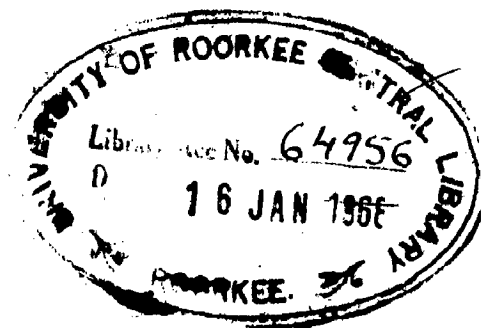


Effect of Boundary Conditions on Uplift Pressure and Exit Gradients Due to Sub-Surface Flow Under Hydraulic Structures

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certified that this dissertation, entitled "Effect of boundary conditions on spatial growth of cell populations due to substrate gradients and its structural organization" in a report of the last year was done by me. It is hereby certified that the course of study of engineering, (civil engineering development) in the academic session 1965-66 and 1966-67. It is hereby certified that our guidance on this dissertation was from October 10, 1966 to October 10, 1967.

To the best of our knowledge this dissertation has not been published anywhere nor submitted for the award of any other degree.

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CHAPTER ONE

CHAPTER I

INTRODUCTION

1.1 HISTORICAL BACKGROUND :

A rational basis to the study of subsoil flow was, for the first time, given in 1856 when Darcy (8) gave his law governing ground water flow in alluvial and sedimentary formations. On the basis of Darcy's work, significant contribution to the subsoil hydraulics were made by Dupuit (9), Forchheimer (10) and Thoin (11) in the later half of the nineteenth century and more recently by Doolittle (7), Leach (12), Horton (13), King and Slichter (14)

In case of hydraulic structure founded on pervious soil, a major force of uplift comes in. One of the first basis for the design of hydraulic structure built on pervious foundations was developed in India, as a result of investigations of Clibborn and Croxford (15). As a result of experiments by Clibborn the hydraulic gradient theory of design of weirs was developed by Cutler and Hagen.

A considerable amount of research on the flow of water through soils, with special regard to the Terzaghi's constant (K) was carried out in 1889 by Slichter and King (16). The main object of this investigation, was to determine the amount of water percolating through various types of soils with a view to ascertain the extent and behaviour of ground waters of United States. It did not deal with the distribution of pressures across structures built on pervious soils.

In 1910 Eligh in his book on "Practical Design of Irrigation Works" enunciated the "Hydraulic Gradient Theory" which later came to be known as the "Eligh Creep Theory". Because of its simplicity, this theory found general acceptance. By a detailed analysis of failure of structures designed by Eligh Theory, Lane (10) introduced his empirical "weighted creep theory".

Pavlovsky, in 1929, introduced the technique of electrical analogy in the investigations of problems of flow of sub-soil water below hydraulic structures.

In 1929 Forzaghi (27) stated that failures occurred by undermining if the hydraulic gradient at exit was in excess of what he called the "Flotation gradient". The concept of flow net was enunciated by Forchheimer (10) and Forzaghi (27). A mathematical treatment of the flow of water through permeable subsoil under dams was evolved by Weaver (29). In 1938 Khosla Doo and Taylor (16) published the method of "Independent variables" which gave a general solution to problem of uplift over various base profiles. These findings were supported by an exhaustive series of laboratory and prototype observations.

1.3 LAPLACE EQUATION AND ITS SOLUTION

Flow can be laminar or turbulent depending upon the Reynolds number of flow. Normally for consolidated sand, the maximum size of the pore spaces is approximately 0.08 mm (28), and the velocity 0.1 cm/sec. and 0.15 cm/sec., the Reynolds number ranges from 0.8 to 0.70. For Reynolds number less than

one it can be assumed that flow is laminar and most of the flows under structures founded on sand fall in this range.

It is well known (17, 18, 21) that flow according to Darcy's Law satisfies the Laplace Equation.

$$V_x = -K \frac{\partial \phi}{\partial x} \quad \dots\dots (1.1)$$

$$V_y = -K \frac{\partial \phi}{\partial y} \quad \dots\dots (1.2)$$

and
$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0 \quad \dots\dots (1.3)$$

where

V_x = Velocity in x direction

V_y = velocity in y -direction

ϕ = potential function

ψ = stream function

K = coefficient of permeability

Laplace Equation is satisfied by the conjugate harmonic function ϕ and ψ and that $\phi(x, y) = \text{constant}$ and $\psi(x, y) = \text{constant}$ are orthogonal curves.

It is seen that a variety of real solution to the Laplace equation can be obtained by taking various functions of z , by separating out real and imaginary parts ϕ and ψ . Which solution to apply in a particular problem depends upon the boundary condition which that particular solution satisfies.

1.3 DIFFERENT METHODS FOR SOLUTION OF LAPLACE'S EQUATION

Following methods are commonly employed for the solution of Laplace's equation.

- (i) Graphical method
- (ii) Hydraulic models
- (iii) Electrical analogy models
- (iv) Analytical methods.

1.3.1 The method of graphical field plotting is a widely used method, more popularly known as "Forschheimer solution". This method is confined to two dimensional cases and is capable of considerable accuracy. The process is begun on a scale drawing of boundaries with assigned potentials by plotting any suitable number of intermediate equipotential lines. Flux lines or stream lines are then drawn to cut the equipotential lines orthogonally and to form curvilinear squares. The errors are corrected by systematic improvements as suggested by Taylor(31). Much experience and foresight are needed for correct representation of flow not by the trial and error method. The greatest advantage of this method is that it requires little equipment and gives quick results. It gives clear idea about the subsoil flow of water.

1.3.2 HYDRAULIC Models may be

- (a) Viscous fluid models
- (b) Sand models

Viscous fluid models are based on the principle that flow through two parallel plates placed at a very small distance apart follows the Laplace Equation. Such a model consists of a small tank constructed with two parallel sides of glass plates-placed a small distance apart. The model is fitted in and the fluid (usually water or glycerine) is made to flow under the model from the upstream to downstream side. When

the flow becomes steady the stream lines can be traced by adding a coloured solution of the same fluid at points and marking or photographing the path of such lines. Then equipotential lines are drawn orthogonal to the stream lines.

In hydraulic scale model sand is filled between two parallel plates. The pressures are observed by means of piezometric pipes introduced at proper points. An ingenious method of tracing the stream lines was followed in Punjab. The flow lines were made visible through sand by injecting potassium dichromate and silver nitrate separately. As a result of chemical reaction red precipitate of silver chromate was produced along the stream lines which became clearly visible.

1.3.3 In the electrical analogy method, the analogy developed between the flow of fluid through porous media which follows Darcy's Law and flow of electricity through an electrolyte which follows Ohm's Law is used. Two dimensional as well as three dimensional problems can be solved by this technique.

The chief advantages of this method are that the laboratory equipment is cheap and handy and flow problems can be analysed speedily and with reasonable degree of accuracy. Many problems which might have been difficult to solve, find an easy solution in electrical analogy method. Models of such smaller scales can be used with electrical analogy set up while capillary forces will vitiate the results if the scale is reduced very much in hydraulic models.

This amongst the various experimental method electrical analogy is the most accurate and quick.

1.3.4 Analytical Method.

Analytical method aims at the solution of Laplace equation mathematically which can be done by help of conformal transformations and conjugate functions. In this method boundary conditions of the problem are expressed by equations and solution obtained mathematically. The method gives exact solution of the problem, although the method becomes involved for complicated boundary conditions. Classical work in this field has been done by A.N. Kholod (16), W. Noor (20), N.N. Pavlovsky (22, 23), N.S. Shklovsky (29). The present work is an extension of Kholod's approach and as a brief review of Kholod's method will be performed.

1.4 KHOLDA'S METHOD

If the Laplace equation could be integrated for a given set of boundary conditions, the mathematical solution of the flow net could be obtained for these conditions. This equation is amenable to mathematical solution for simple boundary condition but becomes too involved for complex boundary conditions. With the help of the method of "Independent variable", evolved by Kholod, an approximate but sufficiently precise results can be arrived at by splitting a complicated four-point profile into several elementary forms. Mathematical solutions for a number of such elementary forms were obtained by Kholod. The basic idea underlying this solution is that

the effect of every additional sheet piling is concentrated mainly in its immediate neighbourhood and dies out gradually as the distance from the pile increases. At a certain given distance from the pile its effect may therefore be estimated with enough accuracy by an empirical formula.

1.5 Limitations of Khosla's pressure and Exit Gradient Diagrams.

The formulae and graphs given by Khosla are applicable for boundary condition when length of pervious reaches on the upstream and downstream are infinite and depth of pervious strata below the foundation of structure is also infinite. In lined canals the boundary conditions are different in that the canal bed is impervious except for downstream inverted filter length. If the lining is effective on the upstream of the structure no water can seep below its foundation and the structure will not be subjected to any uplift. The development of pressure depends upon the extent of assumed pervious bed or width of crack in the lining. The release of pressures downstream of hydraulic structures on lined canals can only be effected through a predetermined filter length. In addition to find cases where depth of pervious strata below the foundations of structure is limited. The formulae and curves given by Khosla are not applicable for such boundaries. Exact solutions for such boundaries are given in the subsequent chapters.

CHAPTER TWO

CHAPTER II

MATHEMATICAL ANALYSIS FOR DIFFERENT BOUNDARY CONDITIONS

2.1 STATEMENT OF THE PROBLEM

If a structure founded on permeable foundation is considered, the flow through porous medium satisfies the two dimensional Laplace equation

$$\frac{\partial^2 H}{\partial x^2} + \frac{\partial^2 H}{\partial y^2} = 0 \quad \dots (2.01)$$

where

H = the difference in upstream and downstream water levels.

In the present study two cases A and B with boundary conditions as described below have been considered for analysis.

A. FINITE PERVIOUS REACHES ON THE UPSTREAM AND DOWNSTREAM OF THE STRUCTURE FOUNDED ON INFINITE PERVIOUS FLOW ZONE (FIGURE 2.1a).

- (i) Impervious floor length, $b_1 + b_2$.
- (ii) Cut off depth below canal bed, d .
- (iii) Length of upstream pervious bed, f_1 .
- (iv) Length of downstream pervious bed or filter length, f_2 .
- (v) Impervious bed HP and GH extending upto infinity on either side.
- (vi) Depth of permeable canal extends upto infinity.

Referring to figure 2.1a, the foundation profile ABCD forms the lower boundary and represents stream line $\psi = 0$. The extreme impervious boundary FE and HC (both extending upto infinity) represents stream line $\psi = q$ where q is the total discharge per unit width infiltrating through the foundation. The other two boundaries are AD at the upstream bed level along which potential $\phi = H_0$, representing the total head and at the downstream end BE, along which potential $\phi = 0$.

B. GENERAL PERVIOUS REACHES ON THE UPSTREAM AND DOWNSTREAM OF THE STRUCTURE FOUNDED ON FINITE PERVIOUS FLOW ZONE (FIGURE 2.2a).

- (1) Impervious floor length, $b_1 + b_2$
- (11) Cut off depth below canal bed, d .
- (111) Length of upstream pervious reach, x_1 .
- (12) Length of downstream pervious reach, x_2 .
- (v) Impervious bed FE and HC extending upto infinity on either side.
- (vi) Depth of permeable subsoil, Z .
- (vii) Impermeable layer of a depth T , extending upto infinity on either side.

Referring to figure 2.2a, the foundation profile ABCD forms the lower boundary and represents stream line $\psi = 0$. The extreme impermeable boundary FEHC represents stream line $\psi = q$, where q is the total discharge per unit width infiltrating through the foundation. The other two boundaries are AD at the upstream bed level along which potential $\phi = H_0$, representing the total head and at the downstream end BE, along which potential $\phi = 0$.

Further for the purpose of making the procedure simple, it is assumed that the coefficient of permeation $k = 1$ and permeable strata is homogeneous.

2.2 METHOD OF SOLUTION

Referring to figures 2.1 and 2.2 both the true profile of the hydraulic structure and the rectangular flow field represented in w -plane are transformed on to the conformal circle of infinite radius $z = \infty$ on the z -plane. The ζ -plane covers the limit joining the two phases of analysis z -plane is transformed into ζ -plane through an auxiliary w -plane.

Using Schwarz-Christoffel transformation, the area in z -plane (see figure 2.1a) bounded by the boundaries mentioned above is transformed to the lower half of the w -plane which in its turn ^{is} transformed into the lower half of the ζ -plane. Thus

$$\begin{aligned} z &= z(w) \\ &= z(\zeta) \end{aligned} \quad \dots\dots(2.02)$$

where

$$z = x + iy, w = u + iv \quad \text{and} \quad \zeta = \xi + i\eta$$

gives the solution of Laplace equation for the z -plane.

The second operation in this process consists in transforming the rectangular flow field ⁱin the w -plane on to the same conformal-infinite plane, giving the relation

$$w = F(\zeta) \quad \dots\dots(2.03)$$

where

$$w = \phi + i\psi$$

and

$$\zeta = F^{-1}(w)$$

on combining the two operations (2.02) and (2.03) into one, we obtain

$$z = f(\zeta) = f \cdot f^{-1}(w) \quad \dots (2.04)$$

and $w = F(\zeta) = F \cdot f^{-1}(z)$

The relation (2.04) holds good for curves defined by their equations, e.g. equipotentials and streamlines. Also for the arrangement of the rectangular flow field shown in Figures 2.1 d and 2.2 d, the potential function is linearly correlated with the path of filtration.

2.3 ANALYTICAL SOLUTION FOR THE BOUNDARY CONDITION 'A'

2.3.1 ANALYTICAL SOLUTION

(A) First Operation

Referring to Figure 2.1 the profile of the hydraulic structure in z-plane, is transformed on to the real axis of the t-plane with

$z = z_A = -b_1$	$t = t_A = -L_1$
$z = z_B = +b_2$	$t = t_B = +L_2$
$z = z_C = +0$	$t = t_C = +0$
$z = z_D = id$	$t = t_D = 0 \quad \dots (2.05)$
$z = z_E = -0$	$t = t_E = -0$
$z = z_F = -(b_1 + \delta_1)$	$t = t_F = -\gamma_1$
$z = z_G = -(b_2 + \delta_2)$	$t = t_G = +\gamma_2$

Schwarz-Christoffel transformation that maps z-plane on to the lower half of t-plane is given by

$$\frac{dz}{dt} = A(t-t_1)^{\frac{1}{2}-1} (t-t_2)^{\frac{1}{2}-1} (t-t_3)^{\frac{1}{2}-1} (t-t_4)^{\frac{1}{2}-1}$$

In this case

$$\frac{dz}{dt} = At (t-\sigma)^{-\frac{1}{2}} (t+\sigma)^{-\frac{1}{2}} \dots\dots(2.06)$$

On integration

$$z = A \int_0^t \frac{k \cdot dt}{\sqrt{t^2 - \sigma^2}} + B.$$

$$z = A \sqrt{t^2 - \sigma^2} + B$$

Constants A and B can be calculated with the help of corresponding known points in Z and t-planes.

At C and E

$$z = 0 \quad \text{and} \quad t = \pm \sigma$$

$$B = 0$$

and at D,

$$z = id \quad \text{and} \quad t = 0$$

$$id = A \sqrt{-\sigma^2} = A i \sigma$$

$$A = \frac{d}{\sigma}$$

Therefore the transformation equation becomes

$$z = \frac{d}{\sigma} \sqrt{t^2 - \sigma^2}$$

$$\text{or } t = \pm \sigma \sqrt{1 + \left(\frac{z}{d}\right)^2} \dots\dots(2.07)$$

This gives

$$L_1 = \sigma \sqrt{1 + \left(\frac{b_1}{d}\right)^2}$$

$$\begin{aligned} \text{or } \zeta_0 &= \sigma \sqrt{1 + \left(\frac{b_1}{d}\right)^2} \\ \zeta_1 &= \sigma \sqrt{1 + \left(\frac{b_1 + \delta_1}{d}\right)^2} \dots\dots (2.09) \\ \zeta_2 &= \sigma \sqrt{1 + \left(\frac{b_2 + \delta_2}{d}\right)^2} \end{aligned}$$

These operation can be made easy by transforming ζ -plane on to ξ -plane such that F and G are placed symmetrically along real axis at -1 and $+1$ respectively. Therefore ζ -plane is transformed on to ξ -plane such that

$$\begin{aligned} \zeta_A &= \beta, \\ \zeta_B &= \delta \beta, \dots\dots (2.00) \\ \zeta_F &= -1 \\ \zeta_G &= +1 \end{aligned}$$

Transformation of ζ -plane to ξ -plane is given by

$$\xi = \Delta \zeta + D'$$

To find the constants consider the values of corresponding coordinates at F and G.

$$\begin{aligned} \text{at F, } \xi &= \gamma_1, \zeta = -1 \\ \text{at G, } \xi &= \gamma_2, \zeta = +1 \\ \therefore \gamma_1 &= D' - \Delta' \\ \therefore \gamma_2 &= D' + \Delta' \end{aligned}$$

Solving these equations and substituting the values of Δ' and D' , the transformation equation becomes

$$\xi = \frac{\gamma_1 + \gamma_2}{2} \zeta + \frac{\gamma_2 - \gamma_1}{2}$$

$$\text{or } \zeta = \left[\xi - \frac{\gamma_2 - \gamma_1}{2} \right] \frac{2}{\gamma_1 + \gamma_2} \dots\dots (2.10)$$

Substituting value of t from equation (2.09),
equation (2.10) becomes

$$\zeta = \left[\sigma \sqrt{1 + \left(\frac{z}{a}\right)^2} - \frac{\gamma_2 - \gamma_1}{2} \right] \frac{2}{\gamma_1 + \gamma_2}$$

Let $\frac{\gamma_1}{\sigma} = \Delta_1 = \sqrt{1 + \left(\frac{b_1 + f_1}{a}\right)^2}$

$$\frac{\gamma_2}{\sigma} = \Delta_2 = \sqrt{1 + \left(\frac{b_2 + f_2}{a}\right)^2}$$

$$\frac{L_1}{a} = \lambda_1 = \sqrt{1 + \left(\frac{b_1}{a}\right)^2} \quad (2.11)$$

$$\frac{L_2}{a} = \lambda_2 = \sqrt{1 + \left(\frac{b_2}{a}\right)^2}$$

$$\frac{t}{a} = \lambda = \sqrt{1 + \left(\frac{z}{a}\right)^2}$$

Substituting equation (2.11) in equation (2.10),

$$\zeta = \left[\lambda - \frac{\Delta_2 - \Delta_1}{2} \right] \frac{2}{\Delta_1 + \Delta_2}$$

$$= 1 - \frac{2(\Delta_2 - \lambda)}{\Delta_1 + \Delta_2}$$

$$\beta_1 = 1 - \frac{2(\Delta_1 - \lambda_1)}{\Delta_1 + \Delta_2} \quad (2.12)$$

and $\beta_2 = 1 - \frac{2(\Delta_2 - \lambda_2)}{\Delta_1 + \Delta_2}$

are obtained

(ID) Second Operation :

Referring to figure 2.1, the rectangular flow field can be transformed on to the same real axis of the ζ -plane noting that

$$\zeta_G > \zeta_B > \zeta_A > \zeta_F$$

Transformation formula is given by

$$\frac{dw}{d\zeta} = \frac{C}{\sqrt{(\zeta - \zeta_G)(\zeta - \zeta_B)(\zeta - \zeta_A)(\zeta - \zeta_F)}}$$

on integration

$$w = C \int_0^{\zeta} \frac{d\zeta}{\sqrt{(\zeta - \zeta_G)(\zeta - \zeta_B)(\zeta - \zeta_A)(\zeta - \zeta_F)}} + D \quad \dots (2.13)$$

Since the uplift pressures on the base of the structure and exit gradient at the downstream end of structure are to be evaluated the integration process is therefore confined to this zone only.

The zone for which the uplift pressures are to be determined corresponds to AB in ζ -plane. Pertaining to this zone

$$\zeta_G > \zeta_B > \zeta > \zeta_A > \zeta_F$$

and the integral equation becomes

$$w = C \int_0^{\zeta} \frac{d\zeta}{\sqrt{(-\zeta + \zeta_G)(-\zeta + \zeta_B)(\zeta - \zeta_A)(\zeta - \zeta_F)}} + D \quad \dots (2.14)$$

Now, substituting the characteristic coordinates of ζ -plane given by equation (2.17) in equation (2.10b) the corresponding values of U are determined. Thus for

$$\zeta = \zeta_B \quad (\text{at } B)$$

equation (2.10b) reduces to (Appendix I equations 8a, 8b, 8c, and 8d)

$$\operatorname{Sn}^2 U = 0$$

or $U = 0$

for $\zeta = \zeta_G \quad (\text{at } G)$

$$\operatorname{Sn}^2 U = \infty$$

or $U = iK$

for $\zeta = \zeta_A \quad (\text{at } A)$

$$\operatorname{Sn}^2 U = 1$$

$$U = K$$

and for $\zeta = \zeta_F \quad (\text{at } F)$

$$\operatorname{Sn} U = \sqrt{\frac{(\zeta_G - \zeta_A)(\zeta_F - \zeta_B)}{(\zeta_B - \zeta_A)(\zeta_F - \zeta_G)}} = 1/k$$

or $U = k + iK'$

Thus the four coordinates of the W -plane are

$$\zeta = \zeta_G \quad U = iK'$$

$$\zeta = \zeta_B \quad U = 0$$

$$\zeta = \zeta_A \quad U = K$$

$$\zeta = \zeta_F \quad U = K + iK'$$

..... (2.10)

in which K and K' are complete elliptic integrals of the first kind with modulus k and k' respectively. The

which on integration, using a bilinear transformation, yields (refer to appendix II)

$$\operatorname{sn}^2 \frac{\mu}{c} W = \frac{(\xi_G - \xi_A)(\xi - \xi_B)}{(\xi_B - \xi_A)(\xi - \xi_G)} \quad (2.15)$$

where

$$\mu = \frac{1}{2} \sqrt{\frac{(\xi_B - \xi_F)(\xi_G - \xi_A)}{1}}$$

and the modulus, k is given by

$$k = \sqrt{\frac{(\xi_B - \xi_A)(\xi_G - \xi_F)}{(\xi_G - \xi_A)(\xi_B - \xi_F)}}$$

from which the co-modulus defined by

$$k' = \sqrt{1 - k^2} = \sqrt{\frac{(\xi_A - \xi_F)(\xi_G - \xi_B)}{(\xi_G - \xi_A)(\xi_B - \xi_F)}}$$

and from equation (2.10) it follows that

$$\xi = \frac{\xi_B(\xi_G - \xi_A) - \xi_G(\xi_B - \xi_A) \operatorname{sn}^2 \frac{\mu}{c} W}{(\xi_G - \xi_A) - (\xi_B - \xi_A) \operatorname{sn}^2 \frac{\mu}{c} W} \quad (2.15b)$$

Further to simplify the procedure, the constant c is assumed equal to μ . This will yield

$$\xi = \frac{\xi_B(\xi_G - \xi_A) - \xi_G(\xi_B - \xi_A) \operatorname{sn}^2 W}{(\xi_G - \xi_A) - (\xi_B - \xi_A) \operatorname{sn}^2 W} \quad (2.16)$$

$$\text{and } \operatorname{sn}^2 W = \frac{(\xi_G - \xi_A)(\xi - \xi_B)}{(\xi_B - \xi_A)(\xi - \xi_G)} \quad (2.16b)$$

Referring to Figure 2.14 the characteristic coordinates of the transformation are

$$\begin{aligned} \xi &= \xi_G = +1 \\ \xi &= \xi_B = +\beta_2 \\ \xi &= \xi_A = -\beta_1 \\ \xi &= \xi_F = -1 \end{aligned} \quad (2.17)$$

four values of U thus calculated determine the layout of the basic rectangular flow field for the second operation.

The values of characteristic coordinates given by the equation (2.17) when substituted in k and k' yield

$$k = \sqrt{\frac{\alpha(\beta_1 + \beta_2)}{(1 + \beta_1)(1 + \beta_2)}} \quad \dots (2.18)$$

$$k' = \sqrt{\frac{(1 - \beta_1)(1 - \beta_2)}{(1 + \beta_1)(1 + \beta_2)}}$$

Similarly the substitution of characteristic coordinate values given in equation (2.17) in equation (2.10b) yield

$$\begin{aligned} S_n^2 W &= \frac{(1 + \beta_1)(\xi - \beta_2)}{(\beta_1 + \beta_2)(\xi - 1)} \\ &= \frac{(1 + \beta_1)(\beta_2 - \xi)}{(\beta_1 + \beta_2)(1 - \xi)} \end{aligned}$$

$$\text{or } S_n W = \sqrt{\frac{(1 + \beta_1)(\beta_2 - \xi)}{(\beta_1 + \beta_2)(1 - \xi)}} \quad (2.20)$$

2.3.2 EVALUATION OF UPLIFT PRESSURES

Along the bottom line of the foundation

$$\begin{aligned} \psi' &= 0 \\ \text{hence } u &= \phi' \end{aligned}$$

therefore equation (2.20) is reduced to

$$S_n \phi' = \sqrt{\frac{(1+\beta_1)(\beta_2-\zeta)}{(\beta_1+\beta_2)(1-\zeta)}} \quad \dots(2.21)$$

The value of β_1 and β_2 are given by equation (2.12). The pressure decreases from maximum head, H at point A in the beginning to 0 at point B in the end. The length of line ABCD in z -plane equals to complete integral of first kind(K) having its modulus equal to k (figure 4).

Potential ϕ for $H = 100$, is found from the relation

$$\phi = \frac{H}{K} \pi \quad \text{or} \quad 100$$

To evaluate the value of potential or pressure at any point along the floor the value of ζ is found from equation (2.12) corresponding to its location in z -plane. Substituting the value of β_1 , β_2 and ζ in equation (2.10) and (2.21), the values of modulus and $S_n \phi'$ are found. Using tables for Jacobian functions (20, 14) the required potential corresponding to any point in z -plane can be evaluated.

2.3.3 DYE GRADIENT

In addition to the characteristic pressure intensities investigated in the last paragraph, it is also important to know the hydraulic gradients at the down stream end of the first percolation trajectory. According to the principle originated by Forzaghi in 1920 and discovered independently by Khoulie a few years later, the stability of the granular soil depended on the limiting value of the hydraulic gradient at the upper surface of the granular

material. Khosla describes this limiting gradient as the
crit gradient. According to Forzaghi's theory of 1928
failure by piping could occur in two different manners

(a) Starting with the lifting and removal of the top
particles of soil from the surface of the carbon bed of the
downstream channel.

(b) Originating in the deeper layers of the granular
material in the case of very steep pressure gradients of the
water filtering through these lower layers were to be
responsible for the local concentration of the uplift force,
capable of lifting the supposed buoyant, weight of the
upper strata.

In so far as failure of type (a) is concerned
Khosla's crit gradient furnishes indeed the significant
design criterion, but in order to ensure safety in regard
to failure of type (b) gradients in the deeper soil levels
should also be investigated. Also however, type (b) failure
is supposed to be a less frequent occurrence (10) there is
enough reason to believe that Khosla's analysis of crit
gradient covers the main part of the problem

Gradient at any point is given by

$$I = \frac{dh}{ds}$$

where h = head at any point along floor or
cut off.

