scheme for the present study	••	45.
3.0	RESULTS	••	47.
4.0	DISCUSSION OF RESULTS	• •	48.
5.0	CONCLUSIONS	••	63.
6.0	SUGGESTIONS FOR FURTHER WORK	••	64.
7.0	REFERENCES	••	65.

•

,

,

.

٠

.

٠

•

· .

**x** 

## LITERATURE - SURVEY

-

#### 1.1 INTRODUCTION

With the beginning of the twentieth century many new phenomenon came into existence. Alfred Wilm<sup>(1,2)</sup>(1903-1911) discovered a new hardening process during his work on Duralumin. He found that this alloy can be hardened quite appreciably by quenching from temperatures below it's melting point and then allowing it to stay at room temperature. The hardness is not produced by quenching alone, but increases during the period of stay at room temperature, and thus the phenomenon got the name of <u>age-hardening</u>. Merica<sup>(3)</sup>(1930) traced it's cause to precipitation of second phase with lapse of time out of the super-saturated solid solution and called it '<u>precipi</u>tation hardening'.

Prom a metallurgical point of view, it is nucleation and growth type of transformation in which atomic re-arrangement is involved mainly by the phenomenon of diffusion, resulting in the break down of the super-saturated solid solution. These changes in the structure and constitution of the alloy during ageing, results in changes in physical, particularly nechanical properties of the alloy.

The above phenomenon has widely been used to improve the properties of various ferrous and non-ferrous alloys. Thus it has been proved very helpful in development of super high otrongth light alloys componial for various industrial purpecce. After stoch, Aluminium alloys are the most important materials used for structural purpose, through the use of the phonomenon of precipitation hardening.

#### 1.2 HISTORICAL REVIEW

As pointed out in the introduction, the carliest corhere on age-herdening core  $\operatorname{Vilm}^{(1,2)}$  and  $\operatorname{Horica}^{(3)}$  the undertook detailed studies of the process.

To other with, it was assumed that the phase procipitating out from the supercaturated colid solution is the equilibrium phase. Hardening was associated with the formation of submissessopic particles of equilibrium phase. However, resistmetric studies of phase changes during precipitation, revealed an anomalous behaviour (4,5). If agoing involved simple precipitation of equilibrium phase, then there should be stoody deerease in resistivity (2) with agoing. But it was observed that there is first an increase in (2, 5) followed by the expected decrease. This behaviour was explained by Gayler and Preston<sup>(6)</sup> on the basis of <u>two alongs precipitation</u> process. The first stage being pro-precipitation of an intermediate or non-equilibrium phase to the equilibrium phase. The internal changes in the lattice of super caturated alloy, resulted in age-hardoning. Resenhain<sup>(7)</sup> and Tammann<sup>(8)</sup> also postulated theories of two stage precipitation.

Horica<sup>(9)</sup> in 1932 suggested that the first stop in precipitation is sogregation of solute atoms on cortain sites within solid solution and referred these regions as <u>knots</u>, which were associated with strain field around them, due to size difforence (miefit) between the solute and the colvent atoms. Proceed of such knots render the process of slip to be difficult leading to hardening. Experimental work of Jenkine and Buokmall<sup>(10)</sup> lent support to Herica's suggestions and the existence of clusters (11,12) was established at a later stage.

Application of small angle X-ray scattering (19,15) lod to the discovery of <u>some</u> by Guinier<sup>(16)</sup> and Proston<sup>(17)</sup>, giving a major break through in the studies on agoing. It is now establiched (18,20) that difference between clusters and somes is only of nomenolature. The hardening mechanics accordated with somes (21,23) has undergone considerable medifications from the one suggested by Morica.

With the cotablishment of <u>non-equilibrium transition</u> <u>phase</u> (0.g. 0<sup>t</sup> in Al-Ou alloys), the concept of <u>cohorency</u> came into existence. Depending upon the <u>nature of interface</u> between

the procipitate and matrix, the procipitation is called cohoront, comi-cohorent or non-cohorent i.e. thether all the lattice planes of the matrix and the precipitate are matching, only come are matching and the none is matching respectively. The matrix and the procipitate have different lattice constants, and hence a cohorent procipitate is associated with cohorency strains, reculting in hardening. Particles of non-equilibrium precipitate to start with, are fully cohorent v.r.t. matrix. With growth, the elactic strain energy (24,25) associated with these increacoulting in lowering of hardness. The credit of explaining the process of precipitation hardening including <u>evergening</u> goes to Mott and Habarre<sup>(26)</sup> on the basis of cohorency strains.

With this bachground the structural changes during agoing can be stated as follows:

> Super saturated colid colution  $\rightarrow$  elustoring  $\rightarrow$ formation of somes and their growth  $\rightarrow$  precipitation of non-equilibrium transitional phase  $\rightarrow$  tranefermation of non-equilibrium phase into equilibrium phase.

Actual number of stops involved varies from system to system, and depends on various factors, like the degree of supercaturation, which in turn depends on colute concentration. Also as agoing is by diffusional process involving colute atoms and

quench-in vacancies, the homogenisation temperature  $(T_{\rm H})$  and the agoing temperature  $(T_{\rm A})$  are most important factors governing the agoing process. Bature of interface between the matrix and the second phase is the most important parameter in controlling the structural properties of ago-hardenable alloy.

#### 1.5 PRESENT UNDERSTANDING OF THE PRECIPITATION PHELOMONON

#### 1.3.1 Introduction

Precipitation of the stable equilibrium phase(s), out of the supersaturated solid solution, takes place in a number of intermediate steps, and their acquence is dependent on various agoing variables like the  $\Sigma_{\rm H}$ ,  $\Sigma_{\Lambda}$ , composition of alloy, and so on. The characteristic steps, and their acquence is dependent on various agoing variables like the  $\Sigma_{\rm H}$ ,  $\Sigma_{\Lambda}$ , composition of alloy, and so on. The characteristic steps, and their acquence of alloy, and so on. The characteristic steps, and the stable driving force for procipitation depends on the degree of supersaturation and decreases as precipitation proceeds. Precipitation follows the path of minisum <u>notivational energy</u> rather than maximum everall loss of free energy. As the primary node of precipitation is by <u>diffusion</u> phonomenon, which amounts to simple interchange of atoms and encodes quench-in vacancies, the factors influencing these like the  $\Sigma_{\rm H}$ , quanching rate,  $\Sigma_{\Lambda}$ , solute concentration etc., govern the structure and sequence of precipitation. The offects of cold working, prior be satisfactorily explained on the process of annealing out of excess quench-in vacancies, diffusion of colute atoms, and the interactions between the vacancies and colute atoms.

In general the precipitation sequence can be classified into following steps, in increasing order of the activational energy:

(1)	Clustoring (Cohorent)	
(11)	Zong formation.	Pro-procipitation

(111) Pornation of intermediate transitional phase.

(iv) Pormation of non-cohorent, equilibrium phace.

At the present study is mainly concerned with the earlier etages (Pro-precipitation) of agoing, we shall limit detailed discussions to clustering and some formation only.

#### 1.3.2 <u>Oluotoring</u>

The first stage in the process () procipitation invariably consists in the collection of solute atoms along certain proferential planes of the matrix, and as this takes place by the ' process of diffusion of solute atoms, the first aggregates of solute atoms are completely coherent with the matrix. The minfit, due to difference in the atomic sizes of solute and colvent atoms, must be accompdated by <u>clastic strains</u> to maintain coherency. Since all the f.c.c., metals have a minimum Young's fieldulus along  $\langle 100 \rangle$  direction, the initial eluctors in f.c.c., (A1) notals form as plates on {100} planes. This is also because A1, has anisotropic elacticity with  $E_{\langle 100 \rangle} = .99 \times 10^{12}$  dynos/cmg and  $E_{\langle 111 \rangle} = 1.14 \times 10^{12}$  dynos/cmg.

There is low activational energy for elustering, and it is almost equal to the activational energy for migration of vacancies, and as the clustering process has no nucleation barrier, the process is of <u>simple growth</u>. Clusters may be formed on agoing, or possibly at  $T_{\rm H}$ , or during the quenching.

The eluctors constitute stable cohorent regions having higher concentration of solute atoms. The elactic strains, due to difference in atomic class of solute and solvent, govern the form of these eluctors. Habarrow<sup>(24)</sup> showed that ,elastic strain energy of the cohorent eluctor increases as it grows with time, and this energy can be reduced, due to enicotropy of elastic constante, thereby causing changes in the initial shape of eluctors. Baker, Braden and Eutting<sup>(27)</sup> explained that elastic strains in and around growing eluctors may be relieved by aggregation of atoms and vacancies; atoms on most spacious planes and vacancies condense most eacily on eleced packed planes [111] in f.e.c., as maximum tensile strains occours in  $\langle 100 \rangle$  direction in cubic motals.

(a) <u>Auston of clustors</u>

Turnbull<sup>(28)</sup> has shown that the minitic law of eluctoring

10 of the form;

$$a_{t} = 1 - \left(\frac{1}{1+bt}\right)^{1/5} \qquad .. \qquad (1)$$

where as, equation for signoidal growth i.e. nucleation of growth type of transformation is;

$$a_{\pm} = 1 - \operatorname{onp}_{\bullet} (-b^{+} b^{\pm}) \qquad \dots \qquad (2)$$

Where,  $a_t = fraction of clustering process completed after$ time t

b = constant containing solute diffusion coofficient.

These are shown in Fig. C. Equation (1) shows, that there is no nucleation barrier for clustering, since the rate of clustering continuously decreases. Activational energy for clustering is low, and is equal to the activational energy for migration of colute atoms or single vacancies. Oritical size of the nucleus for clustering is of the order of one atomic dismeter.

#### (b) Growth and Kinotico of Clustering

With no nucleation barrier, the process consists of <u>com-</u> <u>potitive growth</u> of clusters, and then the rate of growth of a cluster is given by;

 $\frac{dr}{dt} \propto \frac{1}{r^0}$  where, r is radius of cluster at time t.

Abnormally high solute diffusion coefficient of solute in the colvent matrix during clustering was attributed to nonequilibrium concentration of vacancies in the quenched alloy. The excess vacancy theory attributed to Zener<sup>(29)</sup> and Seits<sup>(30)</sup> and developed by Federighi<sup>(31)</sup>, De Sorbo, Treaftic and Turnbull<sup>(11)</sup> and recently by Herman<sup>(17)</sup> explains most of the facts concerning the kinitics of clustering, and also the effects of  $T_{\rm H}$ , quenching rate, step quenching, coefficient of diffusion of colute atoms, reversion etc. The diffusion coefficient of solute (Zn) in Alluminium matrix is given by;

$$\begin{split} \mathbf{D}_{\mathrm{Sn}} &= \mathbf{A} \, \exp_{\mathbf{v}} - \left(\frac{\mathbf{E}_{\mathrm{F}} + \mathbf{E}_{\mathrm{H}}}{\mathbf{KT}}\right) \cdot & \text{where, } \mathbf{E}_{\mathrm{F}} \text{ and } \mathbf{E}_{\mathrm{H}} \text{ are activa-tional energies for formation} \\ & \text{tional energies for formation} \\ & \text{and migration of vacancies} \\ & \text{respectively.} \\ & \mathbf{K} = \text{boltzman's constant}_{=} 8.616 \times 10^{-5} \mathrm{eV/ec.} \end{split}$$

 $\mathbf{A}$  = another constant

Now, as diffusion is governed by a wacancy concentration which would be in equilibrium at  $T_{\rm H}$  rather than  $T_{\rm A}$ , and also to take into account the loss of vacancies to sinks during slow quenching, thus the retained vacancy concentration responding to some temperature value  $T_{\rm H^{+}}$  lower to  $T_{\rm H^{+}}$  the modified equation becomes)

$$D_{Zn} = A \exp(-E_H / KT_A) \exp(-E_p / KT_H)$$

This equation predicts that clustering process has activational energy  $E_{\underline{N}}$ , which is for the migration of vacancies. The second exponential of this equation is constant for the initial stages of egoing but decreases as vacancies are ennihilated at sinks.

Honourod notivational energies are clightly loss than its  $E_{\Box}$  for pure Al, and this is explained in terms of binding energy between colute atoms and & vacancies  $E_{S-V}(32)$ . Such a binding energy arises either due to elastic interaction (cise factor) or electrical interaction (valency factor).

Encode vacancy theory prodicts that the rate of elustering should fall while the vacancy concentration falls from  $(\operatorname{orp.} - \frac{E_P}{W_H})$  to onp.  $(\frac{-E_P}{W_A})$ ; the equilibrium value, as the encode vacancies are annihilated at surface, grain boundaries and other sinks.

Equilibrium to reached in  $A^{\text{punched}}_{A}$  (33) in about 30 min. at room temperature, and we observe that diffusion coefficient deereases fast in a short time and then remains near constant, thus dividing the clustering process into initial <u>fast reaction</u> followed by <u>alow reaction</u> (Fig. C). Cause of slow reaction was traced by Hart<sup>(34)</sup> Turnbull<sup>(12)</sup> and Hermon<sup>(117)</sup> suggesting that clusters and sonce <u>trap vacancies</u> due to binding energy between polute and vacancies.

If eluctore and sence de contain an appreciable volume fraction of vacancies, they night be relatively mobile and migrate amooda like and grow by <u>abcording conclor eluctors</u> or single



FIG.A. FRACTIONAL CHANGES IN P DURING ISOCHRONAL AGEING OF PURE AL QUENCHED FROM TH AS INDICATED

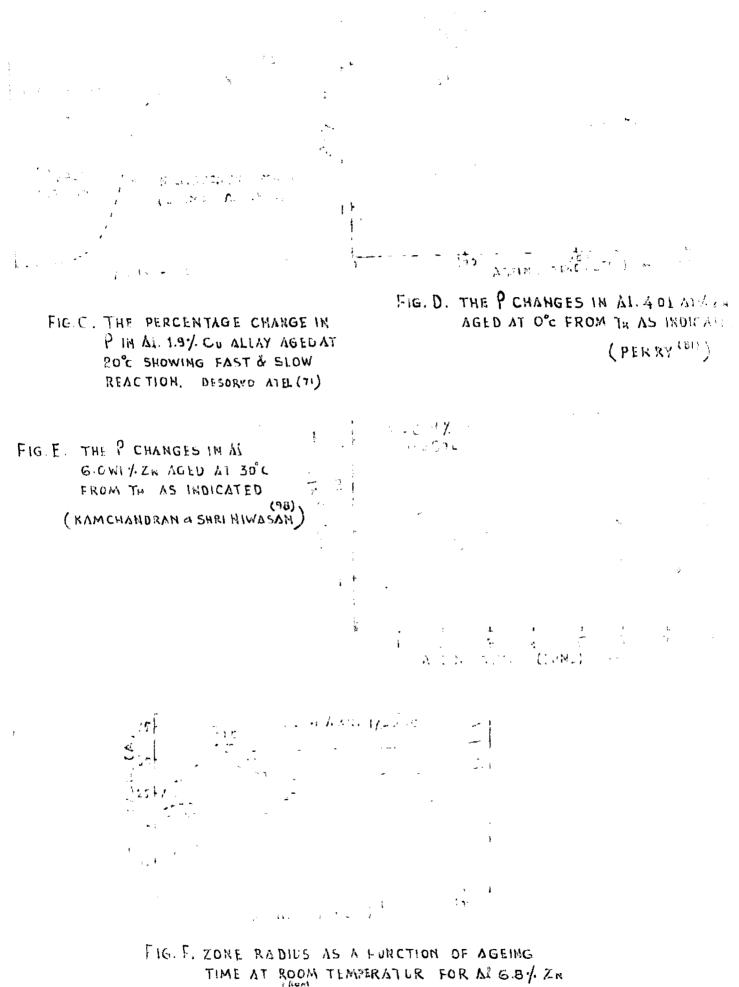
STAGE 2 - 50 TO - 100°C STAGE 1 140 TO 220°C

(PENSARI FEDERIGHI (14))



FIG. B. INDEPEDENT , WIND CURYLS OF ALL 10.2 WHY ZR MITHE NAMES - POOL TO 2000 FROM TH AS INDICATED

( PENSARI FEDERIGHI (84))



QUENCHED ATH AS INDICATED

(GEROLD & SCHEWEIKIER (30)

solute atoms.

The migrational activation energy would approximate to the sum of vacancy migration activation energy and the binding energy between the vacancy and the cluster.

Federighi and Thomas<sup>(35)</sup> modified Hart's mechanism, whereby zones remain static but act as controlled emitters and absorbers of vacancies. The zones are characterised by a binding energy with vacancies. If this is small ( $(AI - Ag, AI - Sn)_i$ ) the zones are poor sinks of vacancies, and there should be no slow reaction. If the binding energy is large (AI = Cu, AI-Mg-SI); the zones contain a high concentration of vacancies and slow reaction is predominant. The nature of interaction between solute and vacancy governs the ageing process by <u>Vacancy-trap mechanism</u>:

#### 1.3.3 Zone formation and their growth

When the clusters attain a definite shape, they are termed as zones (20), and the difference between clusters and zones was only of nomenclature. Clustering, which is over in a short time, is followed by <u>competitive</u> growth of zones. Gerold and Schweizer<sup>(15,36)</sup> studied spherical zones is Al-En system by X-ray studies and found the radius of zones to be a function of  $T_{\rm H}$  and alloy content (Fig. F). Initial rate of growth is very rapid(83), but after some time zone radius reaches almost a constant value, and these stages correspond to fast and slow reaction. Herman<sup>(127)</sup> studied the dependence of the rate of zone growth on vacancies and found emperically;

1 - X = exp. -  $(t/t)^{2}$  Here, X = fraction of solute associated with zones at time t and is equal to  $(C_0 - C_t / C_0 - C_f)$ ,  $C_0$ ,  $C_t$  and  $C_f$  are concentration of solute in matrix initially, at time t, and finally.

In above equation, m is constant depending on geometry of diffusion field, and  $\tau$  is relaxation time given by;  $\tau = a \times C_{\psi}\Gamma$ , in which, N is number of growing zones and  $\Gamma$  is jump frequency.  $\tau$  depends on sink for solute, and  $D_{g}$  which is given by;  $D_{g}= d^{\pm} C_{\psi}\Gamma$ where, d is jump distance.

Gerold and Schweizer<sup>(36)</sup> consider two models of the segregation process. Model I is a conventional distribution of precipitates in a super saturated matrix, where as Model 2 by Walker and Guinier<sup>(37)</sup> consists of clusters of solute atoms surrounded by a shell of denuded in solute with the matrix remaining supersaturated at the original composition of the alloy.

Gerold suggests that the degree of segregation is constant in model 1, whereas it should increase with ageing in model 2. Considering model 1 Gerold<sup>(36)</sup> calculated the concentration of solute inside the sones to be 69, pct. and that outside to be 1.8

:•

per cent. He interprets these values in terms of metaetability gap limits for G.P. sones at room temperature in Aluminium - Zinc alloys.

Considering the competitive growth theory, the initial rate of growth is determined by  $\mathcal{I}_{\mathrm{H}}$ , whilst the final zone radius is a function of the <u>life time of vacancy</u>. The maximum life time is usually obtained at some intermediate value of  $\mathcal{T}_{\mathrm{H}}$ , where the vacancy concentration is just insufficient to nucleate dislocation loops, and thereby reduce the diffusion distance of vacancies (Fig. E and F). The critical value of  $\mathcal{T}_{\mathrm{H}}$  seems to depend on quanching rate and alloy content.

The shape of the sones was determined by the atomic mismatch between the solute and the solvent atoms. Zones were found to be spherical if it was less than 3 pct. as in Al-Zn alloy and are disc like if it is greater than 3 pct. as in Al-Cu alloys.

Plastic deformation after quenching from low  $T_{\rm H}$  causes an increase in the rate of zone growth but the final size is much smaller. The final structure of the alloy is largely determined by the initial stages of ageing.

## 1.3.4 Intermediate Precipitate

Usually the decomposition of supersaturated solid solution to equilibrium phase, takes place through one or more intermediate transition phases e.g.  $\alpha^{\dagger}$  in Al-Zn and  $\theta^{*}$  in Al-Cu system. These phases are obtained by partial loss of coherency of usually larger zones with the matrix during ageing. Priedel<sup>(39)</sup> and Brooks<sup>(40)</sup> discovered condition for an initial coherent precipitate to become partially coherent and finally non-coherent as it grows, and according to them the precipitates become moncoherent when it is so large that the elastic strain energy due to difference in atomic volumes of matrix and precipitate becomes greater than the surface energy ( $\gamma$ ). This criteria then predicts that a precipitate larger in radius, than  $4\gamma / \delta^* G$  will be noncoherent.

Taking  $\gamma = 500 \text{ ergs/cm}^*$  for a non-coherent interface, and  $G = 4 \ge 10^{11} \text{ dynes/cm}^*$ , with  $\delta = 10^{-2}$ , then it predicts that precipitates larger in radius than  $5000^{\circ}$ A will be non-coherent.

#### 1.3.5 Equilibrium Precipitate

This is thermodynamically the most stable state which the structure ultimately attains after any ageing sequence. Thus the final microstructure constitutes the phases, predicted by the phase diagram, vis. the matrix and the second phase. The initial stages of ageing and other ageing variables, determine the marphology (size, shape and distribution) of the final structure. The appearance of intermediate phase is characterized by the peak in hardness while the appearance of equilibrium phase is characterised by a fall in hardness, thus getting the name <u>overageing</u> which basically vorresponds to the loss of coherency. The mode of equilibrium phase, depends on the secondary defects, usually non-coherent precipitate, appears as discontinuous and in inhomogeneous mode of precipitation.

#### 1.3.6. The sequence of precipitation

As pointed out earlier (1.3.1) the ageing process consists of various stages. The precipitation follows initially the path of minimum activational energy rather than the maximum loss of free energy. Thus the G.P. sones are not the only precipitates formed inside the dashed line of metastable phase diagram; they are merely the most rapidly formed precipitate and subsequently they will dissolve in favour of a phase which has a lower free energy but a higher activational energy of formation.

The <u>sequence theory</u> by Guinier<sup>(41)</sup> argued that each stage of precipitation was associated with a definite precipitate type. According to this, the clusters develop into zones by a process of <u>competitive growth</u>. Some of the clusters grow to a critical size and get transformed into intermediate phase, which on further ageing is transformed to the equilibrium precipitate.

Most alloys precipitate at least one intermediate phase on ageing, and the precipitation sequence at any temperature is simply determined by the principle that the alloy decomposes initially along the path of minimum activational energy, and the microstructure finally attains the structure predicted by the phase diagram, though the intermediate phases given by metastable phase diagram. Other ageing variables like  $T_H$  and  $T_A$  etc., may result in predominating or obscuring these successive stages.

#### 1.3.7 Role of quench-in vacancies

Vacancy concentration  ${}^{*}C_{V}{}^{*}$  (fraction of vacant lattice sites n, out of total sites  $n_0$  at  $T_H^0 K$  is given by;  $C_y = n / n_o = \exp_{*}[-E_y / KT_H]$ , which predicts a sharp rise in CV with THE these vacancies are retained in solid solution on drastic quenching. The vacancies introduced by plastic deformation  $(10^{-5}$  to  $10^{-6})$  are much less than that introduced by quenching (10<sup>-4</sup>). In Al-5 pct. Zn, quenched to 20°C from 580°C and 540°C, the Cy values are 6 x 10<sup>-4</sup> and 2 x 10<sup>-4</sup> respectively. In a typical 1. pct. solute alloy,  $C_v$  in as quenched sample is  $10^{-5}$ , with the ratio of solute to vacancy 103, thus most of the vacancies are associated with solute. In this alloy, which form sones, 10<sup>20</sup> solute atoms in 1 Cu.cm. must be transported to 10<sup>18</sup> somes by 10<sup>17</sup> vacancies, thus each vacancy being effectively responsible for transport of at least 10<sup>3</sup> solute atoms in a very short time.

Experically, vacancy decay law (117) is;  $C_{V} = C_{o} \exp_{1} - (\frac{b}{t})^{n}$ , in which  $C_{o}$  and  $C_{V}$  are vacancy concentrations initially and at time t, with a being number of jumps for vacancy annihilation, and b, a constant related to relaxation time. The  $C_V$ value decreases to  $10^{-20}$  from  $10^{-4}$  in about 30 minutes at room temperature in quenched AL. The average number of jumps n, made by a vacancy before annihilation (28) is given by;

 $n = A^{1} Z \in t \exp(-E_{H} / ET_{A})$ , where, e = Debye frequency, and A contains an entropy of activation and is of order 1 to 10, Z is co-ordination number of lattice and t is annealing time. They obtained,  $n = 10^{10}$ , indicating that vacancy sinks are approximately 5 x  $10^{-3}$  cms. apart for a random walk. This indicates that not only surfaces and grain boundaries are the only sinks but dislocations introduced during quenching and formation of secondary defects like voids, vacancy clusters etc. are also effective eink of vacancies.

Pencari and Fedorighi<sup>(33)</sup> and others (96) studied annealing behaviour of pure Al.(Pig. A and B) and found that for  $T_{H^{\pm}}470^{\circ}$ C, the extra resistivity annealed out in one stage at room temperature with an activational energy 0.37 - 0.45 eV., and  $n = 10^{-9}-10^{10}$ ; in good agreement with Bradshaw and Pearcon<sup>(25)</sup>. The process took loss than 30 minutes showing the initial fast reaction followed by the elev reaction. The vacancies tending to cluster on closely packed atomic planes i.e. (111) in Al. For  $T_{\rm H} = 470^{\circ}$ C, the annealing process took place in two stages. Stage I at room temperature, with activational energy 0.58 eV., and  $n = 10^4$ , and stage II in range 140 to 200 °G with activational energy 1.3 eV. These results were interpreted in terms of <u>clustering of vacancies</u> by addition of single vacancies. Stage II is process of celf diffusion as  $E_D = 1.3$  eV. For low  $T_H$ , vacancy supercaturation is insufficient for large clusters to form and annoaling is due to single vacancy and divacancies to sinks. Vacancy clusters chould be unstable and might collapse to form dislocation loops. This is confirmed in pure quenched Al., which shows density of loops  $\simeq 10^{15}/cm^5$ , and  $C_V$  to be one in  $10^4$ .

