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Certified that the attached dissertation on "INVESTIGATIONS OF FRAG. C. P. ROTOR
MOTOR"

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V. K. Jain

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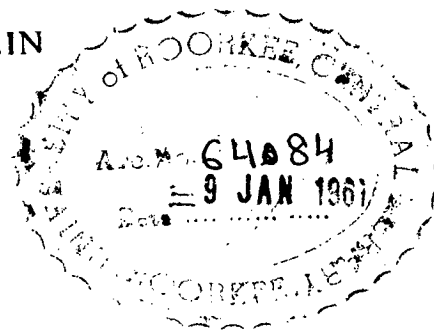
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JAI

INVESTIGATIONS OF DRAG CUP ROTOR MOTOR

A Dissertation
submitted in partial fulfilment
of the requirements for the Degree,
of
MASTER OF ENGINEERING
in
ADVANCED ELECTRICAL MACHINES

by
V.K. JAIN



DEPARTMENT OF ELECTRICAL ENGINEERING
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ROORKEE
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C E R T I F I C A T E.

Certified that the dissertation entitled 'Investigations on Drag Cup Rotor Motor' which is being submitted by Sri V.K. Jai in partial fulfilment for the award of the Degree of Master of Engineering in Advanced Electrical Machines of the University of Roorkee, is a record of student's own work carried out by him under my supervision and guidance. The matter embodied in this dissertation has not been submitted for the award of any other degree or diploma.

This is further to certify that he has worked for a period of 8 months from January '66 to August '66 for preparing dissertation for Master of Engineering Degree at the University.

Roorkee 30th Sept. 1966

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(V.K. Jain)

A B S T R A C T.

In the present work an attempt has been made to develop three dimensional analysis (taking field to be three dimensional) of drag-cup rotor machine. Expressions for combined impedance of airgaps, rotor and central iron core in terms of primary or stator using both rectangular and cylindrical co-ordinate systems have been derived. Radially directed component of magnetic vector potential or rotor currents has been neglected. Effect of web sections of rotor have been neglected. A comparison of results of two dimensional analysis both in rectangular and cylindrical co-ordinates and those of three dimensional analysis in rectangular co-ordinates as well as the experimental ones has been done. It has been shown that for the machine tested that curvature has little effect except at low values of slips. Further the case of mechanical transients has been studied with an eye to non-linearity of torque speed curve.

LIST OF SYMBOLS

Unless otherwise specified MKS system of units has been used.

The symbols defined below have been used without defining them. They have the meaning defined below unless otherwise specified.

A	= magnetic vector potential
	= area under one pole
B	= flux density vector wb/m^2
B_1	= $\frac{1}{2} k_w / D$
B_g	= air gap flux density
D	= average rotor diameter
E	= electric field intensity volt/cm
E_g	= air gap voltage
F_1, F_2, F_3	are function of x, y & z only
H	= magnetic field intensity amp/cm
I_1, I_0	= peak value of stator current
k_w	= $k_d \cdot k_p$ = winding distribution factor
N_1, Z_{ph}	= No. of turns in series per phase
P	= No. of poles pairs
R_g	= reactive part of $Z_{gap-rotor}$
R_1	= stator winding resistance
τ	= pole pitch
T_{xp}, T_x	= turn distribution per phase in x end & s direction respectively.
V_1, V_g	= rotor terminal voltage on reference & control phases respectively.
X_1	= stator leakage reactance
X_g	= reactive part of $Z_{gap-rotor}$

$A_0, I_0, I_2, R_2, C_{11}, R_1, D, \epsilon, A_{20}, I_{20}$ are constants of integration.

$$\alpha = p/D = \pi/T$$

C_1, C_2, C_3 — are constants of integration.

d = air gap thickness meters

f = frequency in c/s

i = current density amp/m²

j = $\sqrt{-1}$

k = π/T

n = no. of air gap

n = order of harmonic

p = no. of air gap

p = no. of poles

ρ = eddy in per unit

τ = time constant

ω = angular velocity

μ = permeability of rotor material

$$= \mu_r \mu_0$$

σ = electrical conductivity of rotor material 1/ohm-cm

ρ = resistivity of rotor material ohm-cm

λ = machine length in axial direction

CHAPTER - 1

INTRODUCTION

The actuator requirements of a.c. servomechanisms are most satisfactorily met by two-phase induction motors. When used as servomotors, one of the phase windings is connected to a fixed a.c. voltage, the reference voltage, and the other winding supplied by a variable control voltage, in space quadrature with the fixed voltage. The direction of rotation is governed by the polarity of the control voltage. Torque developed is a function of the magnitudes of both the voltages.

The two phase servomotor is a specially designed induction motor with a high ratio of rotor resistance to reactance so as to obtain a sloping torque speed characteristic. The two phase servomotors have the unique and most important feature of high torque to weight ratio and quick response. The requirement of high torque to inertia ratio is best met by drag-cup motors.

The drag-cup motor probably is named after the tachometer devices normally used in automatic speed meters. The stator has a standard distributed winding to obtain nearly sinusoidal space distribution. The rotor of such motors consists of two parts; the central iron core, which is stationary and the rotor conductor (drag-cup) which rotates. The drag-cup fits into the air space between the stator winding and the stationary central iron core, the clearances being kept as small as possible. The revolving field is produced by the two-phase winding on the stator and the magnetic circuit is completed through the central core which is independent of moving system. Torque is produced by the interaction of eddy currents flowing in the cup and the flux. Such machines are used as servomotors, tachogenerators and accelerometers.

The analysis of such devices is complicated due to the presence of two air gaps and the consideration of skin effect which cannot be neglected if the cup is thick as in higher rating machines. Furthermore the path of eddy currents is not simple, validity of circuit theory approach is doubtful. An accurate analysis must, therefore, be based on electromagnetic field theory.

1.2. REVIEW OF RESEARCH WORK:

Fillmore⁽¹⁾ (1956) made a study of eddy current paths in the drag-cup induction motor rotors. He assumed cup to be very thin and, therefore, neglected skin effect. He obtained expressions for rotor resistance and average power dissipated in the cup.

When the thickness of the cup is increased as in higher rating machines, skin effect plays an important role and therefore cannot be neglected. Cullford⁽²⁾ (1962) considered skin effect and obtained a rigorous mathematical analysis using the vector potential concept. He neglected the circumferentially directed component of magnetic vector potential on the ground that it makes no contribution to energy flow as measured by Poynting vector. Radial component of vector potential has little effect on power transfer except for fringing effects at ends and has been neglected. He obtained expressions for induced voltage in the stator winding and the developed torque. Cullford's analysis is for a cleave-rotor machine which is a single air gap machine.

Wood and Concordia have analysed the solid rotor machine. They have studied the effect of curvature⁽³⁾ (1959) and indicated that with non-magnetic rotors the results may be deviating from the actual ones. Rectangular co-ordinates give fairly good results when rotor used is magnetic and has a large radius. They have

considered the finite length effect⁽⁴⁾ (1959). They have considered both the circumferential and axial components of vector potential.

In an excellent work by Angst⁽⁵⁾ (1962) solid rotor machine has been analyzed by field theory. He considers fully both the circumferential and axial component of vector potential. He has derived a so-called end effect factor of the machine.

As the tangential components of rotor currents are less producing components, they have been taken into account by Koch⁽⁶⁾ (1964). He neglects both axial and circumferential components of flux density. He derives equivalent circuit with transformer elements of the drag-cup machine.

The characteristic problem in the analysis of drag-cup rotor motors by field theory is that the presence of sleeve and web sections of rotor result in field equations with differential co-efficients and co-ordinates. The effect of web section has been taken into account by several authors. It is well taken into account by Koch⁽⁶⁾. He has solved for both sections separately and synthesized to obtain the result. A classical work of calculating the resistance of the drag-cup in parts and then synthesizing through by circuit theory approach has been done by Fuller and Brzobey⁽⁷⁾ (1962).

In a most recent work by Blackford⁽⁸⁾ (1965) expressions for rotor resistance and reactance, stator to rotor mutual inductance have been derived for a drag-cup machine. Here a comparison of distributed and lumped parameter equations has been done in order to obtain the results. The equivalent circuit representation is simple. He considers field to be two dimensional and magnetic vector potential to be axially directed. He also accounts for

rotor and effects using the Fuller and Trickey's method of calculating the resistance of the cup in two parts.

All the above mentioned work except (3) by Wood and Concordia has been done using rectangular co-ordinate system rather than the natural cylindrical co-ordinate system. As indicated earlier the results may be deviating when non-magnetic rotors with small radius of curvature are employed. With this point in view Mukhopadhyay⁽⁹⁾ (1965) presented an analysis of drag-cup machine in cylindrical co-ordinate system. He considers field to be two dimensional and neglects end effect. An expression for combined ampere turn of the rotor, air gaps and central core has been found. From this the combined equivalent impedance of gaps, rotor and central core can be obtained. Beauty of the approach is that the equivalent circuit of the machine becomes extremely simple making the performance calculation easier.

CHAPTER - 2

2.0 DIMENSIONAL ANALYSIS IN RECTANGULAR CO-ORDINATE
EXTENSION ON ANGLE OF ROTOR COLEPT

2.1.1. Geometry of the Machine-

Fig. 2.1. shows the geometry of the machine with drag-cup rotor. r_1 , r_2 , r_3 and r_4 are radii at stator bore, rotor outer and inner surfaces and rotor core respectively. Y axis dimensions are as indicated in fig.

2.1.2. Assumptions-

Idealised machine shown in fig. 2.1. is used for the purpose of analysis. The following assumptions are made:

- (i) The stator iron and rotor core have infinite permeability.
- (ii) Hysteresis and saturation effects are neglected.
- (iii) Stator winding produces only forward travelling field.
- (iv) Stator and rotor core can be developed into flat, infinitely long bodies.

(v) Laurior's method⁽¹⁰⁾ of representing the stator winding by thin axially directed current sheets against a smooth stator surface is used. Effect of slotting is taken in account by Carter's co-efficient.

(vi) The field is two dimensional, i.e., magnetic intensity is independent of co-ordinate Z and electric intensity has component only in Z direction.

2.1.3. Field Equations

Maxwell's equations for electromagnetic field for space in the medium are, in M.K.S. units:

$$\text{Curl } H = J \quad \dots \quad (2.1)$$

$$\text{Curl } E = -\frac{\partial H}{\partial t} \cdot \mu = -\frac{\partial B}{\partial t} \quad \dots \quad (2.2)$$

$$\text{div } \mathbf{B} = 0 \quad (\text{div } \mathbf{H} = 0) \quad \dots \quad (2.3)$$

$$\text{div } \mathbf{E} = 0 \quad (\text{div } \mathbf{i} = 0) \quad \dots \quad (2.4)$$

In equation (2.1) the displacement current density term is dropped for the induction machine since the magnetic energy storage in the air gap is several orders of magnitude greater than electric storage.

The magnetic vector potential \mathbf{A} is defined by-

$$\text{curl } \mathbf{A} = \mathbf{B} \quad \dots \quad (2.5)$$

In the conducting region ohm's law holds.

$$\mathbf{E} = \rho \mathbf{i} \quad \dots \quad (2.6)$$

Combining (2.1) and (2.5),

$$\frac{1}{\mu} \text{curl} (\text{curl } \mathbf{A}) = \mathbf{i}$$

From the vector identity-

$$\text{Curl} (\text{curl } \mathbf{A}) = \nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A},$$

one obtains-

$$\nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A} = \mu \mathbf{i}.$$

Now since $\text{div } \mathbf{A} = 0$, the above equation simplifies to-

$$\nabla^2 \mathbf{A} = \mu \sigma \frac{\partial \mathbf{A}}{\partial t} \quad \dots \quad (2.7)$$

Equation (2.7) is Poisson's equation applicable to rotor conductor. For air gap, since the conductivity $\sigma = 0$, we have the Laplace's equation-

$$\nabla^2 \mathbf{A} = 0 \quad \dots \quad (2.8)$$

Equations similar to (2.7) and (2.8) hold for other quantities such as \mathbf{B} and \mathbf{E} as well.

where V_s is the stator terminal voltage.