Therefore $I = \frac{1}{L} \cdot \frac{d\sigma^2}{d\theta} \dots\dots(2.28)$

where K = complete Elliptic Integral of first kind with modulus k

It is not possible to find out $\frac{d\sigma^2}{d\theta}$ directly

$$\frac{d\sigma^2}{d\theta} = \frac{d\sigma^2}{d\zeta} \cdot \frac{d\zeta}{d\theta} = \frac{d\sigma^2}{d\theta} \dots\dots(2.29)$$

Now $z = r + iy$

$$dz = dx + i \frac{dy}{\cos \theta}$$

$$\frac{dz}{d\theta} = \frac{dx}{d\theta} + i \frac{dy}{d\theta}$$

$$= \cos \theta + i \sin \theta$$

where θ = angle of stream line with horizontal

at the end

$$\theta = 90^\circ$$

and therefore

$$\frac{dz}{d\theta} = i \dots\dots(2.30)$$

Again $\frac{d\zeta}{dz} = \frac{d\zeta}{d\theta} \cdot \frac{d\theta}{dz}$

From equation (2.30)

$$\frac{d\zeta}{d\theta} = \frac{\sqrt{\Delta_0 \sigma^2}}{\Delta_0}$$

Substituting the value of constant Δ_0 to get

$$\frac{d\zeta}{d\theta} = \frac{\sigma}{d\theta} \sqrt{c^2 - \sigma^2} \dots\dots(2.31)$$

Differentiating Equation (2.10)

$$\frac{d\zeta}{d\theta} = \frac{2}{\gamma_1 + \gamma_2} \dots\dots(2.32)$$

Combining equations (2.25) and (2.26)

$$\frac{d\zeta}{dz} = \frac{2\sigma\sqrt{t^2 - \sigma^2}}{(\zeta_1 + \zeta_2) d.t.}$$

$$= \frac{2}{\Delta_1 + \Delta_2} \cdot \frac{1}{d} \cdot \frac{1}{t/\sigma} \cdot \sqrt{\left(\frac{t}{\sigma}\right)^2 - 1}$$

where Δ_1 and Δ_2 are defined in Equation (2.11)

$$\frac{1}{\sigma} = \sqrt{1 + \left(\frac{z}{d}\right)^2} = \lambda$$

$$\frac{d\zeta}{dz} = \frac{2}{\Delta_1 + \Delta_2} \cdot \frac{1}{d} \cdot \frac{\sqrt{\lambda^2 - 1}}{\lambda} \quad \dots\dots (2.27)$$

$$W = \phi' + i\psi'$$

along the foundation profile $\psi' = 0$

$$W = \phi'$$

$$\frac{dW}{d\zeta} = \frac{d\phi'}{d\zeta}$$

where

$$= \frac{C}{\sqrt{(\zeta - \zeta_G)(\zeta - \zeta_B)(\zeta - \zeta_A)(\zeta - \zeta_F)}}$$

$$C = \frac{1}{2} \sqrt{(\zeta_B - \zeta_F)(\zeta_G - \zeta_A)}$$

$$\frac{d\phi'}{d\zeta} = \frac{1}{2} \sqrt{\frac{(\zeta_B - \zeta_F)(\zeta_G - \zeta_A)}{(\zeta - \zeta_G)(\zeta - \zeta_B)(\zeta - \zeta_A)(\zeta - \zeta_F)}}$$

Substituting characteristic coordinates of transformation.

$$\frac{d\phi'}{d\zeta} = \frac{1}{2} \sqrt{\frac{(1 + \beta_1)(1 + \beta_2)}{(\zeta^2 - 1)(\zeta + \beta_1)(\zeta - \beta_2)}}$$

..... (2.28)

From equation (2.12)

$$\begin{aligned}
 1 + \beta_1 &= \frac{2(\Delta_2 + \lambda_1)}{\Delta_1 + \Delta_2} \\
 1 - \beta_2 &= \frac{2(\Delta_1 + \lambda_2)}{\Delta_1 + \Delta_2} \\
 \zeta - 1 &= \frac{2(\lambda - \Delta_2)}{\Delta_1 + \Delta_2} \\
 \zeta + 1 &= \frac{2(\Delta_1 + \lambda)}{\Delta_1 + \Delta_2} \\
 \zeta + \beta_1 &= \frac{2(\lambda + \lambda_1)}{\Delta_1 + \Delta_2} \\
 \zeta - \beta_2 &= \frac{2(\lambda - \lambda_2)}{\Delta_1 + \Delta_2}
 \end{aligned} \tag{2.29}$$

Substituting equation (2.29) in equation (2.28) and simplifying

$$\frac{d\phi'}{d\zeta} = \frac{\Delta_1 + \Delta_2}{4} \sqrt{\frac{(\Delta_2 + \lambda_1)(\Delta_1 + \lambda_2)}{(\lambda - \Delta_2)(\lambda + \Delta_1)(\lambda + \lambda_1)(\lambda - \lambda_2)}} \tag{2.29a}$$

Combining equations (2.24), (2.27) and (2.29a) and substituting in (2.23), gives after simplification

$$\text{Exit gradient} = G_E$$

$$\frac{H_i}{2kd\lambda} \sqrt{\frac{(\Delta_2 + \lambda_1)(\Delta_2 + \lambda_2)(\lambda - 1)(\lambda + 1)}{(\lambda - \Delta_2)(\lambda + \Delta_1)(\lambda + \lambda_1)(\lambda - \lambda_2)}}$$

With sheet pile at the end of floor, at exit

$$\lambda = 1.0$$

$$\text{and } \lambda_2 = 1.0$$

$$G_E = \frac{H_i}{2kd\lambda} \sqrt{\frac{(\Delta_2 + \lambda_1)(\Delta_1 + 1)(\lambda - 1)(\lambda + 1)}{(\lambda - \Delta_2)(\Delta_1 + 1)(\lambda + \lambda_1)(\lambda - 1)}}$$

On simplification this equation reduces to

$$C_2 = \frac{\pi}{2 K \Delta \lambda} \sqrt{\frac{2(\Delta_2 + \lambda_1)}{(\Delta_2 - \lambda)(1 + \lambda_1)}} = \frac{H}{\Delta K} \sqrt{\frac{\Delta_2 + \lambda_1}{2(\Delta_2 - 1)(1 + \lambda_1)}} \quad (2.30)$$

where

$$\lambda_1 = \sqrt{1 + \left(\frac{b_1}{d}\right)^2}$$

$$\Delta_2 = \sqrt{1 + \left(\frac{b_2 + f_2}{d}\right)^2}$$

d = depth of cut off

K = complete elliptic integral of first kind for modulus $k = \sqrt{\frac{\beta_1 + \beta_2}{(1 + \beta_1)(1 + \beta_2)}}$

H = total head in feet of water.

2.4 ANALYTICAL SOLUTION FOR THE BOUNDARY CONDITION 'B'

2.4.1 ANALYTICAL SOLUTION

(D) First Operation

Referring to figure 2.2 the profile of hydraulic structure in Z -plane is transferred on to the real σ axis of the ζ -plane with

$Z = Z_A = -b_1$	$\zeta = \zeta_A = -l_1$	
$Z = Z_B = +b_2$	$\zeta = \zeta_B = +l_2$	
$Z = Z_C = +0$	$\zeta = \zeta_C = +\sigma$	
$Z = Z_D = id$	$\zeta = \zeta_D = 0$... (2.31)
$Z = Z_E = -0$	$\zeta = \zeta_E = -\sigma$	
$Z = Z_F = -(b_1 + f_1)$	$\zeta = \zeta_F = -\gamma_1$	
$Z = Z_G = +(b_2 + f_2)$	$\zeta = \zeta_G = +\gamma_2$	
$Z = Z_H = -\infty$	$\zeta = \zeta_H = -1$	
$Z = Z_I = +\infty$	$\zeta = \zeta_I = +1$	
$Z = Z_J = i\tau$	$\zeta = \zeta_J = \pm \alpha$	

Chwarz.christeffel transformation that maps Z-plane on the lower half of t-plane is given by

$$\frac{dz}{dt} = A(t - t_1)^{\frac{1}{\alpha}-1} (t - t_2)^{\frac{1}{\beta}-1} (t - t_3)^{\frac{1}{\gamma}-1} (t - t_4)^{\frac{1}{\delta}-1}$$

In this case

$$\frac{dz}{dt} = A(\sigma - t)^{-\frac{1}{2}} (\sigma + t)^{-\frac{1}{2}} (1 - t^2)^{-1} (1 + t)^{-1} (t) \dots (2.32)$$

On integration

$$Z = A \int_0^t \frac{t \cdot dt}{(1-t^2)\sqrt{\sigma^2-t^2}} + B.$$

Let

$$\sigma^2 - t^2 = \lambda^2$$

which gives $-t \cdot dt = 2\lambda \cdot d\lambda$

and $1 - t^2 = \lambda^2 - \sigma^2 + 1 = \lambda^2 + \sigma'^2$

where

$$\sigma'^2 = 1 - \sigma^2$$

$$\begin{aligned} Z &= -A \int \frac{d\lambda}{\lambda^2 + \sigma'^2} + B. \\ &= -\frac{A}{\sigma'} \cdot \tan^{-1} \frac{\lambda}{\sigma'} + B. \\ &= -\frac{A}{\sqrt{1-\sigma^2}} \cdot \tan^{-1} \frac{\sqrt{\sigma^2-t^2}}{\sqrt{1-\sigma^2}} + B. \end{aligned}$$

Constants A and B can be evaluated with the help of corresponding known points in Z and t-planes.

At O and B

$$Z = 0, \quad t = \pm \sigma$$

Substituting we get

$$B = 0$$

At H,

$$Z = iT \text{ and } t = \alpha$$

Substituting we get

$$iT = -\frac{A}{\sigma} \tan^{-1} \alpha = -\frac{A}{2\sigma} \pi$$

or $A = -\frac{2iT\sigma'}{\pi}$

The transformation equation therefore becomes

$$Z = \frac{2iT}{\pi} \tan^{-1} \sqrt{\frac{\sigma^2 - t^2}{\sigma'^2}} \dots\dots(2.33)$$

$$\text{or } \sqrt{\frac{\sigma^2 - t^2}{\sigma'^2}} = \tan \left(\frac{\pi Z i}{2T} \right) \\ = -\tanh \left(\frac{\pi Z}{2T} \right)$$

$$\text{or } t = \sqrt{\sigma^2 + \sigma'^2 \tanh^2 \left(\frac{\pi Z}{2T} \right)} \dots\dots(2.33a)$$

To evaluate σ , consider D,

where $t = 0, Z = id$

$$0 = \sqrt{\sigma^2 + \sigma'^2 \tanh^2 \left(\frac{\pi id}{2T} \right)}$$

$$\text{or } \frac{b}{a} = \tan \frac{\pi d}{2T}$$

$$\sigma = \sin \frac{\pi d}{2T}$$

and $\sigma' = \cos \frac{\pi d}{2T} \dots\dots(2.34)$

Equation (2.33a) can be written as

$$t = \cos \frac{\pi d}{2T} \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \left(\frac{\pi Z}{2T} \right)} \dots\dots(2.35)$$

This gives

$$L_1 = \cos \frac{\pi d}{2T} \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi b_1}{2T}}$$

$$L_2 = \cos \frac{\pi d}{2T} \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi b_2}{2T}}$$

$$\gamma_1 = \cos \frac{\pi d}{2T} \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi(b_1+f)}{2T}}$$

$$\gamma_2 = \cos \frac{\pi d}{2T} \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi(b_2+f)}{2T}}$$

Second operation can be made easy by transforming t -plane on the ζ -plane such that P and H are placed symmetrically at -1 and $+1$ respectively. Therefore t -plane is transformed onto ζ -plane such that

$$\begin{array}{ll} t_A = -L_1 & \zeta_A = -\beta_1 \\ t_D = +L_2 & \zeta_D = +\beta_2 \\ t_P = -\gamma_1 & \zeta_P = -1 \\ t_H = +\gamma_2 & \zeta_H = +1 \end{array}$$

Transformation of t -plane to ζ -plane is given by

$$t = A^2 \zeta + B^2$$

To evaluate the constants consider the values of corresponding coordinates P and H

$$\begin{array}{l} - \gamma_1 = B^2 - A^2 \\ + \gamma_2 = B^2 + A^2 \end{array}$$

Solving these equations and substituting the values of A^2 and B^2 , the transformation equation becomes.

$$t = \frac{\gamma_1 + \gamma_2}{2} \zeta + \frac{\gamma_2 - \gamma_1}{2}$$

or
$$\zeta = \left[k - \frac{\gamma_2 - \gamma_1}{2} \right] \frac{2}{\gamma_1 + \gamma_2} \quad (2.38)$$

let
$$\frac{L_1}{a'} = D_1 = \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi b_1}{2T}}$$

$$\frac{L_2}{a'} = D_2 = \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi b_2}{2T}}$$

$$\frac{\gamma_1}{a'} = D_3 = \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi (b_1 + f_1)}{2T}} \quad (2.39)$$

$$\frac{\gamma_2}{a'} = D_4 = \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi (b_2 + f_2)}{2T}}$$

and
$$\frac{t}{a'} = \mu = \sqrt{\tan^2 \frac{\pi d}{2T} + \tanh^2 \frac{\pi z}{2T}}$$

Combining equations (2.35) and (2.39) and substituting equations (2.39),

$$\zeta = 1 - \frac{2(D_4 - \mu)}{D_3 + D_4}$$

$$\beta_1 = 1 - \frac{2(D_3 - D_1)}{D_3 + D_4} \quad \dots (2.40)$$

$$\beta_2 = 1 - \frac{2(D_4 - D_2)}{D_3 + D_4}$$

(II) Second Operation

Referring to figure 2.2, the rectangular flow field can be transformed onto the same real axis of ζ -plane noting that

$$\zeta_M > \zeta_B > \zeta_A > \zeta_F$$

Transformation formula is given by

$$\frac{dw}{dz} = \frac{c}{\sqrt{(\zeta - \zeta_M)(\zeta - \zeta_B)(\zeta - \zeta_A)(\zeta - \zeta_F)}} + D \quad (2.41)$$

Since the uplift pressures on the base of the structure and exit gradient at the downstream end of structure are to be evaluated the integration process is confined to this zone only.

The zone for which the uplift pressures, are to be determined corresponds to AB in ζ -plane. Pertaining to this zone

$$\zeta_M > \zeta_B \gg \zeta \gg \zeta_A > \zeta_F$$

and the integral equation becomes

$$U = C \int \frac{d\zeta}{(\zeta_M - \zeta)(\zeta_B - \zeta)(\zeta - \zeta_A)(\zeta - \zeta_F)} \quad \dots (2.42)$$

which on integration using bilinear transformation yields (refer to Appendix II)

$$\sin^2 \frac{\mu W}{C} = \frac{(\zeta_M - \zeta_A)(\zeta - \zeta_B)}{(\zeta_B - \zeta_A)(\zeta - \zeta_M)} \quad \dots (2.43)$$

where

$$\mu = \frac{1}{2} \sqrt{(\zeta_B - \zeta_A)(\zeta_M - \zeta_F)}$$

and the modulus k is given by

$$k = \sqrt{\frac{(\zeta_B - \zeta_A)(\zeta_M - \zeta_F)}{(\zeta_M - \zeta_A)(\zeta_B - \zeta_F)}}$$

from which the co-modulus defined by

$$k' = \sqrt{1 - k^2} = \sqrt{\frac{(\zeta_A - \zeta_F)(\zeta_M - \zeta_B)}{(\zeta_M - \zeta_A)(\zeta_B - \zeta_F)}}$$

and from equation (2.43) it follows that

$$\zeta = \frac{\zeta_0(\zeta_M - \zeta_A) - \zeta_M(\zeta_B - \zeta_A) \operatorname{sn}^2 M_2 W}{(\zeta_M - \zeta_A) - (\zeta_B - \zeta_A) \operatorname{sn}^2 M_2 W} \dots\dots(2.43b)$$

Further to simplify the procedure, we assume the constant C equal to 1. This will yield

$$\zeta = \frac{\zeta_B(\zeta_M - \zeta_A) - \zeta_M(\zeta_B - \zeta_A) \operatorname{sn}^2 W}{(\zeta_M - \zeta_A) - (\zeta_B - \zeta_A) \operatorname{sn}^2 W} \dots\dots(2.43a)$$

$$\operatorname{sn}^2 W = \frac{(\zeta_M - \zeta_A)(\zeta - \zeta_B)}{(\zeta_B - \zeta_A)(\zeta - \zeta_M)} \dots\dots(2.44b)$$

Referring to figure 2.2d, the characteristic coordinates of the transformation are

$$\begin{aligned} \zeta_H &= 1 \\ \zeta_D &= \beta_2 \\ \zeta_A &= \beta_1 \\ \zeta_F &= -1 \end{aligned} \dots\dots(2.45)$$

Now substituting the characteristic coordinates of ζ -plane given by equation (2.45) in equation (2.44b) the corresponding values of v are determined. Thus for

$$\zeta = \zeta_B \quad (\text{at D})$$

equation (3.10b) reduces to (refer appendix I equations 5a, 5b, 5c and 5d)

$$\begin{aligned} \operatorname{sn}^2 U &= 0 \\ \text{or } U &= 0 \\ \text{for } \zeta &= \zeta_M \quad (\text{at H}) \\ \operatorname{sn}^2 W &= \alpha \\ \text{or } W &= iK' \end{aligned}$$

for $\zeta = \zeta_A$ (L0 A)

$$\operatorname{Im} w = 1$$

$$w = K$$

and for $\zeta = \zeta_F$ (L0 F)

$$\operatorname{Im} w = \frac{(\zeta_M - \zeta_A)(\zeta_F - \zeta_B)}{(\zeta_B - \zeta_A)(\zeta_F - \zeta_M)} = 1/k$$

or $w = K + iK'$

Thus the four coordinates of the w -plane are

$\zeta = \zeta_M$	$w = iK'$
$= \zeta_B$	$w = 0$
$= \zeta_A$	$w = K$
$= \zeta_F$	$w = K + iK'$

in which K and K' are complete elliptic integrals of the first kind with modulus k and k' respectively. The four values of w thus calculated, determine the layout of the basic rectangular flow field for the second operation.

The value of characteristic coordinates given by the equation (2.45) when substituted in k and k' yield

$$k = \sqrt{\frac{2(\beta_1 + \beta_2)}{(1 + \beta_1)(1 + \beta_2)}} \dots\dots(2.47)$$

$$k' = \sqrt{\frac{(1 - \beta_1)(1 - \beta_2)}{(1 + \beta_1)(1 + \beta_2)}}$$

Similarly the substitution of characteristic coordinate values given in equation (2.45) in equation (2.44b) yields

$$\sin^2 \psi = \frac{(1+\beta_1)(\beta_2-\xi)}{(1+\beta_1)(1+\beta_2)}$$

or

$$\sin \psi = \sqrt{\frac{(1+\beta_1)(\beta_2-\xi)}{(1+\beta_1)(1+\beta_2)}} \quad \dots (2.48)$$

2.4.3 EVALUATION OF UPLIFT PRESSURES

Along the bottom line of the foundation

$$\psi' = 0$$

$$\text{and } \psi = \phi'$$

therefore equation (2.48) is reduced to

$$\sin \phi' = \sqrt{\frac{(1+\beta_1)(\beta_2-\xi)}{(1+\beta_1)(1+\beta_2)}} \quad \dots (2.49)$$

The value of β_1 and β_2 are given by equation (2.40)

The pressure decreases from maximum load H at point A in the beginning to 0 at point B in the end. The length of line $ABCD$ in ψ -plane equals to complete elliptic integral of first kind (K) having its modulus equal to k (Figure 2.21)

Potential for $H = 1$, is found from the relation

$$\phi = \frac{\phi'}{k} \quad \dots (2.50)$$

and in percent as

$$\phi' = \frac{\phi}{k} \times 100$$

To evaluate the value of potential or pressure at any point along the flow the value of ξ is found from equation (2.43) corresponding to its location in Z-plane. Substituting the value of β_1 , β_2 and ξ in equations (2.47) and (2.49) the values of modulus and 2π are found. Using tables for Jacobian functions [14, 20] the required potential corresponding to any point in Z-plane can be evaluated.

2.4.3 BIFT GRADIENT

It has been indicated in para 2.3.3 that it is also important to know the hydraulic gradient at the down stream end of the first percolation trajectory.

Gradient at any point is given by

$$I = \frac{dh}{ds}$$

where h = head at any point

$$= \frac{H}{K} \times \phi'$$

Therefore $I = \frac{H}{K} \frac{d\phi'}{ds} \dots (2.51)$

where K = complete Elliptic integral of First kind with modulus k .

It is not possible to find out $\frac{d\phi'}{ds}$ directly therefore differentiation by stages is adopted as shown below

$$\frac{d\phi'}{ds} = \frac{d\phi'}{d\xi} \cdot \frac{d\xi}{ds} \cdot \frac{ds}{d\theta} \dots (2.51a)$$

$$Z = x + iy$$

or $\frac{dZ}{d\theta} = \frac{dx}{d\theta} + i \frac{dy}{d\theta}$

$$= \cos\theta + i \sin\theta$$

where θ = angle of oblique line with horizontal.
at the end

$$\theta = 90^\circ$$

and therefore

$$\frac{d\eta}{d\theta} = 1 \quad \dots\dots(2.52)$$

Again

$$\frac{d\zeta}{d\theta} = \frac{d\zeta}{d\theta} \cdot \frac{d\theta}{d\theta}$$

From equation (2.52)

$$\begin{aligned} \frac{d\eta}{d\theta} &= \frac{(1-\eta^2)(\sigma^2-\eta^2)^{\frac{1}{2}}}{\Delta\sigma} \\ &= \frac{\pi(1-\eta^2)(\sigma^2-\eta^2)^{\frac{1}{2}}}{2\pi\sigma} \quad \dots\dots(2.53) \end{aligned}$$

Differentiating equation (2.53)

$$\frac{d\zeta}{d\theta} = \frac{2}{\sigma_1 + \sigma_2} \quad \dots\dots(2.54)$$

Combining equations (2.53) and (2.54), we obtain

$$\begin{aligned} \frac{d\zeta}{d\theta} &= \frac{2\pi}{\sigma_1 + \sigma_2} \cdot \frac{(1-\eta^2)(\sigma^2-\eta^2)^{\frac{1}{2}}}{2\pi\sigma} \\ &= \frac{\pi}{(\sigma_1 + \sigma_2)} \cdot \frac{(1-\sigma^2\mu^2)(\sigma^2-\sigma^2\mu^2)^{\frac{1}{2}}}{c.T.\sigma.\sigma.\mu} \quad \dots\dots(2.55) \end{aligned}$$

where σ_1 , σ_2 and μ are defined in equation (2.50)

$$U = \phi' + i\psi'$$

along the foundation profile, $\psi' = 0$

$$U = \phi'$$

$$\frac{dU}{d\theta} = \frac{d\phi'}{d\theta} = \frac{c}{\sqrt{(\sigma-\sigma_M)(\sigma-\sigma_B)(\sigma-\sigma_A)(\sigma-\sigma_F)}}$$

where

$$C = \frac{1}{2} \sqrt{(\xi_B - \xi_F)(\xi_M - \xi_A)}$$

$$\frac{d\phi}{d\xi} = \frac{1}{2} \sqrt{\frac{(\xi_B - \xi_F)(\xi_M - \xi_A)}{(\xi - \xi_M)(\xi - \xi_B)(\xi - \xi_A)(\xi - \xi_F)}}$$

substituting characteristic coordinates of transformation

$$\frac{dq}{d\xi} = \frac{1}{2} \sqrt{\frac{(1+\beta_1)(1+\beta_2)}{(\xi^2-1)(\xi+\beta_1)(\xi-\beta_2)}} \dots (2.56)$$

From equation (2.5)

$$1 + \beta_1 = \frac{B(D_3 + D_2)}{D_3 + D_4}$$

$$1 + \beta_2 = \frac{B(D_3 + D_2)}{D_3 + D_4}$$

$$\xi - 1 = \frac{B(\mu - D_4)}{D_3 + D_4}$$

$$\xi + 1 = \frac{B(D_3 + \mu)}{D_3 + D_4} \dots (2.57)$$

$$\xi + \beta_1 = \frac{B(\mu + D_4)}{D_3 + D_4}$$

$$\xi - \beta_2 = \frac{B(\mu - D_2)}{D_3 + D_4}$$

Substituting values from eqs (2.51) in equation (2.56) and simplifying to obtain

$$\frac{dq}{d\xi} = \frac{1}{2} (D_3 + D_4) \sqrt{\frac{(D_3 + D_1)(D_3 + D_2)}{(\mu - D_4)(\mu + D_3)(\mu + D_1)(\mu - D_2)}} \dots (2.57a)$$

Combining equation (2.52), (2.53) and (2.57a) and substituting in equation (2.51) to get after simplification

$$\text{Exit gradient} = \mathcal{G} = \frac{H\pi(1-\sigma^2)\mu^2}{4kT\sigma^3\mu} \sqrt{\frac{(D_4+\beta_1)(D_3+\beta_2)(\sigma^2-\sigma'^2\mu^2)}{(D_4+\mu)(D_3+\mu)(D_2+\mu)(D_1+\mu)}}$$

with about pile at the end of floor, at unit

$$\mu = \frac{b}{g}$$

$$D_3 = \frac{b}{g}$$

$$\begin{aligned} \theta &= \frac{H \pi}{g \pi^2} \sqrt{\frac{(D_3 + D_1) z \cdot \sigma \cdot \sigma'}{(D_3 + \mu)(D_1 + \mu)}} \\ &= \frac{H \pi}{g \pi^2} \sqrt{\frac{(D_3 + D_1) z \cdot \cos\left(\frac{\pi d}{2T}\right)}{(D_3 - \tan^2 \frac{\pi d}{2T})(D_1 - \tan^2 \frac{\pi d}{2T})}} \end{aligned}$$

where

H = Total head in foot of water

K = Complete elliptic integral of first kind for modulus $k = \frac{2(\beta_1 + \beta_2)}{(1 + \beta_1)(1 + \beta_2)}$

T = Depth of permeable soil in feet

$$D_1 = \sqrt{\tan^2 \frac{\pi d}{2T} + \tan^2 \frac{\pi b_1}{2T}}$$

$$D_3 = \sqrt{\tan^2 \left(\frac{\pi d}{2T}\right) + \tan^2 \frac{\pi (b_2 + f_2)}{2T}}$$

2.8 Application of above Equations for computations

The equations as derived above can be used for computation of uplift pressures and exit gradients. These computations involve the use of elliptic functions which have to be found from tables (14, 20). The computations thus becomes rather cumbersome. In order to facilitate the use of above equations for finding the uplift pressures and exit gradients at key points, values have been computed for different combinations of the variables involved and plotted in the form of curves. These curves can be directly used for practical design purposes. As will be seen interpolation of

values is very easy and accurate and can be safely adopted for design of most of hydraulic structures governed by already mentioned boundary conditions.

These curves have been worked out and plotted for simple elementary profile consisting of a flat floor with a cut off at one end. In any actual problem of hydraulic structure, the foundation profile is more complicated. But it can be reduced to simple elementary profile and pressures at key points worked out. The values have to be corrected by the application of method of "independ^{cut} variables".

The computation of various values and packing of elliptic function tables by linear interpolation for plotting the curves were carried out on an IIM 1620 digital Computer. Computations were carried out for boundary conditions as mentioned below and the curves obtained from each have been shown in figures against corresponding case.