Turnbull et.al.<sup>(21)</sup> discussed clustering with the initial concentration of vacancies and the average life of vacancies. An equation is derived giving the number of atoms n, in the greatest cluster which could form during the vacancy decay period;

 $n = C_{\rm g} \left[ C_{\rm V} \left( T_{\rm H} \right) \right]^{3/2} \cdot \left( L_{\rm V}/d \right)$ . Here  $L_{\rm V}$  is mean spacing for the vacancy sinks and d is atomic diameter. This equation gives, n = 10, for 1 pct, alloy quenched from 550°C and aged at room tomperature; in agreement with Wobb<sup>(43)</sup>. This also predicts that the cluster size will increase with increasing  $T_{\rm A}$ , due to increase in  $L_{\rm V}$ ; with smaller vacancy supersaturation. Turnbull<sup>(12)</sup> working on Al-Ag, deduced that final cluster size is limited for kinetic rather than thermodynamic reasons. There are limitations imposed on cluster size by, the  $C_{\rm V}$ , life time of vacancies, and the diffusion coefficient which increases steadily during agoing. Excess vacancy machanism, is the most important one in explaining ageing processes. De Sorbe<sup>(11)</sup> studied effect of  $T_{\rm H}$ ,  $T_{\rm A}$ , step quanching, step agoing, reversion etc., and gave explanations on the basis of role of vacancies. Step quanching gives time for vacancies to annihilate to sinks or form secondary dofects and thus resulting in slow clustering. <u>Reversion</u> has even greater effect in reducing the rate of clustering. Turnbull<sup>(12)</sup> has shown that light <u>cold working</u> has little effect on rate of eluctoring on specimen rapidly quenched, but there is considerable offect, in enhancing, in specimen which are step quenched and then cold worked.

Addition of small concentration of colute atoms(1 pct.) increases the vacancy concentration (90) compared with pure Al. Thus the activational energy of formation of a vacancy ' $E_p$ ', is reduced, and colute atoms must increase the equilibrium concentration of vacancies at  $T_{\rm H}$ . This is due to emistence of bindingenergy (32, 44, 45) between vacancy and solute atom ,  $E_{\rm S-V}$ . In pure Al.,  $E_{\rm p} = 0.76$  eV., but in Al-2n alloy it is 0.70 eV., thus giving  $D_{\rm S-V} = 0.06$  eV.(Ref.44).

Addition of large concentration of alloying element ( 1 to 10 pct.) generally decreases the number of vacancies which are absorbed on holics or dislocation loops. This is due to decrease in equilibrium number of vacancies at  $T_{\rm H}$ , but it seems unlikely since

it would reverse the trend shown by dilute alloys. Vacancies are concluded to be the most important factor governing the various appets of the agoing process.

1.3.8 General Conclusions on Precipitation Phonomenon

Conclusions concerning the structural changes during agoing are following:

i) Alloy decomposes during or immediately after quonch, and there is formation of coherent clusters with simple growth.

ii) Clusters grow with diffusion controlled process of colute atoms and vacancies, until they attain a shape to be called sones.
The shape of zone is dependent on misfit (5) between solute and colvent atoms; being spherical if it is low, and disc or plates if it is high.

111) Rate of growth of clusters and zones is determined by the motion of vacancies and solute atoms, which is governed by  $C_V$  and  $T_{\Lambda^*}$ 

iv) As the sense grow with agoing it is not possible to maintain cohorency after attaining a critical size and thus development of transitional phase having partial coherency with the matrix.

v) Ultimatoly the non-coherent equilibrium phase is developed out of the transition phase. vi) Sequence of procipitation can be described by a corice of matastable phases in the phase diagram.

vii) The matrix-precipitate interface energy is probably the single parameter of greatest importance in controlling both the microstructural and mechanical properties of the alloy.

Conclusions concerning ageing variables are following:

i) Concentration of vacaneico in an alloy at  $T_{\rm H}$ , their distribution and annihilation controls the low temperature agoing. The  $T_{\rm H}$  and quanching rate determine the quanch-in vacancy concontration and the life time of vacancy is important in determining the final size of zenos on prolonged agoing. There is an optimum  $\Sigma_{\rm H}$ , giving factors rate of precipitation and the meximum zeno size is attained with largest vacancy concentration causing no nucleation of new vacancy sinks. Bature of secondary defects is also closely related to the role of vacancies.

11) As pro-procipitation stages involve interchange of solute atoms and vacaneics, which are diffusional processes, the  $T_A$ , is the other most important variable. Usually with high  $T_A$ , agoing to factor, but the maximum hardoning is obtained at some critical agoing temperature.

111) The agoing temperature determines the degree of supercaturation and is directly related to equilibrium number of sense with a critical size thermodynamically stable at this temperature. iv) Effects of quenching rate, step-quenching, cold working prior to ageing and reversion etc. could be satisfactorily explained on the basis of excess vacancy annihilation mechanism.

v) Einctics of different stages of ageing, is also dependent on the binding energy terms, such as, that between a vacancy and a vacancy  $E_{V-V}$ , between solute atom and vacancy  $E_{S-V}$ , and between a solute and a soluto  $E_{S-V}$ .

vi) Effects due to other variables like the specimen size, concentration of solute etc., are also explainable in terms of  $C_{\rm W}$  and their annihilation mechanicm.

We conclude that there are many inter-dependent factors which govern the <u>thermodynamics</u> and the <u>minetics</u> of ageing process. The nature of interface between the precipitate and the matrix, is most important in governing the mechanical properties, while the role of excess quench-in vacancies being the primary one in governing the structural and marphological changes(kinetics and sequence) during ageing.

#### 1.4 <u>Techniques Employed for Precipitation Studios</u>

Changes in various physical and mechanical properties during agoing must correspond to some quite as definite changes in the structure and constitution of the alloy, conversely, they

can be employed for following the ageing process.

Hardness measurements and X-ray techniques with optical metallography were long before used for deducing precipitate structure and the process. Later on the electron diffraction and microscopic studies were used to reveal details about the pre-precipitation. These were of great advantage as they give information on the number, size, shape and distribution of precipitates. Whereas X-ray diffraction studies showed the existence of G.P. zones, the <u>electrical resistivity</u> measurements being very sensitive to the atomic rearrangement, helped in detection and studies of clusters. Small-angle X-ray diffraction techniques are best provod for the studies of crystal structure of individual particles. Thus the X-ray and electron diffraction studios with electrical resistivity studies are extensively employed for pre-precipitation stages, while the hardnoss and optical metallographic studies for later stages. Sometimes, other techniques such as yield stress, magnetic properties and calorimetric measurements are also found upeful.

In the present study, which mainly deals with pre-precipitation stage, the electrical resistivity measurements were employed for following the ageing process.

# 1.5 Discussion on Resistemetric Studies for Precipitation 1.5.1 Theory

The electrical resistivity (?) of a metal or single phase alloy (82,86,99) is a consitive function of the total number and distribution of point defects, dislocations and solute atoms. The technique therefore offers a powerful method for studying decomposition during againg, since resistivity changes can be attributed to clustering of solute atoms, which basically involves redistribution of solute atoms and vacancies. Introduction of solute atoms and structural defects in a solvent matrix increases resistivity. The resistivity of quenched pure Al., is greater than a well annoaled sample by an emount,  $\Delta C_0$ , given by:

 $\Delta P_0 = A. \exp[-E_F / K T_H]$ . Where,  $E_F$  is the activational onorgy for formation of defect which is responsible for resistivity increase (this is of vacancy type), and a A is a constant, with K boing Boltzman's constant.

In Al-Cu system, resistivity of a single vacancy and a cluster of one hundred vacancies has been found to be as 1.0 and 0.69 µohms cms/ atomic pet., of vacancies respectively, while the resistivity of an array of dislocation, density II lines/cm<sup>9</sup> is approximately 1.1 x  $10^{-14}$  N-µohm cms. Therefore resistivity changes  $\Delta$ (.of the order of  $10^{-2}$  µohms cms., are expected on annealing a pure metal after quenching or cold working. Herman<sup>(117)</sup> pointed out that' (' of an alloy containing spherical sones is sum of two terms:

Contribution of spherical zones to  $\ell'$  was studied by Pensari and Federighi<sup>(44)</sup>, taking unit volume, let  $\Delta \ell_z(r)$  and  $\Delta \ell_s$  to be the contribution of one zone of radius r, and that of an isolated solute atom to resistivity, the total contribution  $\overline{\Delta \ell}$  of these two at any instant t, is given by;

$$\overline{\Delta \rho}(t) = \Pi_{S} \Delta \rho_{S} + \Pi_{Z} \left[ \Delta \rho_{Z}(r) - \frac{16\pi n}{3} (r/a)^{3} \Delta \rho_{S} \right].$$

Where  $\Pi_Z$  and  $\Pi_S$  being number of zones and solute atoms respectively, with a being the lattice parameter. Praction of solutes inside zones (taken constant) is n. Since,  $\Delta \ell_S$ ,  $\Pi_S$  is a constant, the experimental variation of  $\Delta \ell$  is simply given by;

 $\Delta \rho(t) = \Pi \left[ \Delta \rho_{Z}(r) - \frac{16\pi n}{3} (r/a)^{3} \Delta \rho_{S} \right] \equiv N. \varphi(r). \quad \text{There, function}$  $\varphi(r)$  is contribution of  $\rho'$ , of a zone of radius r, computed for deplotion of matrix. The  $\varphi$  (r), increases with r, till a critical value,  $r_1 = 11^{\circ}$ A is reached and is zero for  $r_2 = 39^{\circ}$ A.

Resistivity increase during proprecipitation in (Al-2n, Al-Ag) has been attributed by  $Hott^{(46)}$  to critical scattering by very small zones and by Geislor<sup>(47)</sup> and Fine<sup>(48)</sup> to coherency strains and interface dislocations, respectively. There is good experimental evidence of Hott's theory, while cause of coherency strains was ruled out after studying Al-Zn and Al-Ag systems.

The rate of change and the total change of resistivity with usual peak behaviour observed during precipitation is extremely sensitive to the ageing variables viz.  $T_{H^{0}}$  quenching rate, and  $T_{A^{-}}$ . The great advantage of electrical resistivity measurements is that continuous measurements are possible starting as soon as 5 seconds after quenching.

### 1.5.2 <u>Resistivity increase and the , Cmanimum</u>.

As introduction of defects like the solute atoms, vacancies, dislocations etc., in the ordered solute matrix causes resistivity increase, it is expected that there is an overall increase in P, in the alloy, just after quenching, keeping in view that the concentration of solute atoms is much higher compared to the concentration of quench-in vacancies. The appearance of a peak in resistivity during pre-precipitation was treated by Labusch<sup>(49)</sup> who gave the expression for resistivity as:

 $Q = 1/N = \mu = m^{O}/N = \pi$ , where, N and m<sup>O</sup> are the concontration and the effective mass of the charge carrier whose charge is e.  $\tau$  is relaxation time for scattering and  $\mu$  is carrier mobility which is equal to  $[e \tau / m^{O}]$ .

Labush showed, that  $\tau$  passes through a minimum as solute clustering increases in size, and corresponding to this minimum in  $\tau$  ,there is a maximum in  $\rho$  at the cluster size of  $8^{\circ}A$ . Thus with increasing clustering there is initial decrease in  $\tau$  producing an increase in  $\rho$ , giving minimum  $\tau$  for a wavelength of approximately  $8^{\circ}A$ , following the equation;  $d\rho = -\rho(d \tau / \tau_{\bullet})$  Hence scattering efficiency was largest for this size of clusters.

It has been proposed (44) that,  $\int maximum.corresponds to$  $a definite state in the alloy, which is independent of <math>T_H$  and  $T_A$ . Thus the <u>position and amplitude</u> of resistivity maximum can be used to make important deductions about the structure of alloy. According to them the maximum change in resistivity ( $\Delta \rho_{\rm MOX}$  in Al-10 wt. pct.ZN. alloy is given by,  $\Delta \rho_{\rm MoX}$  388 [1 - 61.2 exp. (-0.13 / K  $T_A$ )].

There are three theories concorning the nature of resistivity maximum, due to  $Mott^{(46)}$ , Geisler<sup>(47)</sup> and Fine<sup>(48)</sup> respectively.

Lott suggested that maximum occurred, when the cluster cize reached a critical value, equal to the wavelength of conduction electrons, to produce strongest scattering. Labusch predicted this to be 8°A, and Matyes<sup>(50)</sup> as 7°A which was also confirmed by Herman and Cohen<sup>(20)</sup>. This was further confirmed by Pensari and Fedrighi<sup>(44)</sup>, and was found to be  $10^{\circ}$ A for Al-5.3 atomic pct., Zn alloy.

Geisler's suggestion explaining, (maximum, in terms of coherency strains occurring during G.P. zones stage was ruled out after studies on Al-Zn and Al-Ag., system.

Fine suggested the '?'increase due to the presence of a large number of dislocations around G.P. zones. But again resistivity peaks in Al-Zn where the misfit with the matrix is low, led to rejection of this theory.

## 1.5.3 Conclusions from resistivity studies.

The resistivity changes ( $\Delta P$ ) in the quenched alloy during ageing is the sum of following effects:

a). An increase caused by growth of G.P. sones. Since maximum in C'occurs in alloys where sones are spheres(117), plates (11 and 12) and needles (51), it is logical to suppose that Mott's theory is most plausible giving a maximum in C'at a critical sone size, signifying strongest scattering.

- b). A decrease due to fall of solute concentration in the matrix, Factors (a) and (b) are interdependent.
- c). A decrease arising from the annihilation of excess quenching vacancies at sinks.
- d). Possibly a decrease in resistivity from the association of vacancies (di-vacancies and voide) and solute atoms.
- ,e). Rate of resistivity increase is controlled by the rate of clustering and zone growth, and shows initial fast and slow reactions. But as zone growth is diffusion controlled, the rate of  $\ell$  increase also becomes dependent on  $T_A$ .
- f). The height of resistivity peak is proportional to the number of gones.
- g). Resistivity peak corresponds to a definite structure during ageing and is independent of  $T_H$  and  $T_A$ , hence its position and amplitude can be used to deduce important deductions like the evaluation of  $E_p$ ,  $E_H$ , and the kinitics of structural changes during pre-precipitation.

We conclude that, (maximum, is not yet completely understood, but there is evidence to show that it corresponds to a definite state of structure of the alloy.

#### 1.6 Aluminium-Zinc System

Extensive studies on ageing in this system have been done by many workers (69, 80, 87, 91, 93, 105, 116). We have evidence of following sequence of precipitation in A1-12 and 15 atomic pot., Zn samples aged at 175°C and 225°C by Simerska and Syneok<sup>(52)</sup>.

Spherical G.P. zones  $\rightarrow$  ellipsoidal G.P. zones  $\rightarrow$  rhombohedral transition phase (R-phase)  $\rightarrow$  a' cubic transition phase  $\rightarrow$  Zn rich stable precipitate.

1)6 Spherical G.P. sones:- The early stages of precipitation consists of formation of Zn rich, coherent spherical G.P. sones, because of misfit between Al and Zn atoms being low (-1.9 pct.). The average size of zones progressively increases with ageing, while solute concentration, within the zones remains constant, and therefore the molar volume of the zones, remains constant during the growth of the zones. This leads to progressive increase of coherency strains. In an 9.4 pct. Zn alloy (53), with  $T_{\rm H}=470^{\circ}$ C, after ageing at room temperature for 8 days the zone diameter is approximately  $50^{\circ}$ A and the density of precipitates 5 x  $10^{16}$  zones/cm<sup>3</sup>.

11) <u>Ellipsoidal G.P.sones</u>:- Progressive increase in the coherency strains during the growth of spherical zones, results in the change of shape towards ellipsoidal ones for the larger zones. Simerska, found that the change of shape from spherical to ellipsoidal, occurred if the zones exceeded the diameter of about  $30^{\circ}$ A. It has been reported (23) that in Al-2n alloys surface energy between coherent boundary between zones and matrix is about  $300 \text{ ergs/cm}^{\circ}$  and this value is exceeded by the elastic strain energy terms when zone reaches a diameter of approximately  $30^{\circ}$ A. It is suggested that the anisotropy of strain field around the spherical G.P.zones, changes their shape to ellipsoids, with their short axis parallel to <111> direction of matrix, and the distance between the closed packed planes (111) within the zones becomes smaller. This mechanism finally results in rhombohedral distortion of matrix, as will be discussed in the next stage.

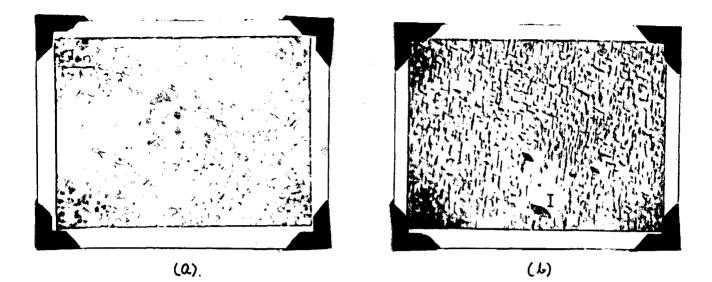
111) Rhombohedral Transition Phase(R-phase):- The development of R-phase on ageing the quenched Al-12.2<sup>st</sup> pot. Zn. samples at  $150^{\circ}$ C, was studied by Simerska and Syneck<sup>(52)</sup>, and was also confirmed by Garf and Lenormand<sup>(54)</sup>, while ageing this alloy at 275°C. They suggested the <u>mechanism</u> for the R-phase formation. Anisotropy of coherency strains around spherical G.P.zones, causing them to become ellipsoidal ones, with further ageing, leads to the contraction of spacing along one of the <111> direction inside the zones. This internal rhombohedral lattice deformation, combined with the loss of coherency of zones with  $\alpha$ -matrix except in (111) habit plane leads to the development of transitional rhombohedral phase. This mechanism was also confirmed (55), on ageing 12 pct. Zn alloy they observed that the lattice parameter of R-phase were,  $a = 3.992^{\circ}A$  with  $a_{\rm R} = 91^{\circ}36^{\circ}$ . He estimated it's 2n content to be about 63 atomic pet. from unit cell volume. The larger interatomic distance between the neighbouring atoms lying in the same atomic plane (111) was found to be 2.862°A, whereas a much shorter average distance of 2.783°A exists between closest neighbours in two succeseive (111) atomic planes. The first of these distance is almost equal to the average inter-atomic distance in  $\alpha$ -matrix, equal to 2.858°A, suggesting that at least a partial coherency of R-phase with the matrix in (111) plane maintains for a long time.

iv) <u>a-transition Phase</u>:- This phase is developed by breaking down of the partial coherency of the R-phase in the (111) habit plane also. Garwood<sup>(56)</sup> was first to detect this f.c.c., transition phase by X-ray techniques in A1-25 pct. In alloy aged at 200°C. The smaller unit cell of this transition phase compared to that of parent f.c.c., (a) phase indicated that it corresponds to a' conjugate solid solution. This In rich a' phase was also observed by Garf<sup>(57)</sup> and Geisler, Barett and Mehl<sup>(58)</sup> in A1-30 pct., alloys aged between 175-225°C. Literature shows that a' phase is f.c.c., with lattice parameter  $a_0 = 3.985^{\circ}A$  (that of A1.is 4.04°A). For A1-9.4 pct., In with  $T_{\rm H} = 540^{\circ}C$ , and ageing for 8 days at room temperature, followed by ageing at 100°C for 1 day, small precipitates of a', parallel to {111} planes were detected. These a' precipitates were about 25°A thick, and about 300°A in diameter, with density of precipitates being  $10^{16}/cm^3$ . The same specimen, if aged at  $100^{\circ}C$  for 5 days, shows weak streaks, which are slightly displaced from f.c.c. matrix spots in electron diffraction studies. The displacement was only 2 pot. in [111] direction perpendicular to the plane of precipitate phase and about 0.5 pct. in the direction in the plane. This was detected by Garwood<sup>(56)</sup> to be elastically distorted f.c.c. with  $a_0 = 3.98^{\circ}A$ , and was doubtless the intermediate a' phase showing granular appearance in optical micrographs. This a' is probably quite stable upto a thickness of  $100^{\circ}A$ , before leading to it's transformation to the equilibrium c.p.h. Zn precipitate. Although G.P.sones can be obtained even below room

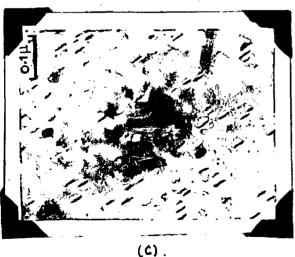
v) Equilibrium Precipitate(o.p.h.Zn):- The stable Zn rich precipitates are formed by continuous precipitation from a'-phase with their basal planes parallel to {111} planes of the matrix. [{111}  $_{\alpha}$ . // (0001) $_{Zn}$ ]. Equilibrium precipitate is also found to be observed by discontinuous mode of precipitation at the defects and at grain boundaries in polycrystalline samples, predominantly at quite low temperatures.

The sequence and mechanism of precipitation in Al-2n alloys can be concluded to be as following:

1) The early stage is the segregation of Zn atoms into



Micrographs of Al-9.4 pct. Zn alloys water quenched from 540°C and aged at room temperature (a) for one day and (b) for five days. (QJ.44)





(d)

Micrographs of A1-25 pct. Zn alloys directly quenched to 200°C and aged (c) for one hour and (d) for four hours. (Ref. 111) .

ophorical G.P. cones, cohoront with a-matrix.

- 11) Anisotropy of cohorency strains around G.P. sones which increases with growing of zones, leads to the contraction of spacing along one of the <111> direction inside the zones, and thus rhombohodral deformation of lattice, resulting in the change of G.P. zones from spherical to ollipsoidal ones.
- iii) The partial loss of cohorency of the senses with the a-matrix except in {111} habit plane, leads to the dovelopment of transitional rhombohedral phase.
  - iv) The breaking down of the partial coherency of the R-phase in the (111) habit plane as well, leads to the development of known non-coherent transitional a'-phase.
  - v) The Zn stable precipitates, are formed by continuous precipitation from a'-phase, with spitaxial growth of hexagonal Zn precipitates, with their basel planes parallel to {111} planes of a-matrix. Also the discontinuous precipitation of Zn on grain boundaries and defects, was found in polycrystalline samples.

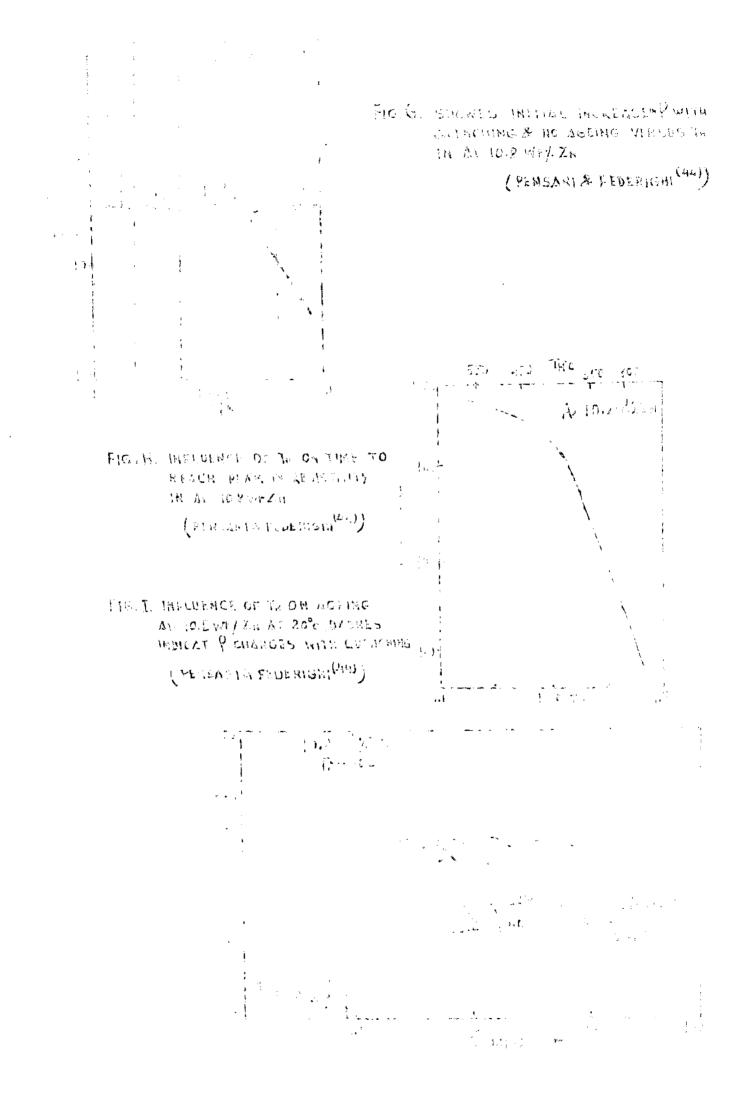
Horo it is to be noted that the process is successive and that the intermediate phaces are not nucleated independently, as has been observed in Al-Cu system.

34.