This combined equivalent impedance is a function of machine constants and slip. The behaviour of the machine can therefore be readily assessed from equivalent circuit given in Fig.2.2.

This equivalent circuit representation though extremely simple is not general since R_2 and X_2 both are functions of slip.

2.2. Rotor Dynamics Analysed in an Elliptical Co-ordinate System

Stathopoulos⁽⁹⁾ carried out an analysis of drag-cup rotor machine using cylindrical co-ordinate system. Assumptions are as in analysis in rectangular co-ordinates. He obtained the following expression of combined rotor turn of the rotor, etc (see also:-

$$R = -j \frac{E_s E_r}{E \mu_0} \cdot \frac{a_{11} E_s^{2p} + 1}{a_{11} E_r^{2p} + 1}$$

$$a_{11} = \frac{E_s^{2p-1} (a_6 a_9 + a_{10}) - \mu_r E_s^{2p-1} (a_6 a_7 + a_8)}{\mu_r E_s^{2p-1} (a_6 a_7 + a_8) + E_s^{2p-1} (a_6 a_9 + a_{10})}$$

$$a_{10} = j \frac{\alpha}{\omega} (Y_{2-1}(\alpha E_s) - Y_{2+1}(\alpha E_s))$$

$$a_9 = j \frac{\alpha}{\omega} (J_{2-1}(\alpha E_s) - J_{2+1}(\alpha E_s))$$

$$a_8 = \frac{\alpha}{\omega} (Y_{2-1}(\alpha E_s) + Y_{2+1}(\alpha E_s))$$

$$a_7 = \frac{\alpha}{\omega} (J_{2-1}(\alpha E_s) + J_{2+1}(\alpha E_s))$$

$$a_6 = \frac{a_5 (E_s^{2p} - a_1) - \mu_r (E_s^{2p} + a_1) a_4}{\mu_r a_2 (E_s^{2p} + a_1) - a_4 (E_s^{2p} - a_1)}$$

a_5, a_4, a_3, a_2 = values of a_{10}, a_9, a_8, a_7 respectively with E_s replaced by E_r .

2.1.4. Boundary conditions-

At the boundary between second air gap and rotor iron H_x vanishes. At the boundaries between conductor and air gap B_y and H_x are continuous. Continuity in H_x is destroyed at the stator surface where H_x is equal to surface current density.

$$\text{At } y = +g_2, H_x = 0 \quad \dots \quad (2.9)$$

$$\begin{aligned} \text{At } y = +g_2, H_x \text{ air} &= H_x \text{ rotor} \\ B_y \text{ air} &= B_y \text{ rotor} \end{aligned} \quad \dots \quad (2.10)$$

$$\begin{aligned} \text{At } y = -g_1, H_x \text{ air} &= H_x \text{ rotor} \\ B_y \text{ air} &= B_y \text{ rotor} \end{aligned} \quad \dots \quad (2.11)$$

$$\text{At } y = 0, H_x = I_0 \quad \dots \quad (2.12)$$

2.1.5. Analysis-

R.M.K.S. system of units is used.

Laplace's equation-

$$\nabla^2 B_x = \frac{\partial^2 B_x}{\partial x^2} + \frac{\partial^2 B_x}{\partial y^2} = 0 \quad \dots \quad (2.13)$$

holds for air gap regions.

Divergence equation-

$$\nabla \cdot B = \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} = 0 \quad \dots \quad (2.14)$$

holds every where.

For rotor Poisson's equation-

$$\nabla^2 E_B = \frac{\partial^2 E_B}{\partial x^2} + \frac{\partial^2 E_B}{\partial y^2} = \mu \sigma \frac{\partial E_B}{\partial t} \quad (2.15)$$

holds good.

Solution of equations (2.13), (2.14) and (2.15) can be effected by the method of separation of variables.

Air gap- In order to solve (2.13) let us assume a

solution of the form-

$$B_x = F_1(x) F_2(y)$$

where,

$F_1(x)$ and $F_2(y)$ are functions of x & y only respectively.

Substituting this in (2.13) one obtains-

$$F_1'' F_2 + F_2'' F_1 = 0$$

or,

$$\frac{F_1''}{F_1} = -\frac{F_2''}{F_2} = -k^2 \quad \text{where } k = \text{constant.}$$

Solution for F_1 and F_2 will be written as-

$$F_1 = c_1 \cos kx$$

$$F_2 = c_2 \text{Exp}(ky) + c_3 \text{Exp}(-ky)$$

since the variation in x direction is sinusoidal and that in y direction is exponential. Further,

$$k = \frac{\pi}{\tau}, \quad \tau \text{ being the pole pitch.}$$

The complete solution of B_x is given by-

$$B_x = \cos kx (A_n \text{Exp}(-ky) + B_n \text{Exp}(ky)) \quad (2.16a)$$

From (2.14)-

$$B_y = - \int \frac{\partial B_x}{\partial x} dy = \sin kx (B_n \text{Exp}(ky) - A_n \text{Exp}(-ky)) \quad \dots \quad (2.16b)$$

n in the above equations represent the number of air gap.

Rotor core - In rotor core flux density should not become infinite at the centre. Solution for rotor core, therefore becomes,

$$B_x = A_3 \cos kx \text{Exp}(-ky) \quad \dots \quad (2.17a)$$

$$B_y = -A_3 \sin kx \text{Exp}(-ky) \quad \dots \quad (2.17b)$$

and $A_1 = a_3 B_1$ where $a_3 = \frac{r-hg}{2lp(2hg_1)(\mu_0 k)}$

and $q = \frac{1+a_2 \cdot \exp(2hg_1)}{1-a_2 \cdot \exp(2hg_1)}$

At stator surface,

$$B_x (\text{air gap}) = (A_1 + B_1) \cos kx$$

$$B_y (\text{air gap}) = \sin kx (B_1 - A_1)$$

If Π represents the peak value of combined ampere turn of air gap, rotor and rotor core, we have-

$$\Delta T = \Pi \sin kx$$

$$\therefore B_x = \mu_0 \frac{\partial (\Delta T)}{\partial x} = \mu_0 \Pi k \cos kx$$

Variation of flux density in y direction is sinusoidal, we have therefore,

$B_y = B_0 \sin kx$, B_0 being the peak value of air gap flux density.

From there one obtains,

$$\Pi = \frac{1+a_3}{1-a_3} \cdot \frac{B_0}{\mu_0 k}$$

Air gap flux density B_0 is related to air gap voltage by the expression-

$$E_1 = 4.44 K_v \int B_0 T_{ph} \cdot A$$

where symbols have usual meaning. Thus-

$$\Pi = \frac{1+a_3}{1-a_3} \cdot \frac{E_1}{4.44 K_v \int T_{ph} \mu_0 k \cdot A}$$

From this combined equivalent impedance of the rotor, air gaps and core in terms of primary can be found as-

$$Z_{oc} = \frac{V_1 \cdot T_{ph}}{\Pi}$$

where V_1 is the stator terminal voltage.

This combined equivalent impedance is a function of machine constants and slip. The behaviour of the machine can therefore be readily assessed from equivalent circuit given in Fig.2.2.

This equivalent circuit representation though extremely simple is not general since R_2 and X_2 both are functions of slip.

2.2. TWO DIMENSIONAL ANALYSIS IN CYLINDRICAL CO-ORDINATE SYSTEM-

Mukhopadhyay⁽⁹⁾ carried out an analysis of drag-cup rotor machine using cylindrical co-ordinate system. Assumptions are as in analysis in rectangular. Co-ordinates. He obtained the following expression of combined impedance of the rotor, air gaps and core-

$$Z = -j \frac{D \cdot E_1}{P \mu_0} \cdot \frac{a_{11} E_1^{2P} + 1}{a_{11} E_1^{2P} - 1}$$

$$\text{where, } a_{11} = \frac{E_2^{P-1} (a_6 a_9 + a_{10}) - \mu_r E_2^{P-1} (a_6 a_7 + a_8)}{\mu_r E_2^{-P-1} (a_6 a_7 + a_8) + E_2^{-P-1} (a_6 a_9 + a_{10})}$$

$$a_{10} = j \frac{\alpha}{\omega} (Y_{P-1}(\alpha E_2) - Y_{P+1}(\alpha E_2))$$

$$a_9 = j \frac{\alpha}{\omega} (J_{P-1}(\alpha E_2) - J_{P+1}(\alpha E_2))$$

$$a_8 = \frac{\alpha}{\omega} (Y_{P-1}(\alpha E_2) + Y_{P+1}(\alpha E_2))$$

$$a_7 = \frac{\alpha}{\omega} (J_{P-1}(\alpha E_2) + J_{P+1}(\alpha E_2))$$

$$a_6 = \frac{a_5 (E_3^{2P} - a_1) - \mu_r (E_3^{2P} + a_1) a_3}{a_2 \mu_r (E_3^{2P} + a_1) - a_4 (E_3^{2P} - a_1)}$$

a_4, a_3, a_2 = Values of a_{10}, a_9, a_8 and a_7 respectively with E_2 replaced by E_3 .

$$a_1 = - \frac{\mu_r - 1}{\mu_r + 1} r_4^{2P}$$

and $\alpha^2 = j\omega \epsilon \mu \sigma$

From this expression for ampere turn performance calculation is done in the manner indicated earlier.

CHAPTER - 3

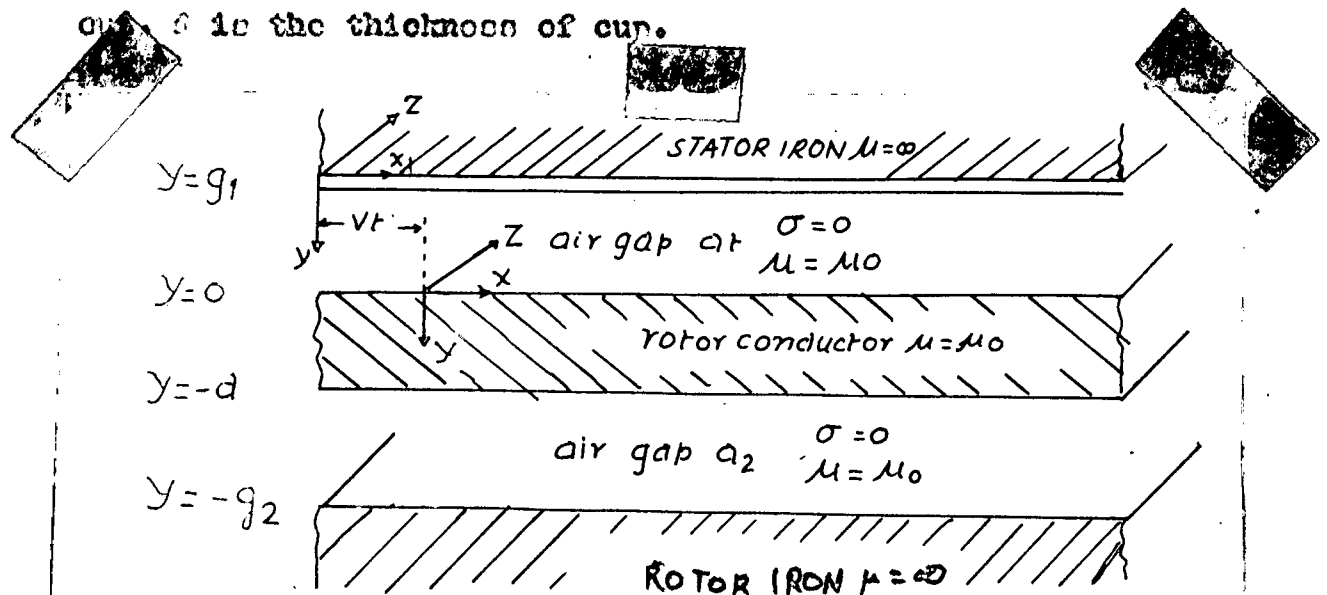
THREE-DIMENSIONAL ANALYSIS IN RECTANGULAR
CO-ORDINATE SYSTEM

3.1.1. INTRODUCTION:

The previous chapter dealt with a two-dimensional analysis of the machine. This chapter deals with the three-dimensional analysis of the machine. The electromagnetic approach is based on vector potential concept. The radial component of vector potential or rotor current has been neglected. It has been shown that the equivalent circuit representation is complex as it requires curvation of a series and the rotor resistance and reactance are both functions of slip.