Boundary condition	Figure no.
A. Structure founded on infinite depth of pervious flow area	
i. uplift pressure at D	
(a) upstream pervious length = 1.00	2.3
(b) upstream pervious length = 2.00	2.4
(c) upstream pervious length = 3.00	2.5
ii. uplift pressure at E	
(a) upstream pervious length = 1.00	2.6
(b) upstream pervious length = 2.00	2.7
(c) upstream pervious length = 3.00	2.8

iii. Ditch positions

- (a) upstream pervious length = 1.00 2.9
- (b) upstream pervious length = 2.00 2.10
- (c) upstream pervious length = 3.00 2.11

D. Structure founded on finite pervious flow zone

i. uplift pressures at D

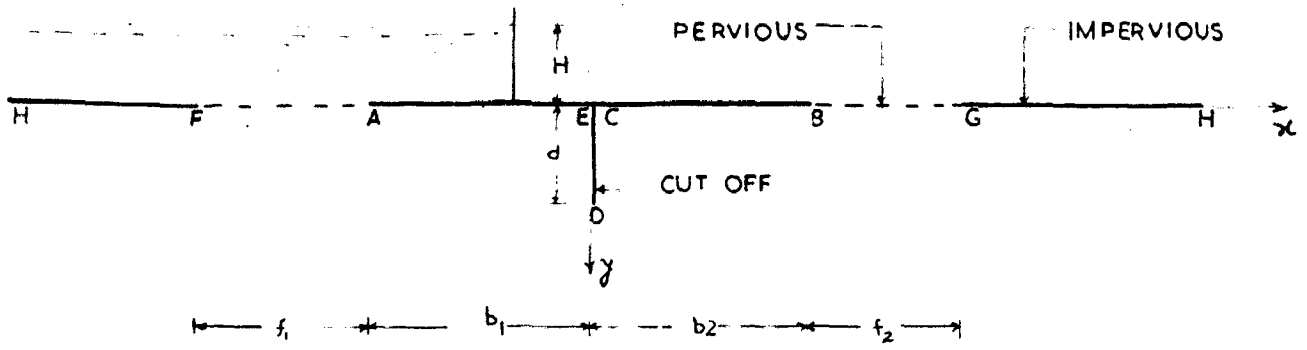
- (a) upstream pervious length = 0.50 2.13
- (b) upstream pervious length = 1.00 2.15
- (c) upstream pervious length = 1.50 2.14

ii. uplift pressures at E

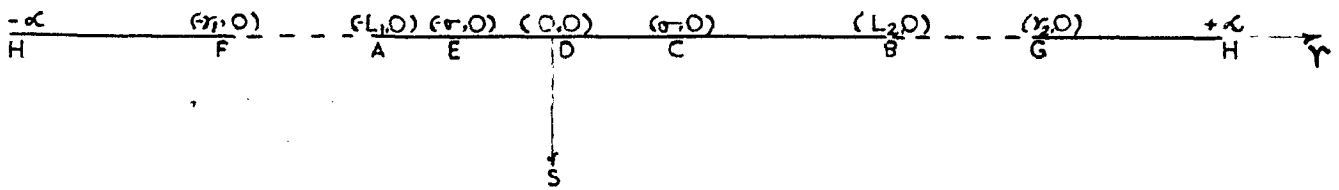
- (a) upstream pervious length = 0.50 2.16
- (b) upstream pervious length = 1.00 2.10
- (c) upstream pervious length = 1.50 2.17

iii. Ditch positions

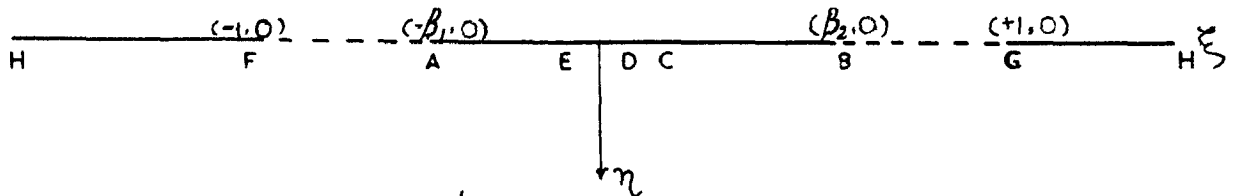
- (a) upstream pervious length = 0.50 2.18
- (b) upstream pervious length = 1.00 2.19
- (c) upstream pervious length = 1.50 2.20



Z - PLANE
(a)



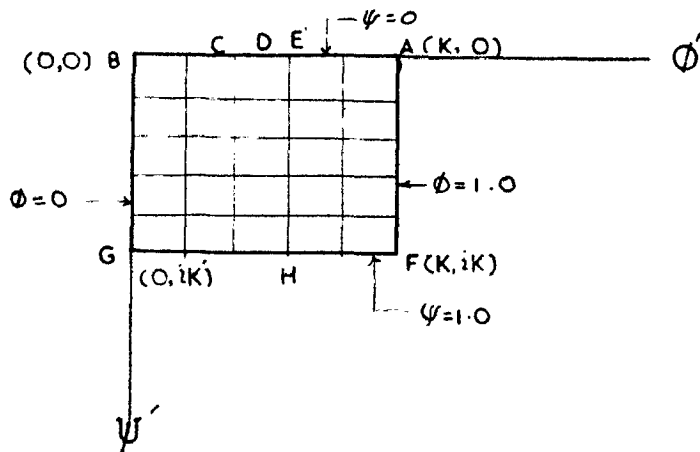
z PLANE
(b)



zeta - PLANE
(c)

$$\phi = \phi/k$$

$$\psi = \psi/k$$



w-PLANE
(d)

FIG. 2-1 SCHWARTZ-CHRISTOFFEL TRANSFORMATION

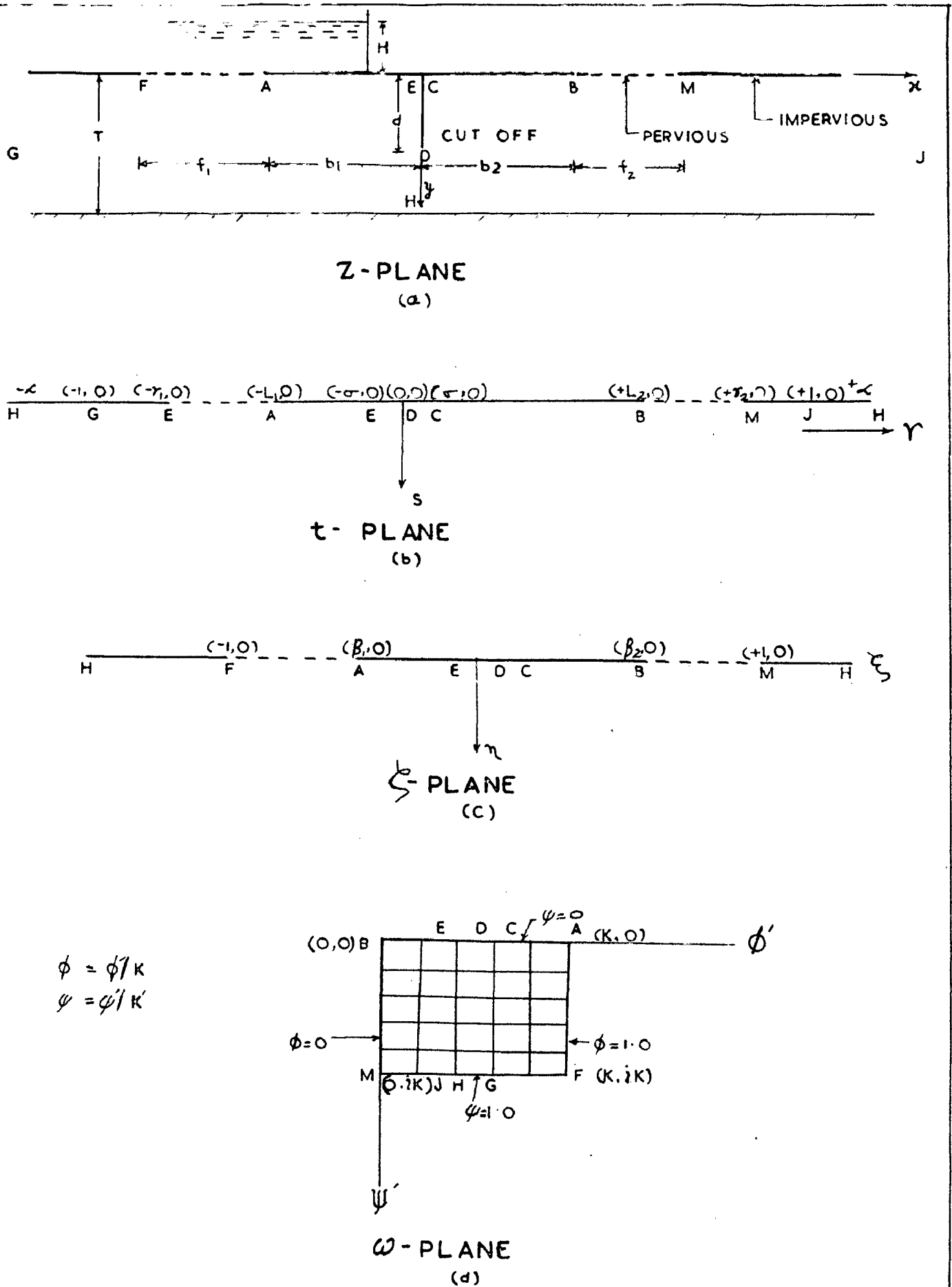


FIG. 2.2. SCHWARTZ-CHRISTOFFEL TRANSFORMATIONS

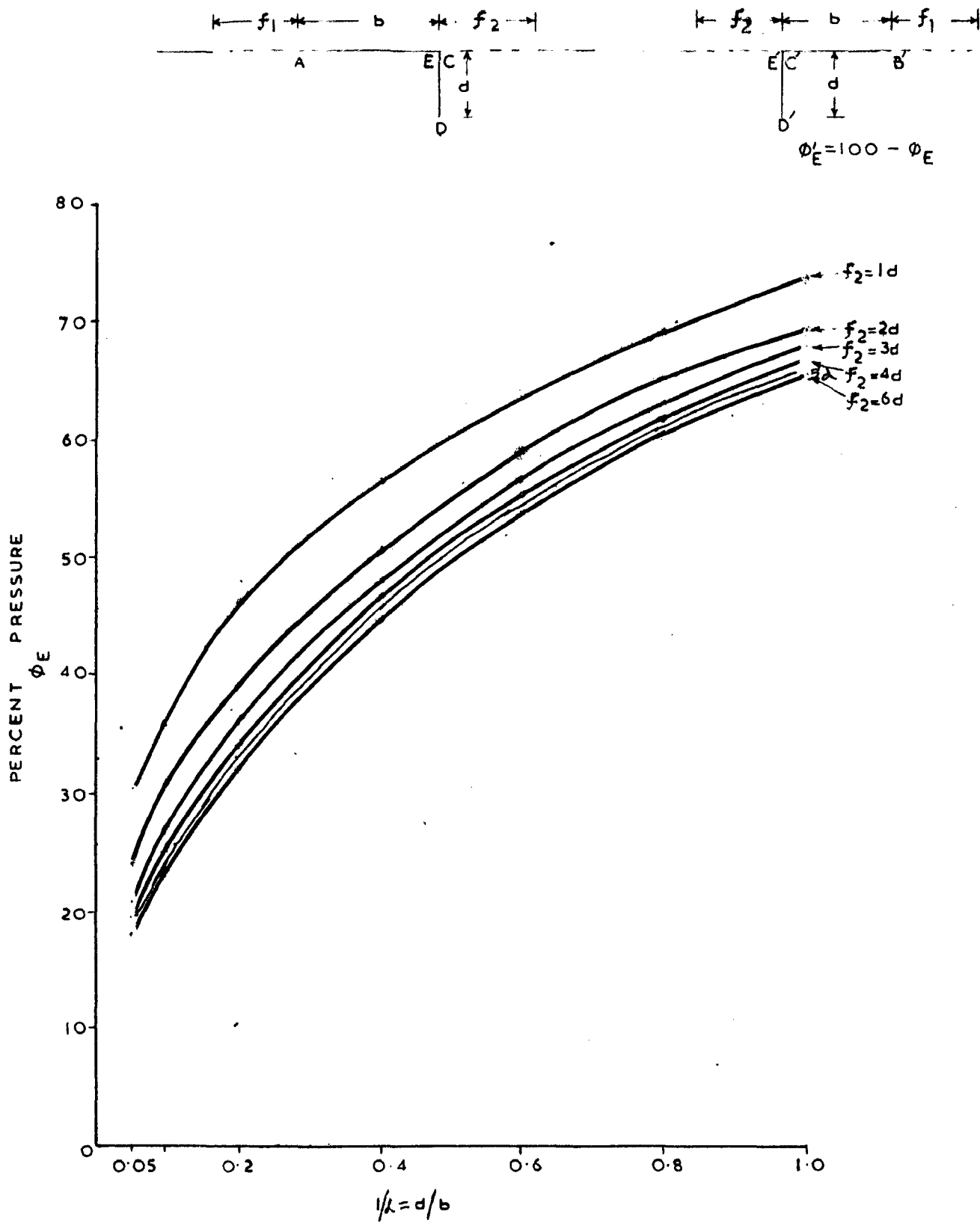


FIG. 2.3 PERCENT PRESSURE AT E FOR $f_1 = 1.0d$.

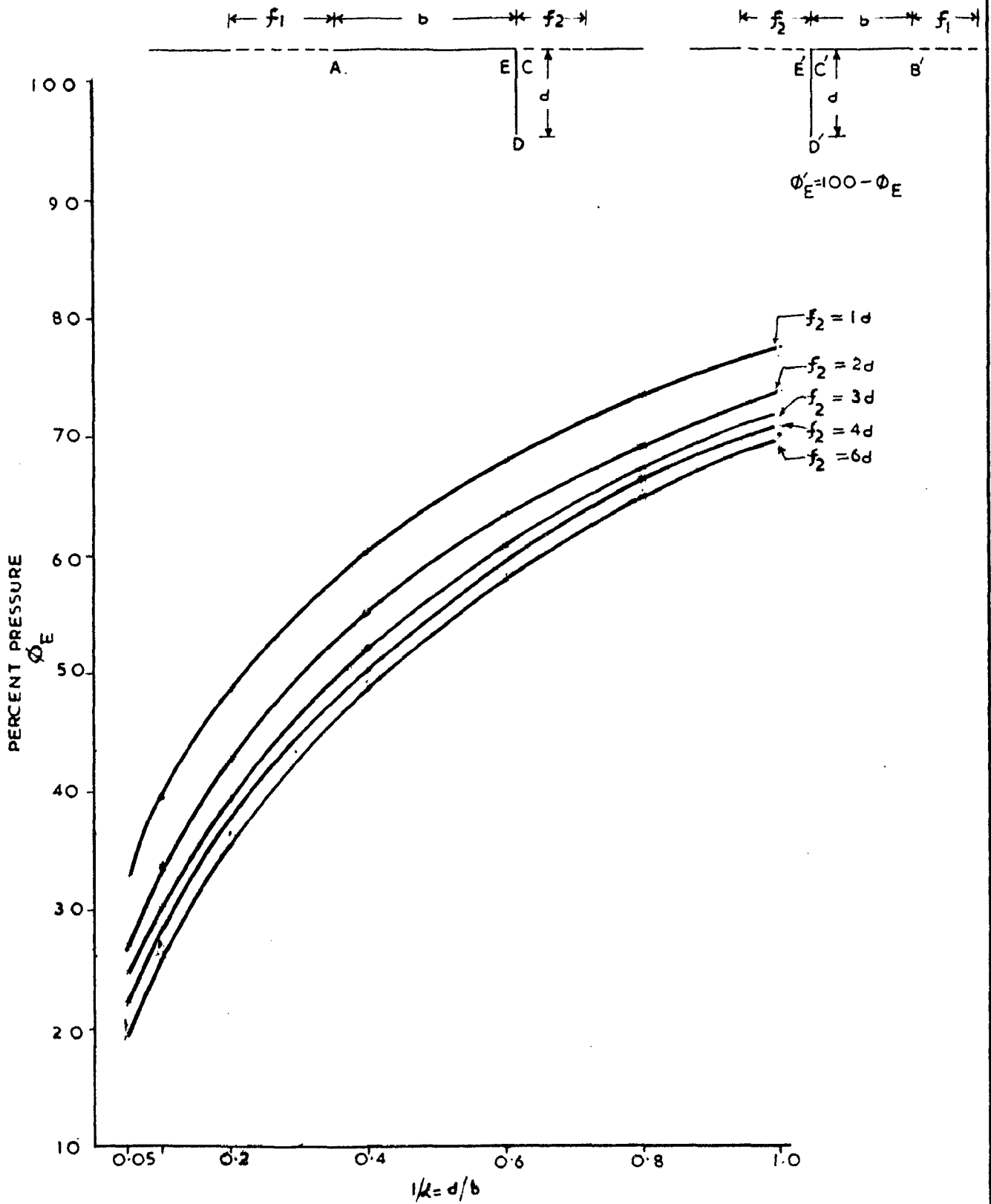


FIG. 2-4. PERCENT PRESSURE AT E FOR $f_1 = 2.0 d$.

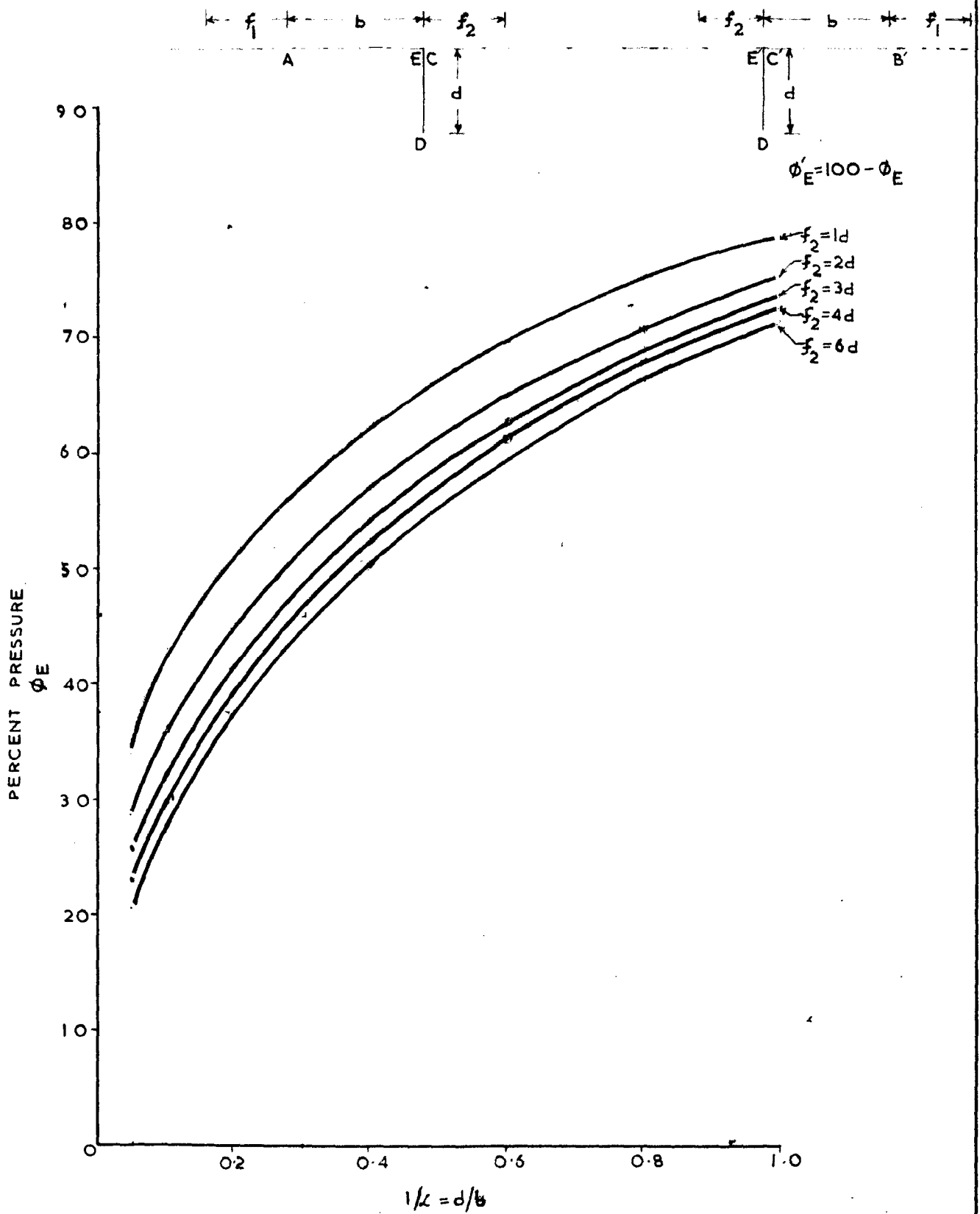


FIG. 2.5. PERCENT PRESSURE AT E FOR $f_1 = 3.0 d$.

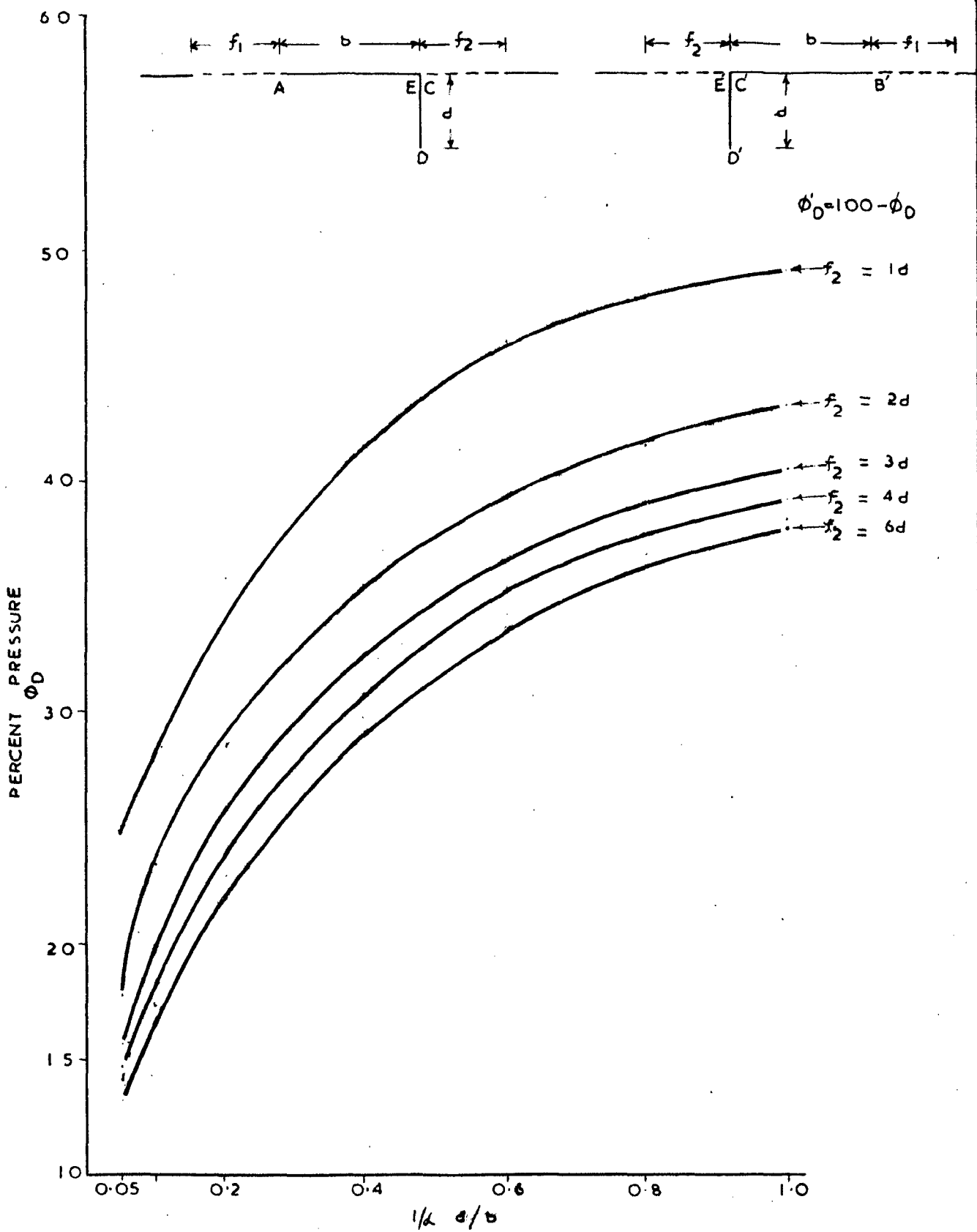


FIG. 2.6. PERCENT PRESSURE AT D FOR $f_1 = 1.0d$.

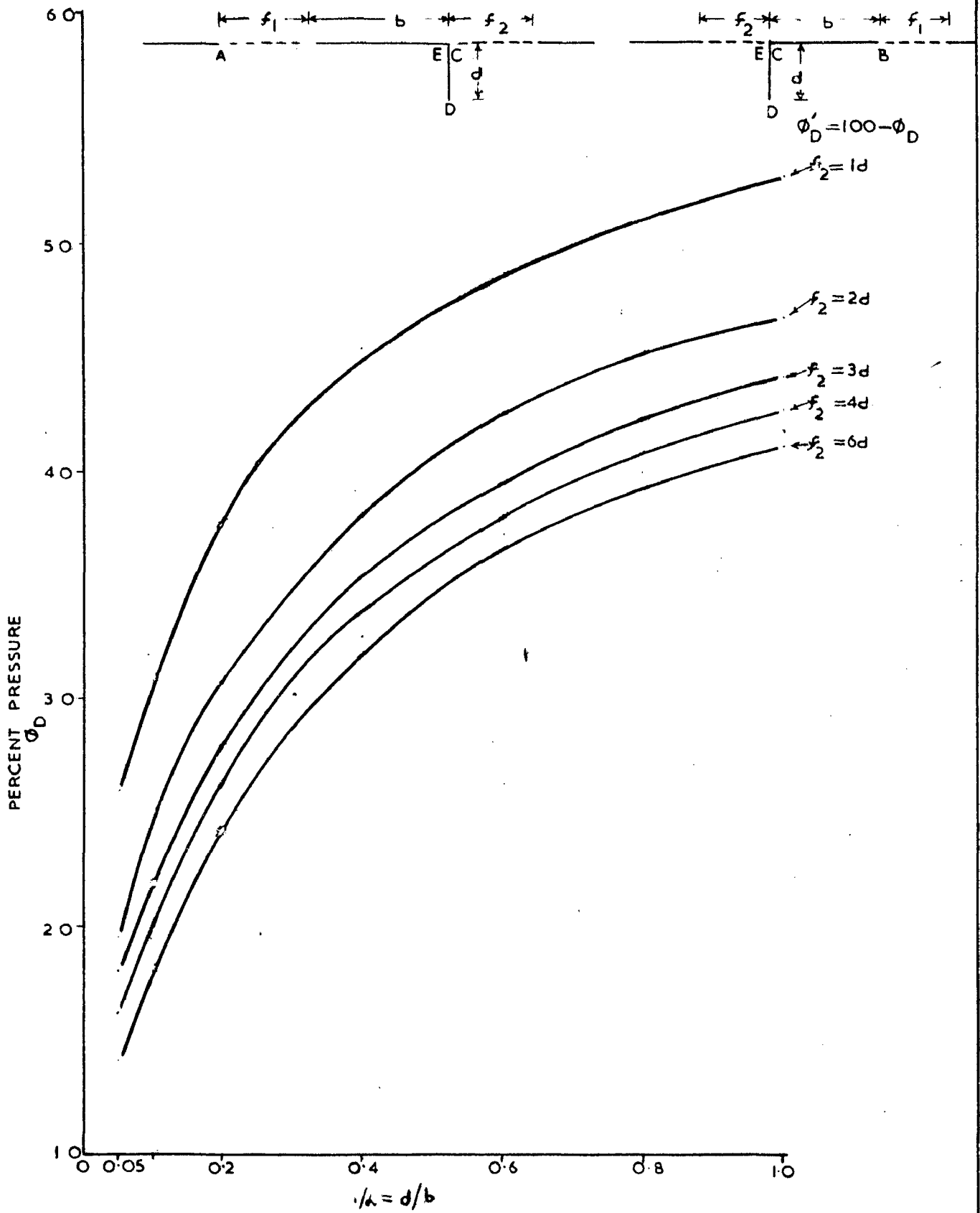


FIG. 2.7. PERCENT PRESSURE AT D FOR $f_1 = 2.0 d$.