#### 1.7 Obcorvations on resistemetric studies on Al-In precipitation

Pansari and Fedorighi<sup>(44)</sup> made thorough invostigation of clustoring in Al-4.5, pot. Zn alloy. They found that the initial (measured just after the quench was greater than the value expected for the super-saturated solid solution. Since the magnitude $of discrepancy increases with increasing <math>T_{\rm H}$ , and decreased with increased quenching rate, the effect was attributed to clustering of Zn atoms during quenching. Other observations are:

- 1) On ageing, the  $\ell$  first increased and then decreased to some constant value, which was dependent on  $T_{H^*}$
- 11) Time to reach maximum in  $\rho$  was less for higher  $T_H$  and  $T_A$ , as is expected because the clustering is caused by vacancy diffucion.
- 111) The height of (maximum, was almost independent of  $T_H$  but strongly dependent on  $T_A$ , being high for low  $T_A$ .
- iv)  $\binom{1}{\text{max.}}$  occurred when mean zone diameter has attained a critical value (10°A), which is independent of  $T_{\text{H}}$  or  $T_{\Lambda}$ .
  - V) The magnitude of resistivity maximum is then proportional to the number of zones, which is dependent on  $T_{A^*}$
- vi) Initial number of somes depends on  $T_A$  and that the number remains constant with time or decreases with the same time law at different temperatures. But since there is no nucleation barrier for clustering this theory is not







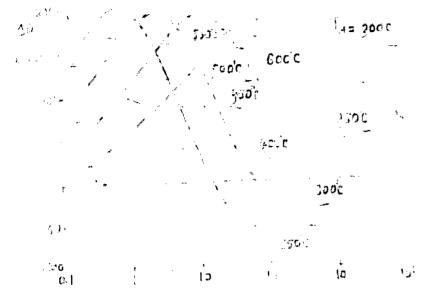
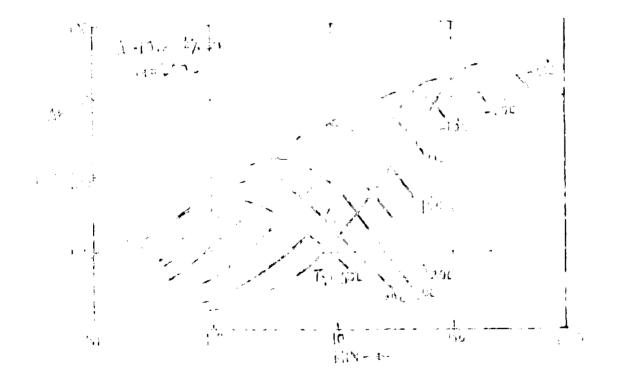
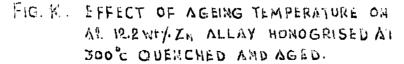


FIG J. RESISTIVITY CHANGES FOR INDICATED THYALVES IN AN - 4 SATIZN ALLAY AGED AT 44°C DASHES INDICATE P CHANGES WITH QUENCHING

(PENSARI & FEDERIGHI





plausible. The number of nuclii would then depend on the degree of supersaturation.

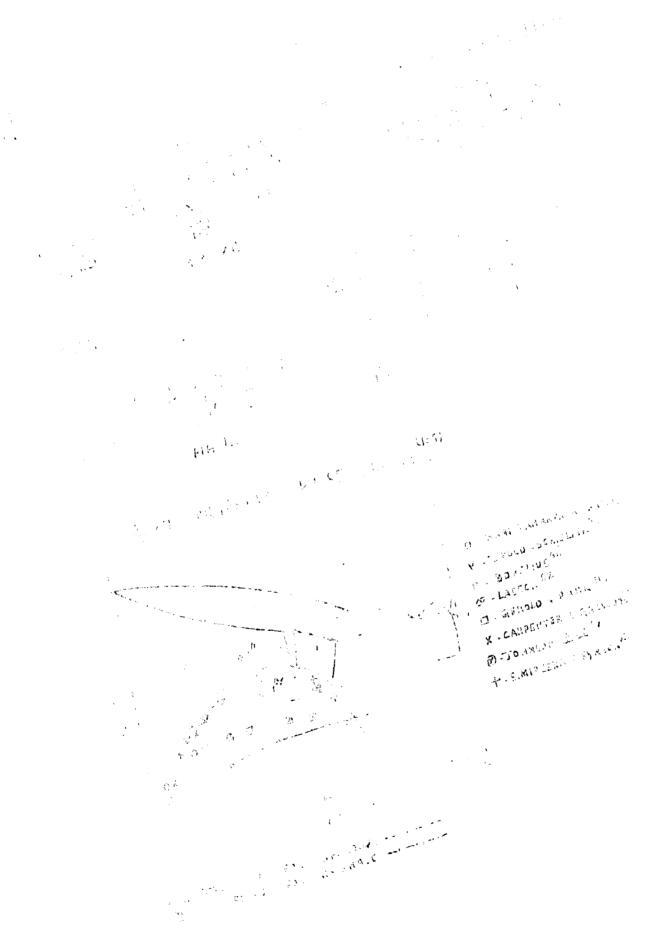
- vii) Pensari and Federighi's results can be explained on the basis of Gerold's (58-60) theory, suggesting the dependence of the degree of segregation simply on the position of metastable miscibility gap, e.g. in Al-4.5/Zn, the phase boundary occurs at 180°C, whereas the solubility limit for G.P. zones is at 94°C, thus giving the equilibrium number of zones (Nz), zero at this temperature. The supersaturation of Zn, and hence the number of Zn atoms available for segregation will be small at temperatures in the range of 50 to 80°C, but will become larger at lower temperatures. Assuming the model of competitive growth for clustering and supposing that the  ${}^{0}$ , is a function of the degree of segregation of the alloy, and hence the number of zones, we would expect the magnitude of resistivity maximum to be small at high  $T_A$  and large at low  $T_A$ . By simply changing the degree of segregation one is simply causing the clusters to grow or dissolve, and this explains the reversion treatments.
- viii) On prolonged isothermal ageing, the final value was only dependent on the kinetics of clustering, and the largest zones were obtained at  $T_{\rm H} = 350^{\circ}$ C. This was explained in terms of concentration of vacancies and their life time.

36.

# 1.8Reversion and Metastable Phase Diagram(Al-Zn).1.8.1Metastable phase diagram

The whole sequence of precipitation during ageing can be described by a metastable phase diagram (Fig. M) showing the solvus line for the coherent (G.P.zones) and the non-coherent precipitates (transitional phases). The solvus line for zones and intermediate precipitate is always displaced towards lower temperature and higher solute concentration compared with equilibrium precipitate, because of the respective activational energy terms for these processes.

Netastable phase diagram can be used to explain formation of G.F. sones, intermediate precipitate and the phenomenon of reversion. Gerold<sup>(15,36)</sup> obtained <u>solvus for G.F.sones</u> by merely <u>extrapolating the high temperature a-a' miscibility gap</u> to lower region, with the limits being 1.8 and 69.0 atomic pct. Zn at room temperature. This view was confirmed by many others (39, 59-63). Dash and Fine<sup>(64)</sup> found no G.F.zones in 1.65 at.pct. Zn alloy aged at room temperature. Pensari and Federighi<sup>(44)</sup> and Garf<sup>(38)</sup>; found, that sones were not stable in Al-4.5 at.pct.Zn alloy above  $100^{\circ}C$  (Gerold gave it to be  $110^{\circ}C$ ). Borelius<sup>(66)</sup>, Johnson<sup>(67)</sup> and Strongin<sup>(68)</sup> indicate the solvus for R-phase some  $25^{\circ}C$  below the monotectoid temperature ( $275^{\circ}C$ ) as shown in fig. M. They indicate that whereas G.P. sones can be formed at even below room temperature, a temperature of  $100^{\circ}C$  or more is required to obtain g-phase.



A study by Wahi and Anantharaman<sup>(69)</sup>, on Al. 10-30 pct. Zn alloys also reveal the presence of R-phase on ageing in the range 200-310°C, and the appearance of R-phase was preceeded by a considerable fall in hardness. They also found that discontinuous precipitation at Grain Boundaries was major mode of decomposition in 10 and 20 at.pct. Zn alloys at and below  $170^{\circ}$ C. Sequential transformation of zones to equilibrium Zn, through intermediate precipitate was found to be negligible or absent in this temperature range. Calculation of Zn concentrations by Ellwood's<sup>(71)</sup> data for lattice parameter of Al-Zn alloy, in the matrix and transition phase, give rise to a <u>second metastable miscibility gap</u>, extended well into high temperature  $\alpha$ - $\alpha$ ' miscibility gap, hence suggesting two miscibility gaps in Al-Zn system.

<u>Pirst</u> extending from lower temperatures to about  $220^{\circ}$ C representing metastable equilibrium between zones and the matrix, and is confirmed by other workers (36, 44, 67).

Second extending into high temperature  $\alpha-\alpha^*$  miscibility gap, and represents metastable equilibrium between R-phase und the matrix. This was in partial agreement with conclusion of many other workers (61, 62, 72, 73).

The two loops overlap around 170-220°C, further complicating the precipitation in this range.

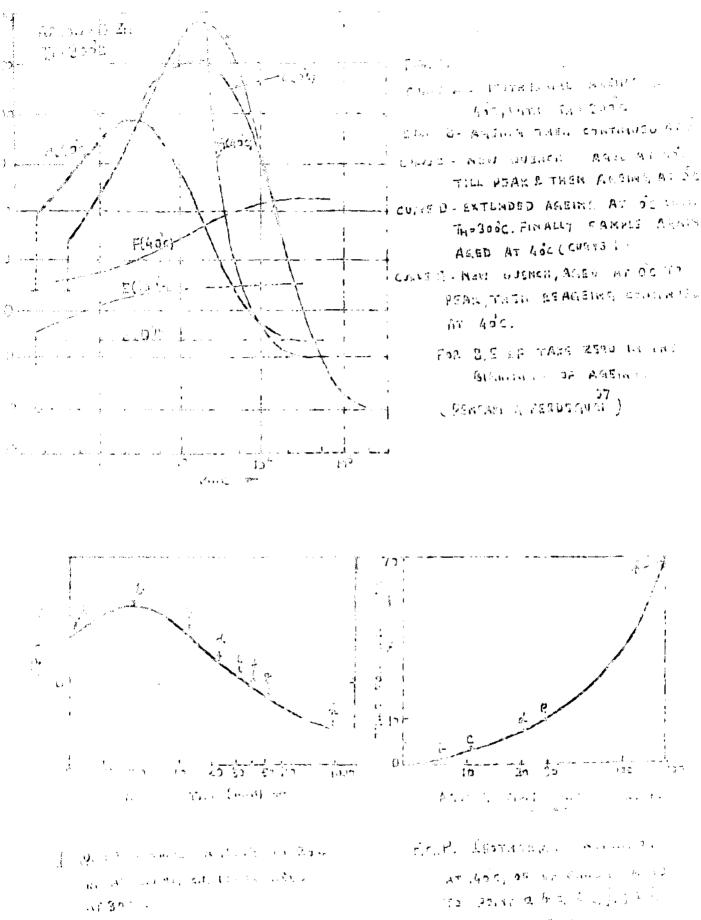
Metastable phase diagram gives clear explanation of

reversion, and determines the stability of transient precipitates.

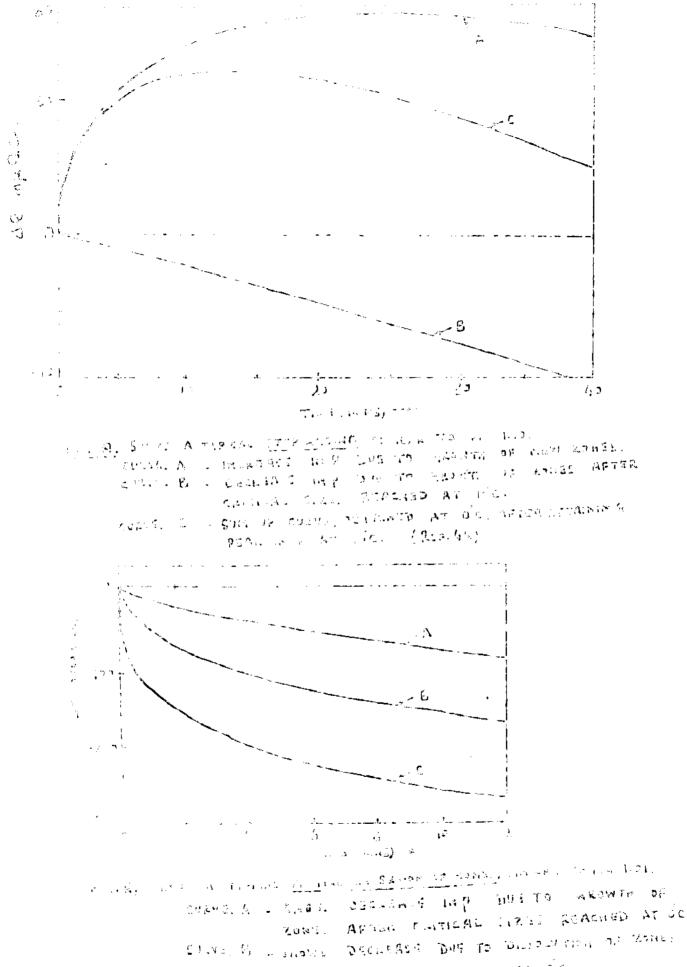
#### 1.8.2 <u>Reversion</u> (or Retrogression)

The phenomenon discovered by Gaylor<sup>(74)</sup> consists of the dissolution of the earlier formed, low temperature ageing products like the G.P. zones or intermediate phase, when the alloy is subjected to a reageing (reversion) treatment at higher temperature given by the solubility limits in the metastable phase diagram. Thus it may be further diversified as reversion of G.P. zones or reversion of intermediate phase, and a partial reversion treatment of any of these. Dehlinger and Knapp<sup>(75)</sup> were first to suggest that reversion was governed by the metastable phase diagram, and this has been confirmed by Silcock et.al<sup>(76)</sup>. Thus we can consider the reversion of G.P. zones and intermediate precipitate at the metastable phase boundary as completely analogous to the dissolution of the equilibrium phase at the solvus line. Reversion involve up-hill diffusion of solute from the zones to the matrix.

On re-heating alloys, which are quenched and aged to produce G.P.sones, some of the zones get dissolved in accordance with the G.P.sones solvus line, and the remaining zones grow further and are able to nucleate the transition phase in accordance with the metastable equilibrium line. In Gerold's view, the degree of segregation of an alloy containing clusters or zones is simply



In priving Parley



- AT LOC PRESSIONAL - POLICE AT SC

The Construction of Arms Ar 150 Artsrather and

dotormined by the position of metastable colubility line. The number of sence is dependent on  $T_A$ . While simply changing the  $T_A$ , above or below the solvus line for G.P.sones, one is simply causing the sence to dissolve or grow. There is a critical size of  $\text{sones}^{(77)}$ which is dependent on  $T_A$ ; the sence smaller than this dissolve, while others grow at this reversion treatment. This study could be termed as <u>partial reversion of sence</u>. The <u>minimum reversion</u> <u>temperature</u>, is that given by the metastable colvus for that phase.

In Garf's<sup>(78)</sup> view <u>true reversion</u> is that in which the low temperature precipitate should not act as nucleus for more stable phase. This would allow simultaneous precipitation of a new phase and dissolution of the intermediate phase, and hence the true reverted state. The process of reversion involves sufficient time to allow the loss of a high concentration of vacancies, and the process of reversion is very slow.

Revorsion techniques were originally employed by Beton ot.al. to determine G.P.sone solvus for Al-Cu, and has been, since then, applied by many for the determination of metastable phase boundaries. Reversion studies also help in better understanding of the mechanicm of various stages of precipitation, involving the formation and discolution of the phase under consideration.

-- 1---

107446

40.

## EXPERIMENTAL

## PROCEDURE

the wire specimen were quenched to water at room temperature. Many such specimen were kept in stock to be used for the reageing studies after necessary heat treatments. Heat treatments were carried out in a vertical homogenization furnace with temperature controlled to  $\pm 5^{\circ}$ C. The homogenized specimen could be dropped out of this furnace through the removable bottom, into the quenching medium, kept directly below the furnace. From here the specimen was quickly transferred for resistivity measurements.

#### 2.2 <u>Resistivity Measurements</u>

2.2.1 <u>Description of Apparatus</u>:- Potentiometric system was employed for following the resistivity changes in the wire specimen, during various ageing and reageing processes. The wire specimen and a standard resistance were connected in series, across an accumulator of 2.1 volts, with a variable resistor in the circuit. The specimen was kept in the ageing beaker, whose temperature was closely controlled through out the experiment by placing it in a bigger bath, fitted with a thermostat. A constant current of nearly 45mA. was made to flow through the circuit, and the potential drops across the wire specimen  $(E_{\underline{X}})$  and that across the standard resistance  $(R_{\underline{S}})$  was measured alternately by balancing this potential drop over a Vernier Portable Potentiometer( Model V-1, Toshniwal). The range of 18 mV. for the full scale deflection and with a least count of  $\cdot$ 01 mV. was selected for this purpose.

#### 2.1 Preparation of Alloy, and the Specimen

2.1.1 <u>Melting and Casting</u>:- The master alloy was prepared from high purity aluminium (99.99 /pure) and high purity zino (99.9 /pure). Weighed quantities of Al. was melted in a graphite cruscible, placed in an electrical resistance,muffled furnace. Required amount of zinc (to account volatilization losses,takem
10 pot. excess) wrapped in Al-foil was added to liquid aluminium and stirred. The molten alloy was homogenized for lo mins, before / into cylindrical cast iron molds of dia. 7/8" and length 6".
2.1.2 <u>Hot forging and wire drawing</u>:- The ingots were hot forged and hot drawn to long wires of 1 mm. diameter, thus removing the microinhomogenity and breaking down of the cast structure. Thereafter a stress relief treatment was given. Fine turnings of the alloy were used for composition estimations.

2.1.3 <u>Specimen preparation for resistivity measurements</u>:- To obtain uniform and smooth wires, they were cold drawn through a fine orifice of a steel die of diameter 0.70 mm. A 55 cms length of this uniform wire was cut out, and marks were made at 50 cms. length, leaving 2.5 cms. on each end, for clamping. This wire was wound in the form of a coil to be used as specimen for resistivity measurements.

2.1.4 <u>Heat treatment</u>:- All wire specimen were given an initial homogenization treatment at  $400^{\circ}$ C, for 15 hrs. After this treatment

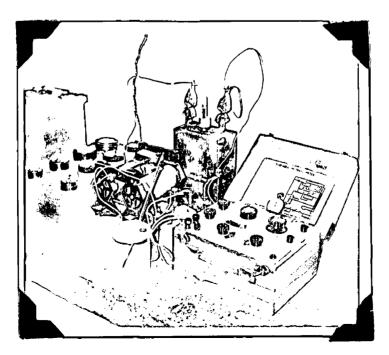
the wire specimen were quenched to water at room temperature. Many such specimen were kept in stock to be used for the reageing studies after necessary heat treatments. Heat treatments were carried out in a vertical homogenization furnace with temperature controlled to  $\pm 5^{\circ}$ C. The homogenized specimen could be dropped out of this furnace through the removable bottom, into the quenching medium, kept directly below the furnace. From here the specimen was quickly transferred for resistivity measurements.

#### 2.2 <u>Resistivity Measurements</u>

2.2.1 Description of Apparatus:- Potentiometric system was employed for following the resistivity changes in the wire specimen, during various ageing and reageing processes. The wire specimen and a standard resistance were connected in series, across an accumulator of 2.1 volts, with a variable resistor in the circuit. The specimen was kept in the ageing beaker, whose temperature was closely controlled through out the experiment by placing it in a bigger bath, fitted with a thermostat. A constant current of nearly 45mA. was made to flow through the circuit, and the potential drops across the wire specimen  $(E_{\rm T})$  and that across the standard resistance  $(R_S)$  was measured alternately by balancing this potential drop over a Vernier Portable Potentiometer( Model V-1, Toshniwal). The range of 18 mV. for the full scale deflection and with a least count of .01 mV. was selected for this purpose.

42.





The lapse of time during agoing and reversion process was measured by a sensitive stop watch.

2.2.2 <u>Calculation of Resistance</u>  $(R_X)$  <u>and Resistivity</u> $({}_{X})$ :- If the potential drop at any instant during ageing, across the wire specimen be  $E_X$  and that across the standard resistance be  $E_S$ , when the same constant current flows through both of them, then the resistance of the wire specimen  $(R_Y)$  can be obtained by :

$$R_{\chi} = E_{\chi} \cdot R_{S} / E_{S}$$
 ohme.

And the resistivity  $\binom{l}{X}$  of the wire of length L cms. and radius r cms. can be obtained by another equation given below :

 $Q_{X} = R_{X} \cdot \pi \cdot r^{2}/L$  ohms. cms.

In our case L = 50 cms. and  $r = \cdot 035$  cms.

## 2.2.3 Mathematical Extrapolation of initial quench resistivity:-

On the findings of Borelius<sup>(79)</sup>, and Do Sorbo at.al.<sup>(11)</sup>, the rate of change of resistance R; (dR / dt) during the clustering and zone formation follows the following relation :

dR / dt = 1 / (a + bt) ... 1. This is analogous to an equation of a straight line in dt/dR and t, with intercept 'a' and slope 'b'. Parameter 1/a is dR/dt at t=0., and is assumed to be proportional to the initial rate of zone growth. Values of 1/a were deduced from the value of R and t measured during isothermal agoing as follows;  $(\Delta t / \Delta R)$ , the quotient of time increase and resistance increment between successive measurements, was set equal to  $(dR/dt)^{-1}$ . This gave data for a plot of  $(dR/dt)^{-1}$  against t, and the values of the intercept 'a' and slope 'b' of the time that best fitted the experimental points was computed, using the method of least squares.  $R_0^-$  the value of the resistance of the specimen immediately after quench was calculated by substituting the measured values of R and t in turn in the integrated form of equation (1).

$$R = R_{a} = 1/b \log (1 + bt/a).$$
 .. 2.

This gave a number of values of  $R_0$  equal to the number of measured values of R and the mean of all these values was quoted as  $R_0$ . With  $R_0$  known, the corresponding resistivity (0,was calculated. A similar method was used by Perry<sup>(81)</sup>.

# 2.2.4 <u>Calculation of $\Delta P$ and the Thormal Coefficient of Electrical</u>

<u>Resistivity ( $\ddot{a}$ ) of the alloy</u> :- The change in resistivity ( $\Delta P$ ) was obtained by the equation,  $\Delta P = P_t(T) - P_0(T)$ , where,  $P_0(T)$ and  $P_t(T)$  are resistivities just after quench and after ageing for time t at  $T^0C$ , respectively.

To account for slight changes in resistivity, due to temperature variations, the correction was employed according to the equation,  $f_t(T) = f_t(0) [1 + \tilde{a}T]$ . The value of  $\ddot{a}$  was calculated from resistivity data at different temperatures of annealed scaples. 2.3 <u>Experimental scheme for the present study</u>:- To obtain oufficient time for pre-precipitation processes in the selected Al-4.17 at. pct. In alloy, three low  $T_H$  values were selected. From each  $T_H$  value, after quenching to  $T_Q^{OC}$ , the ageing was carried out at various temperatures in the vicinity of room temperature as proposed in table I.

After agoing the specimen at  $T_A$  for different values of times called the pre-ageing time  $(t_{PA})$  which are given by times required for a certain percentage of the maximum in resistivity change  $(\% \Delta \rho_{Max.})$  to occur and then the specimen was immediately reaged for 30 mins. At a lower or higher temperature than  $T_A$ , thus constituting the process of step agoing or partial reversion of somes respectively. Table II lists, the required pre-ageing times for different cases, employed in the present study.

Table III summarizes the different ageing and reageing operations employed in the present investigation. The figures in this table indicate the figure numbers given to denote differont operations of reageing. These <u>figure numbers were in a special</u> <u>code</u> as interpreted below :

<u>Momenclature of Pigures</u>:- A special <u>three digit code</u> number was assigned to each figure obtained on reageing operation, according to scheme of Table III giving scheme for agoing and reageing temperature. If the figure number is o.g. <u>Fig. X. YZ</u>, then it is inter-

45.

protod as follows:

1) First digit'X'ropresents the  $T_{\rm H}$  value which is equal to  $300^{\circ}$ C,  $350^{\circ}$ C or  $400^{\circ}$ C for X = 1, 2 or 3 respectively.

11) Second digit i.e. 'Y' with a zero placed after it, gives the temperature of ageing  $(T_A)$  used prior to any reageing. Thus Y = 0,1,2 or 3 means agoing at 0,10,20 or 30°C, respectively.

111) Third digit i.e. '2' with a sore placed after it, gives the reageing temperature used after ageing at  $T_A$ . Thus 2 = 0,1,2 or 3 mean reaging at 0,10,20 or 30 degrees contegrade, respectively.

It is ovident that the various reageing processes in the procent study can be broadly classified into two categories.

- (a) <u>Step Againg</u>:- Then the againg temperature was higher than the reagoing temperature and in such a case reageing temperature was denoted as  $T_{RA}$ . So we have the required condition as  $T_{RA} < T_A$ , signifying Y > Z.
- (b) <u>Partial Roversion of Zonos</u>:- When agoing temperature was lower than the reageing temperature, with reageing temperature donoted as  $\underline{T_R}$  in this case. So we have the required condition as  $\underline{T_R} > \underline{T_A}$ , signifying X < Z.

For a certain rougingp process, the different values of proagoing times as given in Table II was recorded on respective ourves in the figures as 20, 40, 60, 80, and 100 pet. values, respectively.

#### 3.0 Rooulta

The results of primary agoing as proposed in Table I, have been plotted in terms of resistivity changes ( $\Delta Q$ ) versus the ageing time, in figures 1.0, 2.0 and 3.0, corresponding to  $T_{\rm H} = 300$ , 350 and 400 degree contegrade, respectively.

Table II gives the various pro-agoing times and the corresponding resistivity values,  $(t(T_A), \text{for various combinations})$  of  $T_H$  and  $T_A$ . These values were used for different agoing and reagoing operations in the present study.

The results of various reageing processes as proposed in Table III have been plotted in terms of changes in resistivity  $(\Delta \rho)$ versus the reageing time on different figures, which has been clacoified into two categories (step agoing and partial reversion of zones) and are listed systematically in Table III giving the corresponding figure numbers, for these respective operations.

- 1-

THOC value	9 •	Veluo	of	Agoing	tomporaturo	(T <sup>o</sup> C)	Corresponding
300	í	0		10	20	30	1.0
350	ł	0		10	20	30	2.0
400	4	0		10	20	30	3.0

PABLE - III

.