3.1.2. GEOMETRY OF THE MACHINE:

Fig. 3.1 shows the geometry of the machine rolled out slot. The co-ordinate system (x_1, y, z) is attached to the stator. The co-ordinate system (x, y, z) is attached to moving cup. a is the thickness of cup.



IDEALISED DRAG-CUP MACHINE

FIG-31

3.1.3. Assumptions:

Assumptions (i) to (v) of article 2.1.2. are applicable here also.

(vi) Radial component of magnetic vector potential or rotor currents have been neglected. This would in turn mean neglecting the fringing effects which take place at the ends.

3.1.4. Linear Current

The wave due to current sheet at stator surface rotates at an angular velocity ω . Motion of rotor is given by-

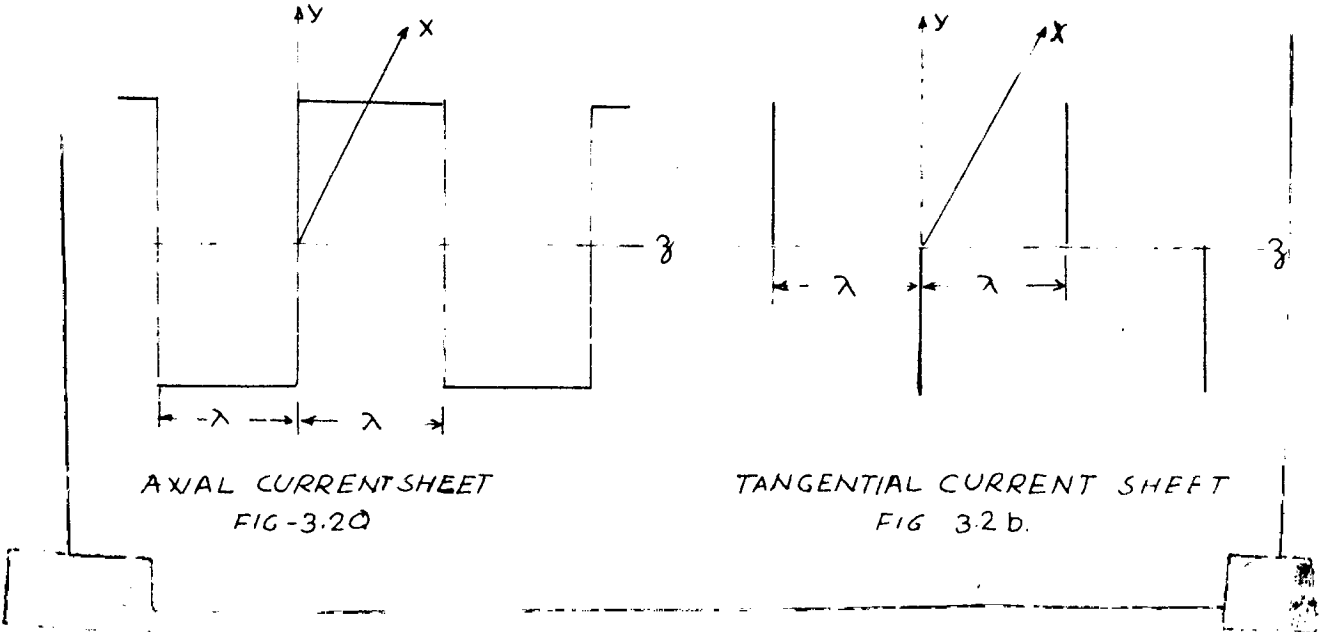
$$x_1 = \int \frac{d(1-c)}{p} dt$$

$$\begin{aligned} \therefore x_1 &= x - \frac{D}{p} \int (1-c) \omega dt \\ &= x - (1-c) \cdot \frac{\omega t}{p} \quad \dots \quad (3.1) \end{aligned}$$

$$\text{here, } \theta = \frac{p}{2} = \frac{\pi}{2}$$

The actual stator winding is replaced by linear current sheet in circumferential direction. This has a sinusoidal variation in circumferential direction and is of constant amplitude in axial direction. For the purpose of analysis the machine may be considered to repeat itself axially with alternate polarity. The current sheet in axial direction may then be expressed by Fourier series. Now since circumferential current sheet is also present it must also have periodic distribution expressible by Fourier series. This can be seen from Kirchhoff's law $\nabla \cdot \mathbf{i} = 0$. The circumferential current sheet is solely confined to the end surfaces of the machine. Fig. 3.2a and Fig. 3.2b show the two current sheets.

More mathematically, axial current sheet,



$$I_{CB} = \frac{4}{n\pi} \left(\frac{3}{2} B_1 I_1 \right) \exp [j(\alpha x_1 + \omega t)] \sum_{n=1,3,\dots} \sin \frac{n\pi x}{\lambda} \dots \quad (3.2)$$

where, B_1 = amplitude of sinusoidal winding distribution
 $= H_1 k_1 / \delta$.

From divergence relation $\nabla \cdot I_0 = 0$, tangential current sheet,

$$I_{ox} = - \int \frac{\partial I_{CB}}{\partial x} dx$$

$$= j \frac{4}{\pi \lambda} \left(\frac{3}{2} B_1 I_1 \right) \exp [j(\alpha x_1 + \omega t)] \left(\sum_{n=1,3,\dots} \cos \frac{n\pi x}{\lambda} + C_n(x) \right) \dots \quad (3.3)$$

where, $C_n(x)$ is a function of x .

Because of symmetry it can be shown that $C_n(x) = 0$.

~~Solution for Δ_{CB}~~ As shown in article 2.1.3., we have for

given the Laplace's equation-

$$\nabla^2 \Delta_{CB} = \frac{\partial^2 \Delta_{CB}}{\partial x_1^2} + \frac{\partial^2 \Delta_{CB}}{\partial y^2} + \frac{\partial^2 \Delta_{CB}}{\partial z^2} = 0 \quad (3.4)$$

$$\nabla^2 \Lambda_{CB} = \frac{\partial^2 \Lambda_{CB}}{\partial x_1^2} + \frac{\partial^2 \Lambda_{CB}}{\partial y^2} + \frac{\partial^2 \Lambda_{CB}}{\partial z^2} = 0 \quad (3.5)$$

$$\Lambda_{CB} = 0 \quad \dots \quad (3.6)$$

A solution to (3.4) can be obtained by the method of separation of variables. Assuming a solution for Λ_{CB}

$$\Lambda_{CB}(x, y, z, t) = \exp(j\omega t) X(x) Y(y) Z(z)$$

Putting it in (3.4),

$$\frac{X''}{X} + \frac{Y''}{Y} = -\frac{Z''}{Z} \quad \dots \quad (3.7)$$

Variation of Λ_{CB} is periodic in z direction and sinusoidal in x direction.

$$\therefore \frac{Z''}{Z} = -\frac{n^2 \pi^2}{\lambda^2} \quad \dots \quad (3.8)$$

$$\frac{X''}{X} = -\alpha^2 \quad \dots \quad (3.9)$$

Solutions to (3.8) and (3.9) are-

$$Z = c_1 \sin \frac{n\pi z}{\lambda}$$

$$\text{and } X = c_2 \exp(j\alpha x_1)$$

From (3.7), (3.8) and (3.9),

$$\frac{Y''}{Y} = \frac{n^2 \pi^2}{\lambda^2} + \alpha^2 = \beta_n^2 \quad (\text{say})$$

Its solution is-

$$Y = c_3 \cosh \beta_n y + c_4 \sinh \beta_n y$$

Complete solution for Λ_{CB} is

$$\Lambda_{CB} = \exp [j(\alpha x_1 + \omega t)] \sum_{n=1,3,\dots} \{ A_{nCB} \cosh \beta_n y + B_{nCB} \sinh \beta_n y \} \sin \frac{n\pi z}{\lambda} \quad \dots \quad (3.10)$$

$$\text{From } \nabla^2 \Lambda_{CB} = \frac{\partial \Lambda_{CB}}{\partial x_1^2} + \frac{\partial \Lambda_{CB}}{\partial y^2} = 0, \text{ we have}$$

$$\begin{aligned}
 A_{xz_1} &= - \int \frac{\partial A_{yz_1}}{\partial z} dz \\
 &= \sum_{n=1,3,\dots} \frac{jn\pi}{\lambda} \exp [j(\alpha x_1 + \omega t)] (A_{nm} \cosh \beta_n y + \\
 &\quad B_{nm} \sinh \beta_n y) \cos \frac{n\pi z}{\lambda} + D_n(x) \quad (3.11)
 \end{aligned}$$

the constant of integration $D_n(x)$ reducing to zero because of symmetry. $m (=1,2)$ in above equations represents a number of airgap.

Solution for rotor: As shown in article 2.1.3. Poisson's and divergence equations are applicable for rotor.

$$\nabla^2 A_{xz} = \frac{\partial^2 A_{xz}}{\partial x^2} + \frac{\partial^2 A_{xz}}{\partial y^2} + \frac{\partial^2 A_{xz}}{\partial z^2} = \mu \sigma \frac{\partial A_{xz}}{\partial t} \quad (3.12)$$

$$\nabla^2 A_{yz} = \frac{\partial^2 A_{yz}}{\partial x^2} + \frac{\partial^2 A_{yz}}{\partial y^2} + \frac{\partial^2 A_{yz}}{\partial z^2} = \mu \sigma \frac{\partial A_{yz}}{\partial t} \quad (3.13)$$

$$\nabla \cdot A_x = \frac{\partial A_{xx}}{\partial x} + \frac{\partial A_{xz}}{\partial z} = 0 \quad \dots \quad (3.14)$$

$$\nabla \cdot A_y = 0 \quad \dots \quad (3.15)$$

A solution to (3.12) may be obtained by the method of separation of variables. Assuming a solution of the form

$$A_{xz}(x, y, z, t) = \exp(j\omega t) X(x) Y(y) Z(z)$$

Substituting in (3.12),

$$\frac{X''}{X} + \frac{Y''}{Y} + \frac{Z''}{Z} = j\omega \mu \sigma \quad \dots \quad (3.16)$$

Variation of A_{xz} in x and z directions is sinusoidal and periodic respectively.

$$\therefore \frac{X''}{X} = -\alpha^2 \quad \dots \quad (3.17)$$

$$\frac{Z''}{Z} = -\frac{n^2 \pi^2}{\lambda^2} \quad \dots \quad (3.18)$$

solutions to (3.17) and (3.18) are-

$$X = c_5 \exp(jax)$$

$$Z = c_6 \sin \frac{n\pi z}{\lambda}$$

From (3.16), (3.17) and (3.18),

$$\frac{Y''}{Y} = j\omega_0 \mu \sigma + a^2 + \frac{n^2 \pi^2}{\lambda^2} = \gamma_n^2 \quad (\text{say})$$

$$\therefore Y = c_7 \cosh \gamma_n y + c_8 \sinh \gamma_n y$$

complete solution for A_{xz} is-

$$A_{xz} = \sum_{n=1,3,\dots} \exp [j(ax + \omega t)] \{ (C_n \cosh \gamma_n y + D_n \sinh \gamma_n y) \times$$

$$\sin \frac{n\pi z}{\lambda} \dots \quad (3.19)$$

$$A_{xz} = - \int \frac{\partial A_{xz}}{\partial z} dz, \text{ from (3.14)}$$

$$= \sum_{n=1,3,\dots} j \frac{n\pi}{\lambda a} \exp [j(ax + \omega t)] (C_n \cosh \gamma_n y + D_n \sinh \gamma_n y).$$

$$\cos \frac{n\pi z}{\lambda} \dots \quad (3.20)$$

The constants of integration A_{nm} ($m=1,2$), B_{nm} ($m=1,2$), C_n and D_n must be determined from boundary conditions.