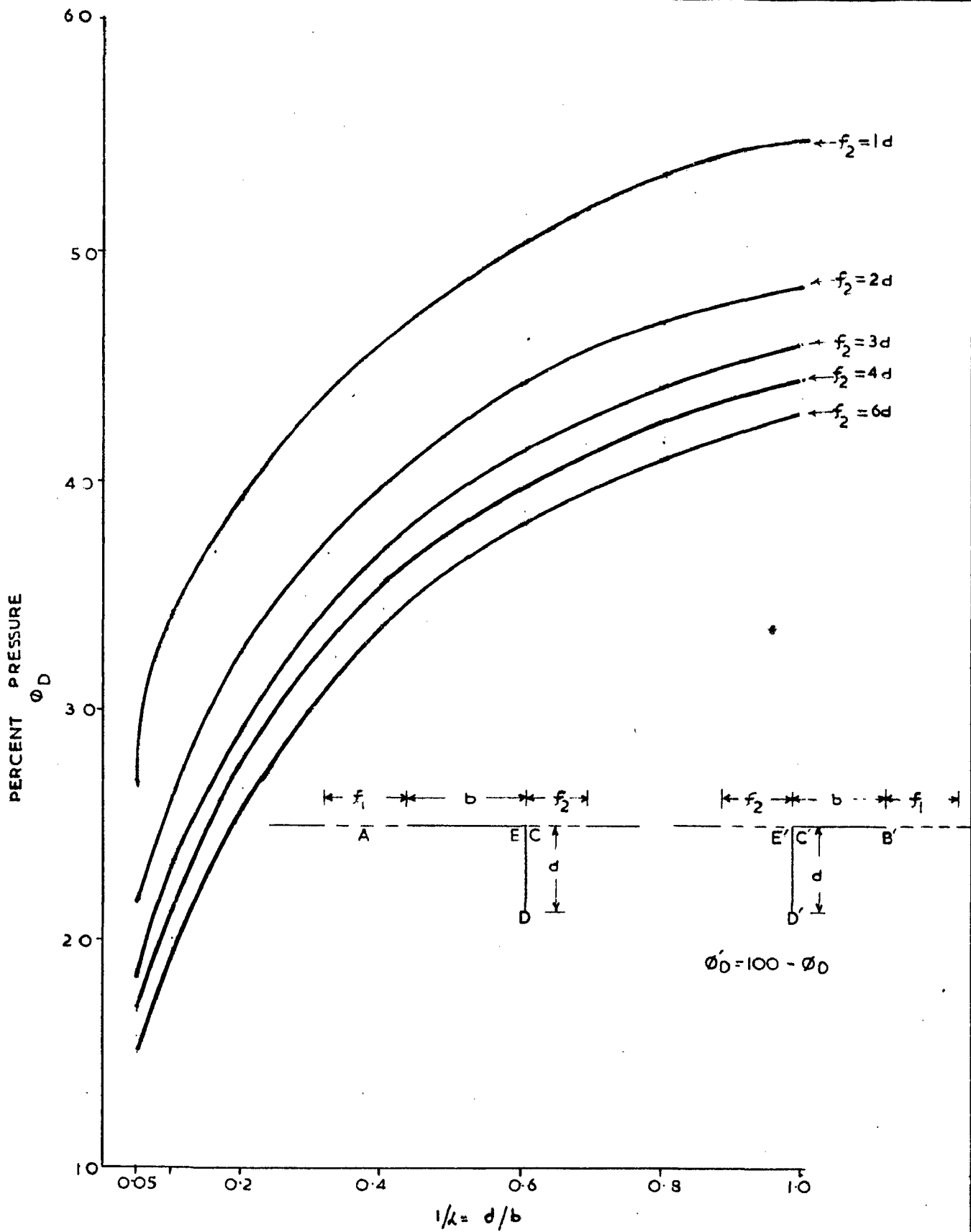


FIG. 2-8. PERCENT PRESSURE AT D FOR $f_1 = 3.0d$.

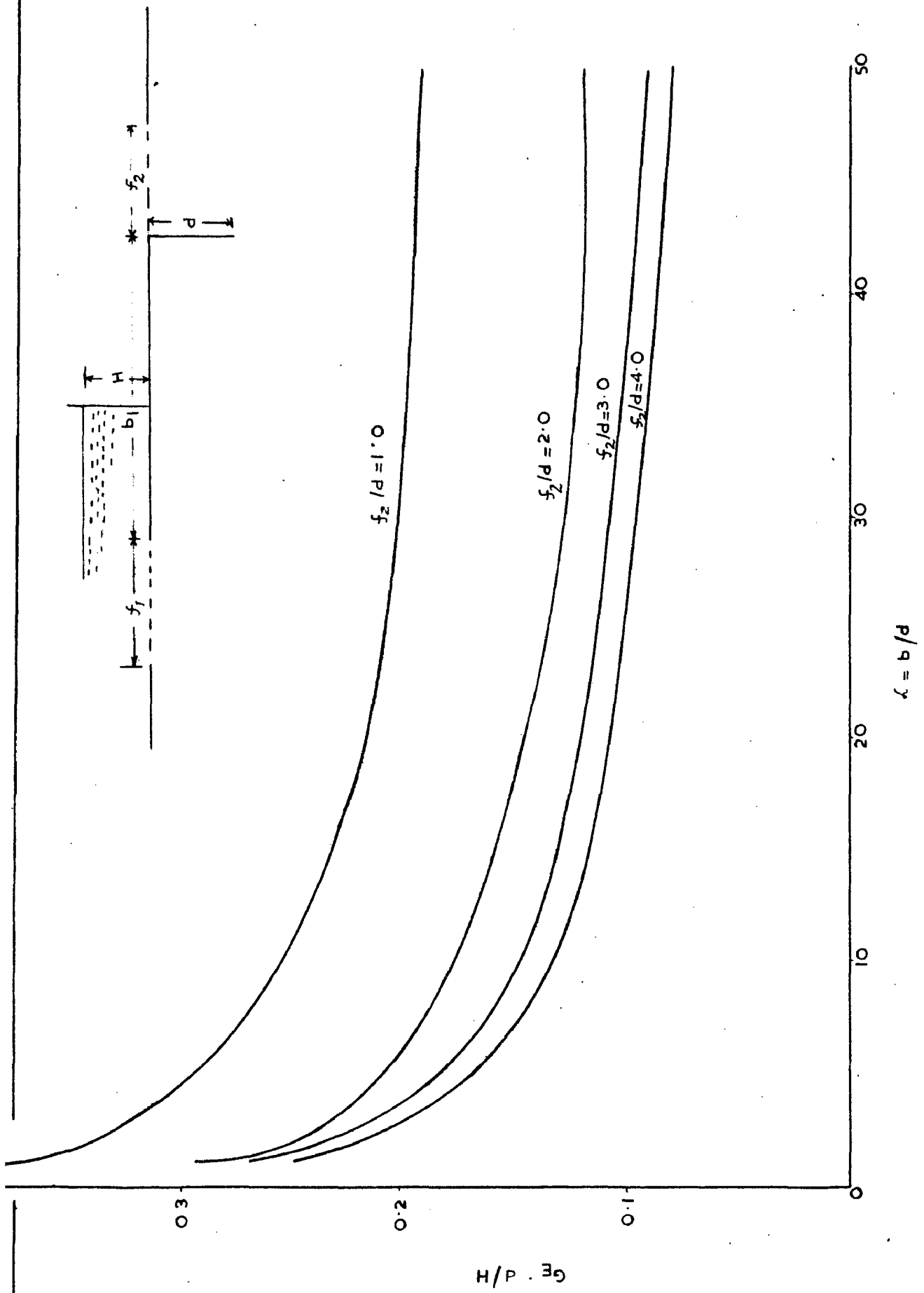


FIG. 2-9. EXIT GRADIENT FOR $f_2/d = 1.0$ d.

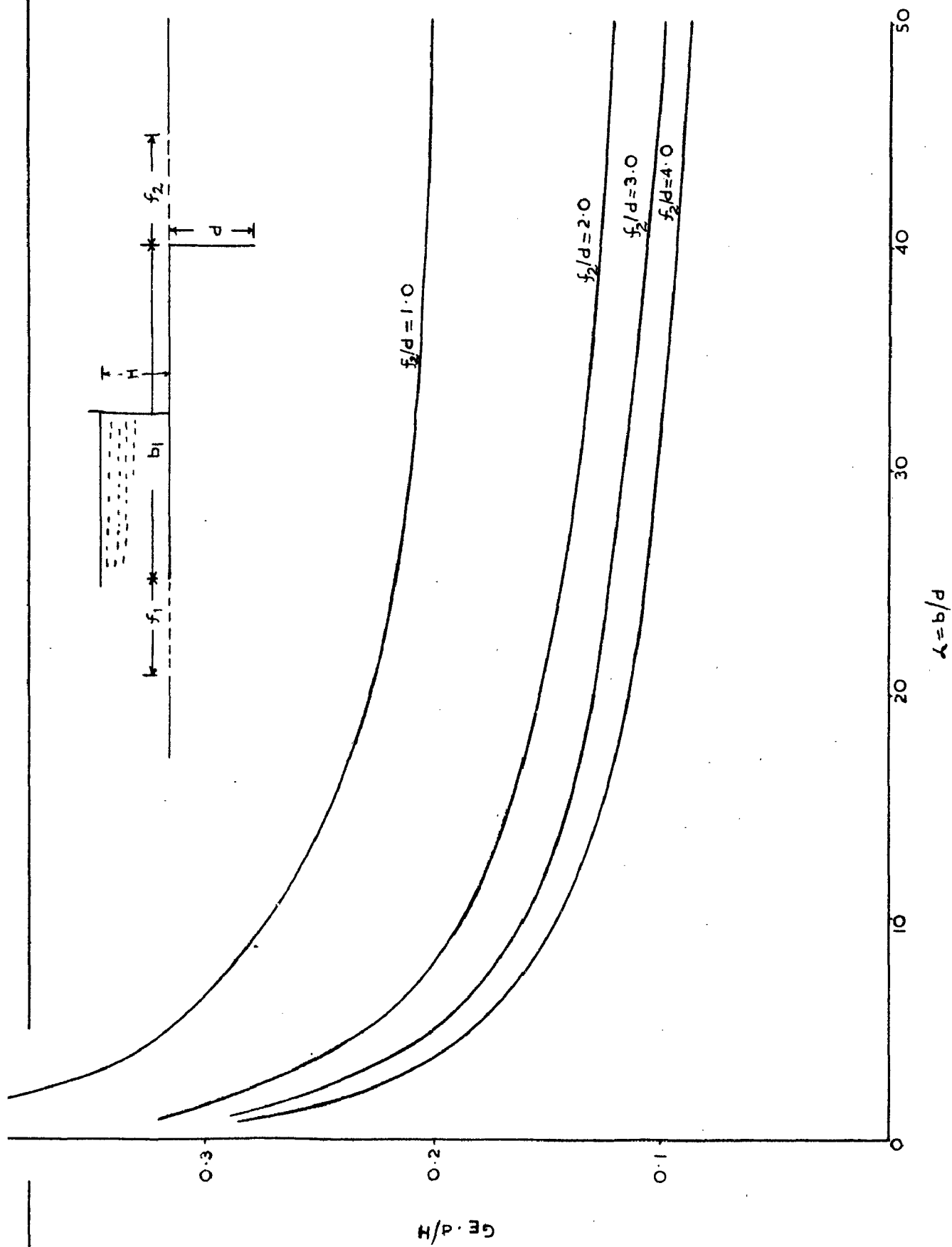


FIG. 2·10. EXIT GRADIENT FOR $f_1 = 2 \cdot 0 d$.

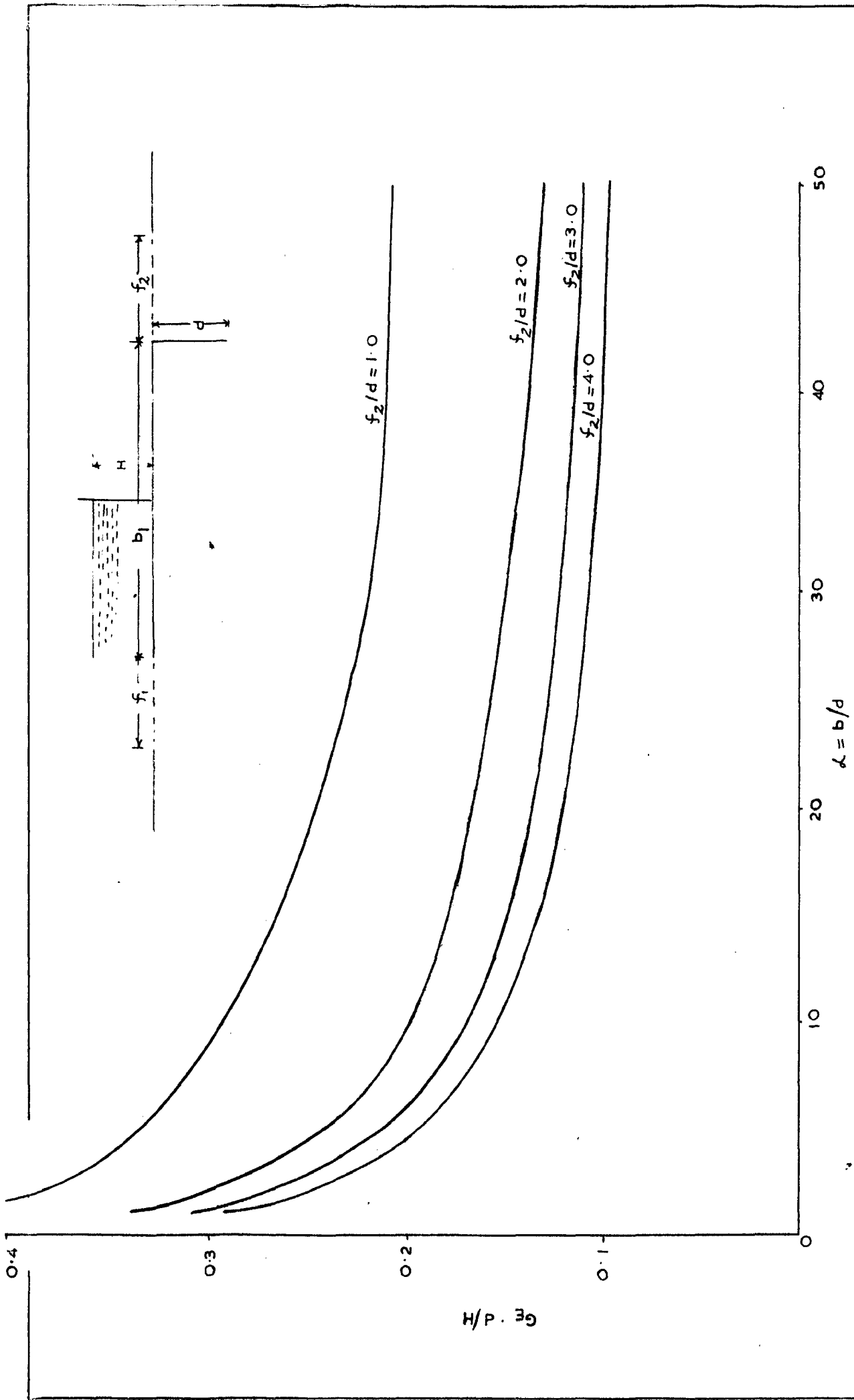


FIG. 2.11. EXIT GRADIENT FOR $f_1 = 3.0d$.

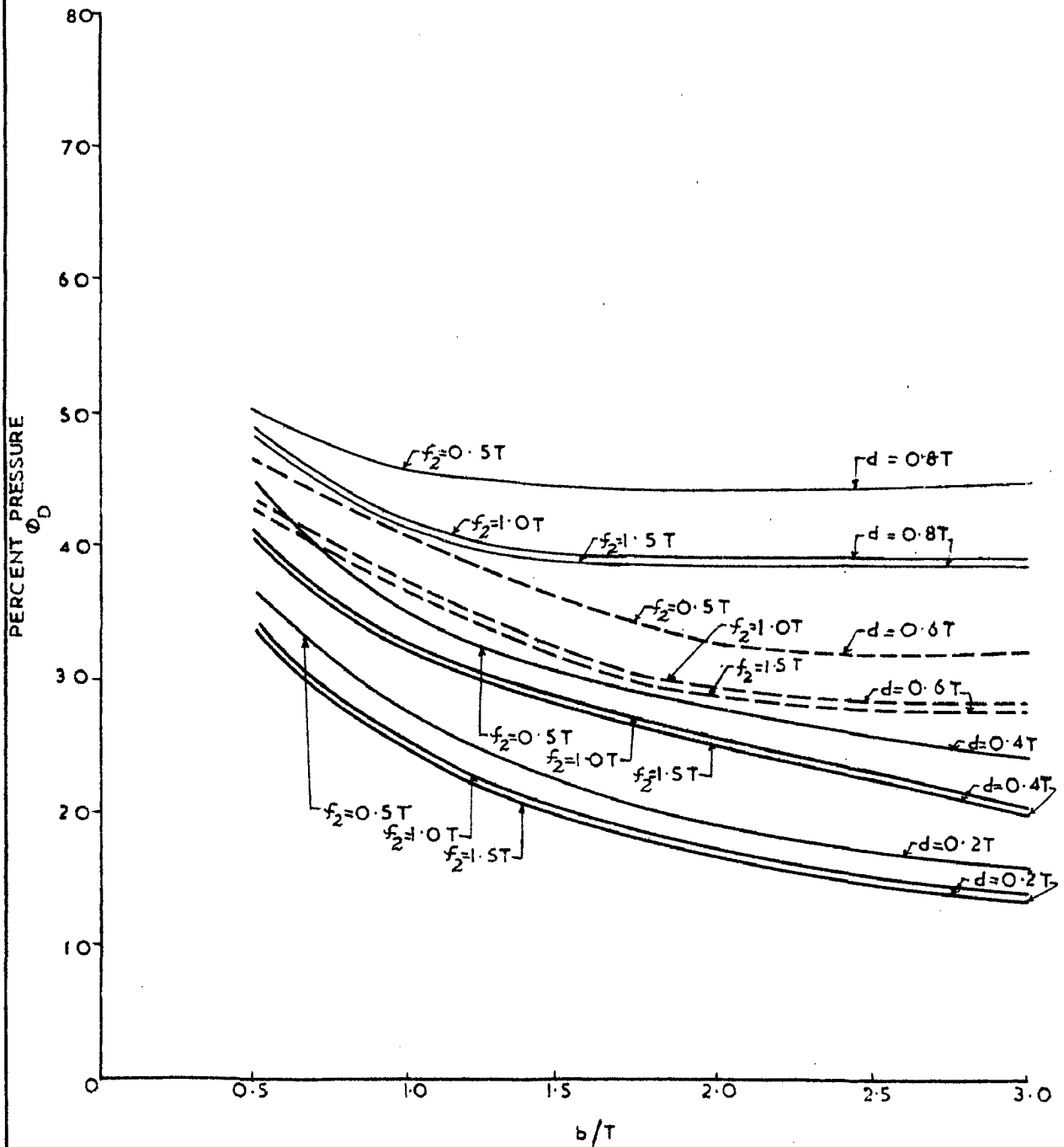
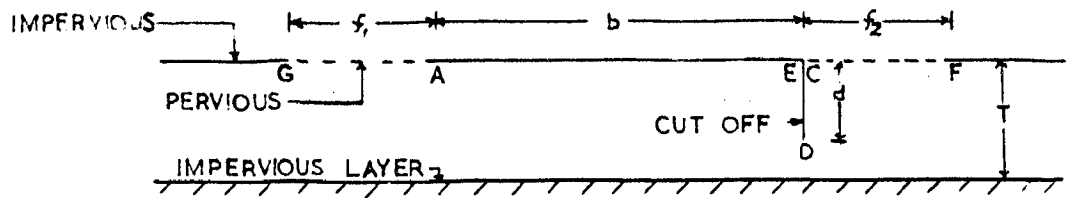


FIG. 2.12. PERCENT PRESSURE AT D FOR $f_1 = 0.5T$.

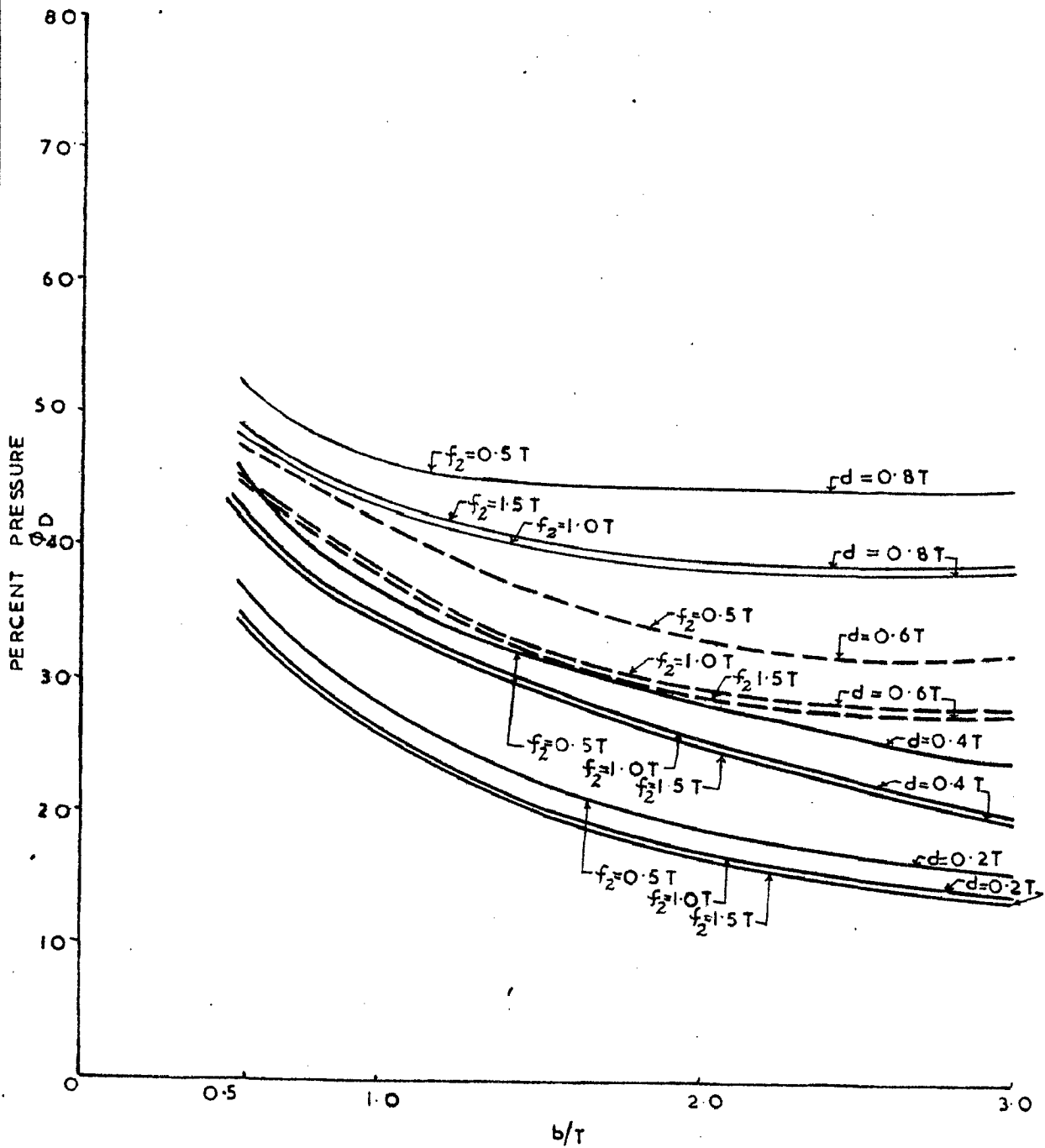
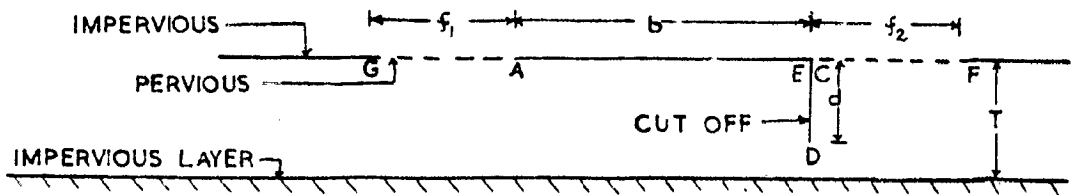


FIG. 2-13. PERCENT PRESSURE AT D FOR $f_1 = 1.0T$.

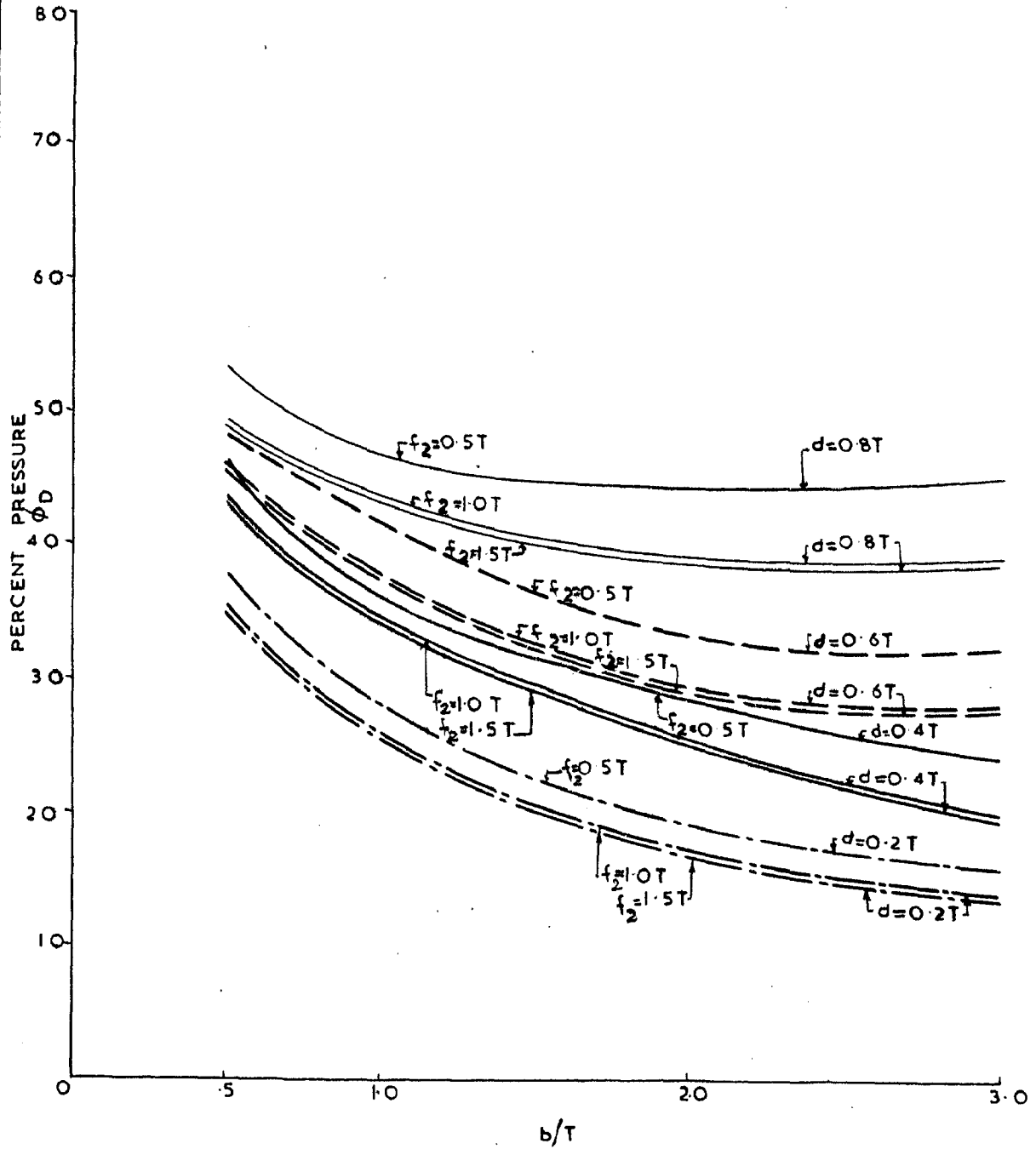
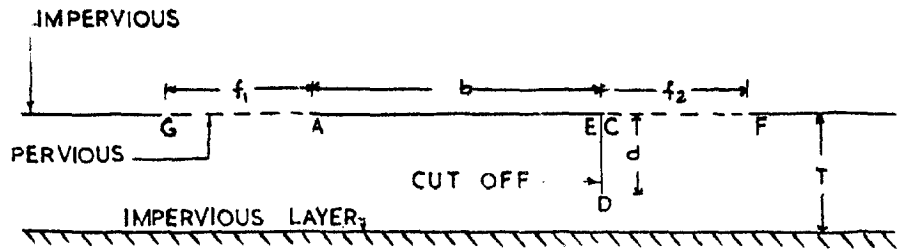


FIG. 2.14 PERCENT PRESSURE AT D FOR $f_1 = 1.5T$

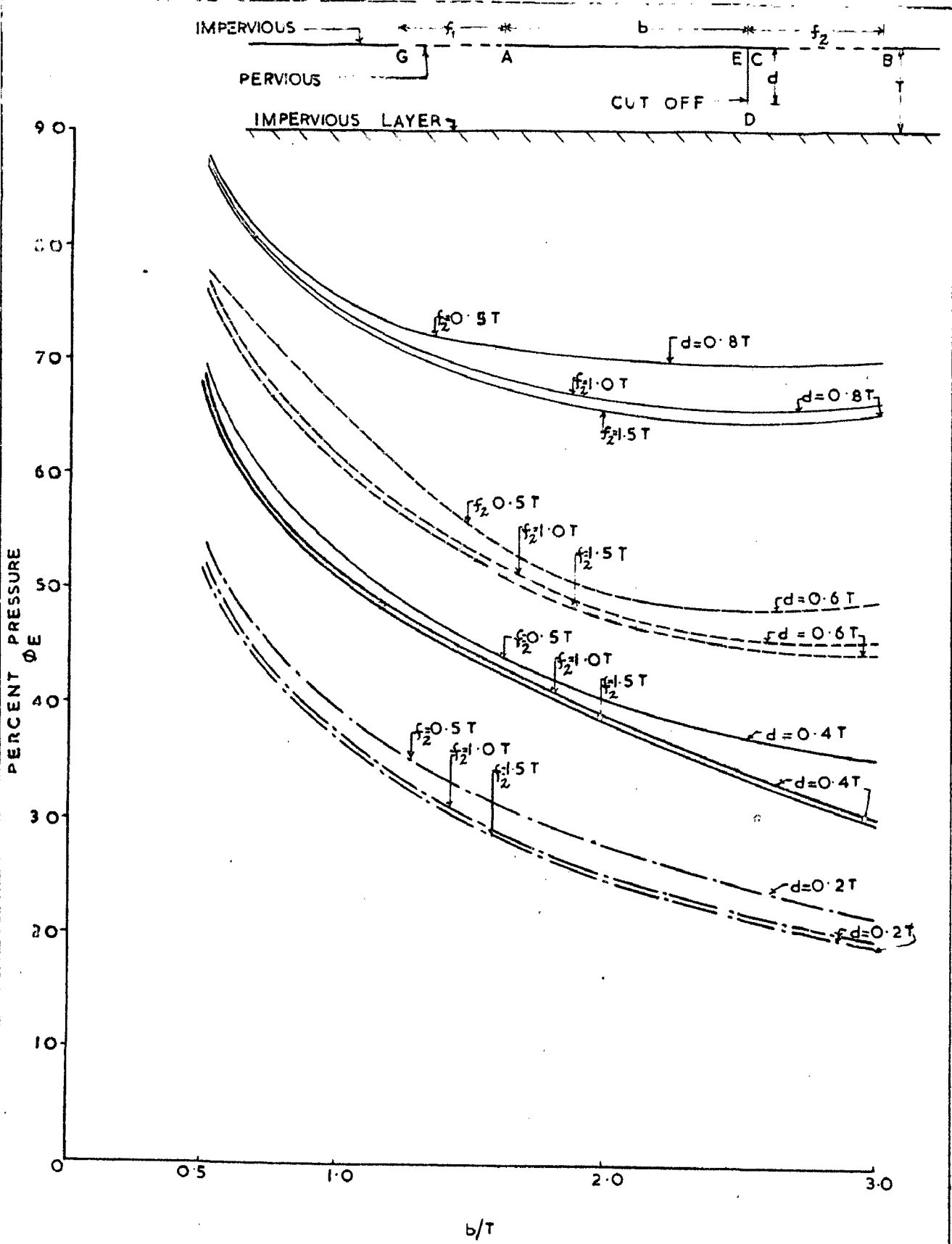


FIG.2-15 PERCENT PRESSURE AT E FOR $f_1 = 0.5T$.