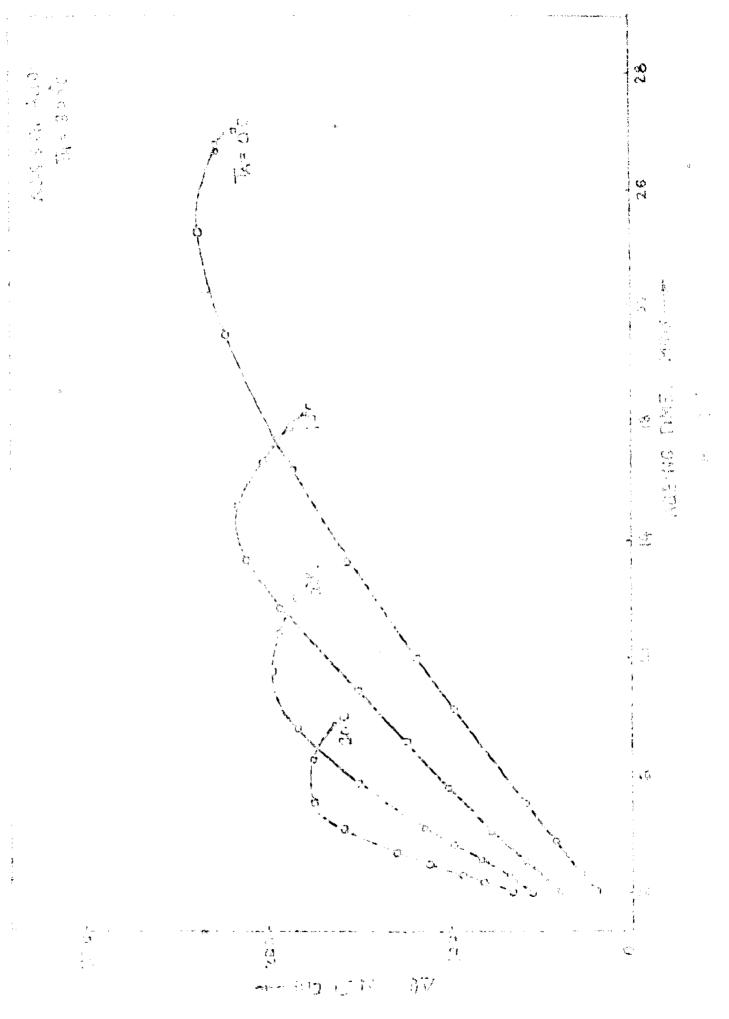
T <sub>H</sub> <sup>O</sup> C	Γ <sub>Δ</sub> ος	Rocaoing tomperature (TR or TRA)						
£1.	25 10 10 10 10 10 10 10 10 10 10 10 10 10	0 <sup>0</sup> 0 .	10 <sup>0</sup> 0 .	20 <sup>0</sup> C	, 30°C			
	0	-	1.01	1.02	1.05			
	10	1.10	•	1.12	1.15			
300	20	1.20	1.21	-	1.23			
	30	1.30	1.51	1.32	-			
	0	-	2.01	2.02	2.03			
	10	2.10		2.12	2.13			
350	20	2.20	2.21		2.23			
	30 2.30 2.31	2.31	2.32					
	0	**	3.01	3.02	3.03			
	10	3.10	<b>4</b> 3	3.12	3.13			
400	20	3.20	3.21	-	3.23			
	30	3.30	3.31	3.52	-			

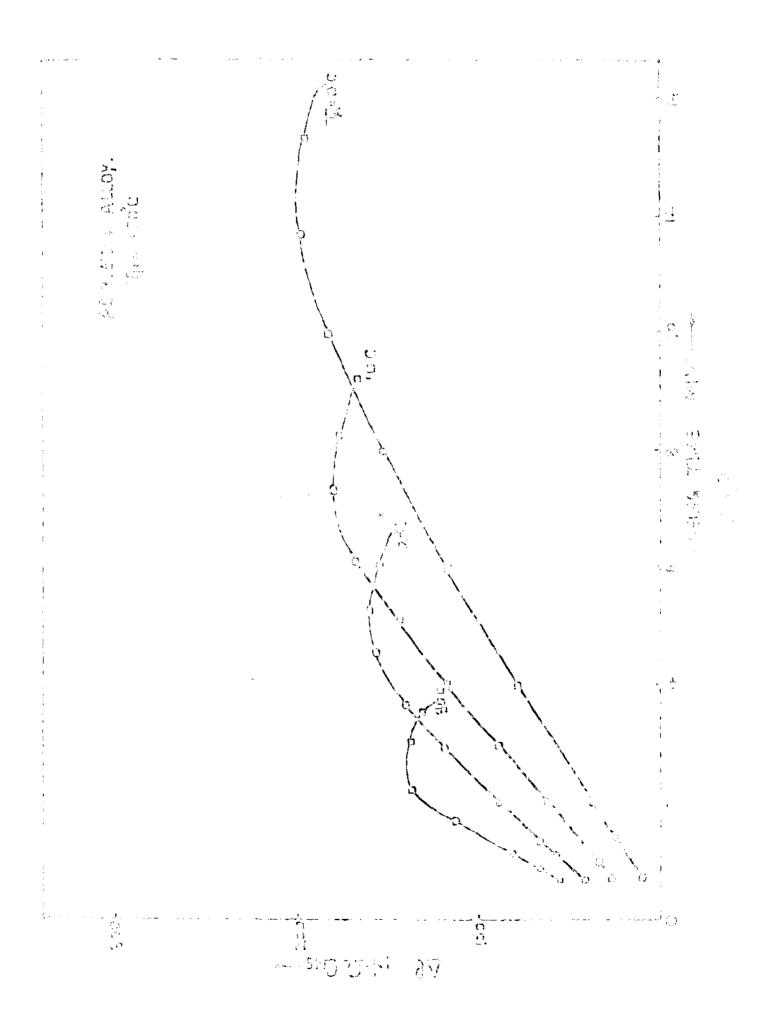
-

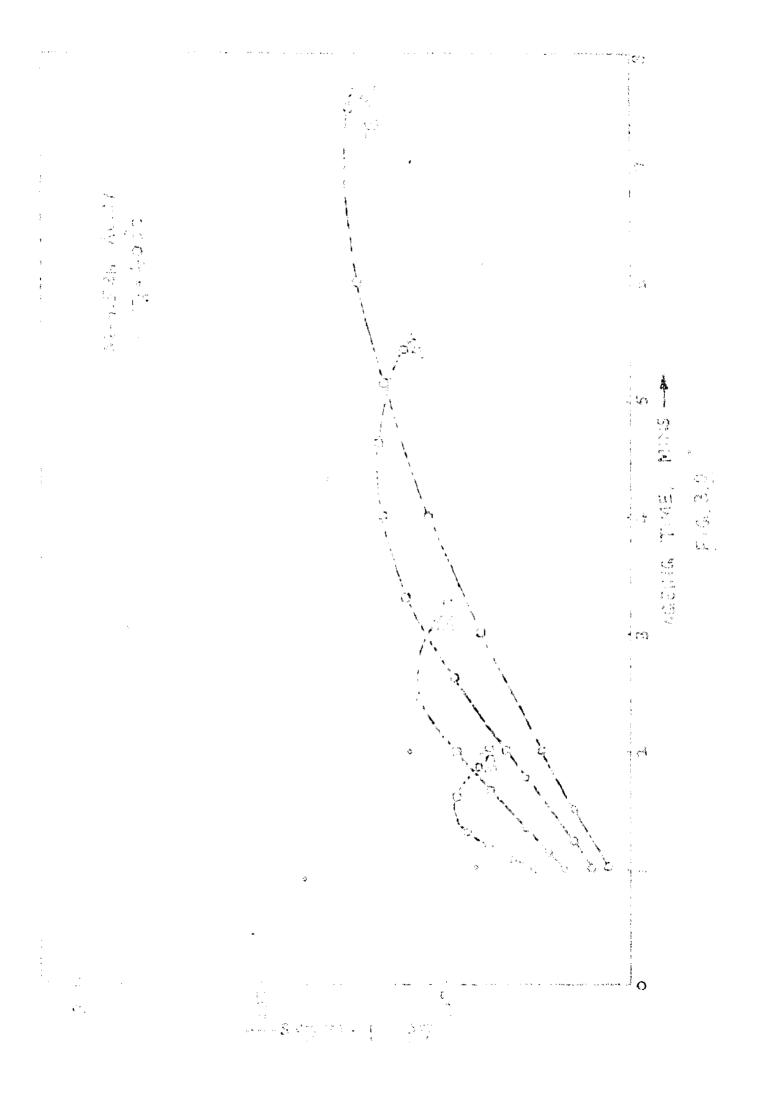
TABLE - II

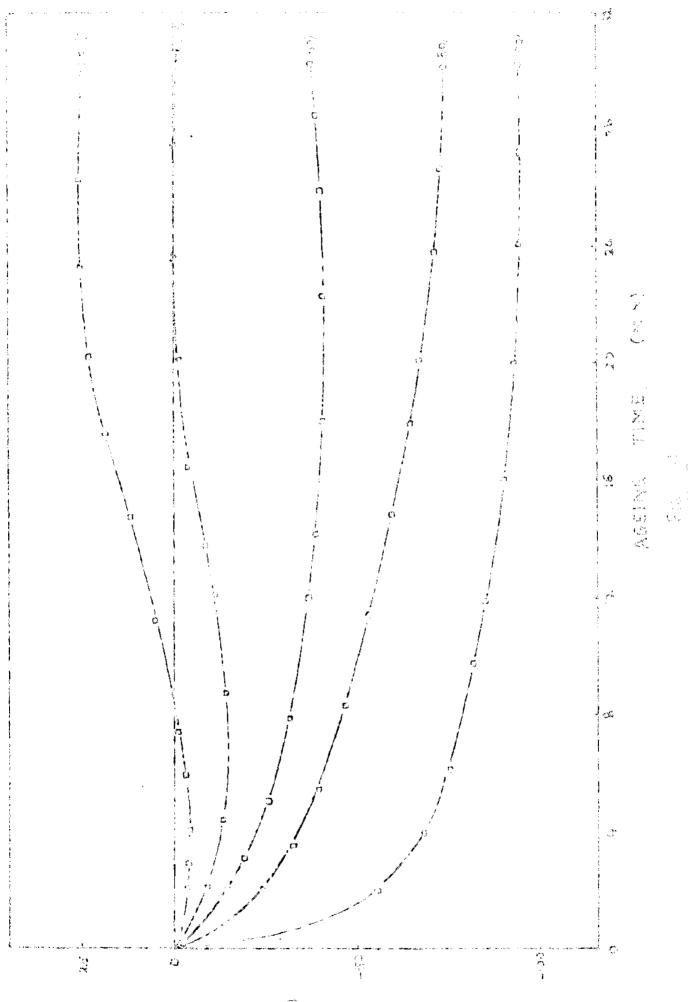
·

TAC	· % APLIC	ax: 300	0 <sup>0</sup> 0		т <sub>н</sub> °с 350°с :		400 <sup>0</sup> C	
	• .	. <sup>t</sup> A	$(t(T_A))$	t <sub>A</sub> .	$(t(T_A))$	. t <sub>A</sub>	$(t(T_{\Lambda}))$	
0	0	000	4.932	0 -00	4.972	0 -00	4.932	
	20	4 -10	4.979	<b>80- S</b>	5.012	4 -10	4.979	
	40	8 -00	5.026	4 -00	5.052	8 -00	5.026	
	60	12 -00	5.073	6 -12	5.092	12 -00	5.073	
	80	15 -30	5.120	8 -30	5.132	15 -30	5.120	
	100	25 -00	5.167	12 -00	5.172	25 -00	5.167	
10	0	0 -00	4.920	0 -00	4.962	0 -00	4.920	
	20	2 -10	4.963	1 -00	5.998	2 -10	4.963	
	40	4 -40	5.006	2 28	5.034	4 -40	5.006	
	60	7 -30	5.049	5 -52	5.070	7 -30	5.049	
	80	10 -25	5.092	5 -12	5.106	10 -25	5.092	
	100	14 -00	5.135	7 -00	5.142	14 -00	5.135	
	. 0	0 -00	4.910	0 -00	4.952	0 -00	4.910	
	. 20	1 -20	4.949	0 -32	4.984	1 -20	4.949	
20	40	3 -00	3,988	1 -20	5.016	3 -00	4.988	
	60	4 -16	5.027	2 -12	5.048	4 -16	5.027	
	80	6 -00	5.066	3 -16	5.080	6 -00	5.066	
	100	9 -00	5.105	5 -00	5.112	9 -00	5.105	
30	0	0 -00	4.900	0 -00	4 <b>• 9</b> 42	0 -00	4.900	
	20	1 -00	4.935	0 -24	4.970	1 -00	4.935	
	40	2 -08		0 -44	4.998	2 -08		
	60	2 -48	5.005	1 -12				
	80	3 -36	5.040	1 -40	5.054	3 -36	5.040	
	100	5 -30	5.075	2 -30	5.082	5 -30	5.075	