Boundary conditions: At the boundaries between conductor and airgap H_x and B_z are continuous. At the rotor core surface H_x vanishes and at the stator surface discontinuity in H_x is equal to stator surface current density.

$$(a) \text{ At } y = -c_2, \quad H_x \text{ air} = 0 \quad \dots \quad (3.21)$$

$$(b) \text{ At } y = -d, \quad H_x \text{ air} = H_x \text{ rotor} \dots \quad (3.22a)$$

$$A_B \text{ air} = A_B \text{ rotor} \dots \quad (3.22b)$$

$$(c) \text{ At } y = 0, \quad H_x \text{ air} = H_x \text{ rotor} \quad \dots \quad (3.23a)$$

$$A_n \text{ air} = A_n \text{ rotor} \quad \dots \quad (3.23b)$$

$$(d) \text{ At } y = G_1, \quad H_x \text{ air} = I_{00} \quad \dots \quad (3.24)$$

Now from relationship- curl $A = D$, we have

$$H_x = \frac{1}{\mu} (\text{curl}_x A) = \frac{1}{\mu} \frac{\partial A_n}{\partial y} \quad \dots \quad (3.25)$$

From (3.21), (3.25) and (3.10),

$$\beta_n (-A_{n2} \sinh \beta_n G_2 + B_{n2} \cosh \beta_n G_2) = 0$$

$$\therefore \frac{A_{n2}}{B_{n2}} = a_1 \text{ (say)} = \coth \beta_n G_2 \quad \dots \quad (3.26)$$

From (3.22a), (3.22b), (3.25), (3.10), (3.19) and (3.26),

$$\frac{C_n}{D_n} = a_3 \text{ (say)} = \frac{a_2 \sinh \gamma_n d + \cosh \gamma_n d}{\sinh \gamma_n d + a_2 \cosh \gamma_n d} \quad (3.27)$$

$$\text{where } a_2 = \frac{\beta_n}{\gamma_n} \cdot \frac{-a_1 \sinh \beta_n d + \cosh \beta_n d}{a_1 \cosh \beta_n d - \sinh \beta_n d} \quad \dots \quad (3.28)$$

From (3.23a), (3.23b), (3.25), (3.10), (3.19) & (3.27),

$$\frac{A_{n1}}{D_{n1}} = a_4 \text{ (say)} = \frac{\beta_n}{\gamma_n} \cdot a_3 \quad \dots \quad (3.29)$$

From (3.24), (3.25), and (3.19) and (3.29),

$$B_{n1} = \frac{\mu}{n \pi \beta_n} \left(\frac{3}{2} B_1 I_1 \right) \cdot \frac{1}{a_4 \sinh \beta_n G_1 + \cosh \beta_n G_1} \quad \dots \quad (3.30)$$

$$A_{n1} = a_4 \cdot B_{n1}$$

Having evaluated all the constants of integration, vector potential at stator surface can be written as-

$$(A_{CE})_{y=G_1} = \sum_{n=1,3,\dots} \exp [j(\alpha x_1 + \omega t)] \left[\frac{A_{n1} \cosh \beta_n G_1 + D_{n1} \sinh \beta_n G_1}{\sin \frac{n\pi d}{\lambda}} \right]$$

$$= \sum_{n=1,3,\dots} \exp [j(\alpha x_1 + \omega t)] \frac{4\mu}{n\pi\beta_n} \left(\frac{3}{2} B_1 I_1 \right) \cdot a_5 \sin \frac{n\pi z}{\lambda},$$

from (3.29) and (3.30)

$$\text{where } a_5 = \frac{a_4 \coth \beta_n G_1 + \sinh \beta_n G_1}{a_4 \sinh \beta_n G_1 + \coth \beta_n G_1} \quad (3.31)$$

and,

$$\begin{aligned} (\Lambda_{(x1)})_{y=G_1} &= \sum_{n=1,3,\dots} \frac{jn\pi}{a\lambda} \exp [j(\alpha x_1 + \omega t)] \frac{4\mu}{n\pi\beta_n} \left(\frac{3}{2} B_1 I_1 \right) \cdot \\ & a_5 \cos \frac{n\pi z}{\lambda} \quad \dots \quad (3.32) \end{aligned}$$

Electric force field is given by,

$$E_G = - \frac{\partial \Lambda_G}{\partial t}$$

$$\therefore (E_{GE})_{y=G_1} = \sum_{n=1,3,\dots} -j\omega \exp [j(\alpha x_1 + \omega t)] \frac{4\mu \cdot a_5}{n\pi\beta_n} \left(\frac{3}{2} B_1 I_1 \right) \cdot$$

$$\begin{aligned} & \sin \frac{n\pi z}{\lambda} \\ (E_{GE})_{y=G_1} &= \sum_{n=1,3,\dots} \frac{n\pi\omega}{a\lambda} \exp [j(\alpha x_1 + \omega t)] \frac{4\mu a_5}{n\pi\beta_n} \left(\frac{3}{2} B_1 I_1 \right) \cdot \\ & \cos \frac{n\pi z}{\lambda} \end{aligned}$$

Taking real part of periodic function in x_1 and

$$(E_{GE})_{y=G_1} = \sum_{n=1,3,\dots} \sin (\alpha x_1 + \omega t) \cdot \frac{4\mu a_5}{n\pi} \cdot \left(\frac{3}{2} B_1 I_1 \right) \sin \frac{n\pi z}{\lambda}$$

$$(E_{GE})_{y=G_1} = \sum_{n=1,3,\dots} \frac{n\pi\omega}{a\lambda} \cos (\alpha x_1 + \omega t) \cdot \frac{4\mu a_5}{n\pi\beta_n} \cdot \left(\frac{3}{2} B_1 I_1 \right) \cdot$$

$$\cos \frac{n\pi z}{\lambda}$$

The real part of turn distribution per phase is, from (3.2) and (3.3),

$$T_G = \sum_{n=1,3,\dots} \frac{4}{a\lambda} B_1 \cos \alpha x_1 \sin \frac{n\pi z}{\lambda}$$

$$T_{H1} = \sum_{n=1,3,\dots} \frac{4}{a\lambda} B_1 \sin \alpha x_1 \cos \frac{n\pi z}{\lambda}$$

Winding voltage: The induced voltage in a conductor is the integral along the conductor length of the total electric field. The conductor voltages accumulated over all the conductors of the winding, yield the winding induced voltage. Voltage induced in a conductor may be written as-

$$v(\text{cond}) = \int_0^{\lambda} \left\{ (T_E E_{(E)})_{y=C_1} + (T_{K1} E_{(K1)})_{y=C_1} \right\} ds$$

Voltage induced in a phase of the stator winding is-

$$V_1 = p \int_0^{\lambda} v(\text{cond}) dx_1$$

It will be observed that $n' = n$ because of the following identity

$$\int_0^{\lambda} \sin \frac{n\pi s}{\lambda} \sin \frac{n'\pi s}{\lambda} ds = \int_0^{\lambda} \cos \frac{n\pi s}{\lambda} \cos \frac{n'\pi s}{\lambda} ds = 0$$

when $n' \neq n$.

$$V_1 = \frac{12 \omega p^2 \mu_0 \lambda B_1^2 I_1}{n \pi \beta_n} \left[\int_0^{\lambda} \left\{ \frac{\sin(\alpha x_1 + \omega t) \cos(\alpha x_1)}{n} + \right. \right.$$

$$\left. \frac{\frac{n\pi}{\alpha^2 \lambda^2} \cos(\alpha x_1 + \omega t) \cdot \sin(\alpha x_1)}{n} \right\} dx_1 \right]$$

$$= \frac{12 \omega p^2 \mu_0 \lambda B_1^2 I_1}{n \pi \beta_n} \left(\frac{1}{n\pi} + \frac{n\pi}{\alpha^2 \lambda^2} \right) \cdot \frac{2}{2} \sin \omega t$$

$$= \frac{6 \omega p^2 \mu \lambda \beta_n B_1^2}{n^2 \alpha^3} \cdot \alpha_5 \cdot I_1 \sin \omega t$$

This is the airgap voltage induced in one phase stator winding. This is in quadrature to the winding referred to. Airgap voltage induced in the same winding is-

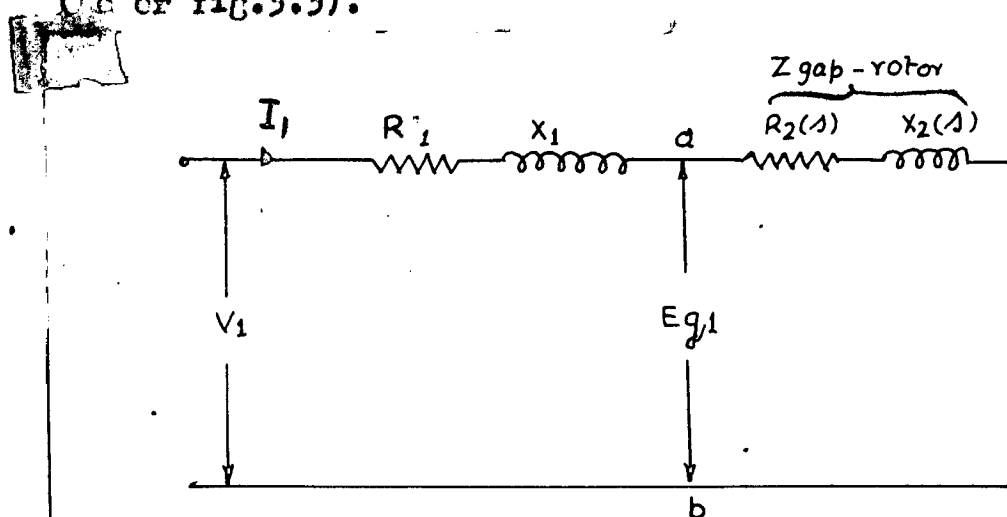
$$V_1 = \sum_{n=1,3,\dots} j \frac{6 \omega p^2 \mu \lambda \beta_n B_1^2}{n^2 \alpha^3} \cdot \alpha_5 I_1 \sin \omega t$$

(3.33)

Equivalent impedance of rotor, airgap and central core in terms of stator is given by-

$$\begin{aligned}
 u_{\text{gap-rotor}} &= \frac{v_1}{\lambda_1 \sin \omega t} \\
 &= \sum_{n=1,3,\dots} 3 \frac{G \omega \beta \mu \lambda \beta_n L_1^2 \sigma_5}{n^2 \sigma^3 \pi} \quad (3.34)
 \end{aligned}$$

Equations (3.33) and (3.34) may be used to represent the equivalent circuit of the machine. Equation (3.34) gives the impedance of the machine as seen from terminals a & b (Refer fig.3.3).



EQUIVALENT CIRCUIT OF DRAG CUP MACHINE.

FIG 3.3.

The equivalent circuit representation as indicated by equation (3.33) is not simple since more than one term of the series must be used.