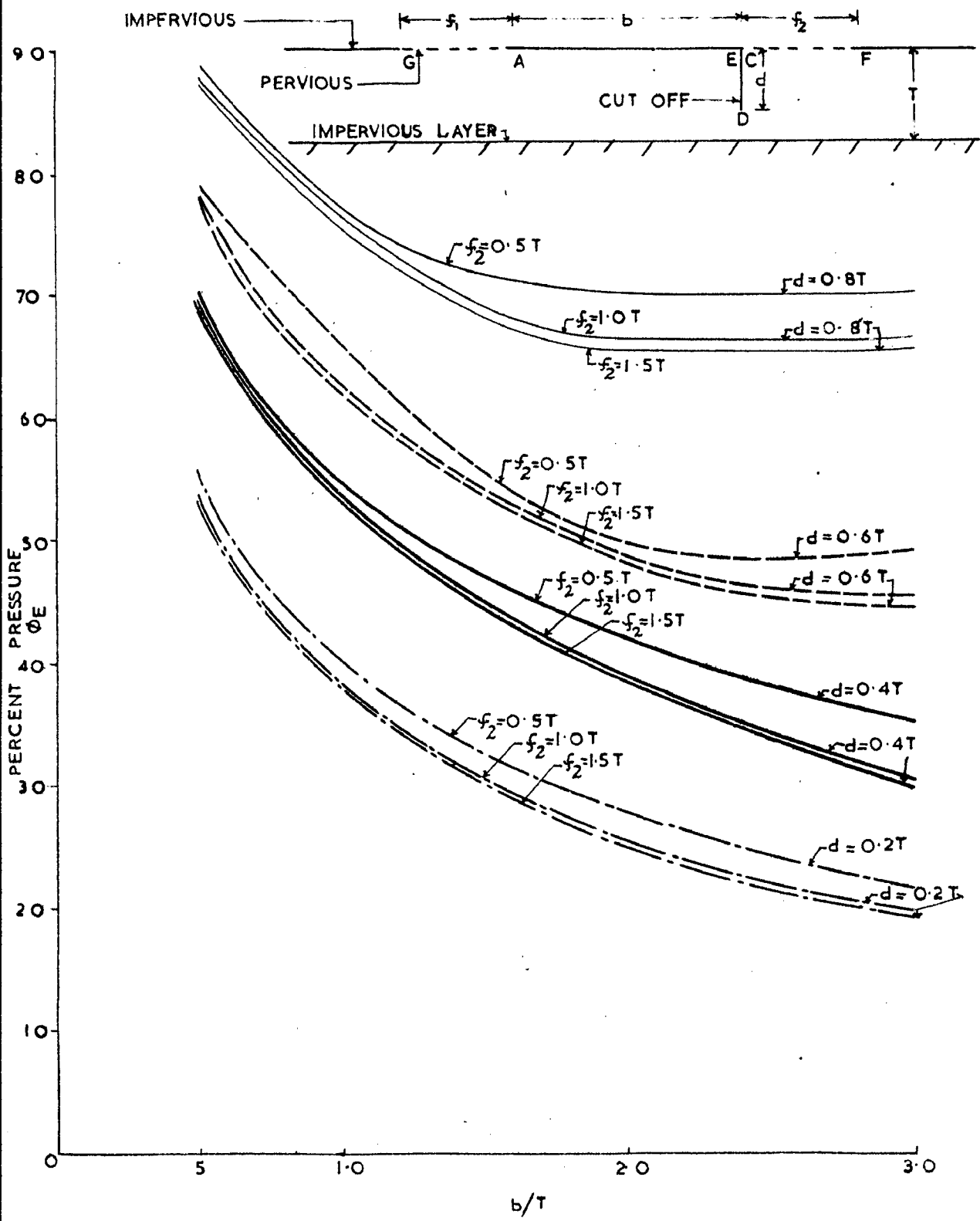


FIG. 2-16. PERCENT PRESSURE AT E FOR $f_1 = 1.0T$.

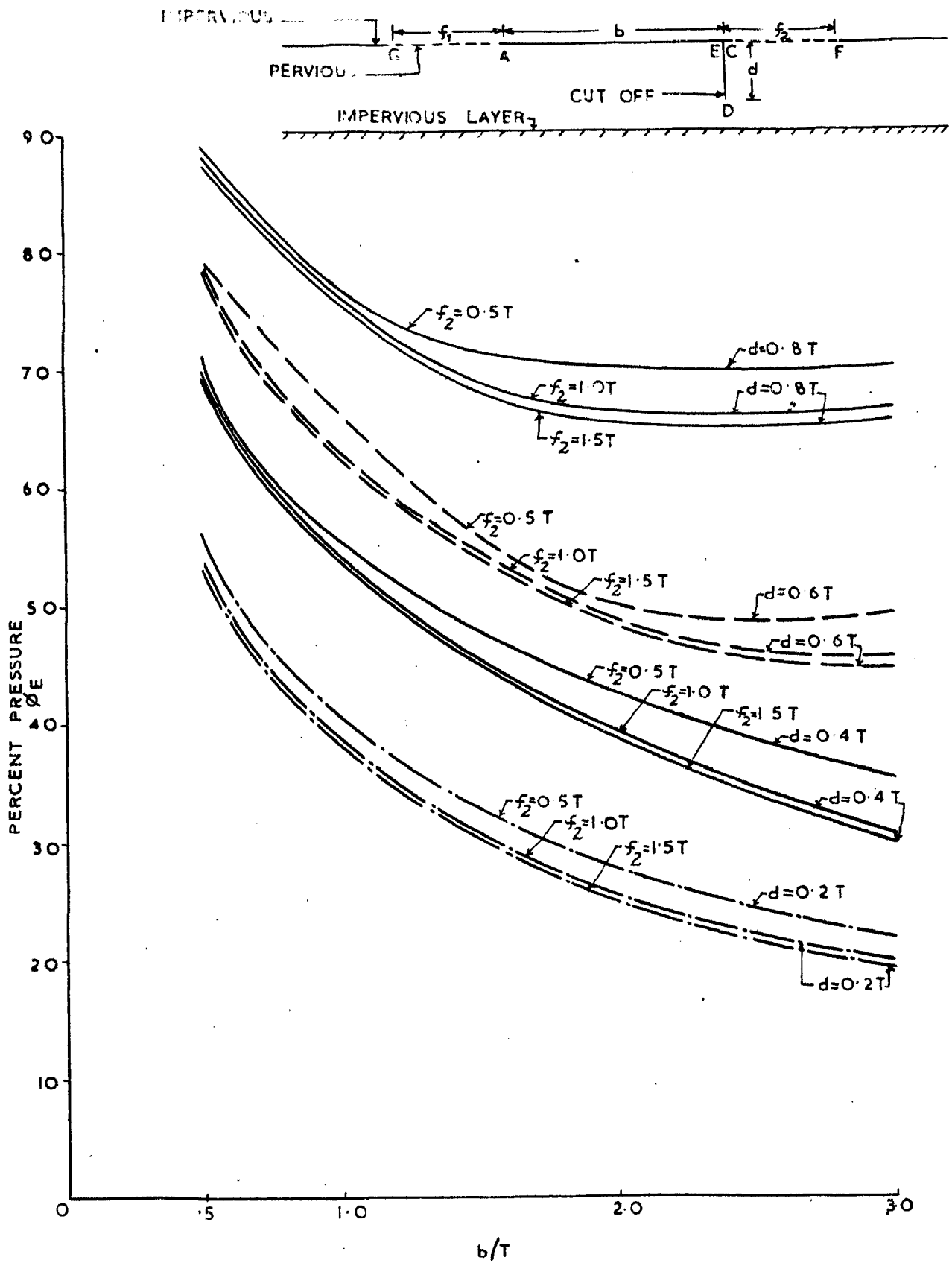


FIG. 2-17. PERCENT PRESSURE AT E FOR $f_1 = 1.5T$.

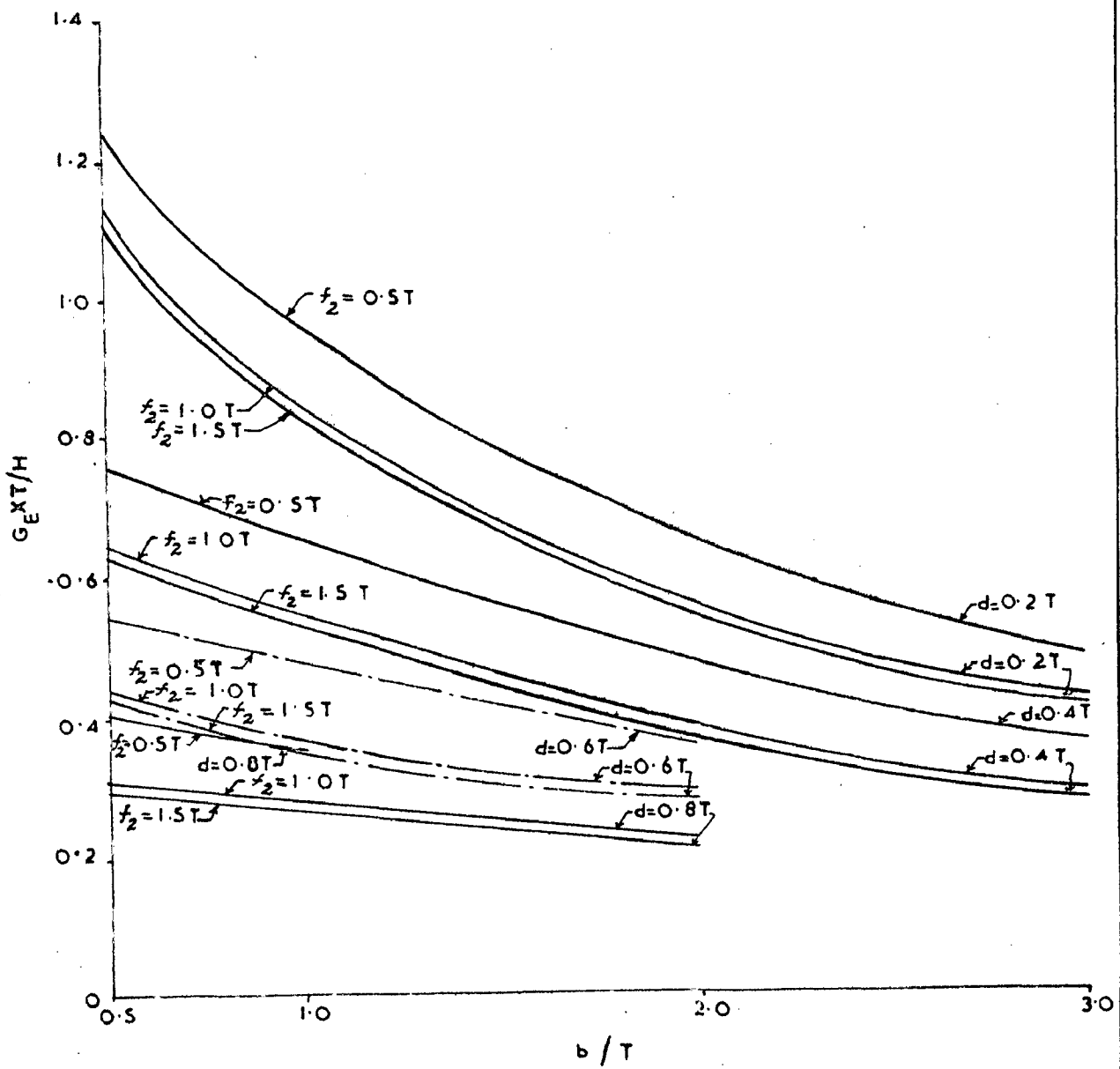
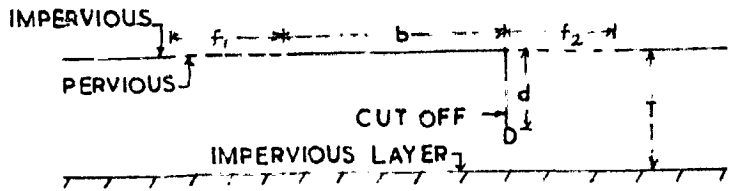


FIG. 2-18. EXIT GRADIENT FOR $f_1 = 0.5 T$.

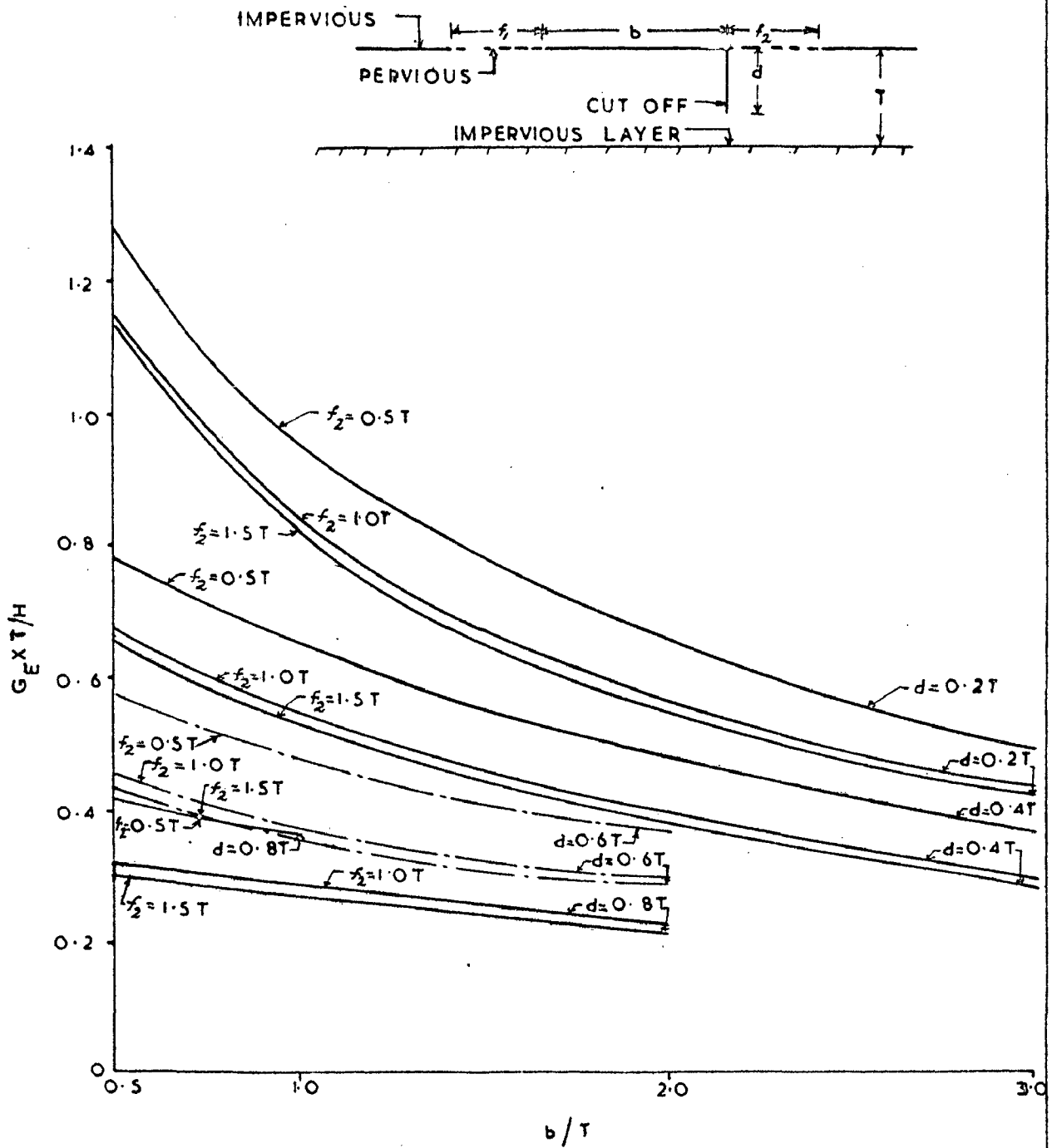


FIG 2.19. EXIT GRADIENT FOR $f_1 = 1.0 T$.

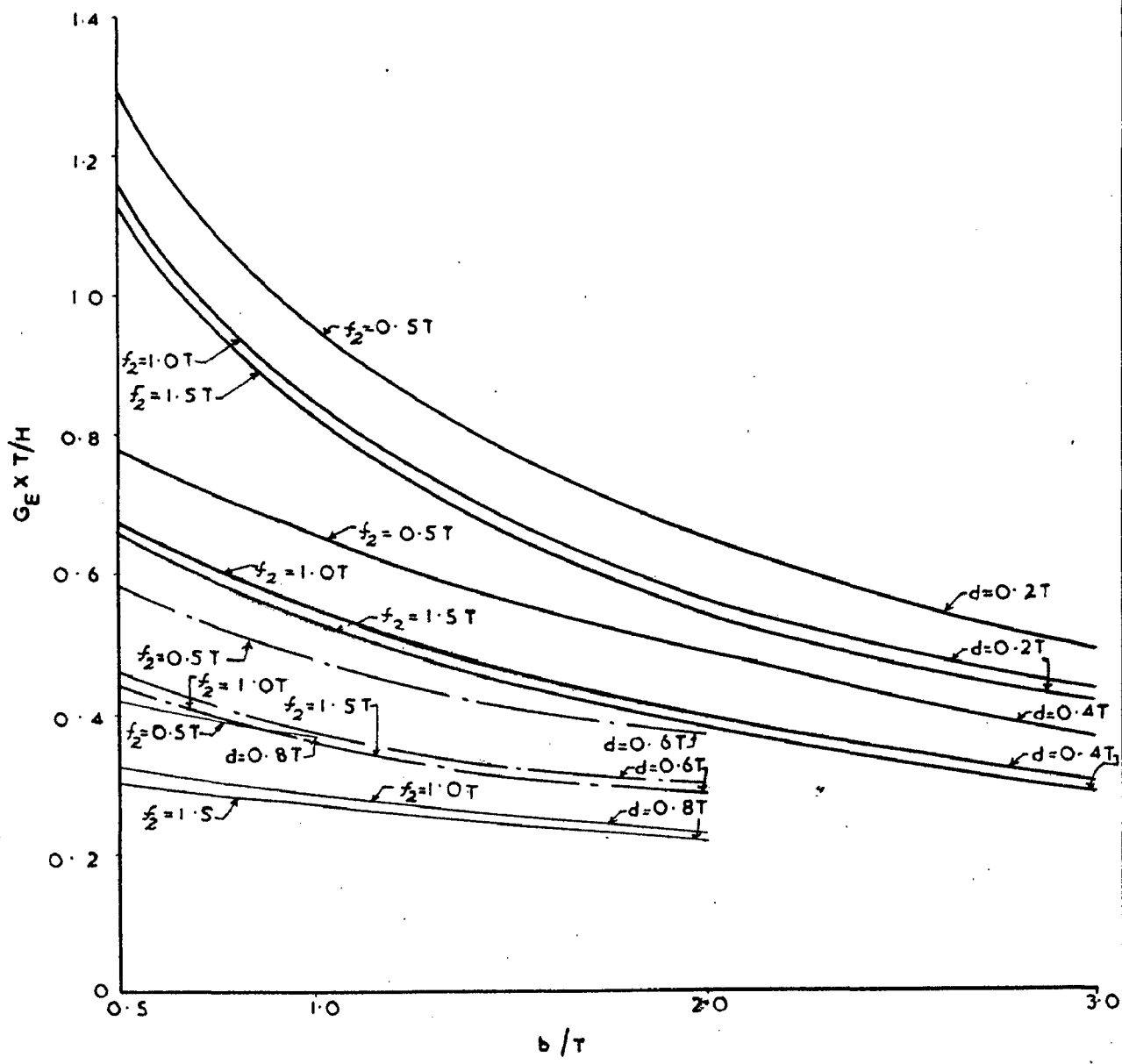
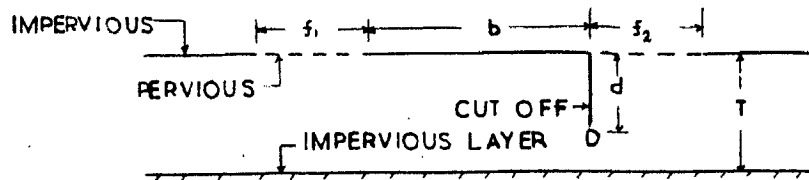


FIG.2-20. EXIT GRADIENT FOR $f_1 = 1.5T$.

CHAPTER THREE

CHAPTER III

EXPERIMENTAL VERIFICATION

3.1 TEST CONDITIONS

Results obtained from exact solutions for the following boundary conditions were verified by two dimensional electrical analogy model technique and analog field plotter.

A. Finite pervious reaches on the upstream and downstream of the structure founded on ^{In-}finite pervious flow zone.

- (i) Depth of cut off = d
- (ii) Impervious floor length upstream of the cut off, b_1 = $5.0 d$
- (iii) Impervious floor length downstream of cut off = b_2 = $5.0 d$
- (iv) Depth of permeable subsoil extends upto infinity
- (v) Length of upstream pervious bed, f_1 = $3.0 d$
- (vi) Length of downstream pervious bed, f_2 = $3 d$

B. Finite pervious reaches on the upstream and downstream of the structure founded on finite pervious flow zone.

- (i) Depth of pervious flow zone = 1
- (ii) Depth of cut off d = 0.4 and 0.6
- (iii) Length of the impervious floor on the downstream of cut off b_2 b_1 = 1.0
- (iv) Length of the impervious floor on the downstream of cut off b_2 = 1.0
- (v) Length of pervious bed on the upstream, f_1 = 0.5
- (vi) Length of pervious bed on the downstream, f_2 = 0.5

3.2 EXPERIMENTAL SET UP

3.2.1 Electrical Analogy Technique

The experimental set up consisted of a tank ~~was~~ filled with electrolytic solution. Ordinary tap water was used as electrolytic solution. The electrodes used to simulate the upstream and downstream porous reaches in the model were $1/8''$ to $1/32''$ thick copper plates. An electrical analogy model showing the foundation profile of the structure to scale was constructed and fitted in the tray. The foundation profile of the structure along with the electrical connections done on the principle of Wheat stone Bridge and the equipment needed to get the accurate mill points are shown vide figure no. 3.1

3.2.2 ANALOG FIELD PLOTTER

An analog field plotter consists of a control unit which houses the power supply, a voltage dividing potentiometer and a sensitive mill detector. Its other accessories are a conducting paper, silver paint, resistance wire, plotting board and test leads. In the sheet of conducting paper, which is the "Current sheet" of the plotter, the electric current analog field pattern is established by means of nullability attached and connected electrodes. The resulting potential drop pattern is marked directly on the conducting paper by means of a moving probe.

64956

3.3.3 The 0.004 inch thick conducting paper is made by uniform spraying carbon or graphite over a uniform sheet of paper. It has an approximate resistance of 4000 ohms per square. For experimental studies, the section of the foundations of structure was drawn to scale on the conducting paper. The previous benches were represented by silver painted electrodes having a resistance of 1 to 4 ohms/square. The sub-cell flow zone was represented by the conducting paper itself and the impervious foundation was simulated by removing the conducting paper from those portions thus rendering it nonconducting to the flow of current. The electrical connections and the model profile are shown vide Drawing no. 3.3

3.3 TEST RESULTS

3.3.1 For test condition 'A' when the structure is founded on the infinite depth of pervious flow zone, the observed uplift pressures along bottom of the floor both by the dimensional electrical analogy model and analog field are given in Table I. The uplift pressures along bottom of the floor calculated from the formulae for above boundary condition are also given in the above table. The results are plotted in Figure 3.3 for comparison.

A perusal of Table I and Figure 3.3 indicates that observed uplift pressures are very close to the calculated values.

3.3.3 For the test condition 'D' when the structure is founded on the full depth of pervious flow zone, the observed uplift pressures along bottom of the floor both by two dimensional electrical analogy model and analog field plotter are given in Table II and III. The uplift pressures, along bottom of the floor, calculated from the formulae for the above boundary conditions are also given in the above tables. The results are plotted in Figure 3.4 and 3.5 for comparison. A perusal of table II and III and Figure 3.4 and 3.5 indicates that observed uplift pressures are sufficiently close to the calculated values.