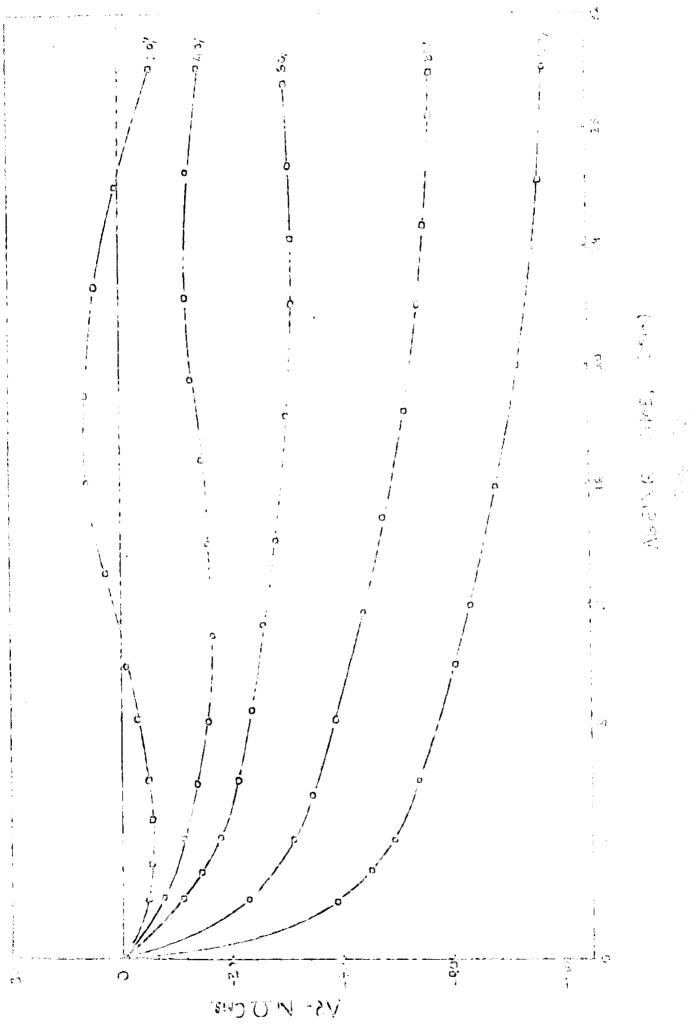


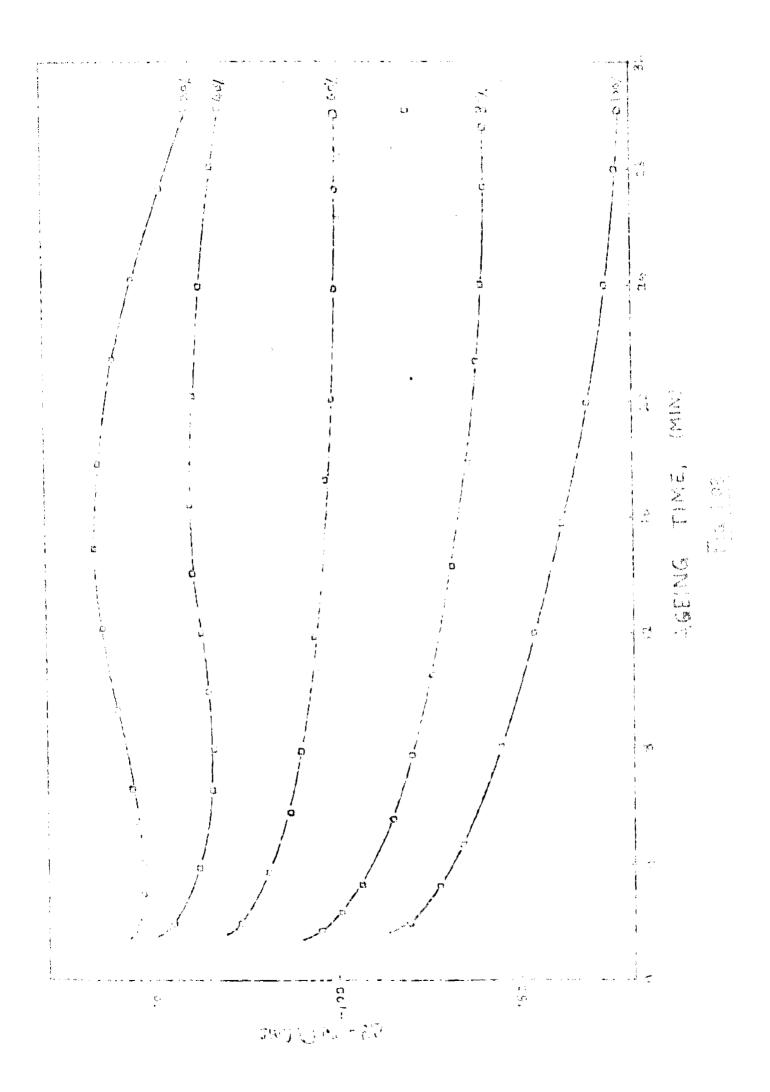




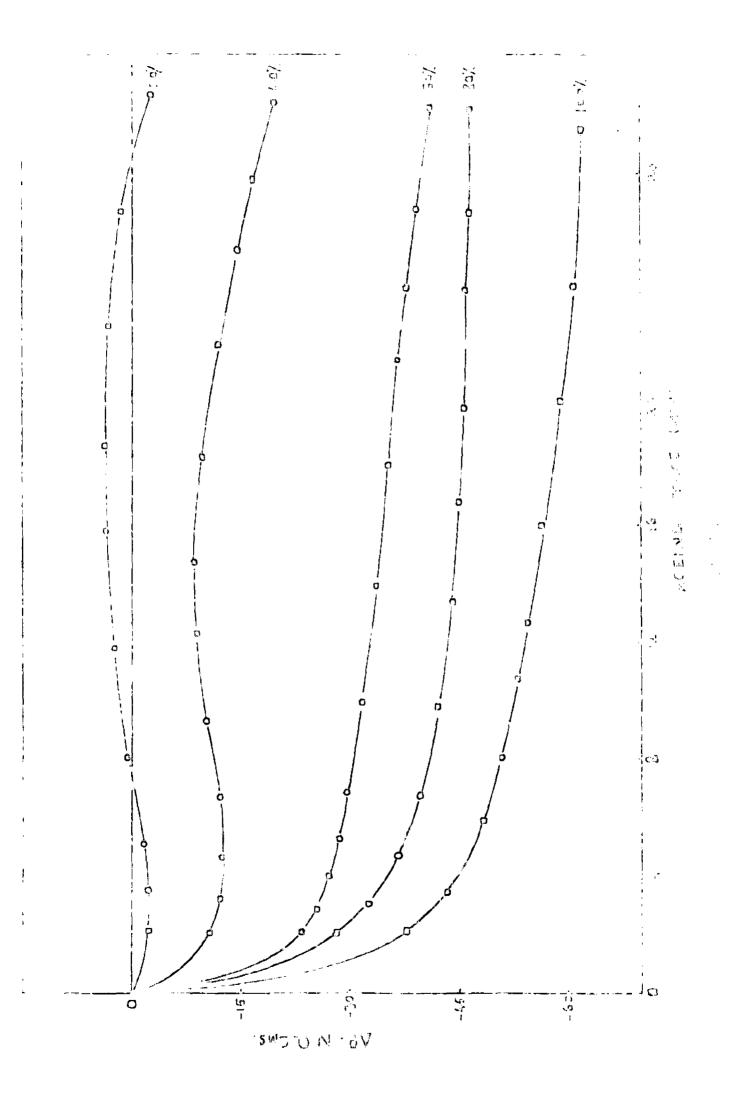


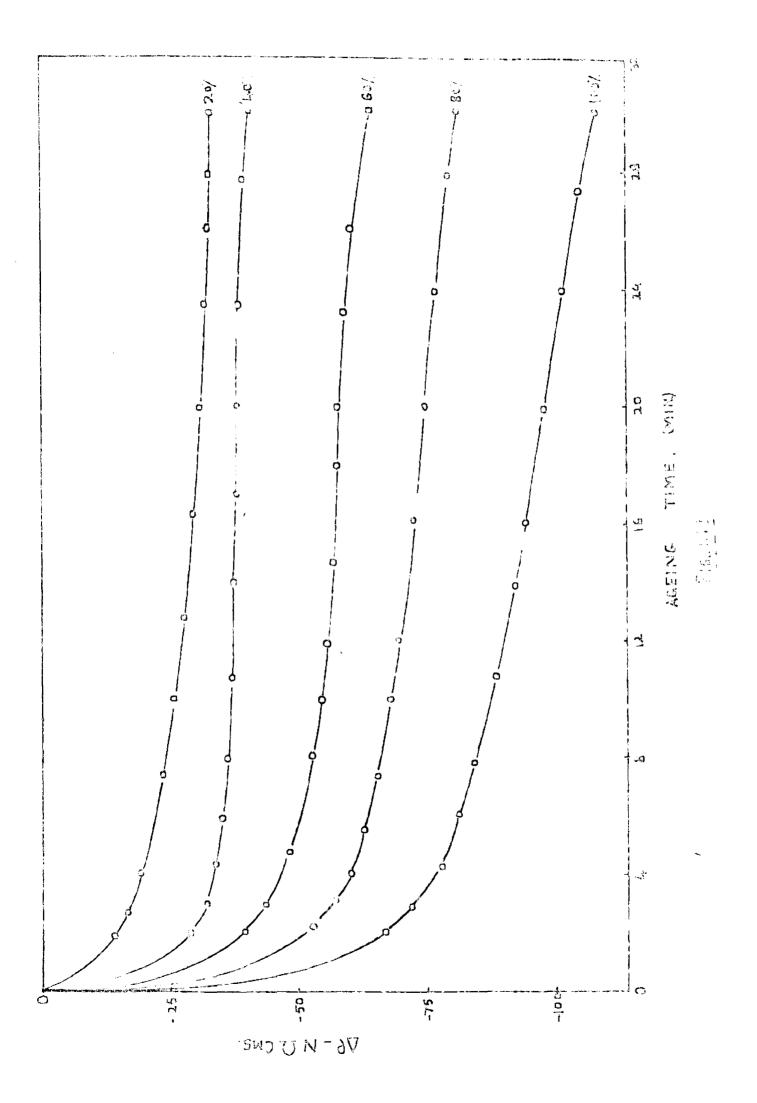
DE MOGNE

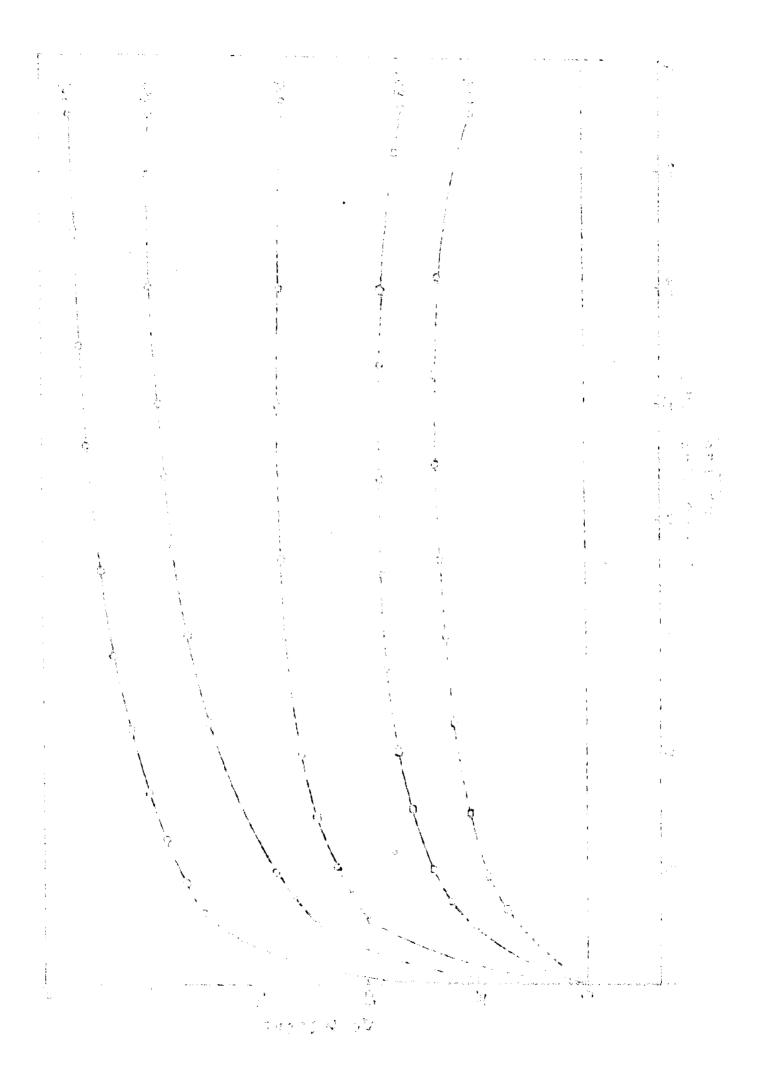


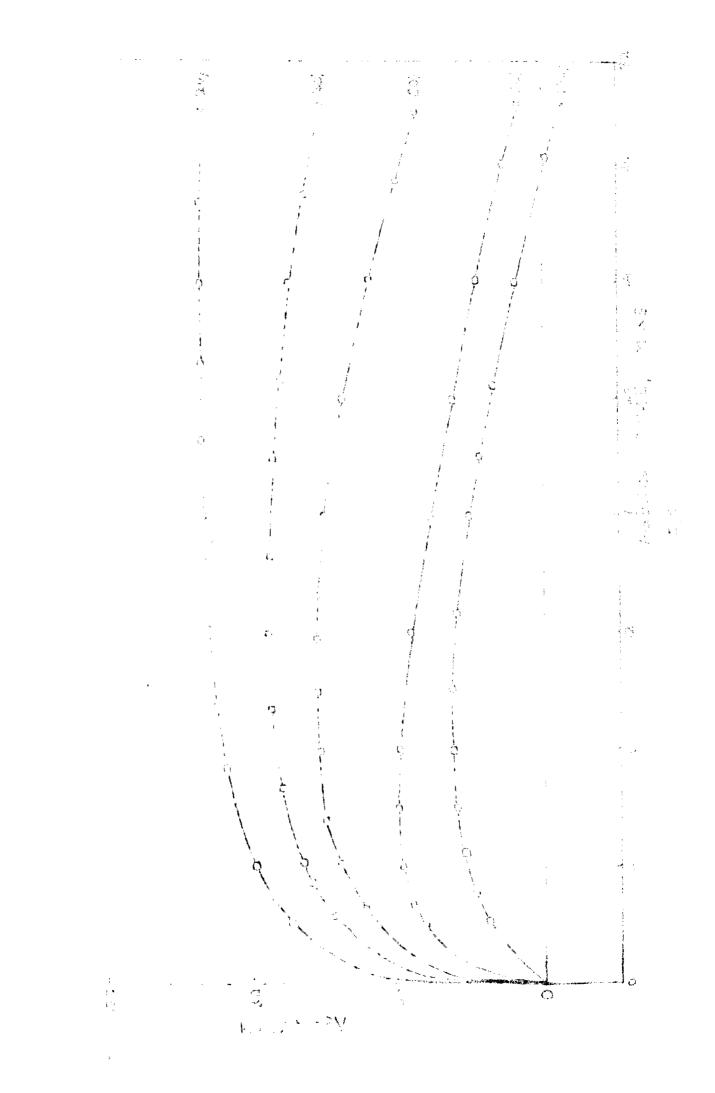


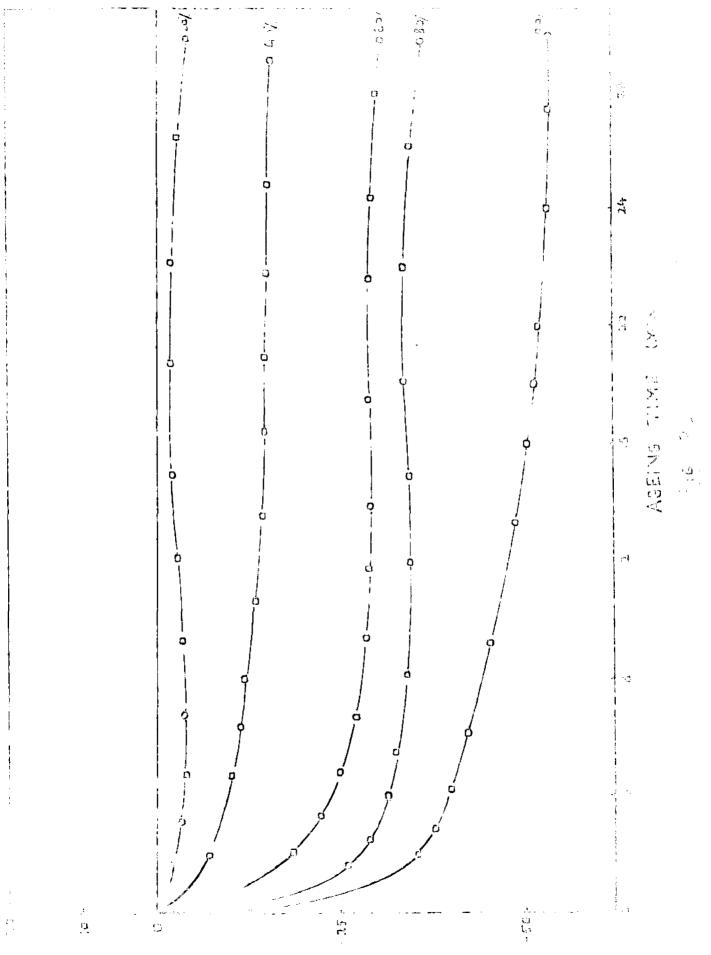








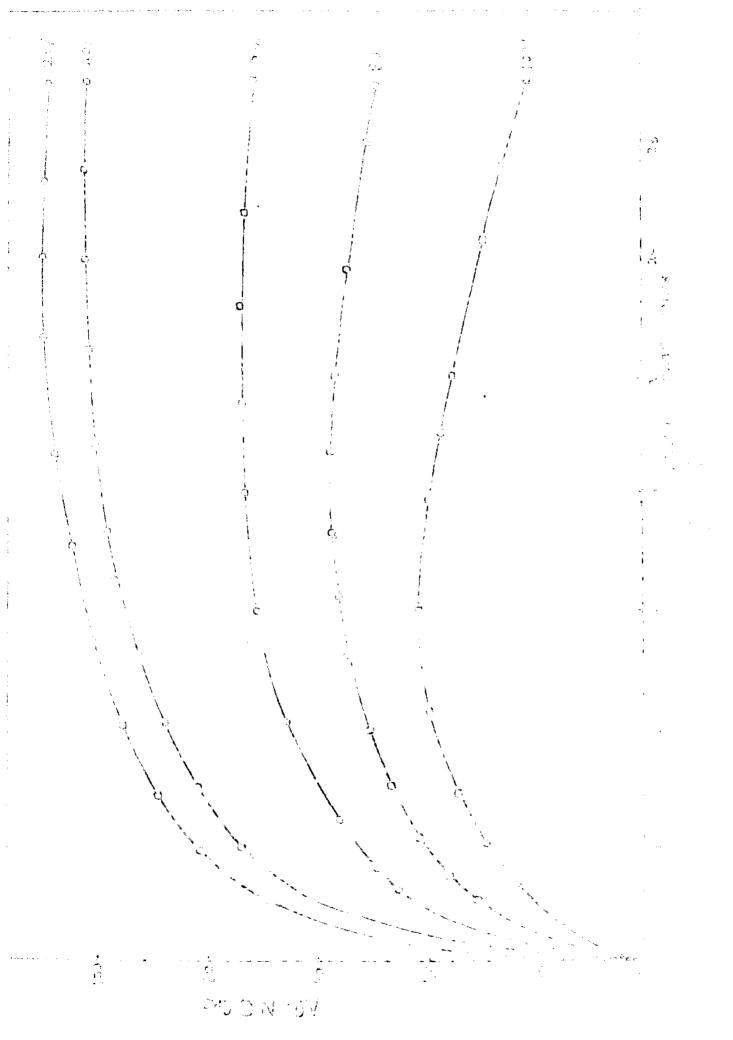


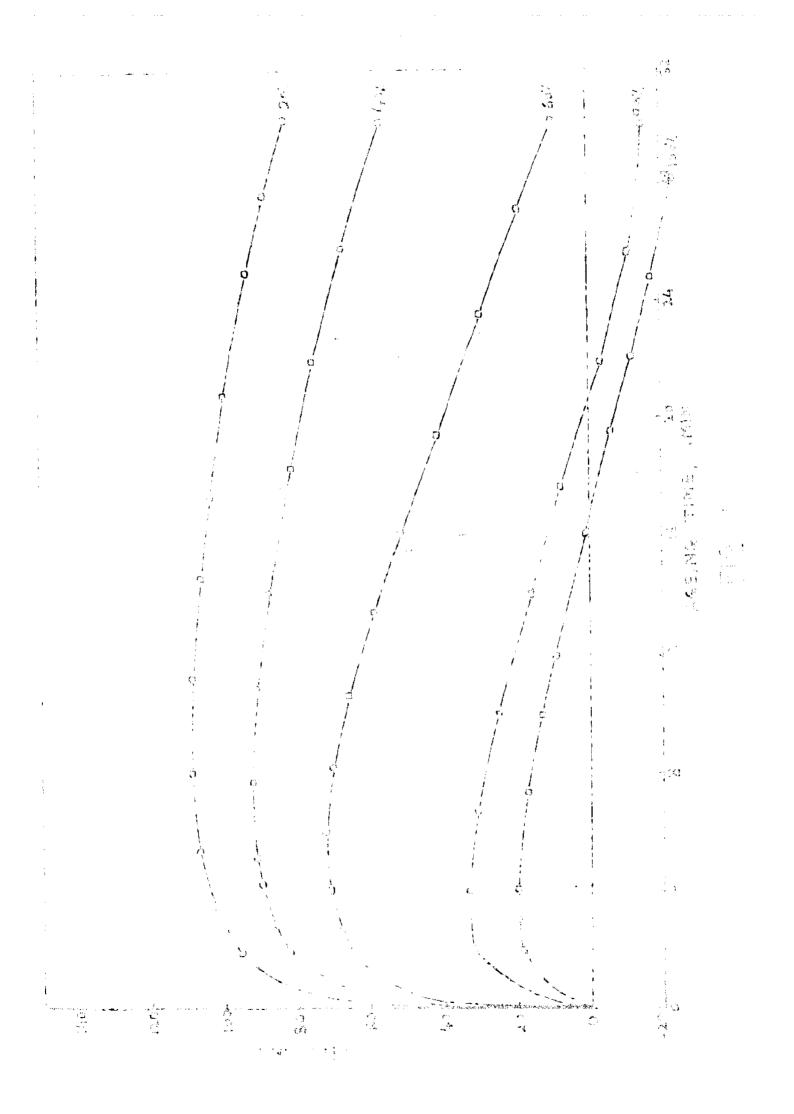


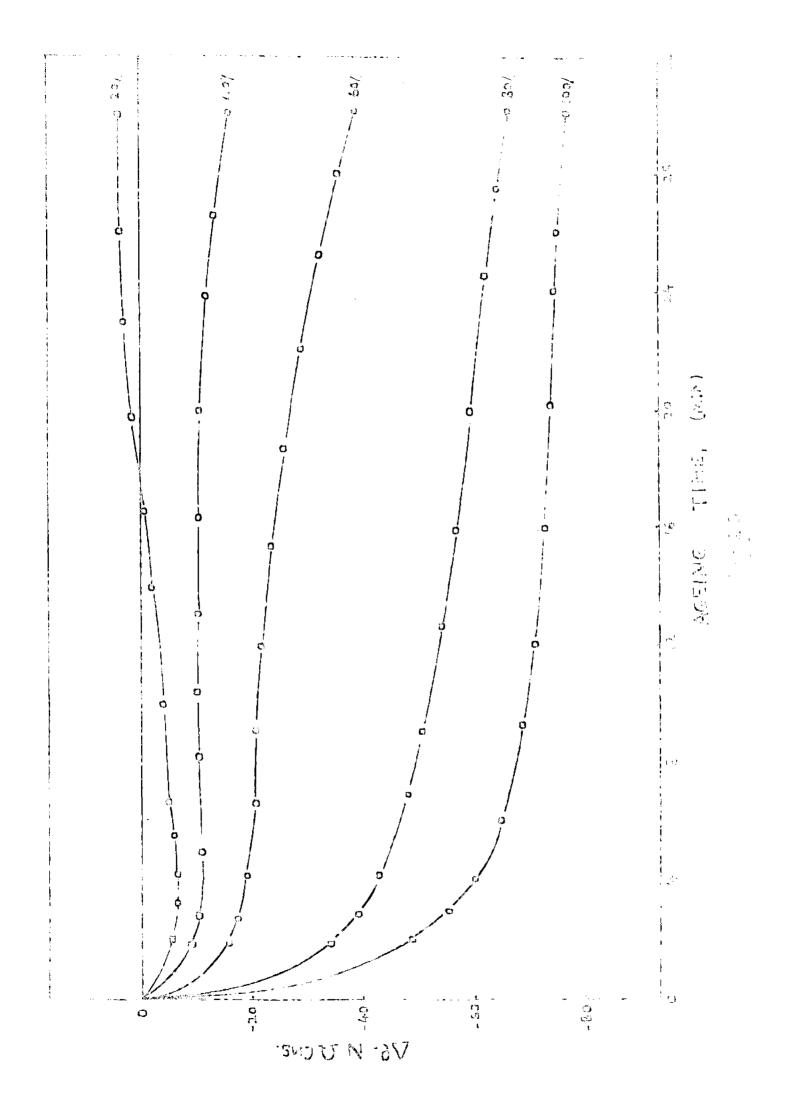
V6- NTTTWE

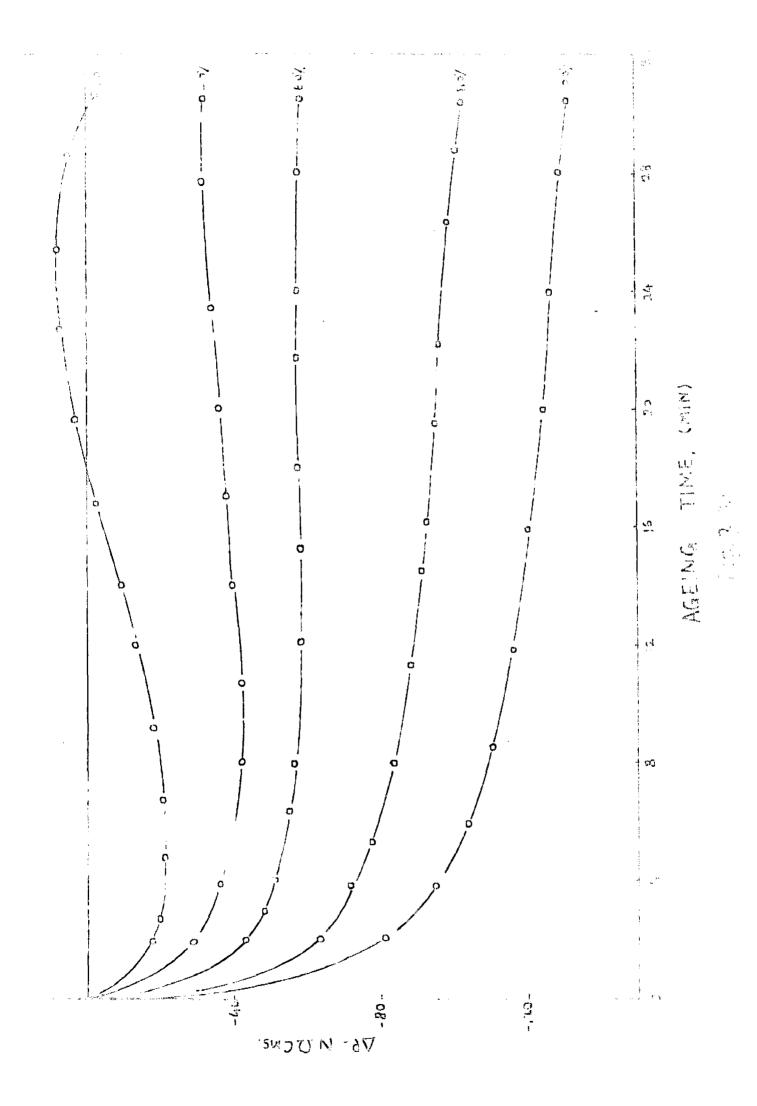


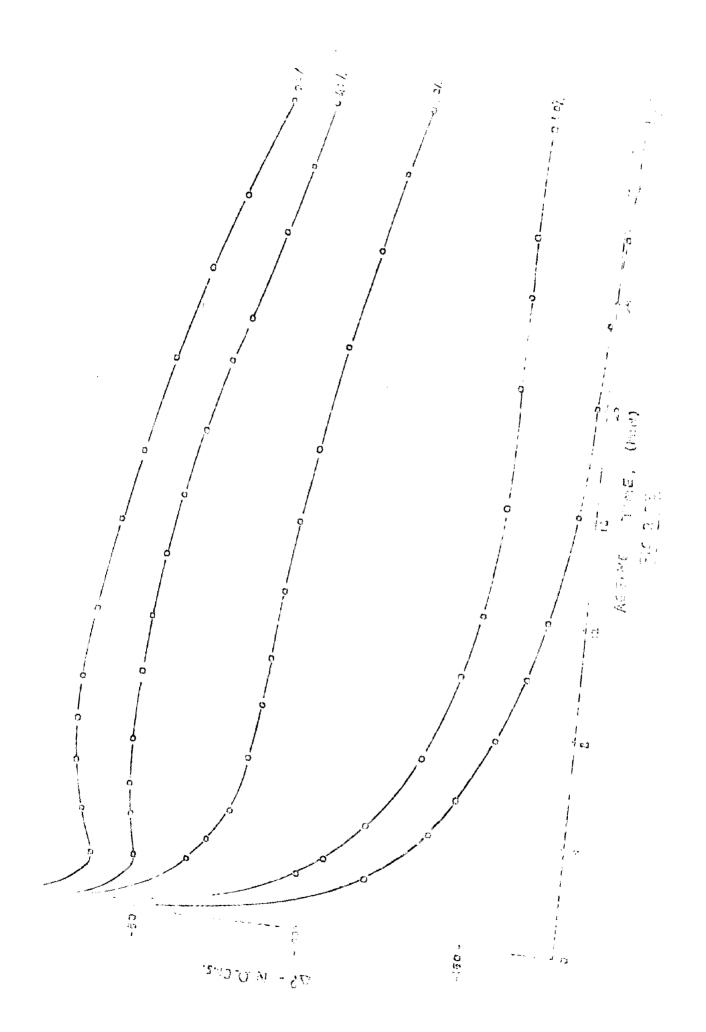
Sugar Charl

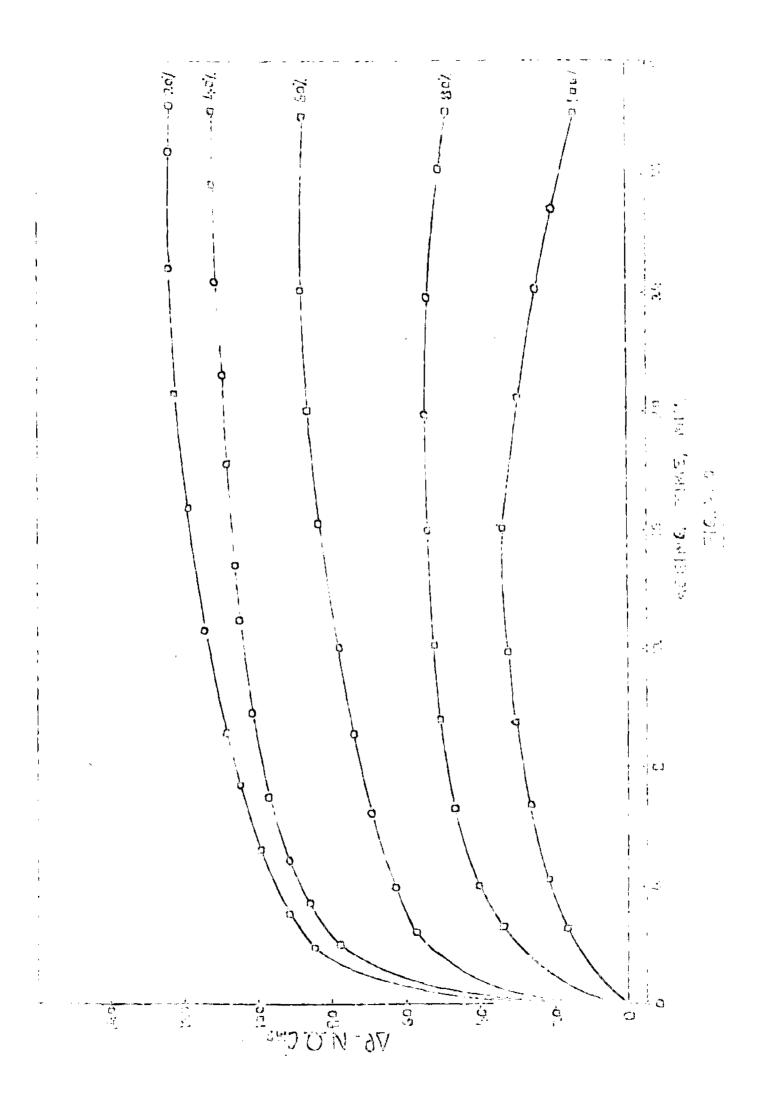


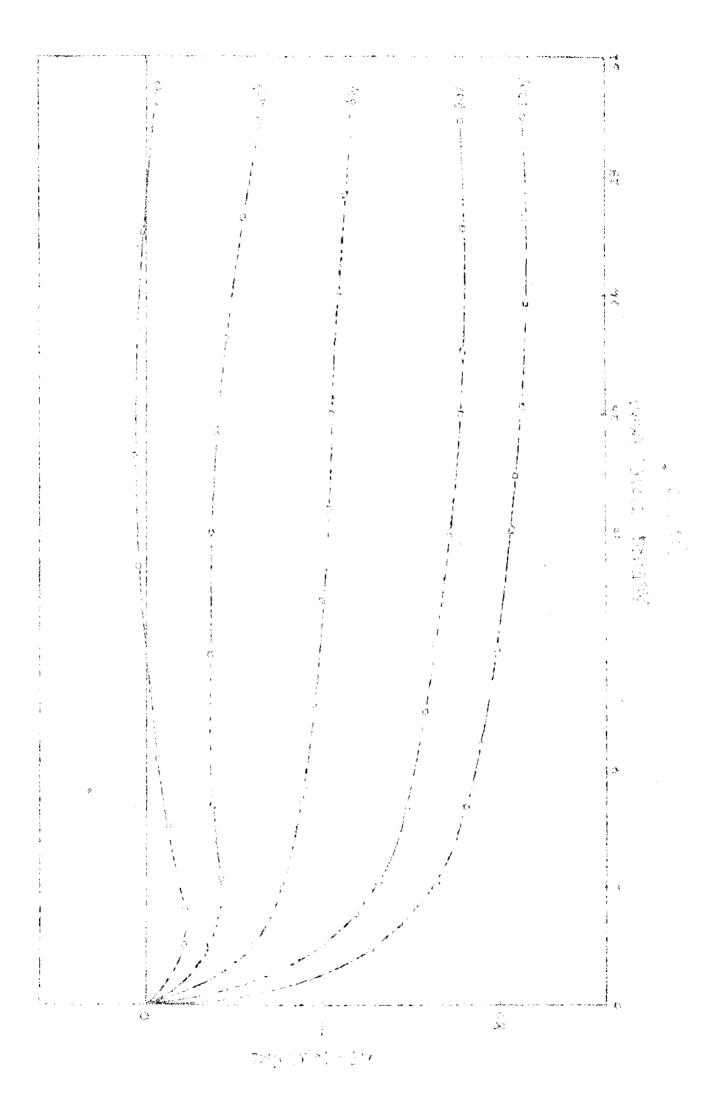


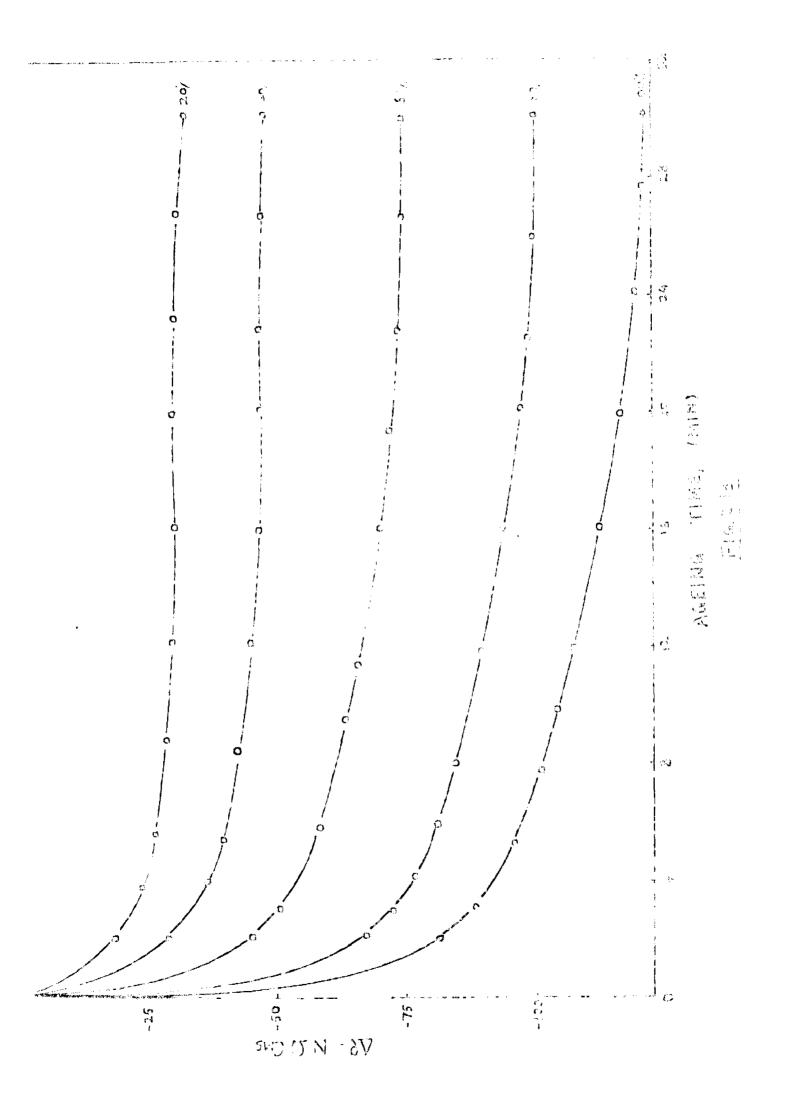


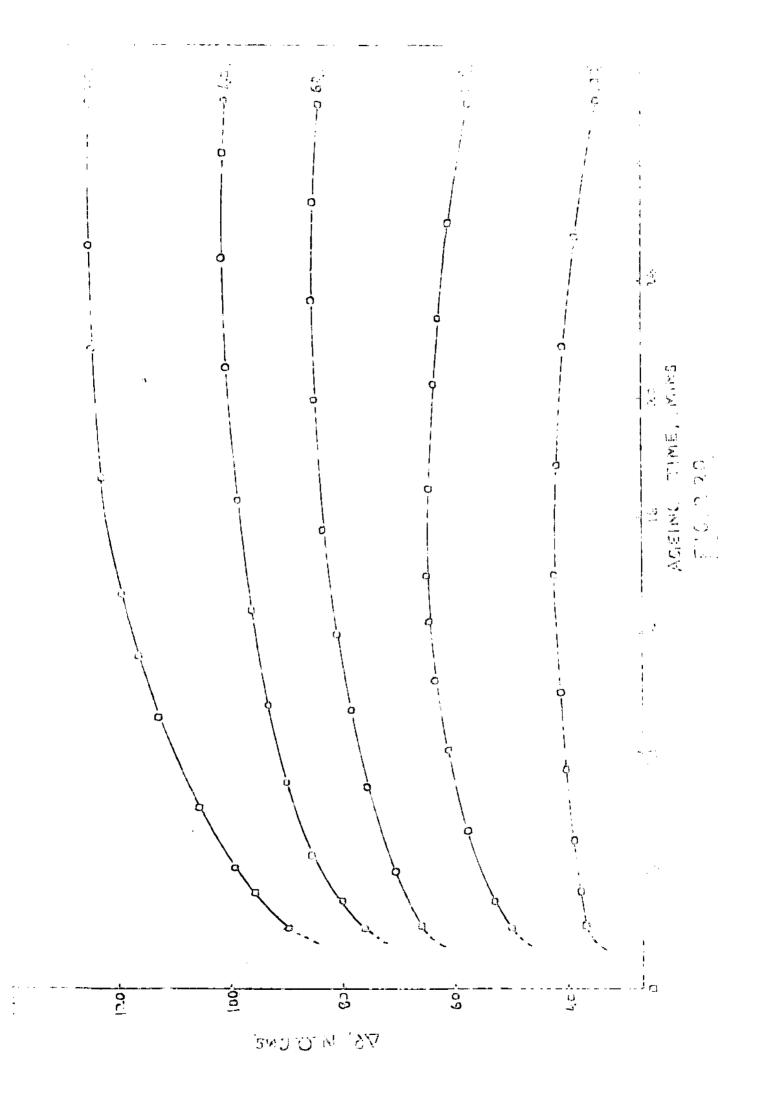


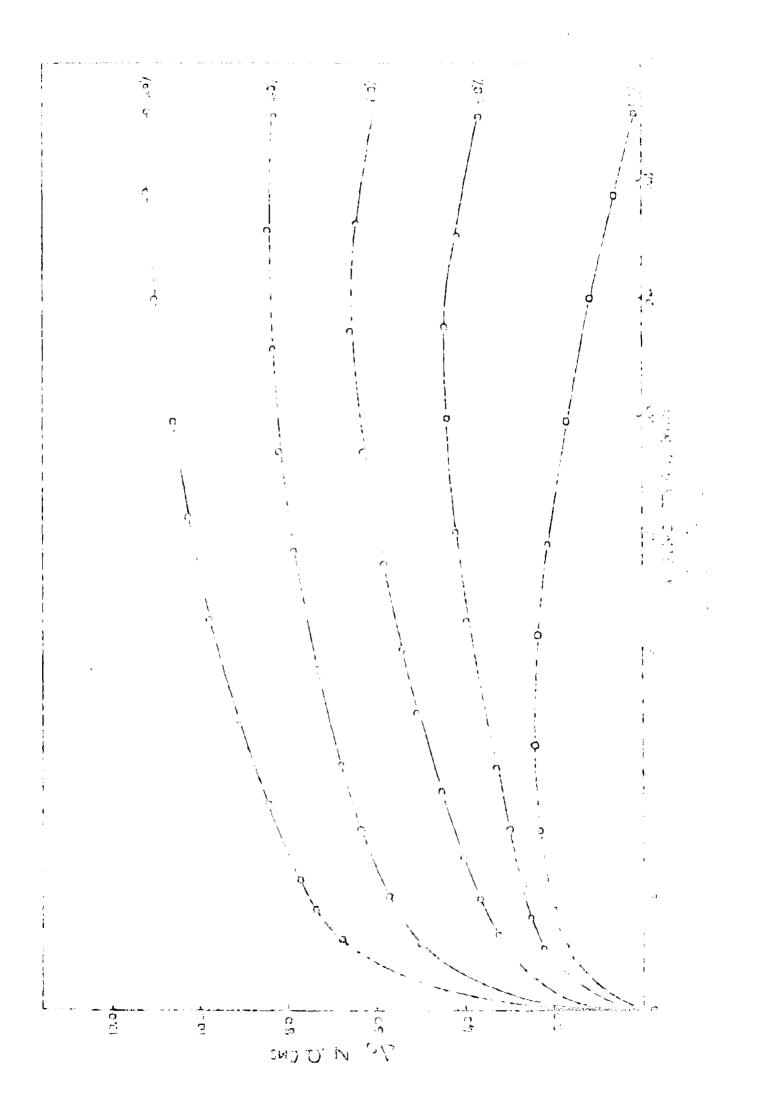


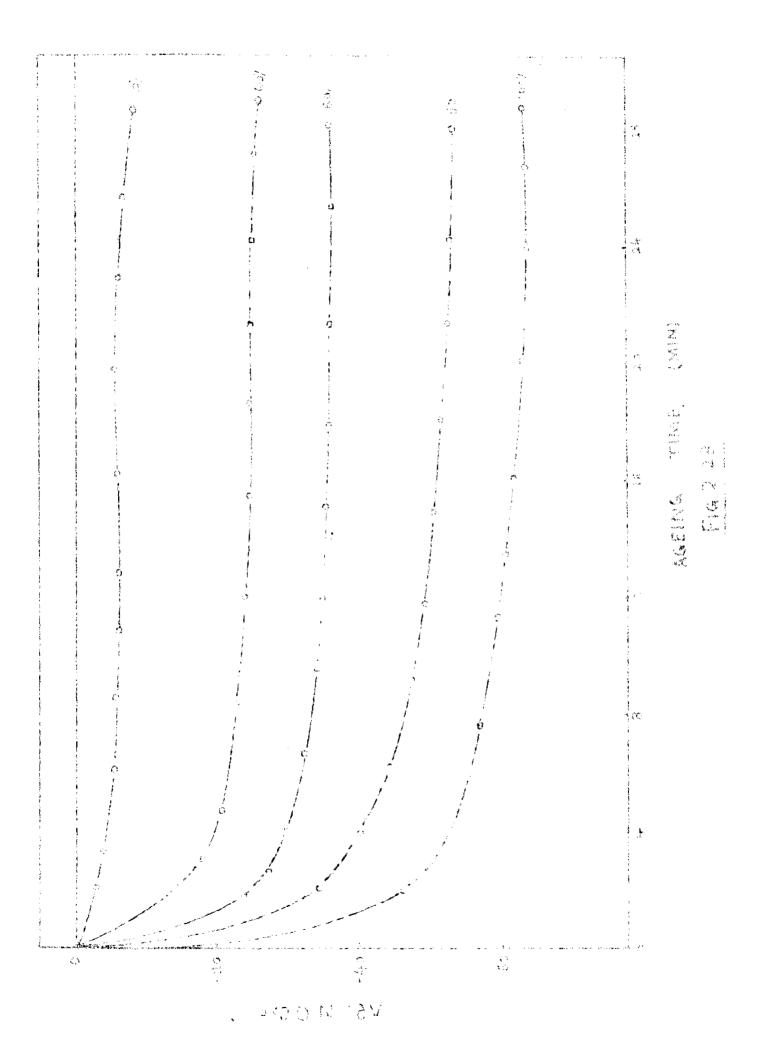


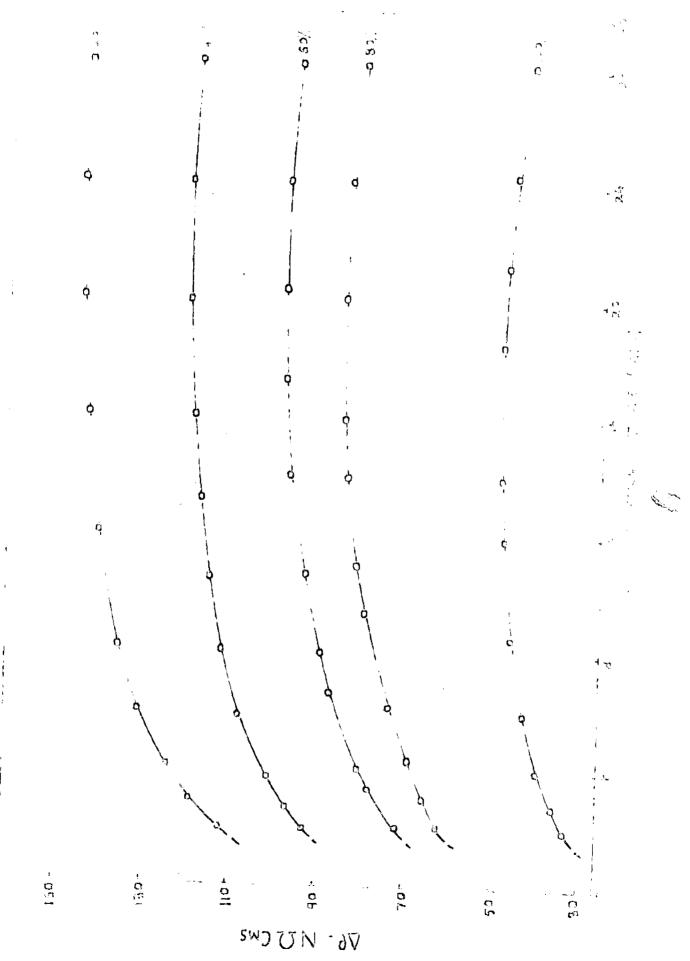


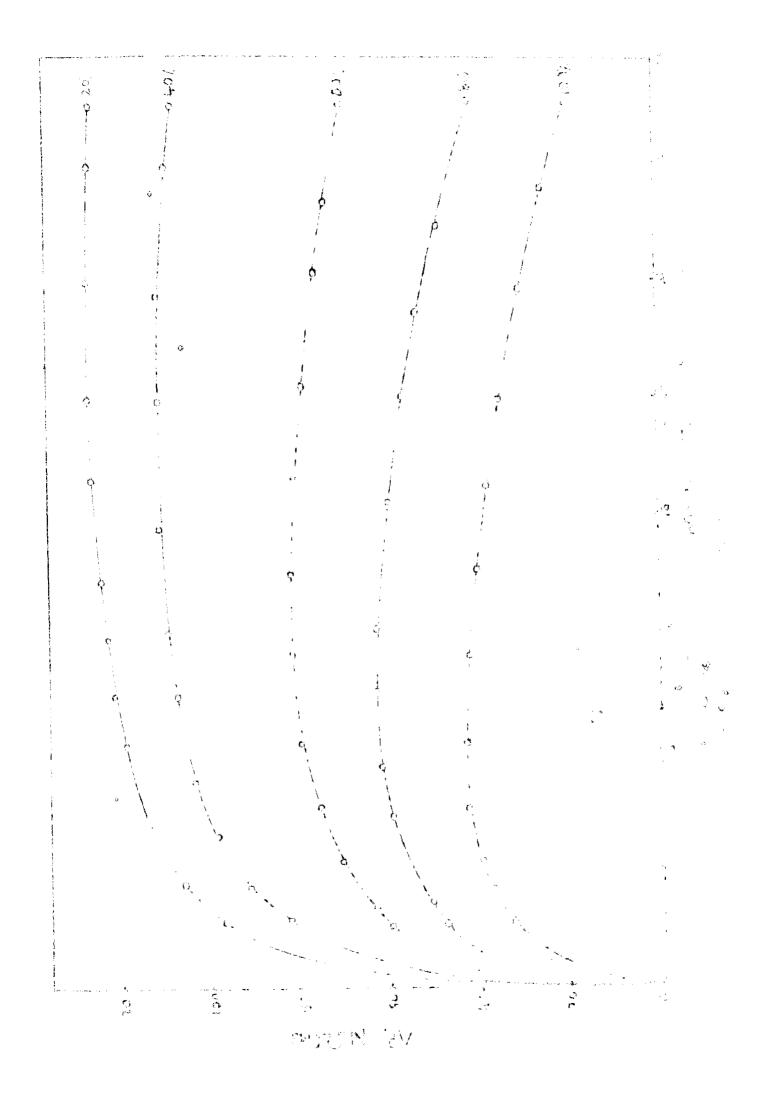


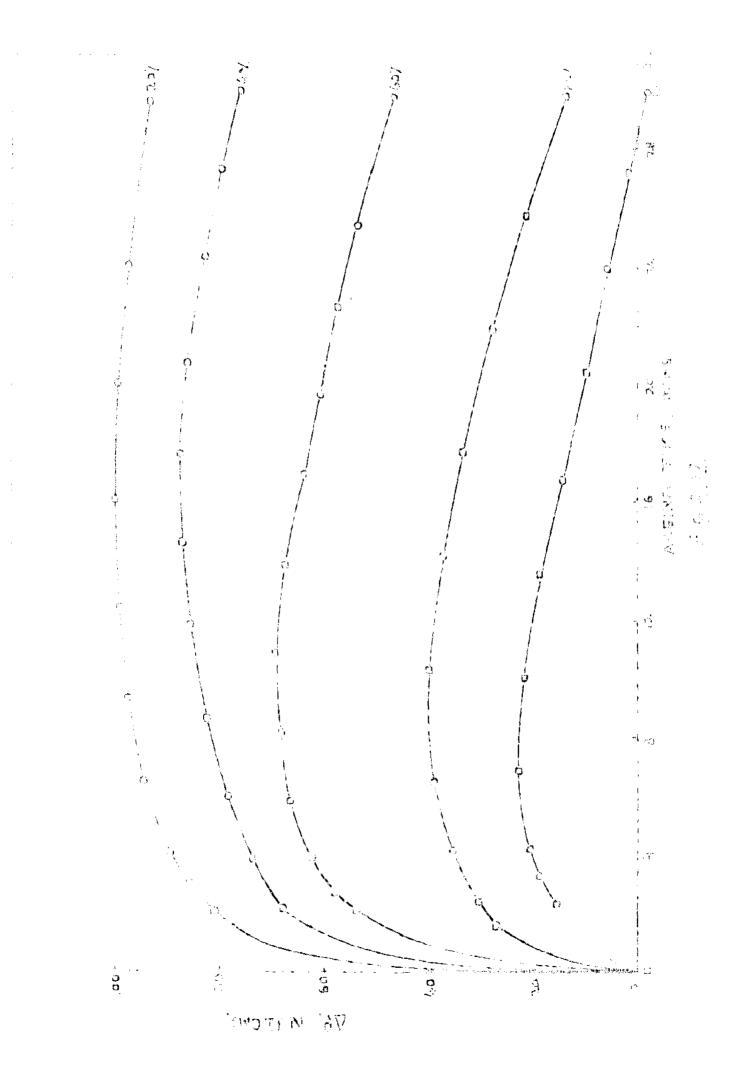


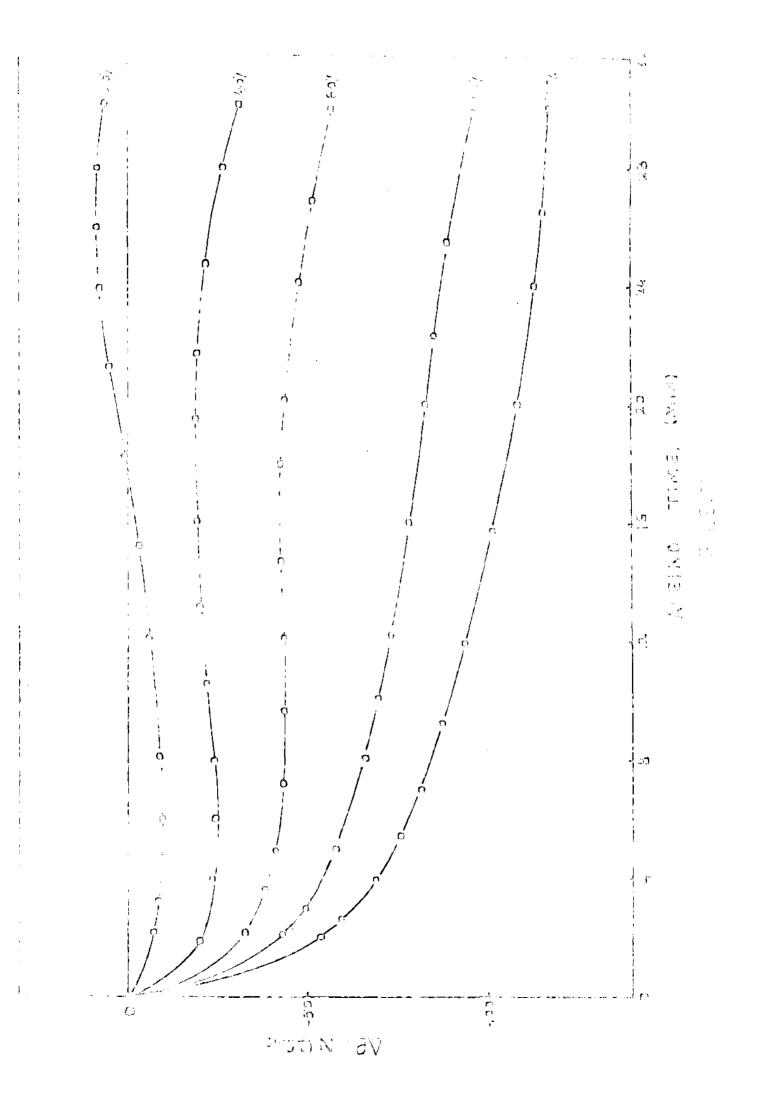


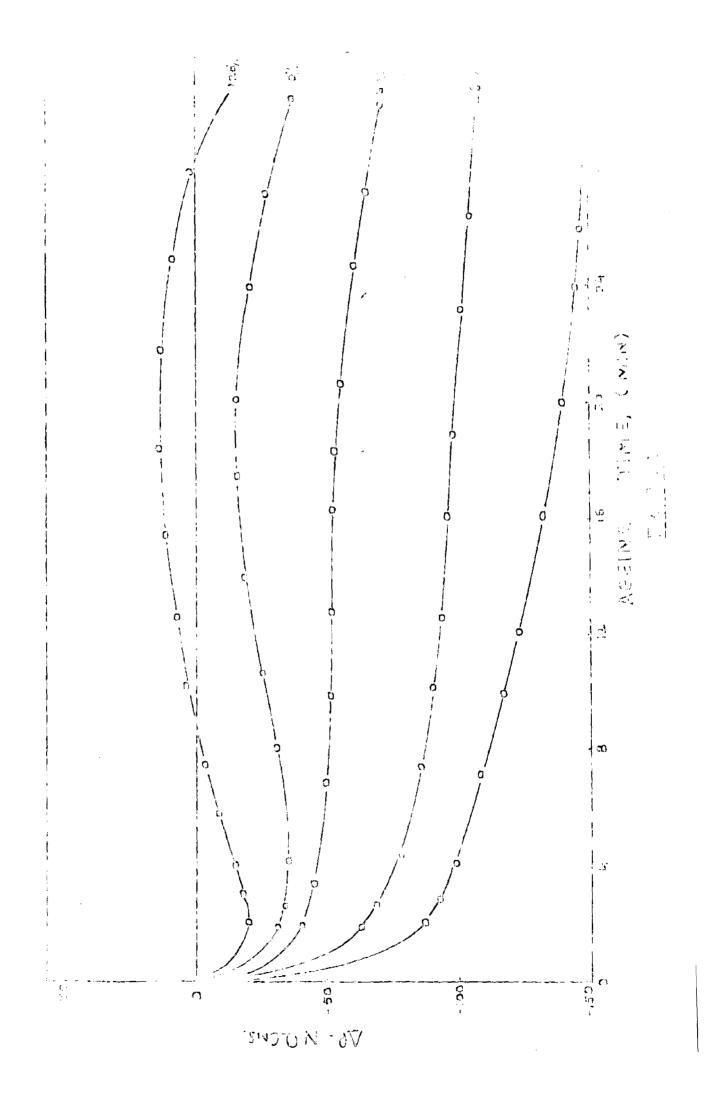


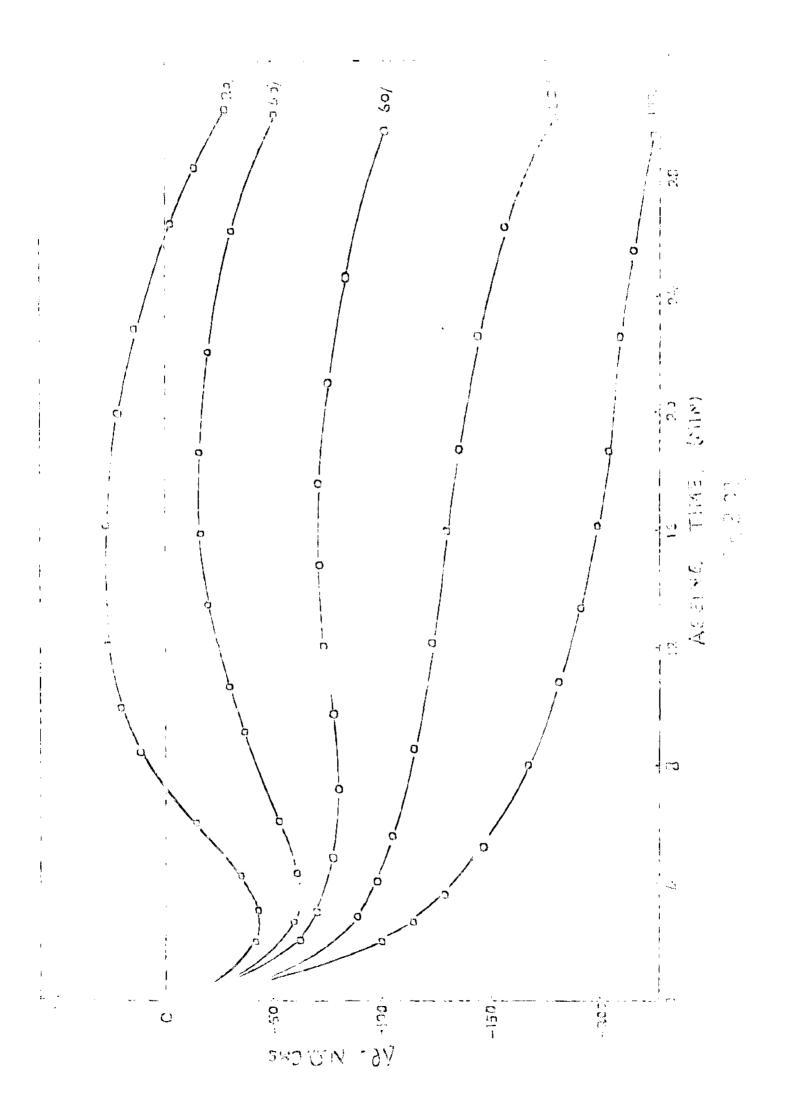


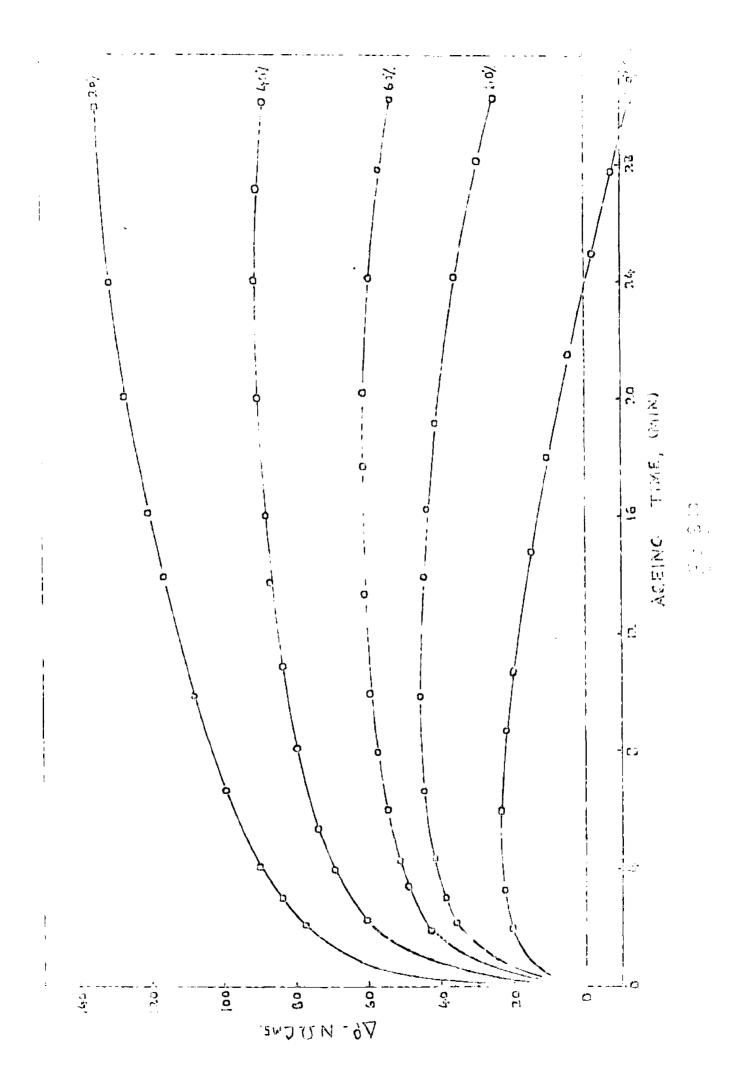


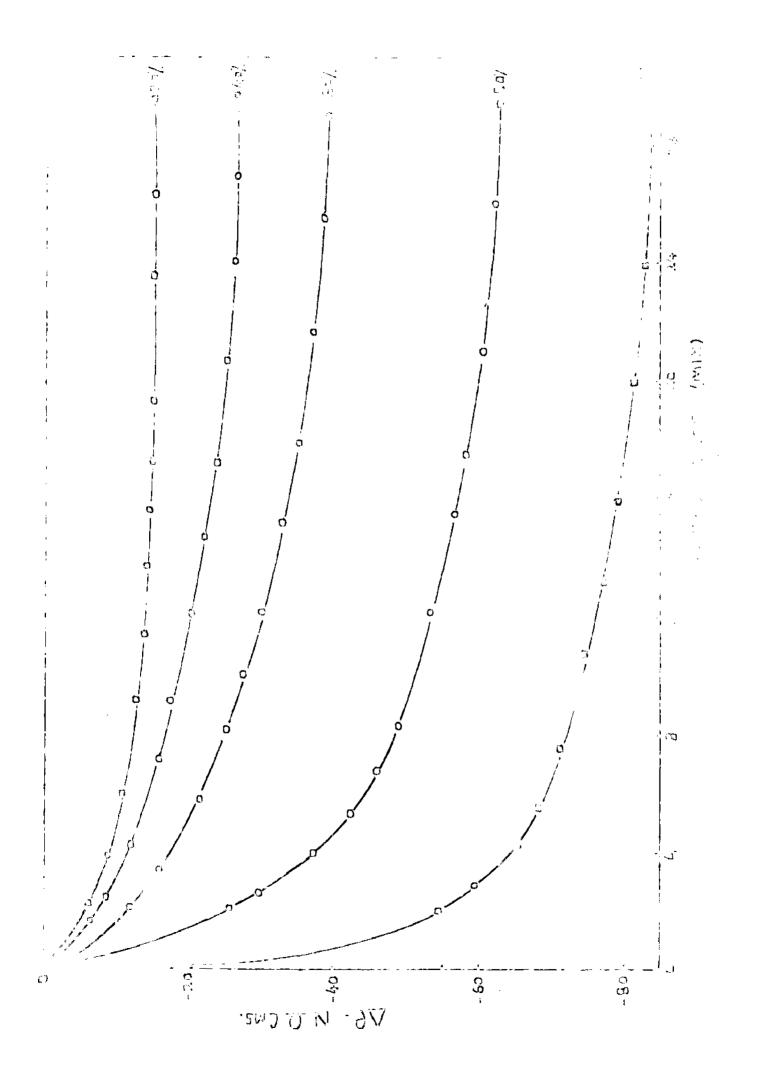


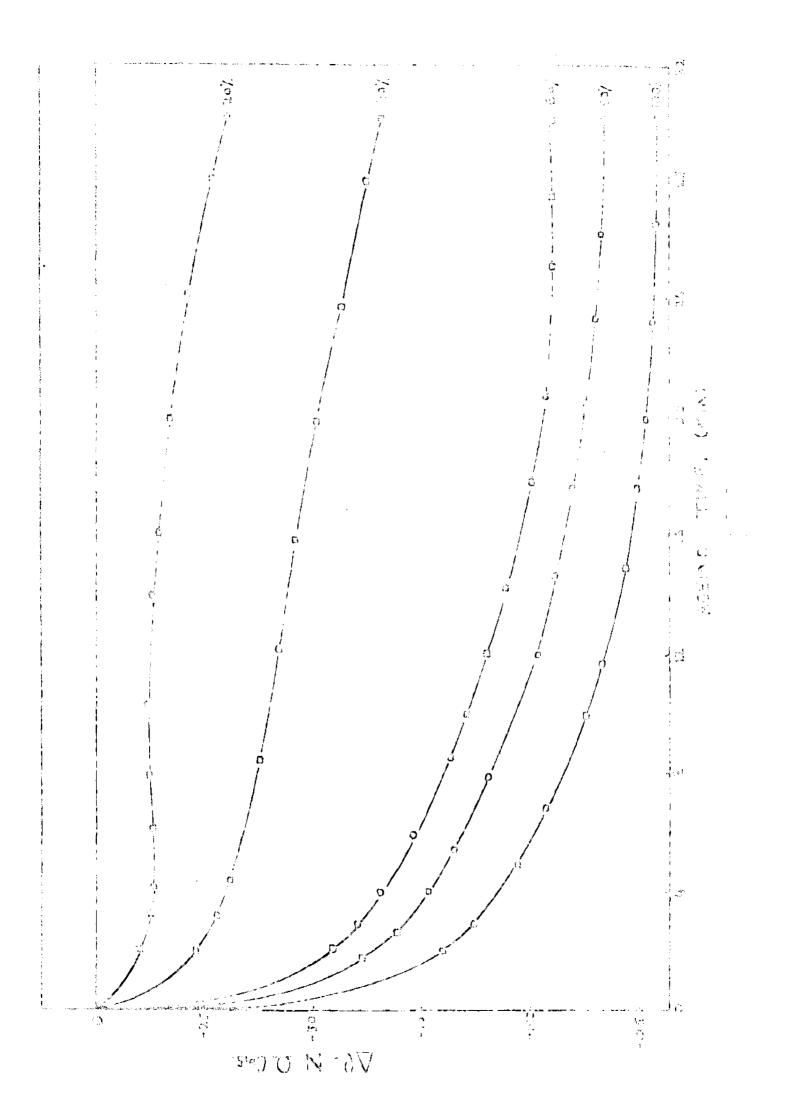


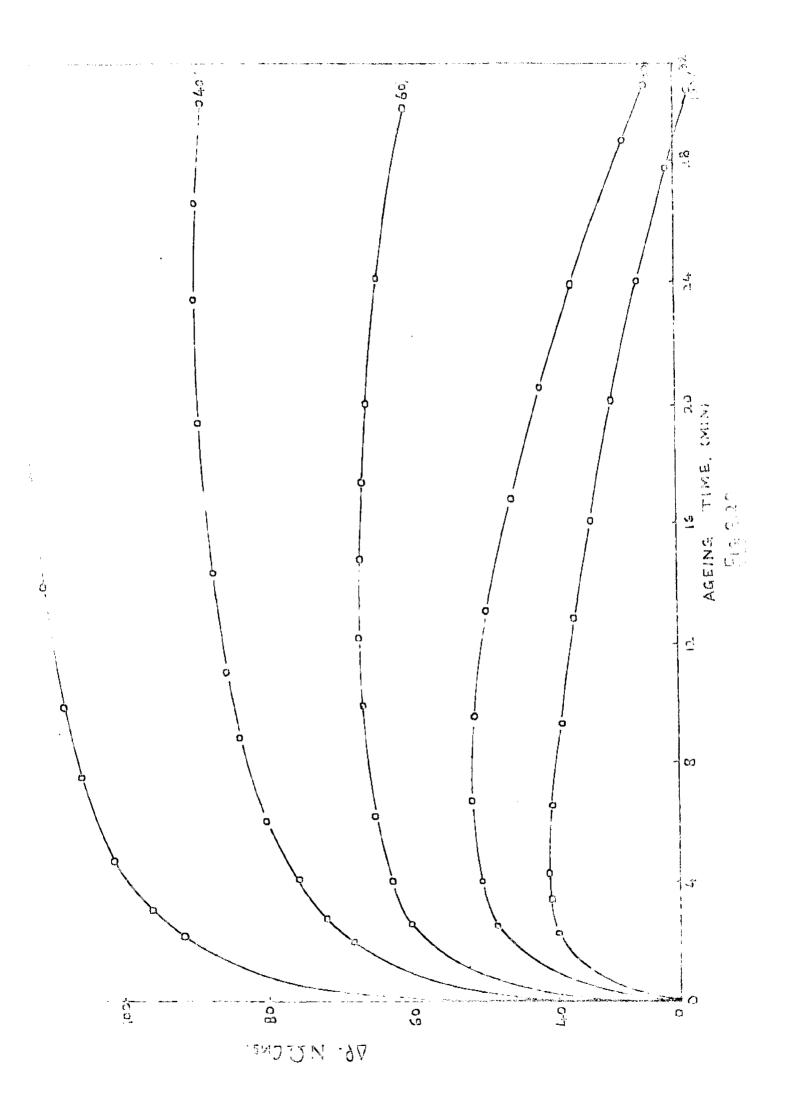


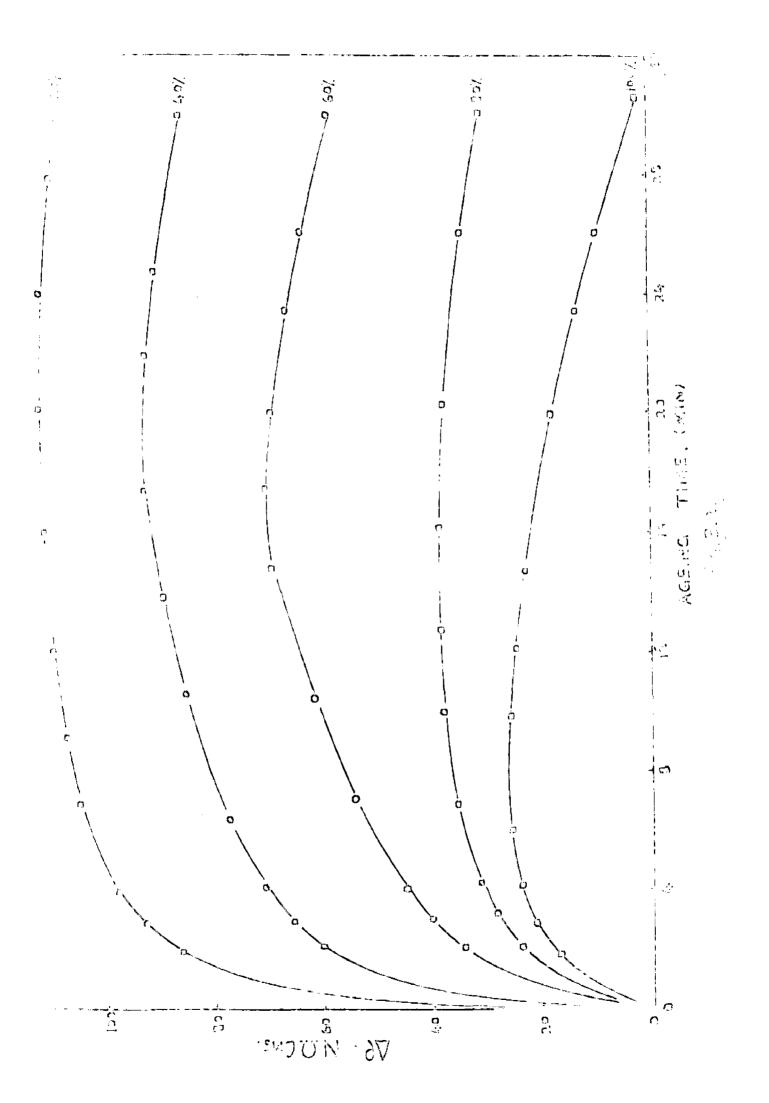


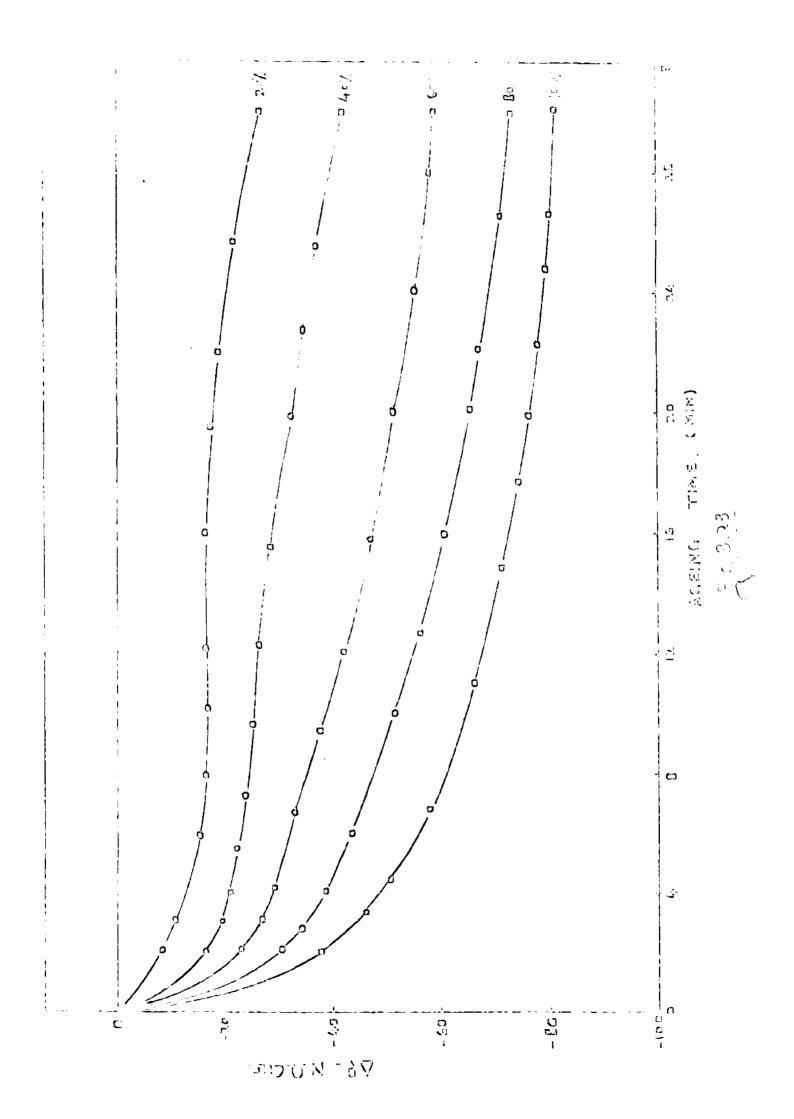


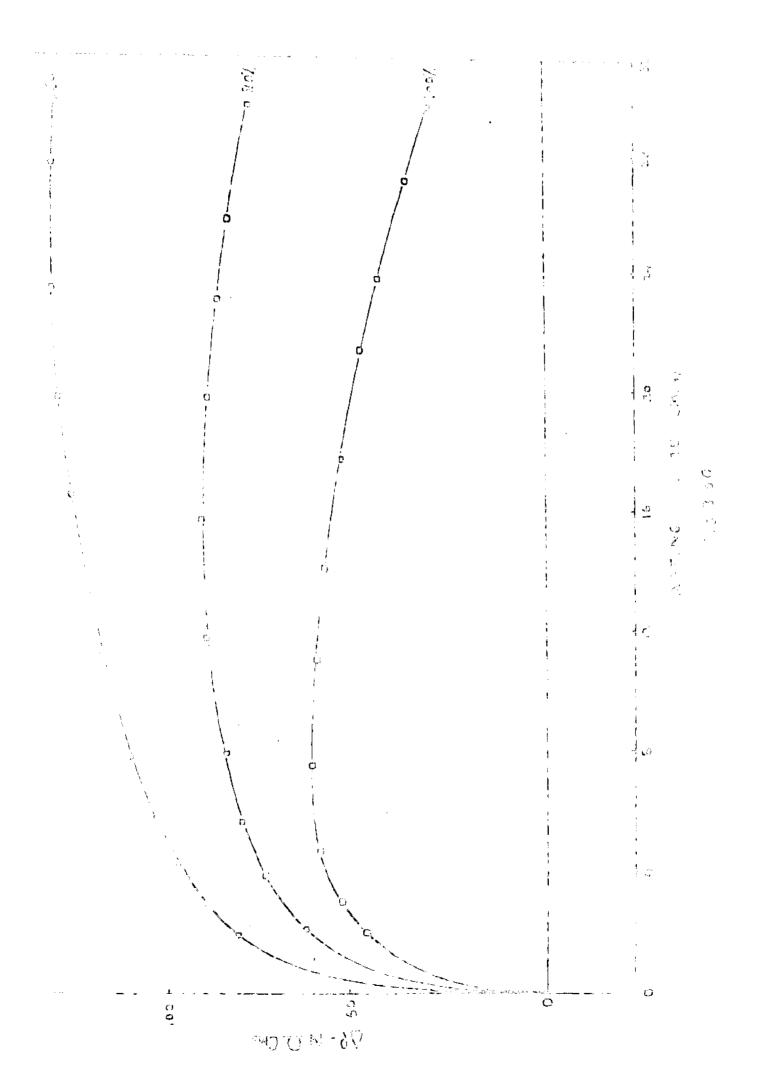


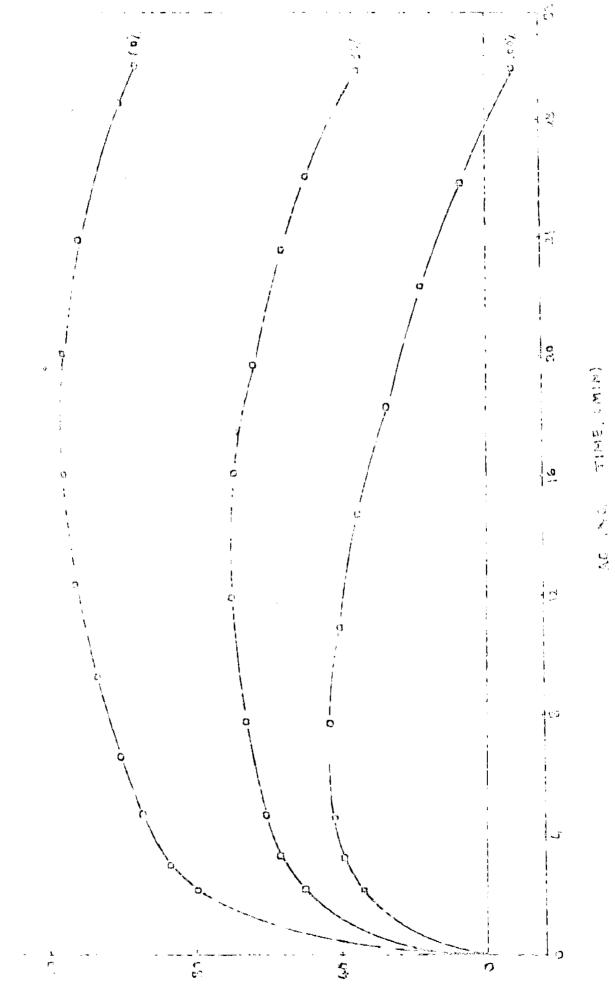






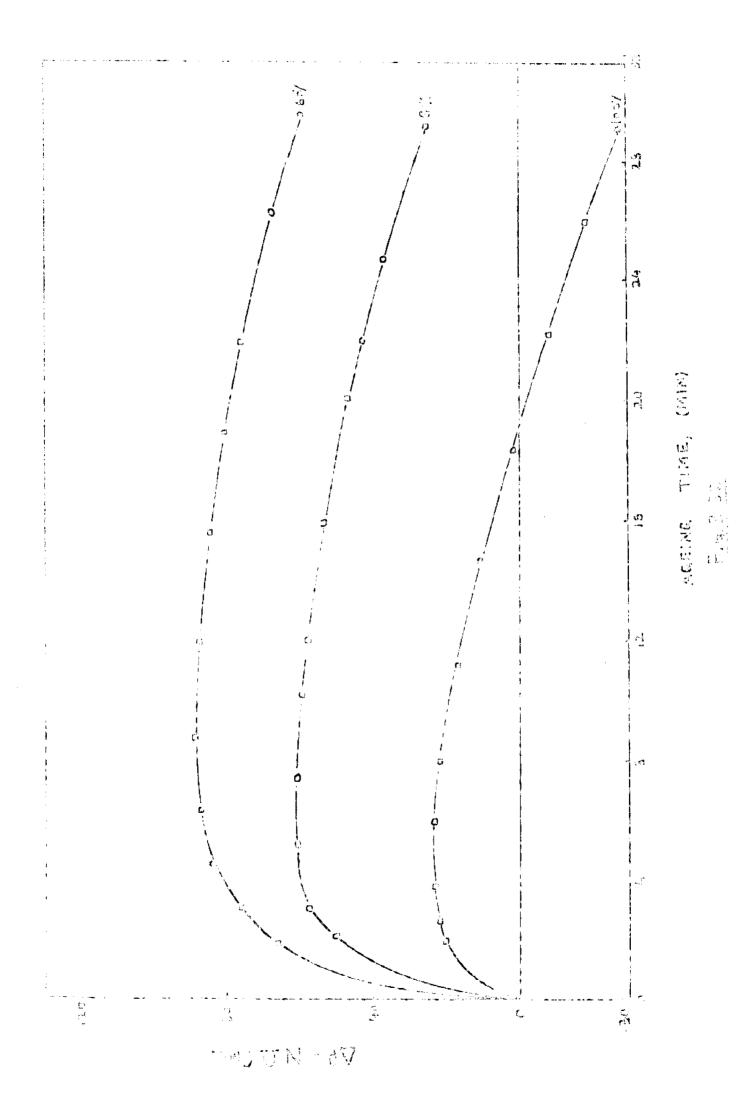






٢Þ

V6 IN UCHE



## DISCUSSION

OF

## RBSULTS

ς,

## 4.0 Discussion of Results

Broadly speaking, the present work can be classified into two categories w.r.t. the ageing and re-ageing temperatures.

- (A) Step ageing
- (B) Partial reversion of G.P.zones

The alloys after homogenization and quenching were subjected to an <u>ageing process</u> at temperature  $T_A$ , for different lengths of time, given by five equal intervals of the resistivity peak, during primary ageing at  $T_A$ , and after the lapse of this time, called as <u>pre-ageing time</u>  $(t_{P_A})$ , they were immediately subjected to a <u>reageing operation</u> at a temperature slightly above or below the ageing temperature  $T_A$ .

Now if, the reageing temperature was lower than the ageing temperature, the process is <u>step ageing</u>, and <u>reageing temperature</u> <u>denoted</u> as  $T_{RA}$ . And if the reageing temperature was above the ageing temperature, the process will result in the dissolution of a fraction of zones already formed at lower ageing temperature, with contemporary growth of the rest of the zones thermodynamically in equilibrium at the new reageing temperature. This process is termed as <u>partial reversion</u> of zones, and the <u>re-ageing temperature</u> <u>denoted as  $T_{R}$ . Thus we conclude that the <u>necessary conditions</u> for for the two operations are :</u>