The above analysis takes into account tangential component of vector potential or rotor currents. It may be interesting to study the effect of neglecting this component of vector potential. A detailed study of this has been done in Appendix I.

3.2. FORCE DYNAMICAL ANALYSIS IN CYLINDRICAL CO-ORDINATE SYSTEM

3.2.1. INTRODUCTION:

Analysis of earlier section neglected the effect of curvature by developing the stator and rotor into flat infinitely long bodies. With non-magnetic rotor having a small radius of curvature the effect of curvature may be pronounced. This section deals with the effect of curvature by considering the natural cylindrical co-ordinate system.

3.2.2. GEOMETRY OF MACHINE:

Fig.3.4 shows the geometry of the machine. The whole machine may be considered to consist of three cylinders namely stator, moving cup and rotor core. The co-ordinate system (r, θ, z) is attached to stator. The co-ordinate system (r, θ, z) is attached to moving cup.

3.2.3. ASSUMPTIONS:

All assumptions except (iv) of article 3.1.3. are applicable here.

3.2.4. ANALYSIS:

Relative Motion: Mag wave due to current sheet at stator surface rotates with an angular velocity ω . Motion of rotor is given by $\int (1-\alpha)\omega dt$.

$$\begin{aligned} \therefore \theta_1 &= \theta - \int (1-\alpha)\omega dt \\ &= \theta - (1-\alpha)\omega t \end{aligned}$$

Current sheets: The stator winding can be replaced by two current sheets, one in circumferential direction and other in axial direction. The axial current sheet is -

$$I_{\theta\theta} = \sum_{n=1,3} \frac{4}{n\pi} \left(\frac{j}{2} B_n I_n \right) \exp(j(\rho\theta_1 + \omega t)) \sin \frac{n\pi z}{\lambda} \quad \dots \quad (3.35)$$

Tangential current sheet is, from divergence equation

$$\nabla \cdot I_{\theta} = 0,$$

$$I_{\theta\theta_1} = - \int r \frac{\partial I_{\theta\theta}}{\partial z} dz$$

$$= \sum_{n=1,3} \frac{j4r_1}{P} \left(\frac{j}{2} B_n I_n \right) \exp(j(\rho\theta_1 + \omega t)) \cos \frac{n\pi z}{\lambda}, \quad \dots \quad (3.36)$$

the constant of integration reducing to zero because of symmetry.

Solution for airgap: For airgap regions we have-

$$\nabla^2 \Lambda_{\theta\theta_1} = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \Lambda_{\theta\theta_1}}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \Lambda_{\theta\theta_1}}{\partial \theta_1^2} + \frac{\partial^2 \Lambda_{\theta\theta_1}}{\partial z^2} = 0 \quad \dots \quad (3.37)$$

$$\nabla^2 \Lambda_{\theta z} = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \Lambda_{\theta z}}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \Lambda_{\theta z}}{\partial \theta_1^2} + \frac{\partial^2 \Lambda_{\theta z}}{\partial z^2} = 0 \quad \dots \quad (3.38)$$

$$\nabla \cdot \Lambda_{\theta} = \frac{1}{r} \frac{\partial}{\partial r} \left(r \Lambda_{\theta r} \right) + \frac{1}{r} \frac{\partial \Lambda_{\theta\theta_1}}{\partial \theta_1} + \frac{\partial \Lambda_{\theta z}}{\partial z} = 0 \quad \dots \quad (3.39)$$

$$\text{and } \Lambda_{\theta r} = 0$$

A solution to (3.38) can be obtained by the method of separation of variables. Assuming a solution for $\Lambda_{\theta z}$,

$$\Lambda_{\theta z}(r, \theta, z, t) = \exp(j\omega t) P_1(\theta) P_2(r) P_3(z),$$

where P_1, P_2, P_3 are purely functions of θ, r & z respectively.

Substituting it in (3.38),

$$\frac{P_3''}{P_3} = - \frac{P_2''}{P_2} - \frac{1}{r} \frac{P_2'}{P_2} - \frac{1}{r^2} \frac{P_1''}{P_1} \quad \dots \quad (3.40)$$

Variation of A_{GB} is periodic in z direction and sinusoidal in θ direction.

$$\therefore \frac{F_3''}{F_3} = - \frac{n^2 \pi^2}{\lambda^2} \quad \dots \quad (3.41)$$

$$\frac{F_1''}{F_1} = - p^2 \quad \dots \quad (3.42)$$

Solutions to (3.41) and (3.42) are -

$$F_3 = c_1 \sin \frac{n\pi z}{\lambda}$$

$$F_1 = c_2 \exp(jP \theta_f)$$

From (3.40), (3.41) and (3.42),

$$r^2 F_2'' + r F_2' + F_2 \left(- \frac{n^2 \pi^2 r^2}{\lambda^2} - p^2 \right) = 0 \quad (3.42)$$

Equation (3.43) is Bessel's equation whose solution may be written as-

$$F_2 = c_3 I_p(\beta r) + c_4 K_p(\beta r)$$

where $\beta = \frac{n\pi}{\lambda}$

and, I_p and K_p are modified Bessel functions of purely imaginary arguments.

Complete solution for A_{GB} is

$$A_{GB} = \sum_{n=1,3,\dots} \exp(j(P\theta_1 + \omega t)) \left\{ \frac{A_{nm} I_p(\beta r) + B_{nm} K_p(\beta r)}{\sin \frac{n\pi z}{\lambda}} \right\} \quad (3.44)$$

From (3.39),

$$A_{GB\theta_1} = - \int r \frac{\partial A_{GB}}{\partial z} dz d\theta_1$$

$$\sum_{n=1,3,\dots} \frac{j n \pi r}{P \lambda} \exp(j(P\theta_1 + \omega t)) \left\{ \frac{A_{nm} I_p(\beta r) + B_{nm} K_p(\beta r)}{\cos \frac{n\pi z}{\lambda}} \right\} \quad (3.45)$$

m in above equations represent the number of gap.

Solution for rotor: For moving conductor we have -

$$\begin{aligned}\nabla^2 \Lambda_{r\theta} &= \frac{\partial^2 \Lambda_{r\theta}}{\partial r^2} + \frac{1}{r} \frac{\partial \Lambda_{r\theta}}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Lambda_{r\theta}}{\partial \theta^2} + \frac{\partial^2 \Lambda_{r\theta}}{\partial z^2} \\ &= \mu\sigma \cdot \frac{\partial \Lambda_{r\theta}}{\partial t} \quad \dots \quad (3.46)\end{aligned}$$

$$\begin{aligned}\nabla^2 \Lambda_{r\theta} &= \frac{\partial^2 \Lambda_{r\theta}}{\partial r^2} + \frac{1}{r} \frac{\partial \Lambda_{r\theta}}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Lambda_{r\theta}}{\partial \theta^2} + \frac{\partial^2 \Lambda_{r\theta}}{\partial z^2} \\ &= \mu\sigma \cdot \frac{\partial \Lambda_{r\theta}}{\partial t} \quad \dots \quad (3.47)\end{aligned}$$

$$\begin{aligned}\nabla \cdot \Lambda_r &= \frac{1}{r} \frac{\partial}{\partial r} (r \Lambda_{r\theta}) + \frac{1}{r} \frac{\partial \Lambda_{r\theta}}{\partial \theta} + \frac{\partial \Lambda_{r\theta}}{\partial z} = 0 \\ &\quad \dots \quad (3.48)\end{aligned}$$

$$\text{and } \Lambda_{r\theta} = 0$$

Solution to Poisson's equation (3.47) can be obtained by the method of separation of variables. Assuming a solution of the form-

$$\Lambda_{r\theta} (r, \theta, z, t) = \exp(j\omega_0 t) F_1(\theta) F_2(r) F_3(z)$$

$$\frac{F_2''}{F_2} + \frac{1}{r} \frac{F_2'}{F_2} + \frac{1}{r^2} \frac{F_1''}{F_1} + \frac{F_3''}{F_3} = \alpha^2 \quad (3.49)$$

$$\text{where, } \alpha^2 = j\omega_0 \mu\sigma$$

Variation in θ direction is sinusoidal and that in z direction is periodic.

$$\therefore \frac{F_1''}{F_1} = -p^2 \quad \dots \quad (3.50)$$

$$\frac{F_3''}{F_3} = -\frac{n^2 \pi^2}{\lambda^2} \quad \dots \quad (3.51)$$

Solutions to (3.50) and (3.51) are,

$$E_1 = c_5 \exp(jP\theta)$$

$$\text{and } E_3 = c_6 \sin \frac{n\pi r}{\lambda}$$

From (3.49), (3.50) & (3.51),

$$E_2'' + \frac{1}{r} E_2' - \left(\alpha^2 + \frac{n^2 \Lambda^2}{\lambda^2} + \frac{P^2}{r^2} \right) E_2 = 0$$

This is a Bessel's equation whose solution may be written-

$$E_2 = c_7 J_P(\gamma r) + c_8 Y_P(\gamma r)$$

$$\text{where } \gamma^2 = -(\alpha^2 + \beta^2)$$

and J_P and Y_P are Bessel's functions of first and second kind of order P respectively.

Complete solution for $A_{r\theta}$ is-

$$A_{r\theta} = \sum_{n=1,3,\dots} \exp(j(P\theta + \omega t)) \left\{ c_n J_P(\gamma r) + D_n Y_P(\gamma r) \right\} \sin \frac{n\pi r}{\lambda} \dots \quad (3.52)$$

From (3.48),

$$A_{r\theta} = - \int r \frac{\partial A_{r\theta}}{\partial r} dr = \sum_{n=1,3,\dots} \frac{jn\pi r}{P\lambda} \exp(j(P\theta + \omega t)) \left\{ c_n J_P(\gamma r) + D_n Y_P(\gamma r) \right\} \cos \frac{n\pi r}{\lambda} \dots \quad (3.53)$$

The constants of integration A_{nm} ($n = 1, 2$), B_{nm} ($n = 1, 2$), C_n and D_n are to be determined from boundary conditions.

Boundary conditions: At the boundaries between conductor and airgap H_θ and A_n are continuous. At the rotor core surface H_θ vanishes and at the stator surface H_r is equal to stator surface current density.

$$(a) \quad \text{At } r = r_4, \quad H_\theta \text{ at } r = 0 \dots \quad (3.54)$$

$$(b) \quad \text{At } r = r_3, \quad H_{\theta \text{ air}} = H_{\theta \text{ rotor}} \quad (3.55a)$$

$$\Lambda_E \text{ air} = \Lambda_E \text{ rotor} \dots \quad (3.55b)$$

$$(c) \quad \text{At } r = r_2, \quad H_{\theta \text{ air}} = H_{\theta \text{ rotor}} \dots \quad (3.56a)$$

$$\Lambda_E \text{ air} = \Lambda_E \text{ rotor} \dots \quad (3.56b)$$

$$(d) \quad \text{At } r = r_1, \quad H_{\theta \text{ air}} = I_{GE} \dots \quad (3.57)$$

From curl relationship $\text{curl } A = D_0$

$$H_{\theta} = \frac{1}{\mu} (\text{curl}_{\theta} \Lambda) = -\frac{1}{\mu} \frac{\partial \Lambda}{\partial r} \dots \quad (3.58)$$

$$\text{and } H_r = \frac{1}{\mu} (\text{curl}_r \Lambda) = \frac{1}{\mu} \left(\frac{1}{r} \frac{\partial \Lambda}{\partial \theta} - \frac{\partial \Lambda_{\theta}}{\partial r} \right) \quad (3.59)$$

From (3.54), (3.58), (3.44),

$$\Lambda_{n2} / B_{n2} = \alpha_1 (\cos \gamma) = - \frac{K_P' (\beta r_1)}{I_P' (\beta r_1)} \quad (3.60)$$