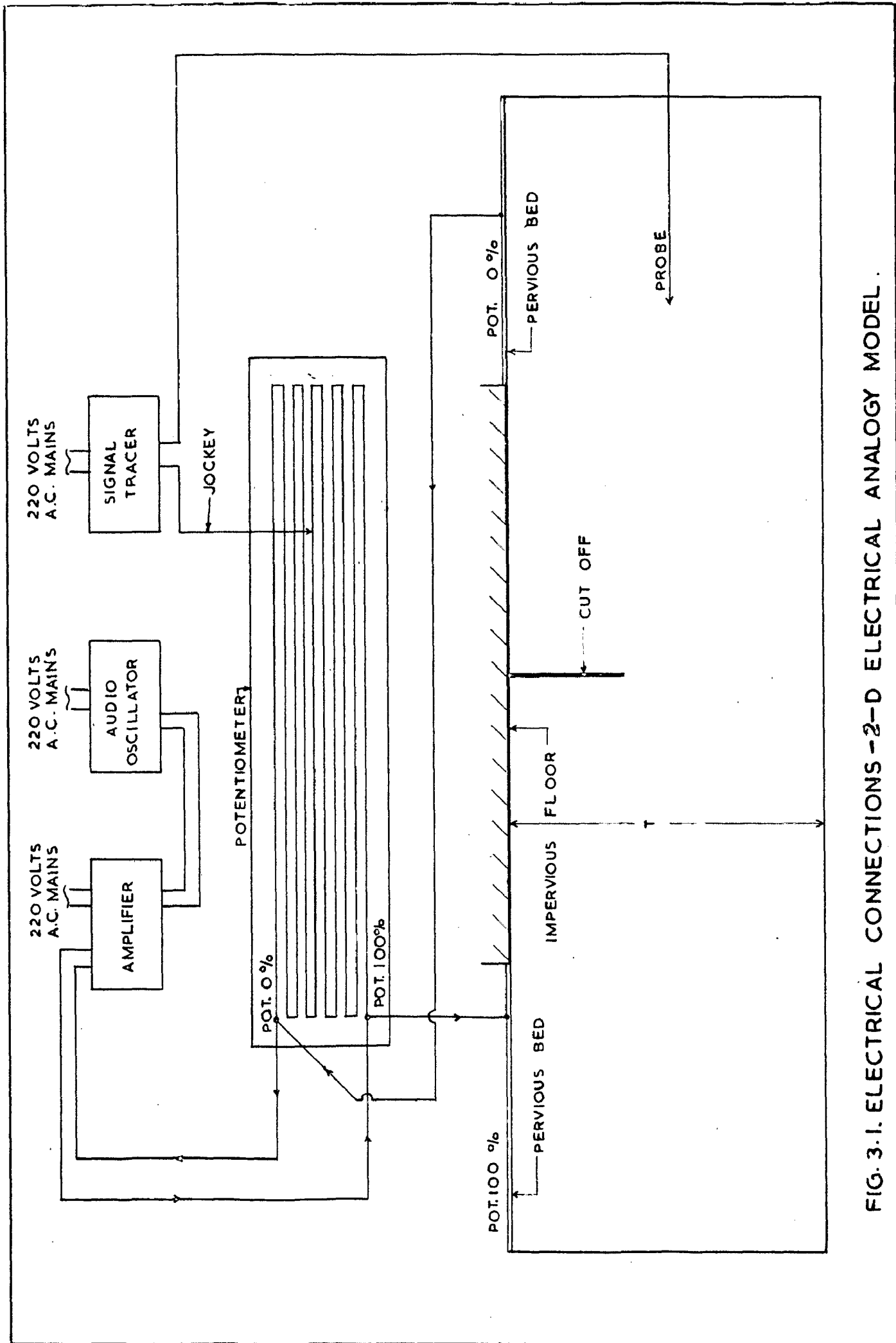


FIG. 3.1. ELECTRICAL CONNECTIONS - 2-D ELECTRICAL ANALOGY MODEL .

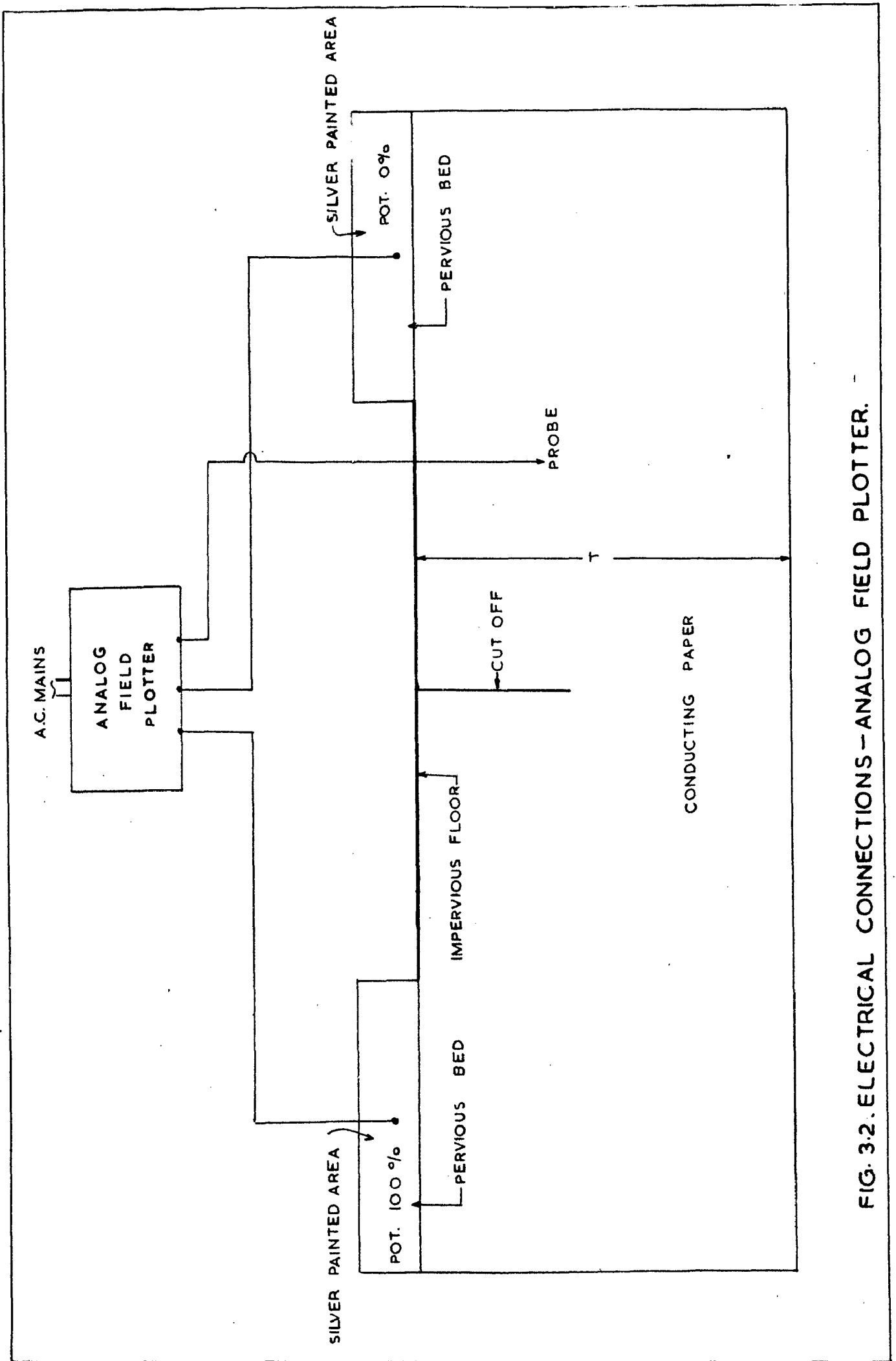


FIG. 3.2. ELECTRICAL CONNECTIONS - ANALOG FIELD PLOTTER.

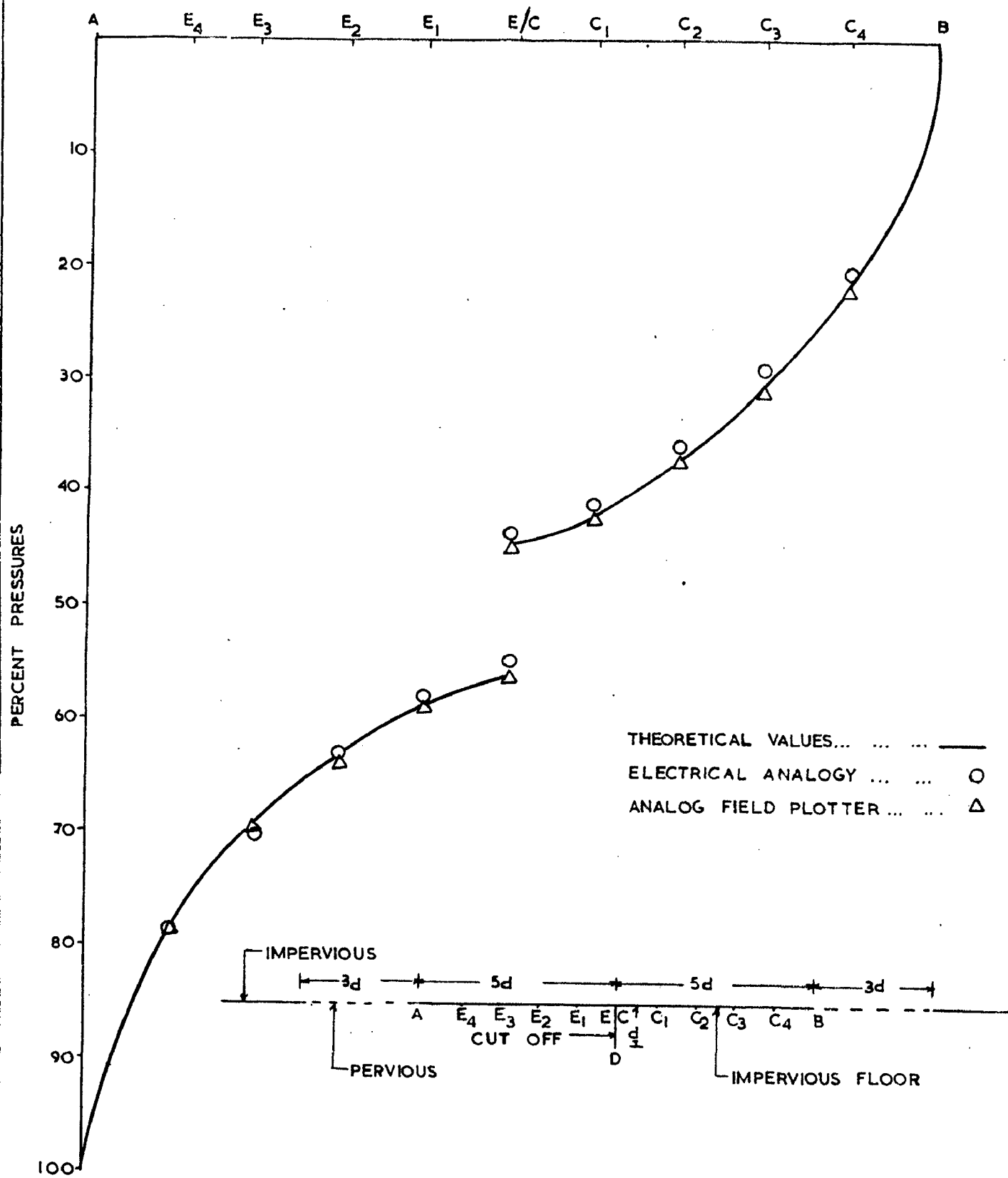


FIG. 3.3 UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATION

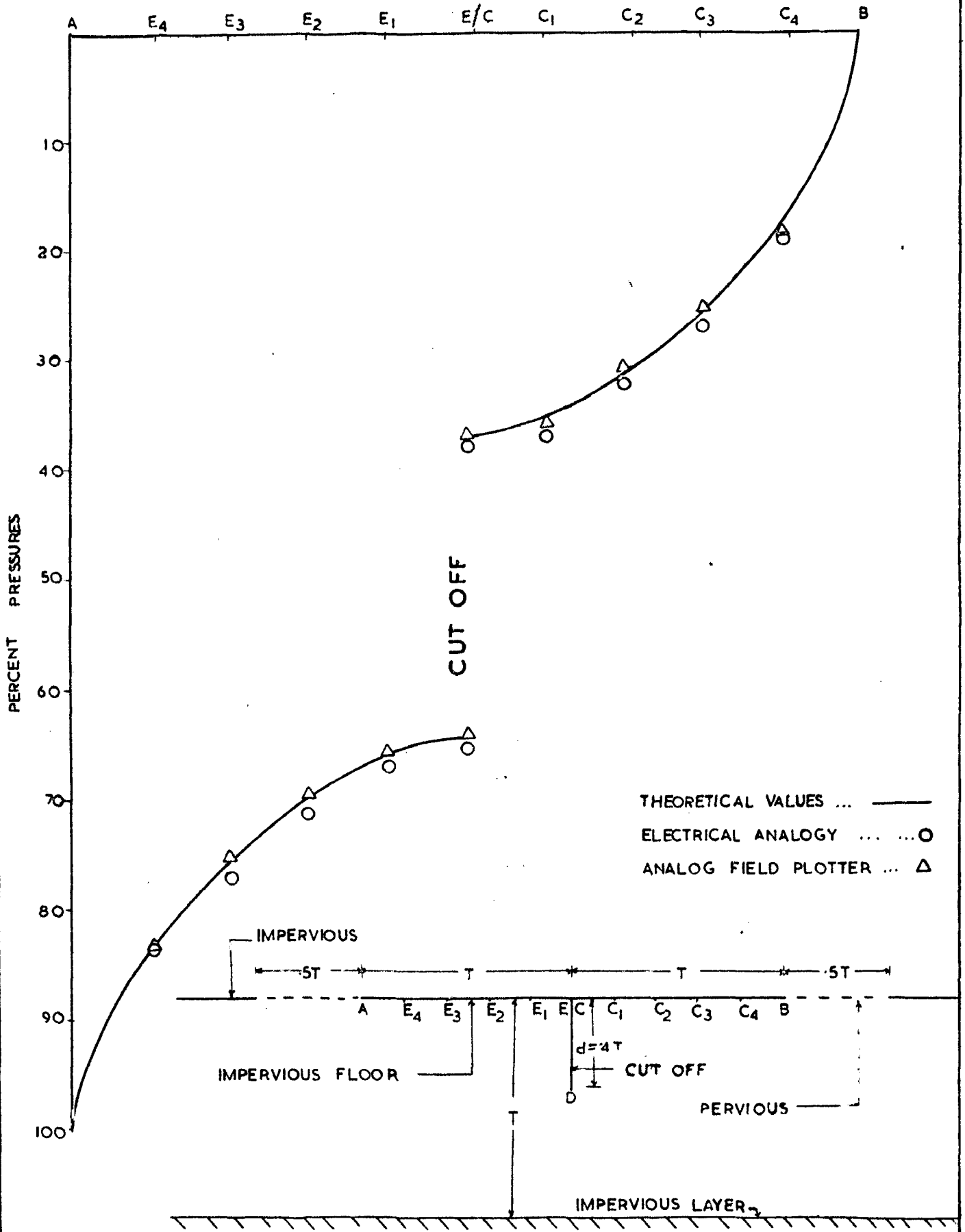


FIG. 3.4. UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATION

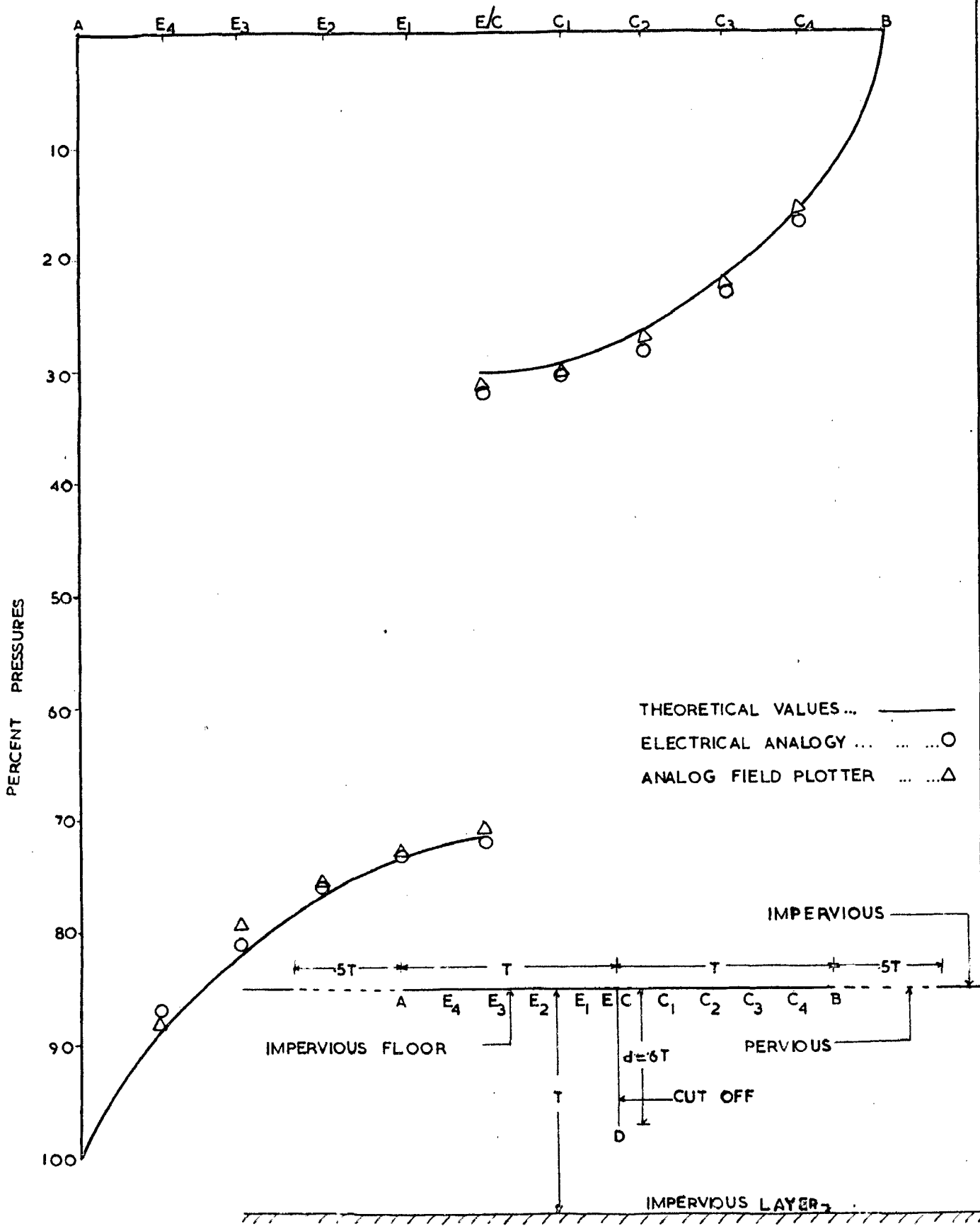


FIG.35. UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATION.

CHAPTER FOUR

CHAPTER IV

DISCUSSION OF RESULTS

4.1 A perusal of Figure 4.1 and 4.2 indicate that uplift pressures increase throughout the floor length of the structure founded on infinite as well as finite pervious flow zone with the increase in the upstream pervious reach. The increase in the uplift pressure due to increase in the upstream pervious reach is slightly more below the floor upstream of the cut off.

4.2 A perusal of Figure 4.3 and 4.4 indicate that uplift pressures decrease throughout the floor length of the structure founded on infinite as well as finite pervious flow zone with the increase in the downstream pervious reach. The decrease in uplift pressures is more on the downstream side of the floor.

4.3 When the pervious reaches on the upstream and downstream are equal, the uplift pressures increase on the floor upstream of the cut off and decrease on the floor downstream of the cut off with the increase in the pervious reaches. This is true for the structure founded on the infinite as well as finite pervious flow zone. This is evident from Figure 4.5 and 4.6. The effect is more pronounced on the floor downstream of the cut off.

4.4 A perusal of figure no. 4.7 indicates that for the same length of pervious reaches both on the upstream and downstream of the structure founded on the finite depth of pervious flow

zone the uplift pressures below the floor upstream of the cut off increase where as those on the floor downstream of cut off decrease with the increase in the depth of cut off.

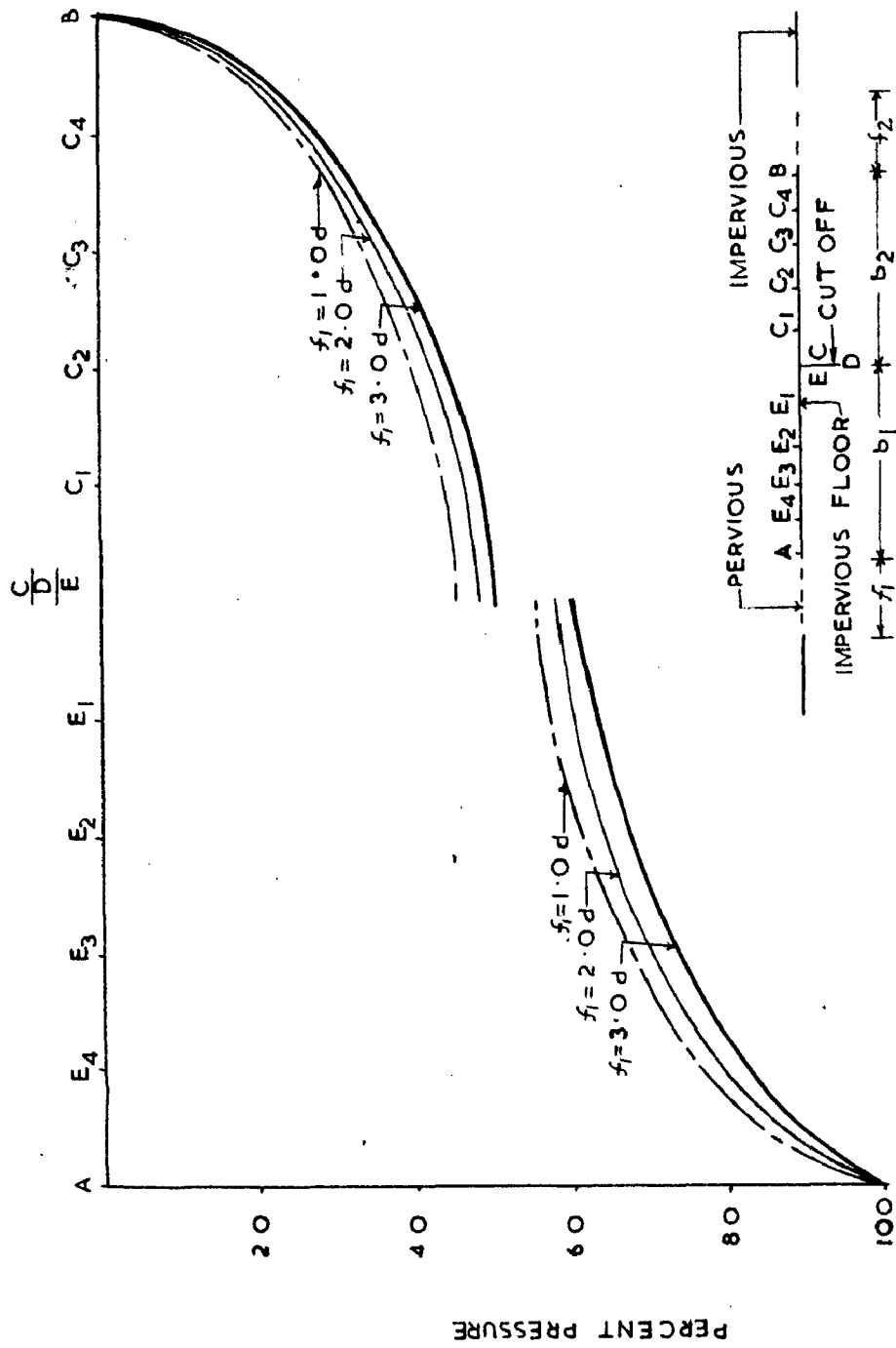


FIG. 4. UPLIFT PRESSURES ON THE BOTTOM OF FOUNDATIONS FOR $f_2 = 1.0 d$.

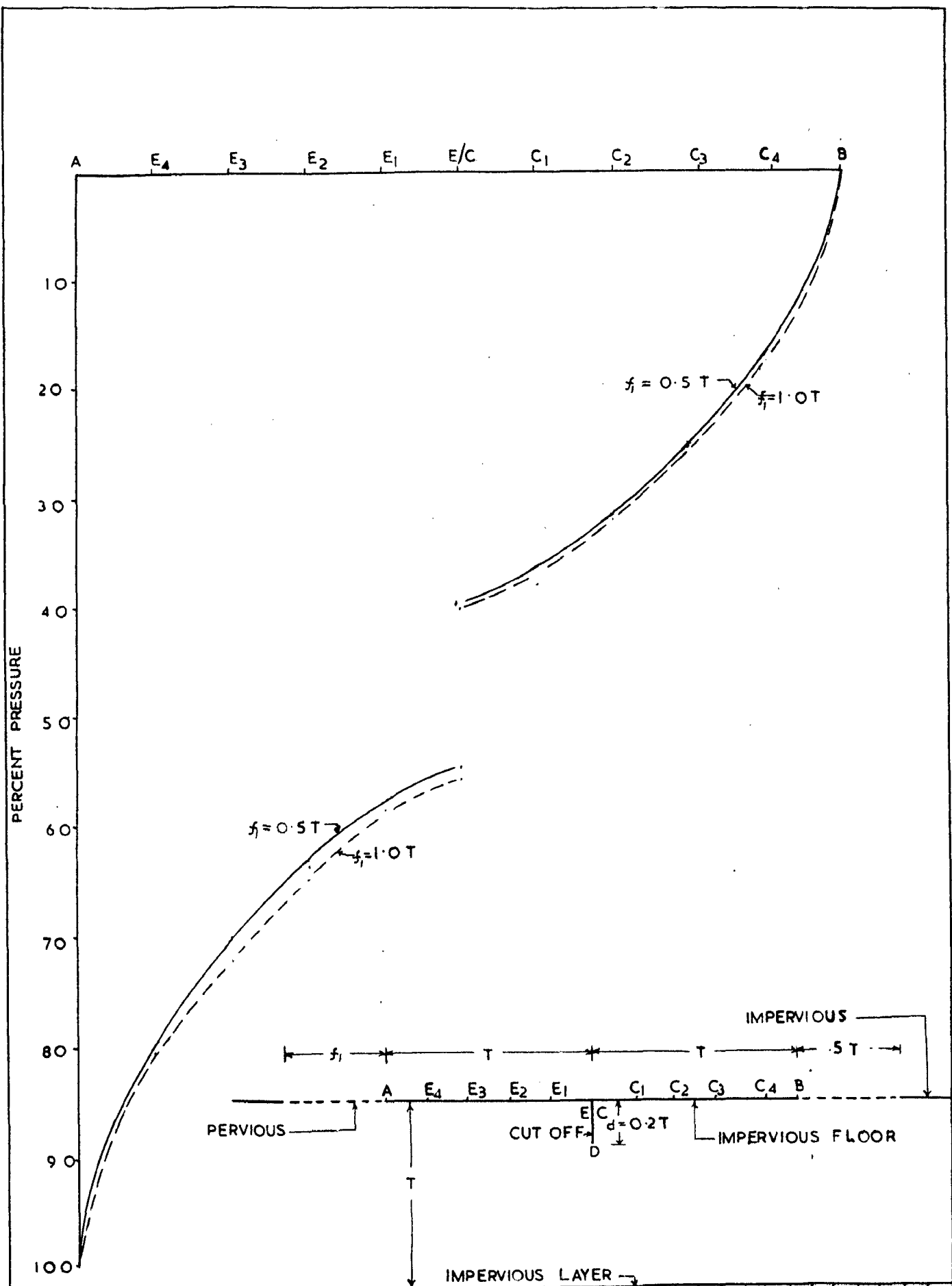


FIG.42.UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATIONS

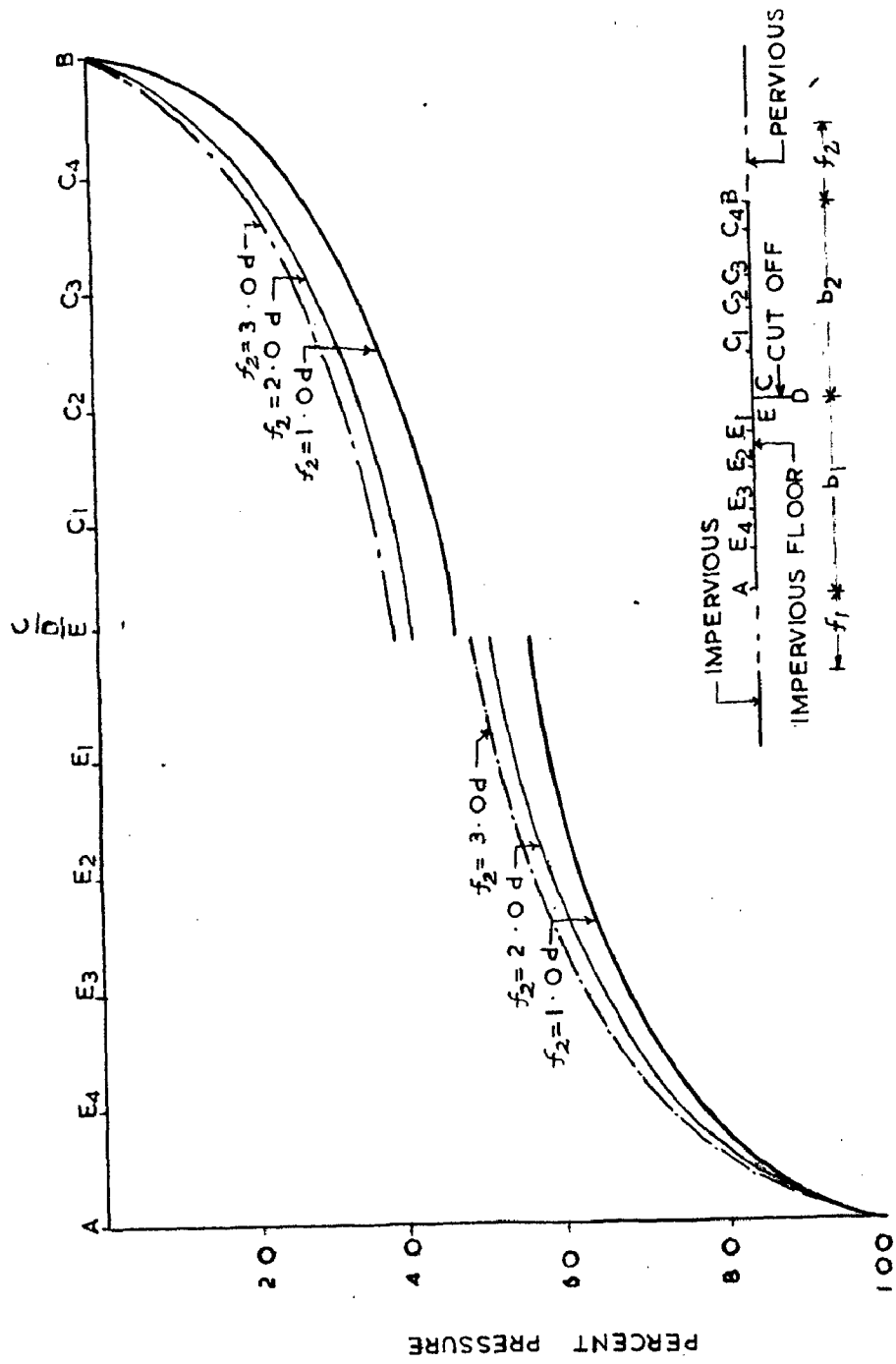


FIG. 4-3. UPLIFT PRESSURES ON THE BOTTOM OF FOUNDATIONS FOR $f_1 = 1.0d$.