(i) Step ageing if,  $T_{RA} < T_A$ (ii) Partial reversion of zones if.  $T_- > T_-$ . The results obtained for these operations with various combinations of ageing and reageing temperatures in the present study wore plotted on various figures as listed in Table III.

The results for step ageing, as well as for partial reversion of zones, are discussed on the basis of some fundamental established principles listed below:

(a) The average size of the zone  $(d_g)$ , reached after the lapse of a certain preageing time  $(t_{PA})$  during ageing at  $T_A$ . The  $d_g'$  of the zones, gradually increases on ageing, and becomes practically constant, after a time almost equal to the <u>average life</u> <u>time of vacancies</u>, which in turn is influenced directly by the prevailing concentration of vacancies and their annihilation.

It has been assumed that by Mott's theory, the peak in resistivity is characterized by a critical scattering, at a critical zone size  $(d_{\rm C})$ , which is approximately  $10^{\rm O}$ A, thus the average size of zones is increasing or decreasing, through a peak in resistivity obtained at size  $d_{\rm O}$ .

Further, it is clear that the time for reversion is directly related to the size of the zone, and thus to the preageing time  $(t_{PA})$  with other conditions being similar (Refer. Fig. 8-1R.).

(b) There is an equilibrium number of zones  $(N_Z)$ , having an average size 'd<sub>T</sub>', which are in thermodynamic equilibrium at the ageing or reageing operational conditions.

This  $\Pi_{Z}$ , is directly proportional to the <u>degree of super-</u><u>anturation</u>, which is determined by the position of the metastable solvus line for G.P. sones. Thus the degree of supersaturation is directly related to the ageing temperature  $T_{A^*}$  (or on  $T_{RA}$  or  $T_R$  for that matter), and hence, the  $\Pi_Z$  also being directly dependent on  $T_A$ . Lower be the  $T_A$ , higher will be the supersaturation, and thus the  $\Pi_Z$  value, and vice versa. In other words a lowering of  $T_A$ , will result in the growth of more number of sones, which are in thermodynamic equilibrium given by the  $T_{RA}$  or  $T_R$ , while raising the ageing temperature would result in discolution of some of the sones, thus establishing a new number of some determined by  $T_{RA}$  or  $T_R$ .

The development of a some with average size  $(d_T)$  which is thermodynamically stable at a cortain temperature, is influenced by the kinetics of the initial stages of agoing, and thus on  $T_H$  and  $T_A$ , which in turn controls the  $C_V$ , their annihilation, and the formation of secondary defects.

(c) At the  $\ell$ , of the alloy is a direct function of the number and close of the sonce (thus their volume fraction), the <u>height of</u> <u>the restotivity</u> peak is directly dependent on the equilibrium number of zones, while its <u>position</u> is characterized by a critical centering of electrons, at sone size  $d_{C}$ .