From (3.55a), (3.55b), (3.58), (3.44), (3.52) and (3.60),

$$\frac{C_n}{D_n} = \alpha_3 (\cos \gamma) = \frac{\alpha_2 Y_P' (\gamma r_3) - Y_P' (\gamma r_3)}{J_P' (\gamma r_3) - \alpha_2 J_1' (\gamma r_3)} \quad (3.61)$$

$$\text{where } \alpha_2 = \frac{\beta \alpha_1 I_P' (\beta r_3) + K_P' (\beta r_3)}{\gamma \alpha_1 I_P' (\beta r_3) + K_P' (\beta r_3)}$$

From (3.56a), (3.56b), (3.58), (3.44), (3.52) and (3.61),

$$\frac{\Lambda_{n1}}{D_{n1}} = \alpha_5 (\cos \gamma) = \frac{\alpha_4 K_P' (\beta r_2) - K_P' (\beta r_2)}{I_P' (\beta r_2) - \alpha_4 I_P' (\beta r_2)} \quad (3.62)$$

$$\text{where } \alpha_4 = \frac{\gamma}{\beta} \cdot \frac{\alpha_3 J_P' (\gamma r_2) + Y_P' (\gamma r_2)}{\alpha_3 J_P' (\gamma r_2) + Y_P' (\gamma r_2)}$$

From (3.57), (3.58), (3.44), (3.55) and (3.62)-

$$B_{n1} = - \frac{\mu K}{n \pi \beta} \left(\frac{3}{2} B_1 I_1 \right) \frac{1}{\alpha_5 I_P' (\beta r_1) + K_P' (\beta r_1)} \quad (3.63)$$

$$\Lambda_{n1} = \alpha_5 \cdot D_{n1}$$

All the constants of integration have now been evaluated. Vector potential at stator surface may be written as-

$$(\Lambda_{G\theta})_{r=r_1} = \sum_{n=1,3,\dots} B_{n1} \exp(j(P\theta_1 + \omega t)) \left\{ a_5 I_P(\beta r_1) + K_P(\beta r_1) \right\} \sin \frac{n\pi R}{\lambda}$$

and $(\Lambda_{G\theta_1})_{r=r_1} = \sum_{n=1,3,\dots} B_{n1} \cdot \frac{j\beta \pi r_1}{P\lambda} \exp(j(P\theta_1 + \omega t)) \left\{ a_5 I_P(\beta r_1) + K_P(\beta r_1) \right\} \cos \frac{n\pi R}{\lambda}$, from (3.44) & (3.45)

$$\text{Let- } a_6 = a_5 I_P(\beta r_1) + K_P(\beta r_1)$$

$$\text{and } a_7 = a_5 I_P'(\beta r_1) + K_P'(\beta r_1)$$

Electric field in airgap is given by-

$$E_G = - \frac{\partial \Lambda}{\partial t}$$

Therefore,

$$(E_{G\theta})_{r=r_1} = - \sum_{n=1,3,\dots} j\omega B_{n1} \exp(j(P\theta_1 + \omega t)) a_6 \sin \frac{n\pi R}{\lambda}$$

$$\text{and } (E_{G\theta_1})_{r=r_1} = \sum_{n=1,3,\dots} \frac{n\pi r_1 \omega}{P\lambda} B_{n1} \exp(j(P\theta_1 + \omega t)) a_6 \cos \frac{n\pi R}{\lambda}$$

Taking real part of periodic function in θ_1 and

$$(E_{G\theta})_{r=r_1} = \sum_{n=1,3,\dots} \omega B_{n1} \cdot a_6 \sin \frac{n\pi R}{\lambda} (\sin(P\theta_1 + \omega t))$$

$$(E_{G\theta_1})_{r=r_1} = \sum_{n=1,3,\dots} \frac{n\pi r_1 \omega}{P\lambda} B_{n1} \cdot a_6 \cos \frac{n\pi R}{\lambda} (\cos(P\theta_1 + \omega t))$$

Real part of turn distribution per phase from (3.35) and (3.36) is-

$$T_E = \sum_{n=1,3,\dots} \frac{4}{n^2 \pi} B_{n1} \cos P\theta_1 \sin \frac{n\pi R}{\lambda}$$

$$T_{\theta_1} = \sum_{n=1,3,\dots} \frac{4 r_1}{P\lambda} B_{n1} \sin P\theta_1 \cos \frac{n\pi R}{\lambda}$$

$$\text{where } B_{n1} = \frac{K_v \cdot T_{ph}}{D}$$

Winding Voltages Voltage induced in a conductor is the integral along the conductor length of the total electric field.

$$\begin{aligned}
 v(\text{cond}) &= \int_0^\lambda \left\{ (E_E \cdot E_{\text{cond}})_{r=r_1} + (E_{\theta_1} \cdot \mu_{\text{cond}})_{r=r_1} \right\} ds \\
 &= \frac{4}{n\pi} \cdot B_1 \cos p\theta_1 \omega B_{n1} \cdot a_6 \sin (p\theta_1 + \omega t) \\
 &\quad \int_0^\lambda \sin \frac{n'\pi s}{\lambda} \cdot \sin \frac{p\pi s}{\lambda} ds + \frac{4r_1}{p\lambda} B_1 \sin p\theta_1 \cdot \\
 &\quad \frac{n\pi r_1 \omega B_{n1}}{p\lambda} \cdot a_6 \cos (p\theta_1 + \omega t) .
 \end{aligned}$$

The integrals $\int_0^\lambda \sin \frac{n'\pi s}{\lambda} \cdot \sin \frac{p\pi s}{\lambda} ds$ and $\int_0^\lambda \cos \frac{n'\pi s}{\lambda} \cdot \cos \frac{p\pi s}{\lambda} ds$ reduce to zero when $n \neq n'$. Hence $n = n'$.

$$\begin{aligned}
 \therefore v(\text{cond}) &= \frac{4}{n\pi} B_1 \omega B_{n1} a_6 \frac{\lambda}{2} \cos p\theta_1 \sin (p\theta_1 + \omega t) + \\
 &\quad \frac{4r_1}{p^2 \lambda^2} B_1 n\pi r_1 \omega B_{n1} a_6 \cdot \frac{\lambda}{2} \cdot \sin p\theta_1 \cos (p\theta_1 + \omega t)
 \end{aligned}$$

Voltage induced in one phase of the stator winding-

$$\begin{aligned}
 V_1 &= D \int_0^\pi v(\text{cond}) d\theta_1 \\
 &= B_{n1} \omega B_1 a_6 \frac{D}{p} \left(\frac{p\lambda}{n} + \frac{n^2 \pi^2 r_1^2}{p\lambda} \right) \sin \omega t \cdot 2p
 \end{aligned}$$

Substituting the value of B_{n1} from (3.63),

$$V_1 = \sum_{n=1,3,\dots} -12 \frac{B_1^2 \mu \omega \lambda r_1^2}{n^2 \pi \beta p} \left(\frac{p^2}{r_1^2} + \beta^2 \right) \cdot \frac{a_6}{a_7} I_1 \sin \omega t \cdot 2p$$

This is the voltage induced in quadrature to winding under consideration. Induced voltage in the same winding, is, therefore,

$$\begin{aligned}
 V_1 &= \sum_{n=1,3,\dots} j \frac{12 B_1^2 \mu \omega \lambda r_1^2}{n^2 \pi \beta p} \left(\frac{p^2}{r_1^2} + \beta^2 \right) \cdot \frac{a_6}{a_7} I_1 \sin \omega t \cdot 2p \\
 &\quad \dots \quad (3.64)
 \end{aligned}$$

Equivalent combined impedance of airgaps, rotor and central core, in terms of primary is,

$$\begin{aligned}
 Z_{\text{gap-rotor}} &= \frac{V_1}{I \sin \omega t} \\
 &= \sum_{n=1,3,\dots} \frac{12 B_1^2 \mu \omega \lambda r_1^2 \cdot 2p}{n^2 \pi \beta p} \left(\frac{p^2}{r_1^2} + \beta^2 \right) \quad (3.65)
 \end{aligned}$$

Equations (3.64) and (3.65) can be used to draw the equivalent circuit. $Z_{\text{gap-rotor}}$ is the impedance as seen from terminals a and b of fig. 3.3.

$$Z_{\text{gap-rotor}} = R_2 + jX_2,$$

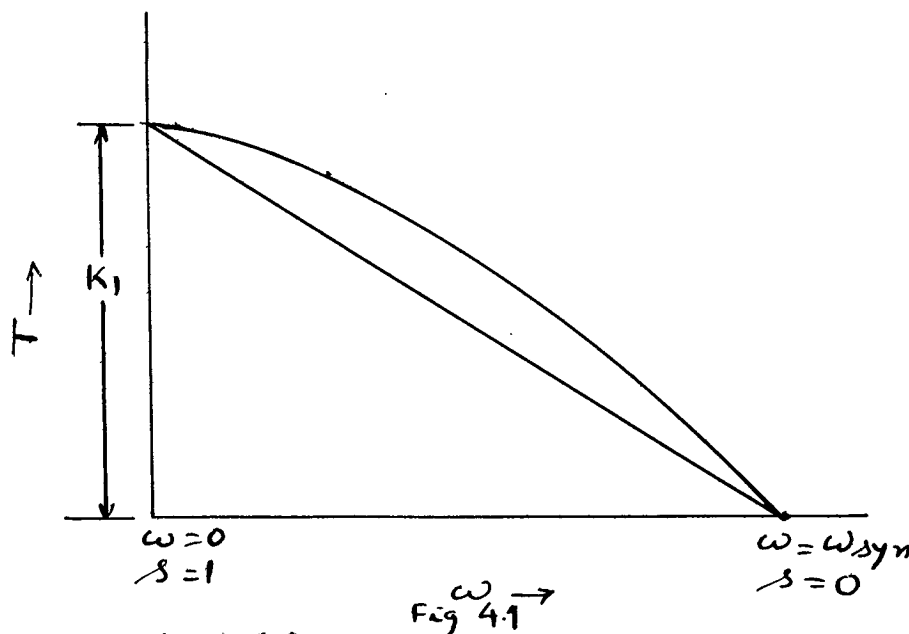
R_2 being the rotor resistance in terms of primary and X_2 being the rotor reactance inclusive of magnetizing reactance in terms of primary.

CHAPTER - 4

M ECHANICAL TRANSIENTS

In this chapter expressions for speed rise and acceleration of a servomotor have been derived taking into consideration the effect of nonlinearity of the torque speed characteristic.

Fig.4.1 shows a typical torque speed curve of the machine. It is nonlinear to take the effect of nonlinearity into account square term of speed will be considered. Torque speed curve can



be approximated by

$$T = K_1 + K_2\omega + K_3\omega^2 \quad \dots \quad (4.1)$$

where T = developed torque

K_1 = stalling torque

and K_2 & K_3 are constants to be determined from the actual torque speed curve.

Writing mechanics equation,

$$J \frac{d\omega}{dt} + K_4\omega = T \quad \dots \quad (4.2)$$

where J = M.I. of rotating parts

K_4 = viscous friction of motor and load.

From (4.1) and (4.2),

$$k_1 + k_2 \omega + k_3 \omega^2 - J \frac{d\omega}{dt} - k_4 \omega = 0$$

$$\begin{aligned} \therefore t &= J \int \frac{d\omega}{k_1 + k_2 \omega + k_3 \omega^2} \\ &= \frac{2J}{q^2} \tanh^{-1} \frac{2k_2 \omega + k_1}{q} + C \end{aligned} \quad (4.3)$$

where, $k_1 = k_2 - k_4$

$$q^2 = k_2^2 - 4k_1k_3$$

and C = constant of integration.