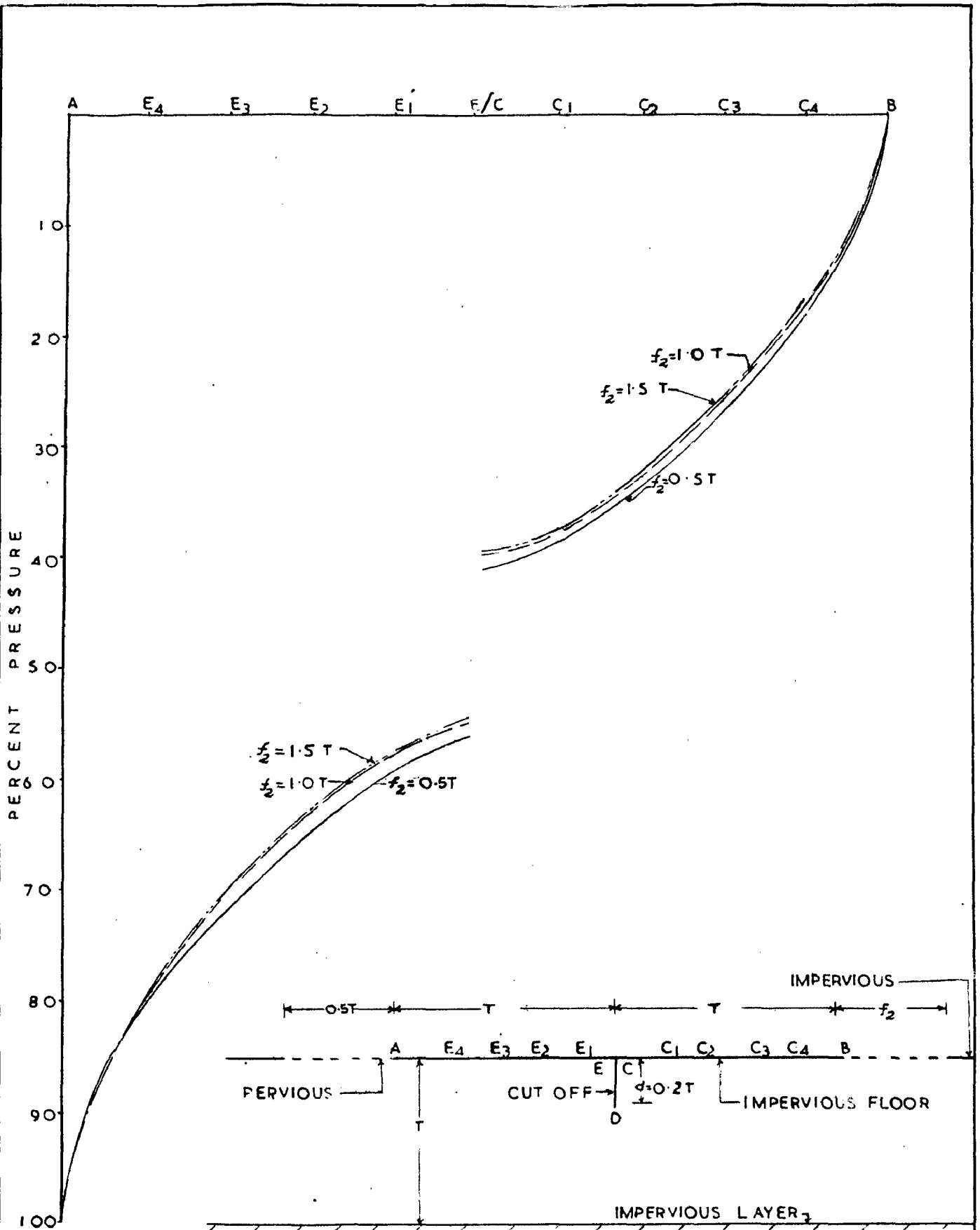


FIG. 4.4. UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATION

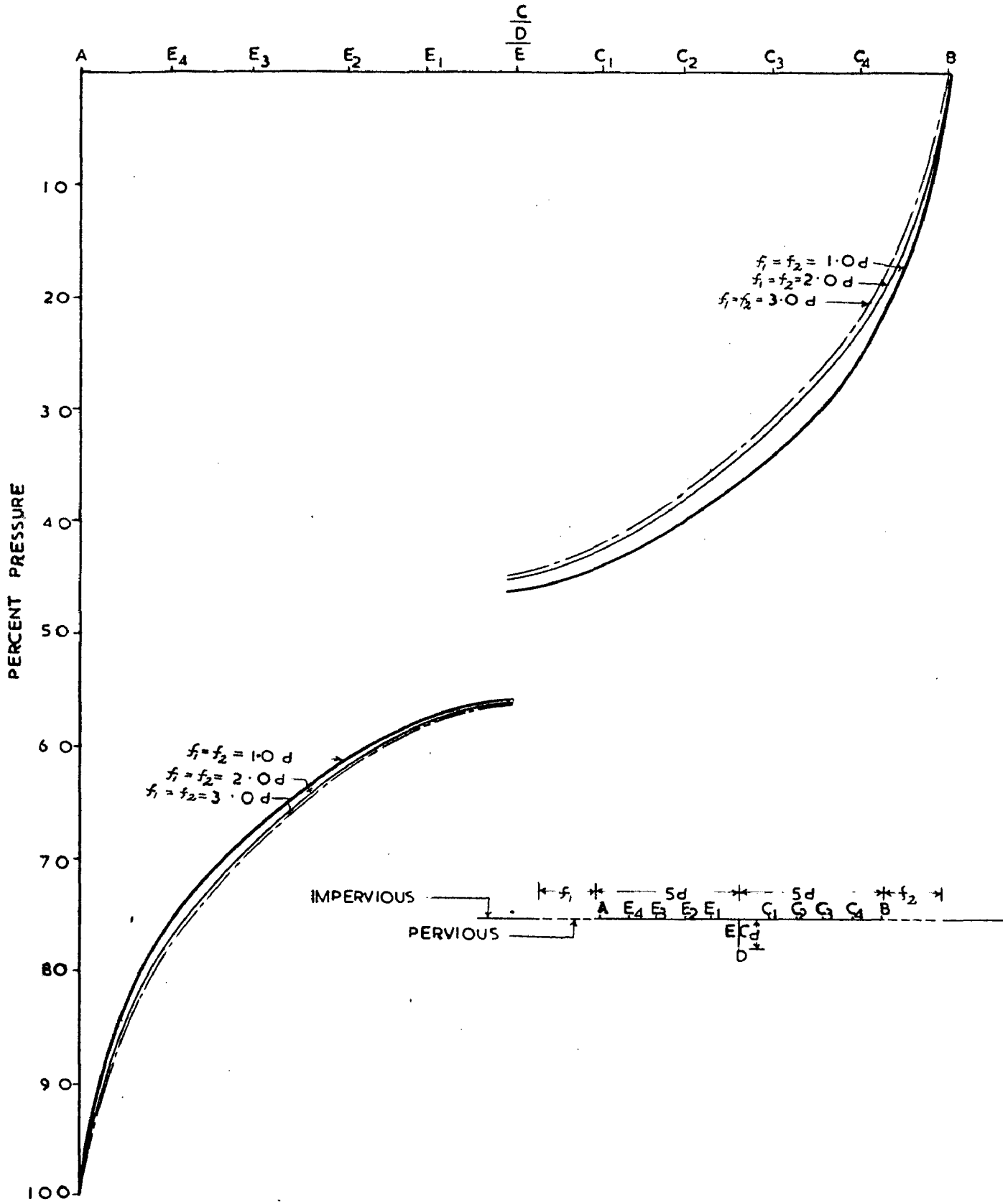


FIG.4.5.UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATIONS

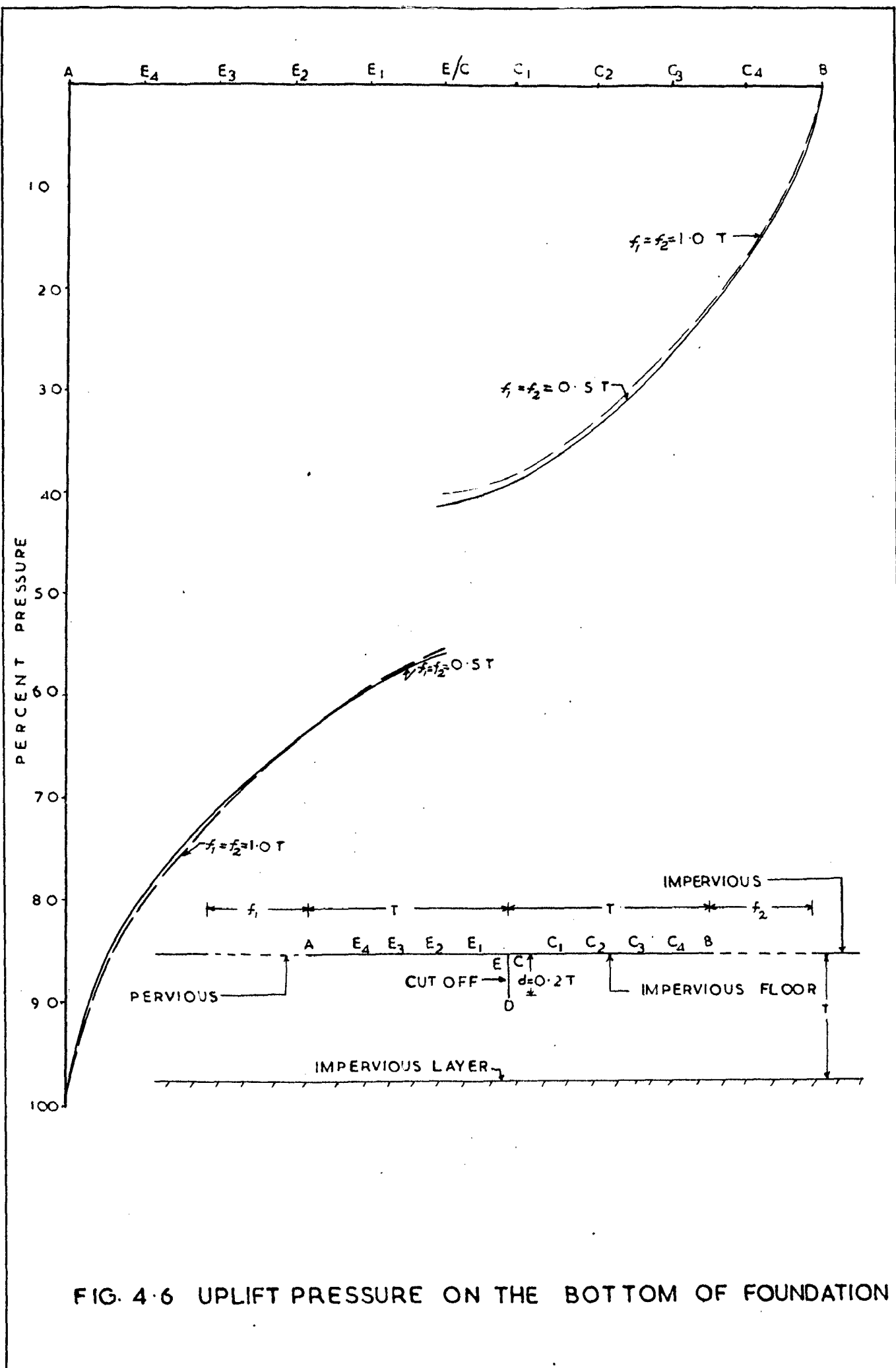


FIG. 4.6 UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATION

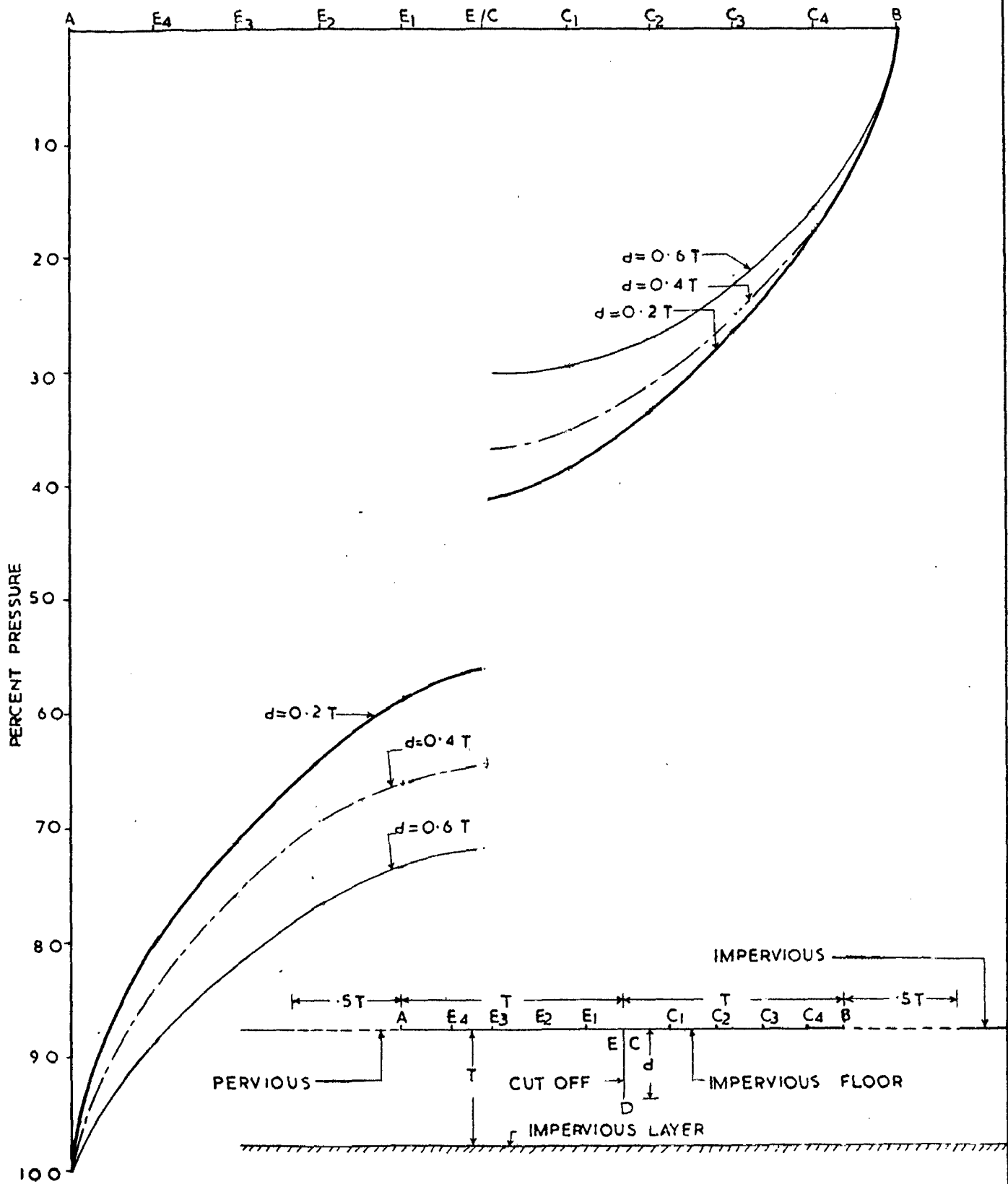


FIG. 4.7. UPLIFT PRESSURE ON THE BOTTOM OF FOUNDATIONS .

CHAPTER FIVE

CHAPTER V

C O N C L U S I O N S

1. Mathematical solution is given to find out exit gradient and uplift pressure on the bottom of foundation of hydraulic structures founded on infinite and finite pervious flow beds.
2. Computations have been made to find out values of exit gradient and uplift pressures at the key points of an elementary profile of a hydraulic structures with various combinations of the variables involved. The computed results have been plotted in the form of curves. These curves can be used for practical design purposes.
3. The computed uplift pressures along the foundation of three simple and elementary profiles of hydraulic structures have been verified by Analog Field Plotter and on two dimensional electrical analogy model. The observed pressures are close to the calculated values.
4. The uplift pressures below the hydraulic structure increase with the increase of upstream pervious reach and decrease with the increase of downstream pervious reach.
5. For the structure founded on finite depth of pervious flow some the uplift pressures below the floor upstream of the cut off increase whereas those on the floor downstream of cut off decrease with the increase in the depth of cut off.

TABLE I

EXPERIMENTAL VERIFICATION
UP-LINE PRESSURES ON THE BOTTOM OF FOUNDATION

Total Conditions :-

- (i) Depth of permeable sub-soil extends upto infinity.
- (ii) Depth of cut off = 4.
- (iii) Impervious floor length upstream of the cut off b1 = 5.04
- (iv) Impervious floor length downstream of the cut off b2 = 5.04
- (v) Length of upstream pervious bed f1 = 3.04
- (vi) Length of downstream pervious bed f2 = 3.04

Sl. No.	Name of the points	Distance from D.S. edge of the structure	Uplift pressures in ft		
			Calculated value	Observed on Electrical analogy model	Analog field plotter
1.	C ₄	4	22.0	20.5	22.0
2.	C ₃	24	30.4	29.0	30.0
3.	C ₂	34	37.2	36.0	37.2
4.	C ₁	44	42.1	41.0	42.1
5.	C	54	44.7	43.5	44.6
6.	D	64	46.7	45.6	46.3
7.	E	64	55.9	54.0	56.1
8.	E ₁	64	60.1	59.0	60.0
9.	E ₂	74	63.2	63.0	63.7
10.	E ₃	84	69.3	70.0	69.0
11.	E ₄	94	78.3	78.0	78.4

TABLE II

EXPERIMENTAL VERIFICATION
UPLIFT PRESSURES ON THE BOTTOM OF FOUNDATION

Test Conditions :-

- (i) Depth of pervious flow zone T
- (ii) Depth of cut off $d = 0.4T$
- (iii) Impervious floor length upstream of the cut off $b_1 = 1.0T$
- (iv) Impervious floor length downstream of the cut off $b_2 = 1.0T$
- (v) Length of upstream pervious bed $f_1 = 0.5T$
- (vi) Length of downstream pervious bed $f_2 = 0.5T$

Sl. No.	Name of the points	Distance from D.S. edge of the structure	Uplift pressures %		
			Theoretically calculated values	By Biological analogy experiments	By Analog field plotter
1.	C_4	$.2$	17.8	19.0	18.2
2.	C_3	$.4$	28.8	27.0	28.2
3.	C_2	$.6$	30.0	32.0	30.0
4.	C_1	$.8$	33.8	37.0	35.8
5.	C	T	36.9	37.0	36.0
6.	D	T	49.8	54.0	49.8
7.	E	T	64.1	65.2	63.0
8.	E_1	1.2	65.8	67.0	65.4
9.	E_2	1.4	69.2	71.0	69.2
10.	E_3	1.6	75.0	77.0	75.8
11.	E_4	1.8	83.8	85.0	83.8

TABLE III

EXPERIMENTAL VERIFICATION
 UPLIFT PRESSURES ON THE BOTTOM FOUNDATION

Soat Conditions :-

- (i) Depth of pervious flow zone = 2
- (ii) Depth of cut off $d = 0.6$
- (iii) Impervious floor length upstream of the cut off $B_1 = 1.0$
- (iv) Impervious floor length downstream of the cut off $B_2 = 1.0$
- (v) Length of upstream pervious bed $L_1 = 0.5$
- (vi) Length of downstream pervious bed $L_2 = 0.5$

Sl. No.	Name of the points	Distance from D.S. edge of the structure	UPLIFT PRESSURE %		
			Theoretically calculated value	By Electrical analogy	By Analog field plotter
1.	C ₄	.2	15.0	17.0	16.0
2.	C ₃	.4	21.7	23.0	22.3
3.	C ₂	.6	27.1	29.1	27.0
4.	C ₁	.8	29.6	30.8	30.0
5.	C	1	30.8	32.0	31.3
6.	D	2	50.9	54.6	51.6
7.	E	2	71.6	71.0	70.0
8.	E ₁	1.2	73.3	73.0	72.8
9.	E ₂	1.4	76.0	76.1	75.8
10.	E ₃	1.6	79.0	81.0	79.2
11.	E ₄	1.8	83.8	87.0	88.2

A P P E N D I X I

APPENDIX I

NOTES ON ELLIPTIC FUNCTIONS

For finding a mathematical solution of problems of seepage below hydraulic structure with complex boundary conditions, use of elliptic integrals is made for solving the differential equation.

The existence of the integral of a continuous function is assured; however, it does not necessarily follow that the integral can be expressed by elementary functions alone. In ground water problems, as a consequence of the Schwarz-Christoffel transformations, such functions are more commonly encountered.

Liouville was the first to demonstrate the existence of genuine elliptic integrals viz. elliptic integrals that can not be expressed in terms of elementary functions. The main classical work on the subject is Legendre's "Traite des fonctions elliptiques" in which the analysis of all elliptic integrals is reduced to the following three normal forms

$$F(x, k) = \int_0^x \frac{dx}{\sqrt{(1-x^2)(1-k^2x^2)}} \quad \dots(1a)$$

$$E(x, k) = \int_0^x \frac{x^2 \cdot dx}{\sqrt{(1-x^2)(1-k^2x^2)}} \quad \dots(1b)$$

$$II(x, k, n) = \int_0^x \frac{dx}{(x^2+n)\sqrt{(1-x^2)(1-k^2x^2)}} \quad \dots(1c)$$

These basic types are termed as first, second and third kinds. For ready made solutions are available for their calculations, as a function of x , k and n . For the modulus, k the limits of the tabulated values are $0 \leq k^2 \leq 1.0$

As regards x , it may have any value (real, imaginary, or complex), but in usual tables real values only are listed, for $0 \leq x^2 \leq 1.0$

If the upper limit of integration is made $x = 1.0$, the elliptic integral is denoted as complete i.e. considering equation (1) it becomes

$$K = F(x, k) = \int_0^x \frac{dx}{\sqrt{(1-x^2)(1-k^2x^2)}} \dots (2a)$$

and $K' = F(x, k') = \int_0^x \frac{dx}{\sqrt{(1-x^2)(1-k'^2x^2)}} \dots (2b)$

where k' is the co-modulus or complementary modulus defined by

$$k' = \sqrt{1 - k^2}$$

The "Jacobi's functions" are defined as follows considering the elliptic integrals of first kind

Let

$$U = F(x, k)$$

$$= \int_0^x \frac{dx}{\sqrt{(1-x^2)(1-k^2x^2)}} \dots (3a)$$

substituting

$$\sin \phi = x \dots (3b)$$

$$U = \int_0^\phi \frac{d\phi}{\sqrt{1 - k^2 \sin^2 \phi}} \dots (3c)$$

Jacobi uses the term "amplitude" to denote the function inverse to this integral and he uses the symbol $\phi = \sin u$, to represent it.

Then, according to equation (3.b)

$$x = \sin \alpha u$$

In referring to elliptic functions, a more common notation is used, viz.:

$$\operatorname{sn} u = \sin \alpha u \quad (\text{sinus amplitudinis } u) \quad \dots(4)$$

In addition to this Jacobi introduced two other functions, which are correlated with $x = \sin \alpha u$ in the following manner

$$\sqrt{1 - x^2} = \cos \alpha u = \operatorname{cn} u$$

$$\sqrt{1 - k^2 x^2} = \Delta \alpha u = \operatorname{dn} u$$

These functions are termed "Jacobi's elliptic functions".

PROPERTIES OF JACOBI'S ELLIPTIC FUNCTIONS

Referring to equation (3.a), since $u = 0$, when $x=0$

$$\begin{aligned} \operatorname{sn} 0 &= 0 \\ \text{or } \operatorname{sn}^{-1} 0 &= 0 \end{aligned} \quad \dots(5a)$$

Considering equation (2.a), since $u = K$, when $x=1$

$$\begin{aligned} \operatorname{sn} K &= 1 \\ \text{or } \operatorname{sn}^{-1} 1 &= K \end{aligned} \quad \dots(5.b)$$

In the same way, $u = iK'$, when $x = \infty$

$$\begin{aligned} \operatorname{sn} iK' &= \infty \\ \text{or } \operatorname{sn}^{-1} \infty &= iK' \end{aligned} \quad \dots(5.c)$$

Similarly $W = K + iK'$, when $x = 1/k^2$

$$\operatorname{Sn}(K + iK') = 1/k^2$$

$$\operatorname{Sn}^{-1}(1/k^2) = K + iK' \quad \dots\dots(5.d.)$$

These four points given by 5.a, 5.b, 5.c, and 5.d are the four corners in order of the rectangle in W plane (see main text)

A P P E N D I X I I

APPENDIX II

REPRESENTATION OF A RECTANGLE UPON A HALF PLANE

The w -plane represented in figure 4.4 is transformed into the ζ -plane as indicated in fig. 4c. Corresponding points in the w and ζ -planes are indicated by the same letters.