Dow,  $\Pi_Z$  is directly related to the  $\Upsilon_A$ ; the height of the resistivity peak is proportional to this  $\Pi_Z$  value. In such a case,

the difference in the  $\binom{max}{max}$  at agoing and reagoing temperatures is directly proportional to the difference in these temperatures. We summarize by saying, that peak heights of resistivity changes  $(\Delta \ell)$  versus reaging time, curves will be directly proportional to the <u>difference</u>  $(T_A - T_{RA})$  for step agoing case and to the <u>differ</u>-<u>ence</u>  $(T_R - T_A)$  for partial reversion case.

(d) The <u>slope</u> of the recistivity variation curves at instant t; (d  $\ell$  / dt), during reagoing is a <u>kinotic term</u>, giving a hint about the rates of the process. Paster process will have greater slopes. Also a zero slope does not always means no reaction, but it may signify a balance in the two opposing effects contributing to the recistivity, as shall be discussed later.

(c) The ageing process is diffusion controlled, and hence, the migration of solute atoms, the concentration of vacancies and their danihilation (activational anorgies for formation and migration of vacancies) the interactions between the solute and the vacancies etc. are the factors governing the kinetice of the process. These factors govern the <u>alops</u> of the ageing curve, and also the <u>position</u> of the maximum.

The concentration of vacancies, the preageing time, is important for kinetics of the reageing process. With the lapse of  $t_{PA}$  the  $C_V$  is greatly reduced, and hence the <u>rates of the step</u> <u>acting</u>, as well as reversion process is bound to be slover, with

-51-

its effects very clear during later stages of reageing, giving more or less horizontal curves. Explanations of the curves for various,  $t_{PA}$  values (from 20 pct. to 100 pct. resistivity peak values), can be advanced in these terms. The times to attain resistivity peaks on step ageing process is much higher compared to that for ageing process.

(f) The variation in resistivity during the step ageing and specially during the reversion process, is a complex phenomenon and is the <u>sum up of various factors</u> contributing to the resistivity, as discussed below, taking the case of partial reversion:

i) In actual practice there is a <u>distribution of somes</u> of varying sizes in the vicinity of the average zone size,  $d_{z}$ . Thus,  $\rho$  is the sum up of the individual contributions of the zones of different sizes.

ii) There is a decrease in resistivity due to dissolution of a fraction of the zones lower than a temperature dependent critical size  $(d_T)$ , which are thermodynamically unstable at reageing conditions. Thus a decrease in volume fraction of zones.