Applying the initial condition $\omega = 0$ at $t = 0$, one obtains

$$0 = \frac{2J}{q^2} \tanh^{-1} \frac{k_1}{q} \quad \dots \quad (4.4)$$

$$\therefore t = -\frac{2J}{q^2} \left(\tanh^{-1} \frac{2k_2 \omega + k_1}{q} - \tanh^{-1} \frac{k_1}{q} \right),$$

From (4.3) & (4.4)

$$\begin{aligned} \text{or, } \omega(t) &= \frac{q}{2k_2} \tanh \left(-\frac{q^2}{2J} t + C \log_0 \frac{q^2 + k_1}{q^2 - k_1} \right) - \frac{k_1}{2k_2} \\ &\dots \quad (4.5) \end{aligned}$$

This is the expression for speed rise similarly an expression for speed fall can be obtained by applying the initial condition $\omega = \omega_{\max}$ at $t = 0$.

Eqn.(4.5) may be applied to determine the speed response of a servomotor provided its torque speed characteristic and moment of inertia are known.

ACCELERATION

Since acceleration is time derivative of speed a curve for acceleration may be obtained either by ^{exact} or numerical differentiation.

Differentiating (4.5) the expression for speed with respect to time there results an expression for acceleration-

$$\text{Acceleration} = \frac{d\dot{\omega}(t)}{dt}$$

$$= \frac{-q}{4Jk_3} \operatorname{sech}^2 \left(-\frac{q^{\frac{1}{2}} t}{2J} + \frac{1}{2} \log_e \frac{q^{\frac{1}{2}} + k_5}{q^{\frac{1}{2}} - k_5} \right)$$

CHAPTER - 5

EXPERIMENTAL TESTS & RESULTS

Tests were performed on two motors which were built in different sizes $\frac{1}{2}$ h.p. and 3 h.p. The smaller motor was built with two rotors one with copper cup and the other having aluminium cup. The copper cups of two different thickness were used. Design consideration and principal design data of these motors are given in Appendix III.

TORQUE SPEED CURVE:

A few torque-speed curves both under balanced and unbalanced applied voltages were obtained experimentally by means of a d.c. motor coupled to test motor. Theoretical curves were obtained by (a) two dimensional analysis both in rectangular and cylindrical co-ordinates of Chapter 2 and (b) three dimensional analysis of Chapter 3. Use of a digital Computer was made to obtain the theoretical curves. The Computer programmes are given in Appendix IV. However in three dimensional analysis the axial component of vector potential was neglected.

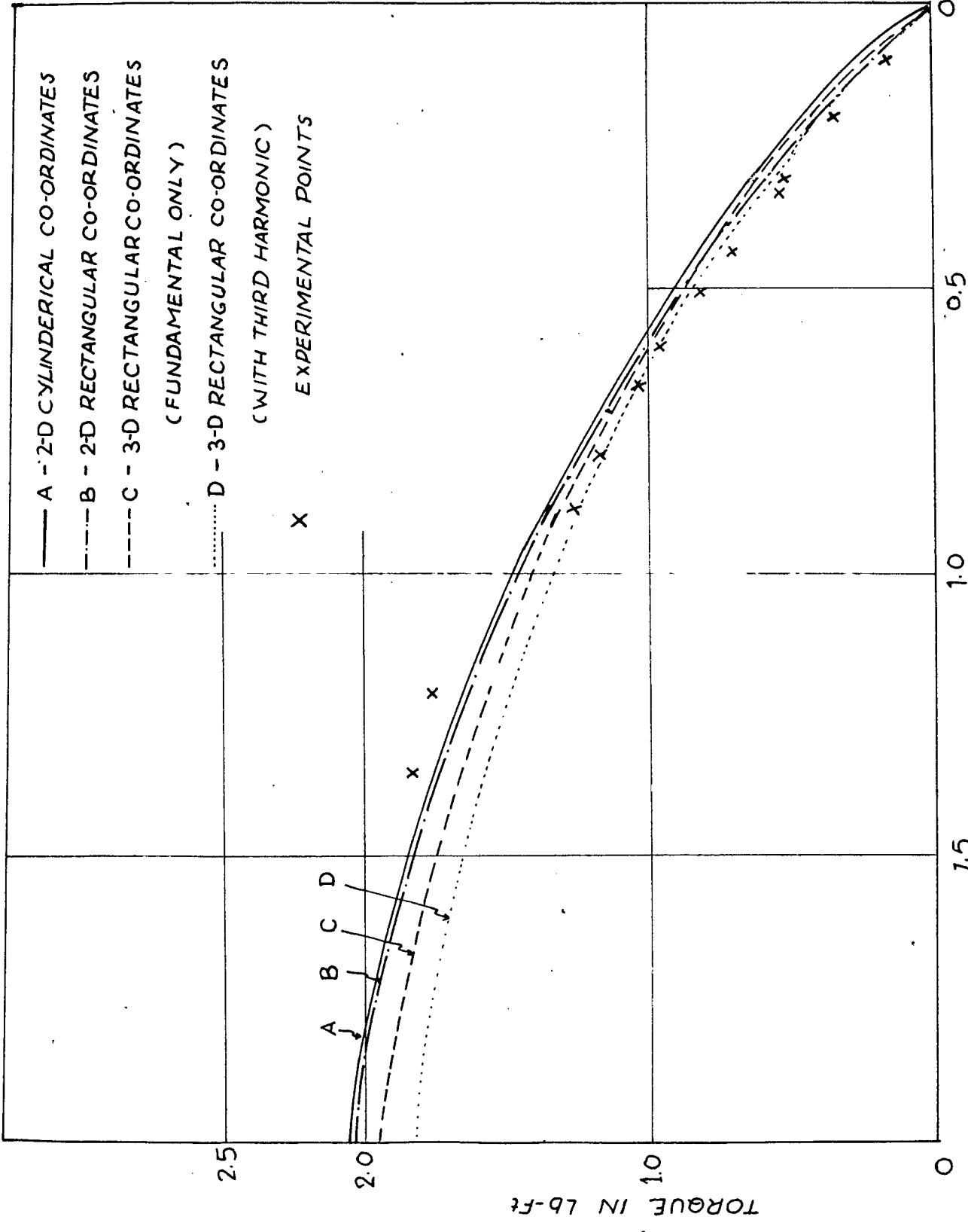
Some relations of Bessel functions for use in analysis of sections 2.2 & 3.2 are given in Appendix II.

Figs. 5.1, 5.2, 5.3 and 5.4 show the torque speed curves both for balanced and unbalanced applied voltages for various rotors: Figs. 5.1, 5.2, 5.3 and 5.4 refer to copper thin cup, copper thick cup, Aluminium cup and large rotor. The first three figs. refer to smaller machines

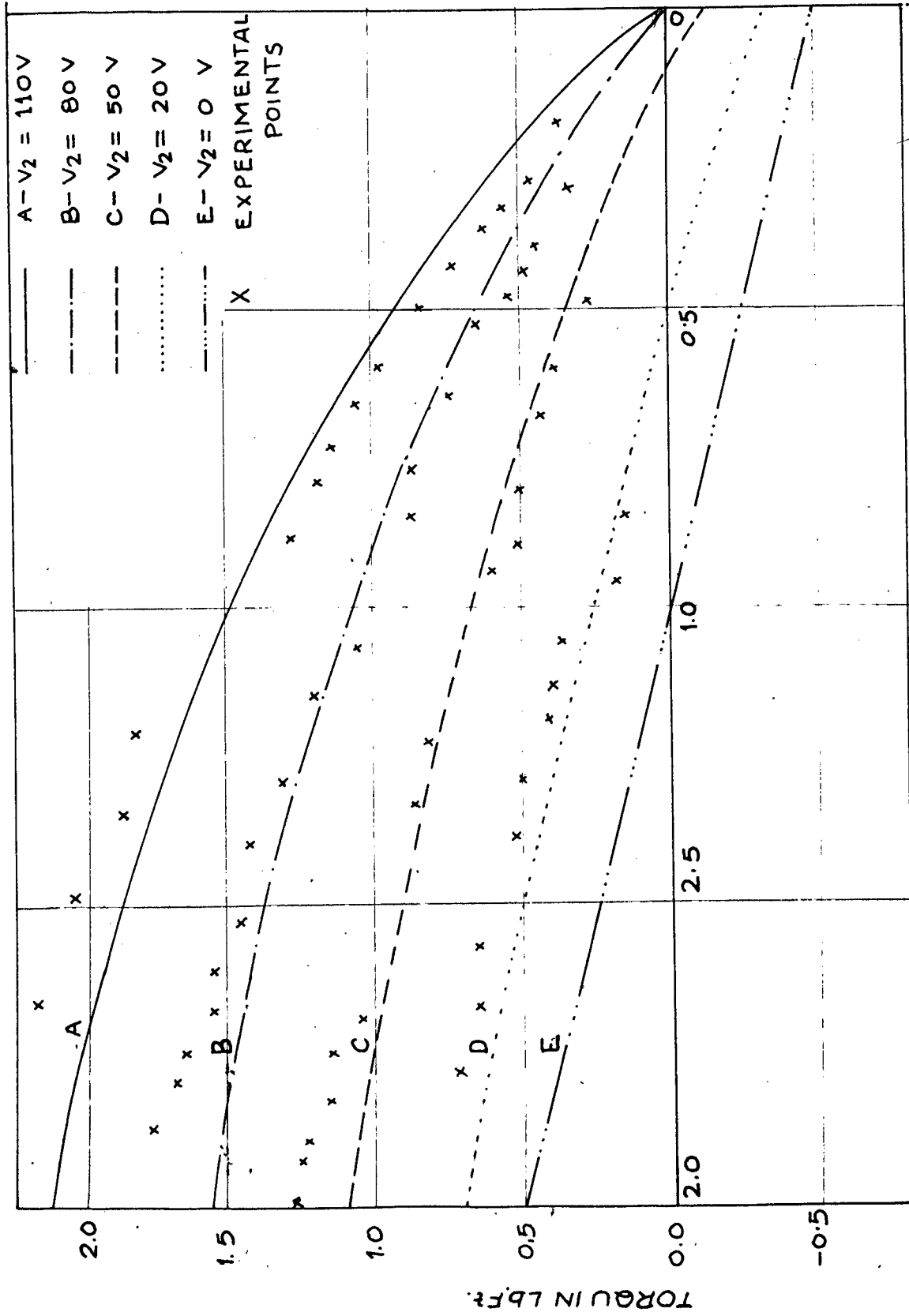
MECHANICAL TRANSIENTS:

A small tachogenerator was coupled to test motor for obtaining oscillograms of speed rise. The current transients were also recorded. Figs. 5.5(a), 5.5(b), 5.6(a), 5.6(b) and

5.7, ^(a&b) show these curves for various rotors. Figs. 5.5(a) and 5.5(b) refer to thick copper cup. Figs. 5.6(a), 5.6(b) refer to thin copper cup and Fig. 5.7, ^(a&b) refers to aluminium cup rotor. It was observed that the aluminium cup rotor gives much faster response than the copper cup of the same size. The aluminium cup rotor when run at no load rises to full speed in 5 to 10 cycles of 50 c/s wave whereas the electrical transients die out in about 3 cycles. For this case the electrical time constant is quite comparable to mechanical one and therefore it is not justifiable to neglect it in transfer function. The copper thick takes much more to attain the full speed. When started at no load the time taken for speed to gain steady value is large and may be more than 40-50 cycles of a 50 c/s wave. As the servomotor is always connected to load the electrical time constant will be small in comparison to mechanical time constant and can therefore, be neglected in transfer function of the motor.