The integral equation (2.10) corresponding to zone AB in the ζ -plane as given in the text, is

$$W = A \int \frac{d\zeta}{(\zeta_G - \zeta)(\zeta_B - \zeta)(\zeta - \zeta_A)(\zeta - \zeta_F)} \quad \dots(1)$$

$$\text{where } \zeta_G > \zeta_B > \zeta > \zeta_A > \zeta_F$$

In order to integrate this we first make a bilinear transformation of the ζ -plane into t -plane. There is just one bilinear transformation that maps three given distinct points t_1 , t_2 and t_3 in w -plane respectively. The transformation is given by

$$\frac{(t - t_1)(t_2 - t_3)}{(t - t_3)(t_2 - t_1)} = \frac{(\zeta - \zeta_1)(\zeta_2 - \zeta_3)}{(\zeta - \zeta_3)(\zeta_2 - \zeta_1)} \quad \dots(2)$$

The corresponding points in t and ζ -plane in our case are as below

$$\begin{array}{ll} t_1 = 0 & \zeta_1 = \zeta_B \\ t_2 = 1 & \zeta_2 = \zeta_A \\ t_3 = \infty & \zeta_3 = \zeta_G \\ t_4 = 1/t_3 & \zeta_4 = \zeta_F \end{array} \quad \dots(3)$$

The transformation then, can be written

$$t = \frac{(\zeta - \zeta_B)(\zeta_A - \zeta_G)}{(\zeta - \zeta_G)(\zeta_A - \zeta_B)}$$

and $1-t = 1/k^2 t = \frac{(\zeta_F - \zeta_B)(\zeta_A - \zeta_G)}{(\zeta_F - \zeta_G)(\zeta_A - \zeta_B)}$

from which it follows

$$1-t = \frac{(\zeta_A - \zeta)(\zeta_B - \zeta_G)}{(\zeta - \zeta_G)(\zeta_A - \zeta_B)}$$

and $1-k^2 t = \frac{(\zeta_G - \zeta_B)(\zeta - \zeta_F)}{(\zeta_F - \zeta_B)(\zeta - \zeta_G)}$

and $dt = \frac{(\zeta_G - \zeta_B)(\zeta_F - \zeta_G)}{(\zeta_F - \zeta_B)(\zeta - \zeta_G)} d\zeta$

and hence after little reduction equation (1) becomes

$$u = \frac{BA}{\sqrt{(\zeta_B - \zeta_F)(\zeta_G - \zeta_A)}} \int_0^t \frac{dt}{\sqrt{4t(1-t)(1-k^2 t)}} + B \dots\dots(4)$$

where integration constant B is zero in this case

If we put $M = \frac{1}{2} \sqrt{(\zeta_B - \zeta_F)(\zeta_G - \zeta_A)}$

$$M/A u = \int_0^t \frac{dt}{\sqrt{4t(1-t)(1-k^2 t)}} \dots\dots(5)$$

If we introduce intermediate variable x,

such that $x = \sqrt{t}$, the above equation (5) becomes

$$M/A u = \int_0^x \frac{dx}{\sqrt{(1-x^2)(1-x^2 k^2)}}$$

Since the above equation (6) is of the form of elliptic integral of the first kind, using the results of equation (6) of Appendix I

we can write

$$\sin \frac{M}{A} W = x$$

or $\sin^2 \frac{M}{A} W = x^2$

On substituting the value of $x^2 = t$, we arrive at

$$\sin^2 \frac{M}{A} W = \frac{(S - S_B)(S_A - S_G)}{(S - S_G)(S_A - S_B)} \dots\dots(7)$$

The above equation (7) is the same as equation (2.15) and equation (2.43) of the text.

APPENDIX III

APPENDIX III

NOTATIONS

The following symbols are used in this dissertation:

- D_1, D_2 = upstream and downstream impervious floor lengths respectively.
- d = depth of cut off
- L_1, L_2 = upstream and downstream pervious lengths respectively.
- k = modulus of the elliptic integral
- k' = co-modulus defined as $\sqrt{1-k^2}$
- q = seepage discharge per unit width
- t = $x + iy$, complex variable representing auxiliary semi-infinite plane
- U_x, U_y = velocity components in x and y directions respectively
- U = complex potential, i.e. $\phi + i\psi$, complex plane representing rectangular flow field.
- Z = complex potential, i.e. $x + iy$, complex plane representing physical plane
- A, C = Magnification factors in Schwarz-christoffel transformation
- B, D = constants of integration
- H = Head difference between upstream and downstream water levels.
- L_1, L_2 = transformation parameters.
- K, K' = complete elliptic integrals of the first kind, with modulus k and k' , respectively.
- T = Depth of permeable flow zone
- σ, β, γ = transformation parameters.
- ζ = $\xi + i\eta$, complex variable representing auxiliary semi-infinite plane
- ξ, η = real and imaginary axes in ζ -plane
- K = permeability coefficient

ϕ' = potential function

ϕ = reduced value of ϕ'

ψ' = stream function

ψ = reduced value of ψ' , and

Sn = elliptic sine

A N N E X U R E I

```

C C A S CHAWLA 7 IRI UPLIFT PRESSURES FOR BOUNDARY CONDITION A
DIMENSION AM(20),SN(11,20),U(11),SNFX(11)
READ 10,(AM(J),J=1,20)
10 FORMAT (10F7.5)
READ 20,((SN(I,J),J=1,20),I=1,11)
20 FORMAT (10F7.5)
READ 30,(U(I),I=1,11)
30 FORMAT (11F6.4)
C UFOR=B1/D
C DFOR=B2/D
C EL IS THE LOCATION OF KEY POINT IN T-PLANE
200 READ 700,EL,UFOR
700 FORMAT(2F10.4)
C UFIR=F1/D
C DFIR=F2/D
40 UFIR=1.0
50 DFIR=1.0
ELON=SQRTF(1.0+(UFOR)**2)
ELT=1.0
DON=SQRTF(1.0+(UFOR+UFIR)**2)
60 DT=SQRTF(1.0+(DFIR)**2)
BON=1.0-2.0*(DON-ELON)/(DON+DT)
BT=1.0-2.0*(DT-ELT)/(DON+DT)
T=1.0-2.0*(DT-EL)/(DON+DT)
AMX=SQRTF(2.0*(BON+BT)/(1.0+BON+BT+BON*BT))
SNFU=SQRTF((BT+BON*BT-T-T*BON)/(BON+BT-BON*T-BT*T))
IF (AMX-.9999) 11,11,12
12 J=20
GO TO 2
11 DO1 J=1,20
IF (AMX-AM(J)) 3,2,1
1 CONTINUE
2 K=J
DO 4 I=1,11
IF (SNFU-SN(I,K)) 6,5,4
4 CONTINUE
5 M=I
PHI=U(M)
GO TO 100
6 M=I
PHI=U(M-1)+(U(M)-U(M-1))/(SN(M,K)-SN(M-1,K))*(SNFU-SN(M-1,K))
GO TO 100
3 K=J
DO 7 I=1,11
A=AMX-AM(K-1)
SNFX(I)=SN(I,K-1)+(SN(I,K)-SN(I,K-1))/(AM(K)-AM(K-1))*A
7 CONTINUE
8 DO 9 I=1,11
IF (SNFU-SNFX(I)) 14,5,9
9 CONTINUE
14 M=I
PHI=U(M-1)+(U(M)-U(M-1))/(SNFX(M)-SNFX(M-1))*(SNFU-SNFX(M-1))
100 PUNCH 300,EL,UFOR,UFIR,DFIR,AMX,SNFU,PHI
300 FORMAT (7F10.7)

```



```

IF(DFIR-6.0)70,80,80
70 DFIR=DFIR+1.0
GO TO 60
80 IF(UFIR-3.0)90,200,200
90 UFIR=UFIR+1.0
GO TO 50
END

```

00175	00872	01737	02588	03420	04226	05000	05736	06428	07071	
07660	08192	08660	09063	09397	09659	09848	09962	09999	099999	
000000										
0000000										
01564	01567	01576	01591	01612	01640	01676	01720	01773	01838	
01917	02012	02128	02272	02456	02702	03054	03655	04956	05947	
03090	03096	03112	03139	03178	03229	03293	03373	03469	03585	
03724	03891	04093	04340	04649	05047	05595	06452	07958	08787	
04540	04547	04568	04602	04652	04717	04800	04900	05021	05166	
05336	05539	05779	06066	06415	06846	07403	08184	09262	09674	
05878	05887	05907	05944	05997	06066	06152	06257	06382	06529	
06701	06701	07134	07405	07723	08099	08552	09121	09745	09917	
07071	07078	07098	07132	07180	07243	07321	07414	07525	07654	
07802	07972	08165	08384	08632	08913	09231	09591	09914	09978	
08097	08095	08112	08138	08176	08225	08285	08357	08442	08539	
08649	08773	08912	09065	09234	09417	09613	09817	09971	09995	
08910	08914	08924	08942	08966	08998	09037	09083	09137	09198	
09267	09343	09427	09517	09615	09717	09822	09924	09991	099986	
09511	09512	09517	09526	09538	09554	09573	09595	09621	09650	
09683	09719	09758	09799	09843	09887	09932	09973	099975	099997	
09877	09877	09879	09881	09884	09888	09893	09899	09906	09914	
09922	09931	09941	09952	09963	09974	09985	09994	099993	099999	
10	10	10	10	10	10	10	10	10	10	
10	10	10	10	10	10	10	10	10	10	
0	1	2	3	4	5	6	7	8	9	10

0000000	1000
00000	1250
0000	16667
0000000	2500
00000	5000
000	10000
00000	20000
-10	10000
-100	125
-1000	16667
-100	2500
-100	5000
-10	10000
-100	200000

```

C C A S CHAWLA IRI Z EXIT GRADIENT- BOUNDRY CONDITION A
    DIMENSION AM(10),CK(10)
    READ 10, (AM(J),J=1,10)
10  FORMAT (10F7.5)
    READ 10, (CK(J),J=1,10)
C    UFOR=B1/D
C    DFOR=B2/D
    200 READ 700,UFOR
    700 FORMAT (F10.4)
C    UFIR=F1/D
C    DFIR=F2/D
    90  UFIR=1.0
    40  DFIR=1.0
    ELON=SQRTF(1.0+(UFOR)**2)
    ELT=1.0
    DON=SQRTF(1.0+(UFOR+UFIR)**2)
60  DT=SQRTF(1.0+(DFIR)**2)
    BON=1.0-2.0*(DON-ELON)/(DON+DT)
    BT=1.0-2.0*(DT-ELT)/(DON+DT)
    AMX =SQRTF(2.0*(BON+BT)/(1.0+BON+BT+BON*BT))
    DO 1 J=1,10
    IF (AMX-AM(J)) 3,2,1
1    CONTINUE
2    K=J
    CKX=CK(K)
3    K=J
    CKX=CK(K-1)+(CK(K)-CK(K-1))/(AM(K)-AM(K-1))*(AMX-AM(K-1))
    S=SQRTF((DT+ELON)/(2.0*(DT-1.0)*(ELON+1.0)))
    SKR=S/CKX
300  FORMAT (7F10.7)
    PUNCH 300,UFOR,UFIR,DFIR,AMX,CKX,S,SKR
    IF(DFIR-4.0) 70,80,80
70  DFIR=DFIR+1.0
    GO TO 60
80  IF(UFIR-3.0)90,100,100
90  UFIR=UFIR+1.0
    GO TO 40
100 GO TO 200
    END
08192 08660 09063 09397 09659 09848 09962 09986 099985 099993

```

```

20347 21565 23088 25046 27601 31534 38317 433P7 56349 57914
20000
3000
4000
50000
10000
5000
1000

```

```

C C A S CHAWLA Z PRESSURE VARIATION FOR BOUNDRY COMDITON A
DIMENSION AM(20),SN(11,20),U(11),SNFX(11)
READ 10,(AM(J),J=1,20)
10 FORMAT (10F7.5)
READ 20,((SN(I,J),J=1,20),I=1,11)
20 FORMAT (10F7.5)
READ 30,(U(I),I=1,11)
30 FORMAT (11F6.4)
C EL IS THE LOCATION OF KEY POINT IN T-PLANE
200 READ 700,EL
700 FORMAT (F10.6)
C UFOR=B1/D
C DFOR=B2/D
C UFIR=F1/D
C DFIR=F2/D
UFOR=5.0
DFOR=5.0
UFIR=1.0
60 DFIR=1.0
50 ELON=SQRTF(1.0+(UFOR)**2)
ELT=SQRTF(1.0+(DFOR)**2)
DON=SQRTF(1.0+(UFOR+UFIR)**2)
DT=SQRTF(1.0+(DFOR+DFIR)**2)
BON=1.0-2.0*(DON-ELON)/(DON+DT)
BT=1.0-2.0*(DT-ELT)/(DON+DT)
T=1.0-2.0*(DT-EL)/(DON+DT)
AMX =SQRTF(2.0*(BON+BT)/(1.0+BON+BT+BON*BT))
ENU=BT+BON*BT-T-T*BON
DENU=BON+BT-BON*T-BT*T
SNFU=SQRTF(ENU/DENU)
IF (AMX-.9999) 11,11,12
12 J=20
GO TO 2
11 DO1 J=1,20
IF (AMX-AM(J)) 3,2,1
1 CONTINUE
2 K=J
DO 4 I=1,11
IF (SNFU-SN(I,K)) 6,5,4
4 CONTINUE
5 M=I
PHI=U(M)
GO TO 100
6 M=I
PHI=U(M-1)+(U(M)-U(M-1))/(SN(M,K)-SN(M-1,K))*(SNFU-SN(M-1,K))
GO TO 100
3 K=J
DO 7 I=1,11
A=AMX-AM(K-1)
SNFX(I)=SN(I,K-1)+(SN(I,K)-SN(I,K-1))/(AM(K)-AM(K-1))*A
7 CONTINUE

```

```

8   DO 9 I=1,11
   IF (SNFU-SNFX(I)) 14,5,9
9   CONTINUE
14  M=I
   PHI=U(M-1)+(U(M)-U(M-1))/(SNFX(M)-SNFX(M-1))*(SNFU-SNFX(M-1))
100 PUNCH 300,EL,UFIR,DFIR,AMX,SNFU,PHI
300 FORMAT (6F10.7)
   IF(DFIR-3.0) 15,16,16
15  DFIR=DFIR+1.0
   GO TO 50
16  IF(UFIR-3.0) 17,18,18
17  UFIR=UFIR+1.0
   GO TO 60
18  GO TO 200
   END

```

```

00175 00872 01737 02588 03420 04226 05000 05736 06428 07071
07660 08192 08660 09063 09397 09659 09848 09962 09999 099999

```

```

00000000
00000000

```

01564	01567	01576	01591	01612	01640	01676	01720	01773	01838	
01917	02012	02128	02272	02456	02702	03054	03655	04956	05947	
03090	03096	03112	03139	03178	03229	03293	03373	03469	03585	
03724	03891	04093	04340	04649	05047	05595	06452	07958	08787	
04540	04547	04568	04602	04652	04717	04800	04900	05021	05166	
05336	05539	05779	06066	06415	06846	07403	08184	09262	09674	
05878	05887	05907	05944	05997	06066	06152	06257	06382	06529	
06701	06901	07134	07405	07723	08099	08552	09121	09745	09917	
07071	07078	07098	07132	07180	07243	07321	07414	07525	07654	
07802	07972	08165	08384	08632	08913	09231	09591	09914	09978	
08090	08095	08112	08138	08176	08225	08285	08357	08442	08539	
08649	08773	08912	09065	09234	09417	09613	09817	09971	09995	
08910	08914	08924	08942	08966	08998	09037	09083	09137	09198	
09267	09343	09427	09517	09615	09717	09822	09924	09991	099986	
09511	09512	09517	09526	09538	09554	09573	09595	09621	09650	
09683	09719	09758	09799	09843	09887	09932	09973	099975	099997	
09877	09877	09879	09881	09884	09888	09893	09899	09906	09914	
09922	09931	09941	09952	09963	09974	09985	09994	099993	099999	
10	10	10	10	10	10	10	10	10	10	
10	10	10	10	10	10	10	10	10	10	
0	1	2	3	4	5	6	7	8	9	10

```

0000000000
-1000000
-1414214
-2236068
-3162278
-4123106
+1000000
+1414214
+2236068
+3162278
+4123106

```

```

C C A S CHAWLA Z IRI UPLIFT PRESSURES FOR BOUNDARY CONDITION B
DIMENSION AM(20),SN(11,20),U(11),F(5),X(5),D(5),SNFX(11)
RFAD 10, (AM(J),J=1,20)
10 FORMAT (10F7.5)
RFAD 20, ((SN(I,J),J=1,20),I=1,11)
20 FORMAT (10F7.5)
RFAD 30, (U(I),I=1,11)
30 FORMAT (11F6.4)
C F(1)=R1/T
C F(2)=R2/T
C UFIR=F1/T
C DFIR=F2/T
C CUT=D/T
C EMU IS THE LOCATION OF KEY POINT IN T-PLANE
F(1)=0.5
114 F(2)=0.0
CUT=0.2
54 UFIR=0.5
53 DFIR=0.5
52 F(3)=F(1)+UFIR
F(4)=F(2)+DFIR
S=SINF(1.5708*CUT)
C=COSF(1.5708*CUT)
DO 51 I=1,4
X(I)=3.1416*F(I)
D(I)=SQRTF(S*S/(C*C)+((EXPF(X(I))-1.0)/(EXPF(X(I))+1.0))**2)
51 CONTINUE
BN=1.-2.*(D(3)-D(1))/(D(3)+D(4))
BT=1.-2.*(D(4)-D(2))/(D(3)+D(4))
115 EMU=-S/C
50 T=1.-2.*(D(4)-FMU)/(D(3)+D(4))
AMX=SQRTF(2.*(BN+BT)/((1.+BN)*(1.+BT)))
SNFU=SQRTF(((1.+BN)*(BT-T))/((BN+BT)*(1.-T)))
IF(SNFU) 62,62,63
62 PHI=0.000
GO TO 100
63 IF(AMX-.99999) 11,12,12
12 J=20
GO TO 2
11 DO1 J=1,20
IF(AMX-AM(J)) 3,2,1
1 CONTINUE
2 K=J
DO 4 I=1,11
IF (SNFU-SN(I,K)) 6,5,4
4 CONTINUE
5 M=I
PHI=U(M)
GO TO 100
6 M=I
PHI=U(M-1)+(U(M)-U(M-1))/(SN(M,K)-SN(M-1,K))*(SNFU-SN(M-1,K))
GO TO 100
3 K=J
DO 7 I=1,11

```

```

A=AMX-AM(K-1)
SNFX(I)=SN(I,K-1)+(SN(I,K)-SN(I,K-1))/(AM(K)-AM(K-1))*A
7 CONTINUE
8 DO 9 I=1,11
  IF (SNFU-SNFX(I)) 14,5,9
9 CONTINUE
14 M=I
  PHI=U(M-1)+(U(M)-U(M-1))/(SNFX(M)-SNFX(M-1))*(SNFU-SNFX(M-1))
100 PUNCH 60,FMU,F(1),F(2),UFIR,DFIR,CUT,BN,AMX,SNFU,PHI
60 FORMAT (F10.3,4F6.3,4F8.5,F6.3)
  IF(DFIR-1.5) 31,32,32
31 DFIR=DFIR+0.5
  GO TO 52
32 IF(UFIR-1.5) 33,34,34
33 UFIR=UFIR+0.5
  GO TO 53
34 IF(CUT-0.8) 35,36,36
35 CUT=CUT+0.2
  GO TO 54
36 IF(F(1)-1.0) 111,112,112
111 F(1)=F(1)+0.5
  GO TO 114
112 IF(F(1)-3.0) 113,49,49
113 F(1)=F(1)+1.0
  GO TO 114
49 STOP
END
00175 00872 01737 02588 03420 04226 05000 05736 06428 07071
07660 08192 08660 09063 09397 09659 09848 09962 09999 099999
00000000
00000000
01564 01567 01576 01591 01612 01640 01676 01720 01773 01838
01917 02012 02128 02272 02456 02702 03054 03655 04956 05947
03090 03096 03112 03139 03178 03229 03293 03373 03469 03585
03724 03891 04093 04340 04649 05047 05595 06452 07958 08787
04540 04547 04568 04602 04652 04717 04800 04900 05021 05166
05336 05539 05779 06066 06415 06846 07403 08184 09262 09674
05878 05887 05907 05944 05997 06066 06152 06257 06382 06529
06701 06901 07134 07405 07723 08099 08552 09121 09745 09917
07071 07078 07098 07132 07180 07243 07321 07414 07525 07654
07802 07972 08165 08384 08632 08913 09231 09591 09914 09978
08090 08095 08112 08138 08176 08225 08285 08357 08442 08539
08649 08773 08912 09065 09234 09417 09613 09817 09971 09995
08910 08914 08924 08942 08966 08998 09037 09083 09137 09198
09267 09343 09427 09517 09615 09717 09822 09924 09991 099986
09511 09512 09517 09526 09538 09554 09573 09595 09621 09650
09683 09719 09758 09799 09843 09887 09932 09973 099975 099997
09877 09877 09879 09881 09884 09888 09893 09899 09906 09914
09922 09931 09941 09952 09963 09974 09985 09994 099993 099999
10 10 10 10 10 10 10 10 10 10
10 10 10 10 10 10 10 10 10 10
0 1 2 3 4 5 6 7 8 9 10

```

```

C C A S CHAWLA Z IRI EXIT GRADIENT FOR BOUNDRY CONDITION B
DIMENSION X(5),F(5),AM(14),CK(14),D(5)
READ 10, (AM(J),J=1,14)
10 FORMAT (7F10.7)
READ 10, (CK(J),J=1,14) /
C F(1)=B1/T
C F(2)=B2/T
C UFIR=F1/T
C DFIR=F2/T
C CUT=D/T
F(1)=0.5
114 F(2)=0.0
CUT=0.2
54 UFIR=0.5
53 DFIR=0.5
52 F(3)=F(1)+UFIR
F(4)=F(2)+DFIR
S=SINF(1.5708*CUT)
C=COSF(1.5708*CUT)
DO 51 I=1,4
X(I)=3.1416*F(I)
D(I)=SQRTF(S*S/(C*C)+((EXPF(X(I))-1.0)/(EXPF(X(I))+1.0))**2)
51 CONTINUE
BN=1.-2.*(D(3)-D(1))/(D(3)+D(4))
BT=1.-2.*(D(4)-D(2))/(D(3)+D(4))
55 AMX=SQRTF(2.*(BN+BT)/((1.+BN)*(1.+BT)))
IF (AMX-1.000) 5,6,6
6 K=14
GO TO 2
5 K=1
11 IF(AMX-AM(K)) 3,2,1
1 K=K+1
IF(K-14) 11,11,2
2 CKX=CK(K)
GO TO 4
3 CKX=CK(K-1)+(CK(K)-CK(K-1))*(AMX-AM(K-1))/(AM(K)-AM(K-1))
4 ENU=2.0*C*(D(4)+D(1))/S
DENU=(D(4)-S/C)*(D(1)+S/C)
S=0.7854*SQRTF(ENU/DENU)
G=S/CKX
PUNCH 60,F(1),F(2),UFIR,DFIR,CUT,BN,AMX,CKX,S,G
60 FORMAT (5F5.2,2F8.5,F12.5,2F8.5)
IF(DFIR-1.5) 31,32,32
31 DFIR=DFIR+0.5
GO TO 52
32 IF(UFIR-1.5) 33,34,34
33 UFIR=UFIR+0.5
GO TO 53
34 IF(CUT-0.8) 35,36,36
35 CUT=CUT+0.2
GO TO 54
36 IF(F(1)-1.0) 111,112,112
111 F(1)=F(1)+0.5
GO TO 114
112 IF(F(1)-5.0) 113,49,49
113 F(1)=F(1)+1.0
GO TO 114
49 STOP
END

```

.7660	.8192	.8660	.9063	.9397	.9659	.9848
.996195	.99863	.99985	.99993	.9999756	.9999985	1.0000000
1.9356	2.0347	2.1565	2.3088	2.5046	2.7681	3.1534000

```

C C A S CHAWLA Z PRESSURE VARIATION FOR BOUNDRY CONDITION P
DIMENSION AM(20),SN(11,20),U(11),F(5),X(5),D(5),SNFX(11)
READ 10, (AM(J),J=1,20)
10 FORMAT (10F7.5)
READ 20, ((SN(I,J),J=1,20),I=1,11)
20 FORMAT (10F7.5)
READ 30, (U(I),I=1,11)
30 FORMAT (11F6.4)
C F(1)=R1/T
C F(2)=B2/T
C UFIR=F1/T
C DFIR=F2/T
C CUT=D/T
C EMU IS THE LOCATION OF KEY POINT IN T-PLANE
F(1)=1.0
F(2)=1.0
CUT=0.2
54 UFIR=0.5
53 DFIR=0.5
52 F(3)=F(1)+UFIR
F(4)=F(2)+DFIR
S=SINF(1.5708*CUT)
C=COSF(1.5708*CUT)
TAN=S/C
DO 51 I=1,4
X(I)=3.1416*F(I)
D(I)=SQRTF(S*S/(C*C)+((EXPF(X(I))-1.0)/(EXPF(X(I))+1.0))**2)
51 CONTINUE
BN=1.-2.*(D(3)-D(1))/(D(3)+D(4))
BT=1.-2.*(D(4)-D(2))/(D(3)+D(4))
L=1
Z=0.8
56 Y=3.1415927*Z
EMU=SQRTF(TAN*TAN+((EXPF(Y)-1.0)/(EXPF(Y)+1.0))**2)
55 T=1.-2.*(D(4)-EMU)/(D(3)+D(4))
AMX=SQRTF(2.*(BN+BT)/((1.+BN)*(1.+BT)))
SNFU=SQRTF(((1.+BN)*(BT-T))/((BN+BT)*(1.-T)))
IF(AMX-.99999) 11,12,12
12 J=20
GO TO 2
11 DO1 J=1,20
IF(AMX-AM(J)) 3,2,1
1 CONTINUE
2 K=J
DO 4 I=1,11
IF (SNFU-SN(I,K)) 6,5,4
4 CONTINUE
5 M=I
PHI=U(M)
GO TO 100
6 M=I
PHI=U(M-1)+(U(M)-U(M-1))/(SN(M,K)-SN(M-1,K))*(SNFU-SN(M-1,K))
GO TO 100
3 K=J
DO 7 I=1,11
A=AMX-AM(K-1)
SNFX(I)=SN(I,K-1)+(SN(I,K)-SN(I,K-1))/(AM(K)-AM(K-1))*A
7 CONTINUE
8 DO 9 I=1,11
IF (SNFU-SNFX(I)) 14,5,9
9 CONTINUE
14 M=I

```



```

PHI=U(M-1)+(U(M)-U(M-1))/(SNFX(M)-SNFX(M-1))*(SNFU-SNFX(M-1))
100 PUNCH 60, Z, F(1), F(2), UFIR, DFIR, CUT, BN, AMX, SNFU, PHI
60  FORMAT (5F6.3, 4F8.5, F8.3)
    IF(L-4) 120, 121, 122
120  L=L+1
    Z=Z-0.2
    GO TO 56
121  L=L+1
    Z=0.0
    EMU=+S/C
    GO TO 55
122  IF(L-6) 123, 124, 125
123  L=L+1
    Z=CUT
    FMU=0.0
    GO TO 55
124  L=L+1
    Z=-0.0
    EMU=-S/C
    GO TO 55
125  IF(L-10) 126, 126, 128
126  L=L+1
    Z=Z-0.2
    Y=3.1415927*(-Z)
    EMU=-SQRTF(TAN*TAN+((EXPF(Y)-1.0)/(EXPF(Y)+1.0))**2)
    GO TO 55
128  IF(DFIR-1.5) 31, 32, 32
31  DFIR=DFIR+0.5
    GO TO 52
32  IF(UFIR-1.5) 33, 34, 34
33  UFIR=UFIR+0.5
    GO TO 53
34  IF(CUT-0.8) 35, 36, 36
35  CUT=CUT+0.2
    GO TO 54
36  STOP
    END

```

00175	00872	01737	02588	03420	04226	05000	05736	06428	07071	
07660	08192	08660	09063	09397	09659	09848	09962	09999	099999	
00000000										
00000000										
01564	01567	01576	01591	01612	01640	01676	01720	01773	01838	
01917	02012	02128	02272	02456	02702	03054	03655	04956	05947	
03090	03096	03112	03139	03178	03229	03293	03373	03469	03585	
03724	03891	04093	04340	04649	05047	05595	06452	07958	08787	
04540	04547	04588	04602	04652	04717	04800	04900	05021	05166	
05336	05539	05779	06066	06415	06846	07403	08184	09262	09674	
05878	05887	05907	05944	05997	06066	06152	06257	06382	06529	
06701	06901	07134	07405	07723	08099	08552	09121	09745	09917	
07071	07078	07098	07132	07180	07243	07321	07414	07525	07654	
07802	07972	08165	08384	08632	08913	09231	09591	09914	09978	
08090	08095	08112	08138	08176	08225	08285	08357	08442	08539	
08649	08773	08912	09065	09234	09417	09613	09817	09971	09995	
08910	08914	08924	08942	08966	08998	09037	09083	09137	09198	
09267	09343	09427	09517	09615	09717	09822	09924	09991	099986	
09511	09512	09517	09526	09538	09554	09573	09595	09621	09650	
09683	09719	09758	09799	09843	09887	09932	09973	099975	099997	
09877	09877	09879	09881	09884	09888	09893	09899	09906	09914	
09922	09931	09941	09952	09963	09974	09985	09994	099993	099999	
10	10	10	10	10	10	10	10	10	10	
10	10	10	10	10	10	10	10	10	10	
0	1	2	3	4	5	6	7	8	9	10

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