An increase in resistivity, due to the growth of zones,
 with size above than those dissolving on reageing, as pointed out
 above. This gives rise in volume fraction of zones.

We conclude, that a net variation of redictivity during step agoing or reversion, will be the sum up of these factors. The individual contribution of these factors will determine the sum up curve, thus leading to an increase in  $\ell$ , if the factors contributing to the  $\ell$  increase over ride the factors contributing to the decrease in  $\ell$ , and vice versa. A horizontal portion of agoing curves may either signify the balance of the two opposing tondencies or a vory slow reaction because a major part of vacancies has alroady been lost by that time because of lapse of  $t_{\rm PA}$ .

Thus a complox phenomenon of step agoing or partial revercion can be comfortably explained in terms of above discussed and well established principles. In any of these operations, the tendency for the reaction is such so as to obtain the new equilibrium state which is thermodynemically most stable under these conditions.

It will be more systematic to discuss the present results under two separate headings of stop ageing and partial reversion.

## Stop Agoing

Taking the case(similar to fig.1.10)where reageing is done at  $T_{RA}$ , after a specimon has been pro-aged at  $T_A$ ,  $(T_{RA} < T_A)$ , till peak in resistivity, we expect the following :

1) The average size of zones in the alloy pro-aged till peak  $(t_p)$ , at  $T_A$  has attained the critical scattering value,  $d_C = 10^{\circ}A$ 

Now reageing of this at a lower temperature  $T_{RA}$ , will result in further growth of these zones, thus contributing a decrease in  $\ell$  (curve B of fig. Q), which is actually the part of the ageing curve after the peak at  $T_A$ .

11) An increase in l, due to increase in the equilibrium number of zones, which are expected to grow, due to lowering of the ageing temperature from  $T_A$  to  $T_{RA}$ , thus increasing the degree of supersaturation. This is shown by curve A.of fig. Q, and represents the growth of new zones at reageing temperature.

111) As the equilibrium number of zones is dependent on ageing temperature in this case the height of peak in curve A, is directly proportional to the difference in the equilibrium number of zones at  $T_{RA}$  minus those at  $T_A$ . Greater be the difference  $(T_A - T_{RA})$ , greater is the degree of supersaturation, and correspondingly the difference in the equilibrium number of zones, hence resulting in higher peaks for curve A.

iv) Times to peak in curve A is appreciably greater than compared to the time to peak for primary ageing at  $T_A$ , because of low  $C_v$  value, in the former case.

v) The net effect of these two factors, responsible for ourve A and B, should be the curve expected on step ageing, which has been shown by curve C in fig. Q. The cum up curve C, chows the poculiar nature of a rise, a flattened peak and finally a drop, corresponding to initially pre-dominating nature of curve  $\Lambda$ , while at later stages, the curve B pre-dominates.

vi) The reaction rates, as expected are much slower, because of considerable loss of vacancies during the pro-ageing stay time.

Now with this back ground, we shall take individual cases of step agoing, with their features and explanations.

Fig. 1.10 :- It may be observed that :

(1) For all values of  $t_{PA}$  from 20 pct. to 100 pct. there is an increase in resistivity, with least distinct peak in 20 pct. case. This is expected, because of the increase in volume fraction of zones, due to lower  $T_{RA}$ , compared to the  $T_A$ .

(11) Rate of stop agoing is factor in the case with lesser  $t_{PA}$ , because of high  $C_y$  left over with lesser pro-agoing.

(111) A net increase in  $\ell$ , shows that the decrease of  $\ell$ , due to growth of somes, boyond size  $d_{C}$ , on re-agoing process is far less, than the increase due to growth of new number of somes.

(1v) The height of peak in 20 pct. case is greatest, because of maximum difference in the number of zones and the size of zones formed after 20 pct. pre-ageing time at  $10^{\circ}$ C and those given by the peak at  $0^{\circ}$ C. <u>Fig. 1.20 and 1.21</u> :- Basic arguments are similar as for fig.1.10 with a few new observations :

(i) Comparing the curves in fig.1.20 and 1.21, it is observed that the peaks are attained earlier, while reageing at  $10^{\circ}$ C, simply because of enhanced diffusion at  $10^{\circ}$ C, compared to that at  $0^{\circ}$ C.

(11) A dropping of curves for 100 pct. case in later stages could be explained on the basis, that the zones which were previously growing have also attained resistivity critical size;  $\int_{h}^{d_c} by$ that time and they also contribute a decrease in resistivity. This also explains the appearance of earlier peaks in 100 pct.case compared to the 20 pct. cases in the same figure.

Pig. 1.30, 1.31 and 1.32 :- The nature of curves is similar to those obtained for step ageing in fig. 1.20. Here we note :

(1) In the case of fig. 1.32, the peaks are obtained earlier when we move from 100 pct. case towards 20 pct. case. This trend is also visible for curves in fig. 1.31 and 1.30. This is simply because of the earlier setting-in of the factors contributing to a decrease in (, in the case where the specimen has been already pre-aged so as to give high, everage zone size, prior to step-ageing. This is clear with the fact that the longer be the  $t_{pA}$ , larger will be the average zone size prior to step-ageing.

(11) A comparison of times to reach peak in 100 pct. cases, for

the figures 1.30, 1.31 and 1.32 shows that peaks are carlier with higher ro-agoing temperature, which is because enhanced diffusion with higher re-ageing temperature.

(111) Slopes of the curves will lessor preageing time is steeper in each figure, and this is because of high concentration of vacancies with lessor preageing times.

(iv) Sloping down of curves with higher preageing times, is prominent in all these figures, which is still clear in figures with higher reageing temperatures. These could be explained on the quicker attainment and growth of somes, which had already attained sizes near to resistivity peak, when they are subjected to higher reageing temperatures.

(v) The crossing of zero reference line at an earlier time in 100 pct. case than in 80 pct. case, is due to larger prior zero clise in 100 pct. case. A value lower than the zero reference line indicates that the zeros have grown to such a size, where their contribution to resistivity is lessor than that just after proagoing. <u>Effect of  $T_{\rm H}$ </u>:- Upto now the only variables were the agoing tomperature and the reageing temperature with different preagoing times. The nature of the curves obtained with different  $T_{\rm H}$  values is similar to those obtained in figures with  $T_{\rm H} = 300^{\circ}$ C, already discussed above. For the offect of  $T_{\rm H}$ , we shall, as an onempto compare the figures 1.10, 2.10 and 3.10:

(1) For a certain pre-ageing time, for the above three cases, there is only a slightly higher peak with higher  $T_H$  value, but this effect is not very appreciable. This probably due to higher concentration of quench-in vacancies with higher  $T_H$  value.

(11) For the initial rates of reageing in the above three cases, it can be observed that the rates of stepageing are faster in the case which has highest  $T_{H}$ . Again this is probably due to higher concentration of vacancies with higher  $T_{H}$  values.

The effect of  $T_{\rm H}$  can also be similarly explained while comparing the figures 1.30, 2.30 and 3.30.

# Partial Reversion of Zones

Taking a case similar to figure 1.01, where an alloy has been aged at  $T_A$  and reaged at  $T_{RA}$ ,  $(T_R > T_A)$ , we observe : (1) A decrease in resistivity due to growth of those sones at  $T_R$  which had already attained a resistivity critical size at  $T_A$ . This has been shown by curve A of fig.R; and is actually, the part after the peak, of the ageing curve at 0°0.

(11) A decrease in resistivity due to decreasing volume fraction of somes, because of equilibrium number of somes at  $T_R$  is much less than that at  $T_A$ , thus resulting in dissolution of a fraction of sones. This is given by curve B of fig. R.

(iii) On prolonged reageing at  $T_R$ , there is possibility that the zones with size larger than those thermodynamically stable at  $T_R$ , may grow, thus contributing an increase in resistivity.

(iv) Thus the net sum of these factors will give the resultant, partial reversion curve 'C', which shows a gradual decrease in resistivity on reageing.

(v) As should be expected, the maximum drop in resistivity will be observed in the reversion operation in which the degree of supersaturation is greatly reduced, i.e. greater be the $(T_R-T_A)$ value, greater the depth of reversion curves, as this is directly related to the ratio of the equilibrium number of zones given by the peaks of the ageing curves at  $T_P$  and  $T_A$  respectively.

(vi) One should expect the rates of reaction for lower  $t_{PA}$  to be faster because of high concentration of vacancies with lesser  $t_{PA}$ .

<u>Fig. 1.01:-</u> We observe following features in 100 pct. preageing:

Now we shall discuss individual cases of reversion.

(1) A decrease in resistivity due to growth of those zones at  $T_R$  which have already attained critical scattering size during preageing, till the peak in resistivity.

(11) A decrease in resistivity due to dissolution of a fraction of zones, because equilibrium number of zones at  $10^{\circ}$ C is lower than that it o'c

Both of these factors contribute in decreasing the f, hence the sum up curve also shows a continuous decrease on reageing.

In 20 pct. preageing case, we observe a decrease followed by an increase and finally not very distinct peak, which is explained in following terms:

Probably, before crossing the zero reference line, the
 dissolution of zones is the predominating reaction, resulting in
 a net decrease in resistivity.

(11) After some time, the increase in volume fraction of zones, due to their growth at 10°C, overrules the decrease in volume fraction of the zones dissolving, thus giving a net increase.

It is seen that the greatest decrease is shown by the 100 pct. case, in which both the factors, discussed earlier, contribute a decrease in resistivity whereas in cases with lesser preageing stay, dissolution of already formed zones takes lesser time, and the curves cross zero line after some time, which is quicker for 20 pct. case compared to the 40 pct. case.

Fig. 1.02 and 1.03:- These are explainable on the similar lines as for fig.1.01 with additional features as:

(1) Comparing the curves for fig. 1.01, 1.02, and 1.03, it is observed that the magnitude of decrease in resistivity, is highest for reageing at  $30^{\circ}$ C, obviously because of higher difference in  $(T_{\rm R} - T_{\rm A})$ , as we raise the reageing temperature from  $10^{\circ}$ C to  $30^{\circ}$ C. (11) Further, the crossing of the zero line for curves with lesser  $t_{PA}$ , is earlier, because of earlier superseeding of the factors contributing an increase in  $\ell$ . (i.e. growth of zones at  $T_R$ ). (iii) The peaks after the initial decrease, are observed earlier, with a rise in reageing temperature, because of increased diffusional properties with higher reageing temperatures.

Fig. 1.12 to 1.13:- Here again curves are almost similar in nature to those in fig. 1.02. In the case of fig. 1.13 even with 20 pct. case, the initial decrease followed by an increase has been eliminated, because of the fact that the degree of supersaturation is lesser for reageing at 30°C compared to, reageing at 20°C. Here the flatness of almost all the curves is Fig.1.23 :probably due to loss of a considerable fraction of vacancies during  $t_{PA}$ , in the alloys preaged at a higher temperature ope.j.20°C, compared to the case of preageing at 10°C or 0°C. Here the flatness of the curves is more likely due to concentration of vacancy effect, rather than the balance of two opposing factors contributing to  $\rho$  :  $_{\phi}$ It may be observed that there is no major Effect of T<sub>H</sub> 1change with respect to the nature of curves, due to variations in T<sub>H</sub>, if we compare the respective figures for reversion, with different  $T_H$  values e.g. comparing the figures 1.03, 2.03 and 3.03, it may be observed that :

(1) The rates of reageing are somewhat slower, in the cases,with lesser homogenization temperature. This effect is due to

the variation of concentration of quench-in vacancies.

(11) Because of the higher concentration of quench-in vacancies with higher homogenization temperatures, thus resulting in faster reaction, it is observed, that almost each stage of the curves becomes more prominent. This will be clear, if we compare the curves for 20 pct, preageing time, in the figures 1.03,2.03 and 3.03.

It is concluded, that the explanations for various stepageing and partial reversion of zones, employed in the present study, can be satisfactorily advanced in terms of the established principles enumerated in the beginning of the discussions.

Further, the discussions of the step-ageing processes in greater detail taking the case similar to fig. 1.10, makes clear all relevant features of such a process.

Similarly a detailed discussion of the partial reversion process, taking the case similar to fig.1.01, can be used for the understanding of the other reversion processes employed in the present study.

62.(0).

The apparant activational energies for the migration of vacancies E<sub>11</sub>, were calculated, from the data of fig.1.0, 2.0 and 3.0. The results are listed below:

Homogenization temperature	(THOC)	Value of En
300°C	đ	0.42 01.
350°C		0.40 ov.
400 <sup>0</sup> 0		0.35 eV.

The increase in the  $E_{\rm H}$  values, with increasing  $T_{\rm H}$ , can be explained on the basis of the effects of  $T_{\rm H}$ , on the concentration and role of vacancies. With increasing the homoginization temporature, the proportion of di-vacancies, which have a lower activational energy for migration, compared to that for a single vacancy increases. Hence, lesser proportion of single vacancies, with increasing homoginisation temperature is the probable cause for decrease in values of  $E_{\rm H}$ .

### 5.0 CONCLUSIONS

There are many aspects of the agoing mechanism which are confirmed by the present study.

The role of quench-in vacancies appears to be the most important factor in governing the various aspects of ageing. The much slover rates during re-ageing are explainable in terms of the prevailing concentration of vacancies, the diffusional controlled processes of migration of solute atoms, and annihilation of vacancies at the temperature of ageing or re-ageing.

It has been possible to explain the results of the present work on the cotablished understanding about the size of the sones, the dependence of their equilibrium number on the ageing temporature, and the observance of a peak in resistivity at a critical sone size.

Purther, the results have been explained on the proposed theory that there is a temperature dependance of the size of the cones with respect to their thermodynamic stability at this temperature.

The apparant, activational energy,  $E_{\rm H}$  for the migration of vacancies, was calculated to be, 0.42 eV, 0.40 eV and 0.35 eV for the homogenization temperatures, 300°C, 350°C and 400°C respectively.

#### 6.0 SUGGESTIONS FOR FURTHER WORK

Keeping in view the significance of the present work, in better understanding of the pre-precipitation stage of ageing, a further advanced work is suggested on the following guidelines:-1. Electrical resistivity studies should be supported by electron diffraction,X-ray intensity, and small angle X-ray diffraction studies for still better understanding of the **EXAMPLE EXAMPLE**, EXAMPLE EXAMPLE, size, shape, number and distribution of clusters and zones.

2. The effect of solute concentration can be studied.

- Similarly the effect of ternary additions (trace additions) can be undertaken.
- 4. The study can be extended to sub-zero temperatures, so as to make the processes slower, which can be more conveniently studied.
- 5. There is the unexplored aspect with respect to the determination of the activational energies of the reageing processes (step-ageing and reversion of zones) as has been hinted in Ref.(44).

- 1. Vilm, A., Notellurgie, 8, (1911) 225.
- 2. Vilm, A., Detallurgie, 8, (1911) 650.
- Merica, P.D., Valternberg, R.G., and Scott, H., Trans. Met. Soc.
   A.I.M.E., 64, (1921) 3.
- 4. Knno, S., Science Reports, Tohuku Imperial University, 11(1922) 269.
- 5. Franknch, V., and Schweizer.E., Z. Letallk, 14 (1922) 49.
- 6. Gaylor, M.L.V., and Preston, G.D., J. Inst. Metals, 41 (1929) 191.
- 7. Rosonhain, V., J. Inst. Dotals, 44 (1930) 226.
- 8. Tommann, G., Z. Letallk, 22, (1930) 365.
- 9. Derica, P.D., Trans. Not. Soc. of A.I.D.E., 99 (1932) 131.
- 10. Jonkins, C.H.H., and Buchmall, E.H., J. Inst. Metals 57 (1935) 141.
- 11. Do Sorbo, U., Treaftis, H.N. and Turnbull, D., Acta. Let. 6 (1958) 401.
- 12. Turnbull, D., Roconbaum, H.S., and Troaftic, H.N., Acta. Mot. 8 (1960) 277.
- 13. Guinior, A., Acta. Not. 3 (1955) 510.
- 14. Guinior, A., X-ray diffraction, V.H. Procean, San-Prencisco, 1963.
- 15. Gorold, V., Phys. Status Solidi, 1, (1951) 57.
- 16. Guinior, A., Ann. Phycique, 12, (1939) 161.
- 17. Proston, G.P., Proc. Royal Society A 161, (1938) 526.
- Borolius, G., Defects in Crystallin Solids, Physical Society p.169, (1955).
- 19. Dorolius, G., Larsson, L.E., The mochanism of phase transformation in metals., J. Instituto of Metals. p 285, (1956).
- 20. Horman, H., and Cohon, J.B., Haturo, 191, (1951) 63.
- 21. Haborkron, H., and Gorold, V., Phys. Status Solidi, 15,(1966) 167.
- 22. Hirsch, P.B., and Kolly, A., Phil. Mag. 12, (1965) 881.
- 23. Harknoso, S.D. and Hron, J.J., Trans. Met. Soc. A.I.M.E.

- 24. Habarro, P.R.H., Proc. Roy. Soc. A52 (1940) 90.
- 25. Habarro, P.R.H., 101d., A175 (1940) 519.
- 26. Mott, N.F., and Habarro, F.R.N., Proc. Roy. Soc. 52 (1940) 86.
- 27. Baker, R.G., Bradon, D.G., and Nutting, J., Phil. Mag.[8] 4 (1959) 1339.
- 28. Turnbull, D., Solid Stato Physics, 3, (1956) 226.
- 29. Zonor, C., Proc. of Int. Conf. on Physics of Metals. Amestordom(1948), p 117.
- 30. Soits, P., L'Etat Solide, Inst. Int. Phys. Solray, (1952) 401.
- 31. Podorighi, T., Acta. Met. 6, (1958) 379.
- 32. Lomor, V.M., Vacancies and other defects in Motals and Alloys J. Inst. of Metals, (1958) 79.
- 33. Pensari, C, and Pederighi, T., Phil. Mag. [8] 3 (1958) 1225.
- 34. Hart, E.V., Acta. Let. 6, (1958) 553.
- 35. Federighi, T., and Thomac, G., Phil. Mag. [8] 7 (1962) 127.
- 36. Gorold, V., and Schubisor, W., Z. Detallt 52, (1961) 76.
- 37. Walkor, C.B., and Guinier, A., Acta. Mot. 1, (1953) 568.
- 38. Garf, R., Compton Rendus 249 (1959) 1110.
- 39. Friedol, J., Les Dislocations, Gauthoir Villars, 1956.
- 40. Brooks, H., Detal intorfaces, Trans. Amer. Soc. of Dotals. (1952) 20.
- 41. Guinior, A., Acta. Cryst. 5 (1952) 51.
- 42. Bradshaw, P.S., and Pearcon, S., Phil. Hag. [8] 2 (1957) 94.
- 43. Nebb, L.B., Acta. Let. 7, (1959) 748.
- 44. Ponsari, C., and Federighi, R, Acta. Met. 8 (1950) 217.
- 45. Westmacott, K.H., Barnes, R.S., Hull, D., and Smallman, R.E., Phil. Mag.[8] 6 (1961) 929.
- 46. Lott, N.P., J.Inot. Metals, 60 (1937) 267.
- 47. Goiolor, A.H., Phase transformations in colids, (1951) 387.
- 48. Pino, U.E., Acta. Lot. 7 (1959) 228.
- 49. Labusch. Phys. Status. Solidi 3 (1963) 1661.
- 50. Uatyos, Z., Phil. Uag. [7] 40 (1949) 324.
- 51. Woodlos, ibid, [7] 40 (1950) 31.
- 52. Simorska and Synock, Acta. Elet. (1957) 223.
- 53. Fink, V.L., and Smith, D.V., Trans. A.I.U.E. 124 (1937) 162.
- 54. Garf, R. and Lonormond, Acad. Sci., Paris (1964) 3494.
- 55. Horts, Anantharcman, T.R., and Gorold, V., Phys. Status Solidi 8 (1955) 5.

- 56. Garwood, R.D., Devies, A.L., and Richards, G.L., J. Inst. Motals 88, (1959-60) 311 and 375.
- 57. Gorf, R., Journal Physics, Paris 22 (1962) 819.
- 58. Geisler, A.H., Barrtt, C.S., and Hohl, R.F., Trans. A.I.U.E. 152 (1943) 201.
- 59. Baur, R., and Gerold, V., Z. Metallk, 52 (1961) 671.
- 60. Baur, R., and Gorold, V., Acta. Let., 10 (1962) 637.
- 61. Carpenter, A., and Garwood, R. D., J. Inst. Metals 94, (1966) 301.
- 62. Larsson and Lasok, 1bid., 95 (1967) 320.
- 63. Nicholson, R.B., and Kelly, A., Progress in Metal Phys. VOL.10, Edited by Bruce Chalmors.
- 64. Dash, J., and Fino, H.E., Acta. Not. 9 (1961) 149.
- 65. Katz, K. I. and Herman, H., Mat. Sc. Engg. 1, (1966) 162.
- 66. Borclius, G., J. Inst. Lotals. 3 (1951) 477.
- 67. Johnson, A.A., Huges and Barton., ibid., 94 (1966) 186.
- 68. Strongin, B.G., Phys. Hot. Letallography, 23(2), (1967) 55.
- 69. Wohi, R. P., and Anantharaman, T. R., Current Science, 28, (1969) 1.
- 70. Rao, K., Herman, H., and Parche, Uat. Sc. Engg., 1, (1966) 162.
- 71. Ellwood's, A., J. Inst. Motals. 80, (1951) 217.
- 72. Cohen, J.B., Trans. A.I.W.E., 242 (1968), 166.
- 73. Rudman and Hillard, H, Acta. Met. 15, (1967) 1023.
- 74. Gaylor, M.L.V., Jr. Inst. Metals, 28, (1922) 213.
- 75. Dohlinger, V., and Knapp, H., Z. Cetallk, 43, (1952) 223.
- 76. Silcock, J. H., Hoal, and Hardy, J. Inst. Detals. 82, (1953)239.
- 77. Knobeevski, S.T., 1514., 69 (1943) 397.
- 78. Garf, R., 1bid., 86,(1957-58) 534.
- 79. Borelius, G., Vestensk, Acad. Handl, 169, (1943) 177.
- 80. Richards and Garwood, R.D., J. Inst. of Motals, (1965) 393.
- 81. Porry, P., Acta. Mot., (1966) 1143.
- 82. Asdonte, II., Acta. Not. (1961) 587.
- 85. Entwisttle, and Perry., J. Inst. Hetals (1966) 24.
- 84. Eurakani, Kawano and Eurakani, Acta. Hot. (1969) 29.
- 85. Smugosky, Herman and Pollack, 1bid., (1969) 883.
- 86. Vilkes, ibid., (1958) 863.
- 87. Larsen, Acta. Net. (1967) 35.
- 88. Horman, H., 1bid., (1967) 154.

Hormon, H., Cohon, J.B., ibid., (1963) 43 89. Poderighi, T., and Ponsari, C., Aota. Dot. (1959) 422. 90. Lurakani, Kveno, end Kurakani, Trans. A.I.M.E. (1969) 815. 91. 92. Herkons, S. D., Hron, J. J., Net. Trans, (1970) 43. Lipsit, H., Lotallurgical Transactions, (1971) 1739. 93. Ciach. R., Scripta Ectallurgica, (1970) 39. 94. 95. Ramon. R. S., ibid., (1970) 197. 96. Bohn.G., and Gerold.V., 1bid., (1970) 269. Chattorico.A. and Pabian.D.J., ibid., (1970) 285. 97. 93. Remchandran, V., and Srinivasan, N.K., 1914., (1967) 103. Remon, K.S., 1bid., (1971) 791. 99. K.Krishna Reo, Katz, V.I., and Herman, H., Mct. Sc. Engg., 28, 100. 1, (1966-67) 263. Reman, K.S., Das, E.S.D. and Vasu, K.I., J. Met. Sc. (1970) 5. 101. 102. Shee, S.K., Vchi, R.P., and Liera, S., J. Indian Institute of Letals, (1967) 148. 103. Rots, L. E., Rao, K.K., and Herman, H., 1bid., (1966) 95. 104. Vahi, R. P., and Ananthoraman. T. R., ibid., (1956) 41. Wahi, R. P., and Anantharaman, T. R., 1bid., (1971) 61. 105. 106. Vahi, R. P., Kutumba Rao, V.U. P., and Anantharaman, ibid., (1970) 20. 107. Carpentor and Garwood, R. D., Not. Sc., Engg., 1 (1967) 202. 108. Hardy and Heal, Report on precipitation, Prog. Metal Phys. V-5. 109. Smith. G.C., Age Hardening of Lictals, Prog. Lietal Phys. V-1. 110. Kimura, H. Hasiguti, R.R., Acta, Mot. 9 (1961) 1076. 111. Gould, R.V., J. Inst. Lietals, (1971), 1. 112. Murakani, Kawano, and Murakani, ibid., (1971) 160. 113. Rao, K., and Harman, H., ibid., (1966) 180, 420. 114. Garf.R., ibid., (1950) 282. 115. Price and Kolly, 1bid., (1950) 425. 116. Rao, and Herman, J. Inst. Metals, (1966), V-94, p320. 117. Herman, H., Lietallurgical Transactions, (1971), 16.

# 

- 2<sub>17</sub> Nonogenisation Temperature
- T<sub>A</sub> Ageing Temperature
- T<sub>RA</sub>- Re-agoing temperature for stop-ageing case
- T<sub>R</sub> Re-agoing temperaturo for revorsion case
- tp Timo required to reach peak in resistivity
- tpa- Pro-ageing tino
- E<sub>M</sub> Activational Energy for migration of vacancies
- E<sub>F</sub> Activational Energy for formation of vacancies
- D<sub>S</sub> Diffusion Coefficient of soluto atom in solvent
- $E_{B_{-}V}$  Interaction energy between solute and solvent
  - Atomic minimatch between solute and solvent atoms.
- Cy Concentration of vacancies .
  - $\gamma$  Interfacial energy between zones and matrix
  - Electrical resistivity of alloy
- $\begin{pmatrix} (T) & \text{Resistivity at instant t after quenching} \\ this agains at TC$ 
  - G. modulus ] rigi duty