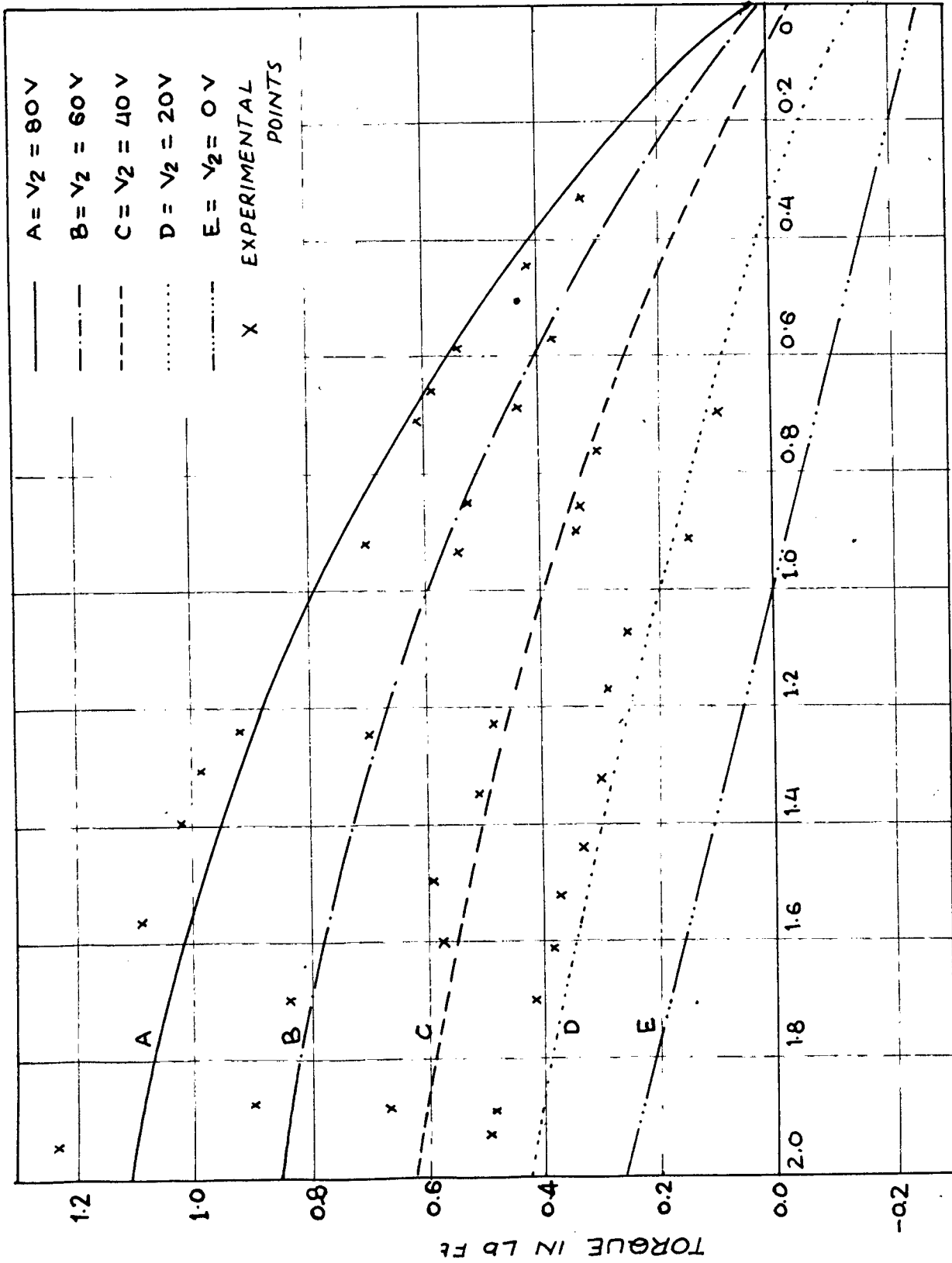
SLIP IN P.U.
 TORQUE SPEED CURVE (COPPER THIN CUP)
 WITH BALANCED APPLIED VOLTAGES $V_1 = V_2 = 110$



SLIP IN P.U.

TORQUE-SPEED CURVE (COPPER THIN CUP)
 WITH UNBALANCED APPLIED VOLTAGES $V_1=110V$

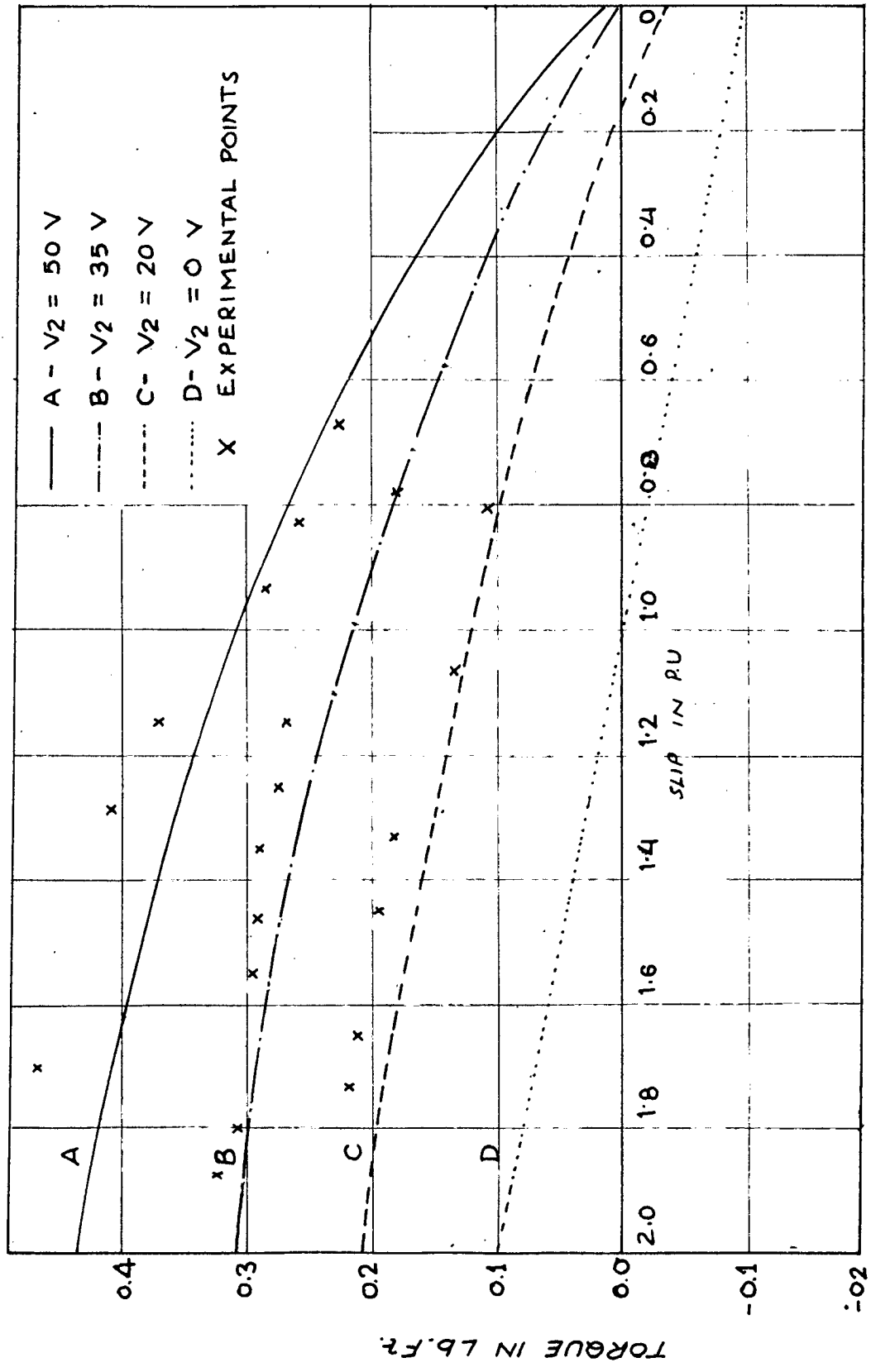
Fig 5.16



SLIP IN PU.

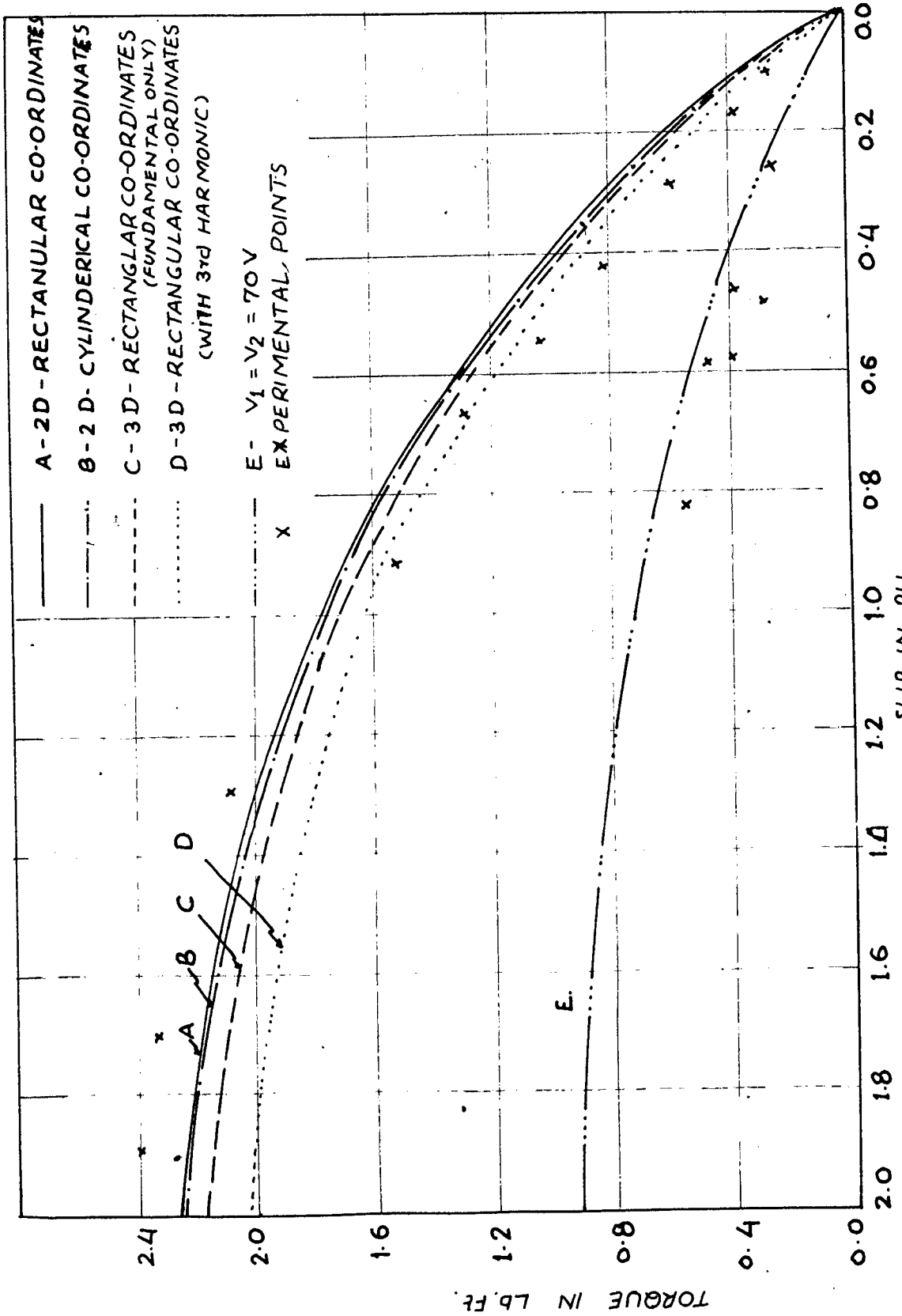
TORQUE - SPEED CURVE (COPPER THIN CUP)
 WITH UNBALANCED APPLIED VOLTAGE $V_1 = 80\text{ V}$

Fig. 5.1C



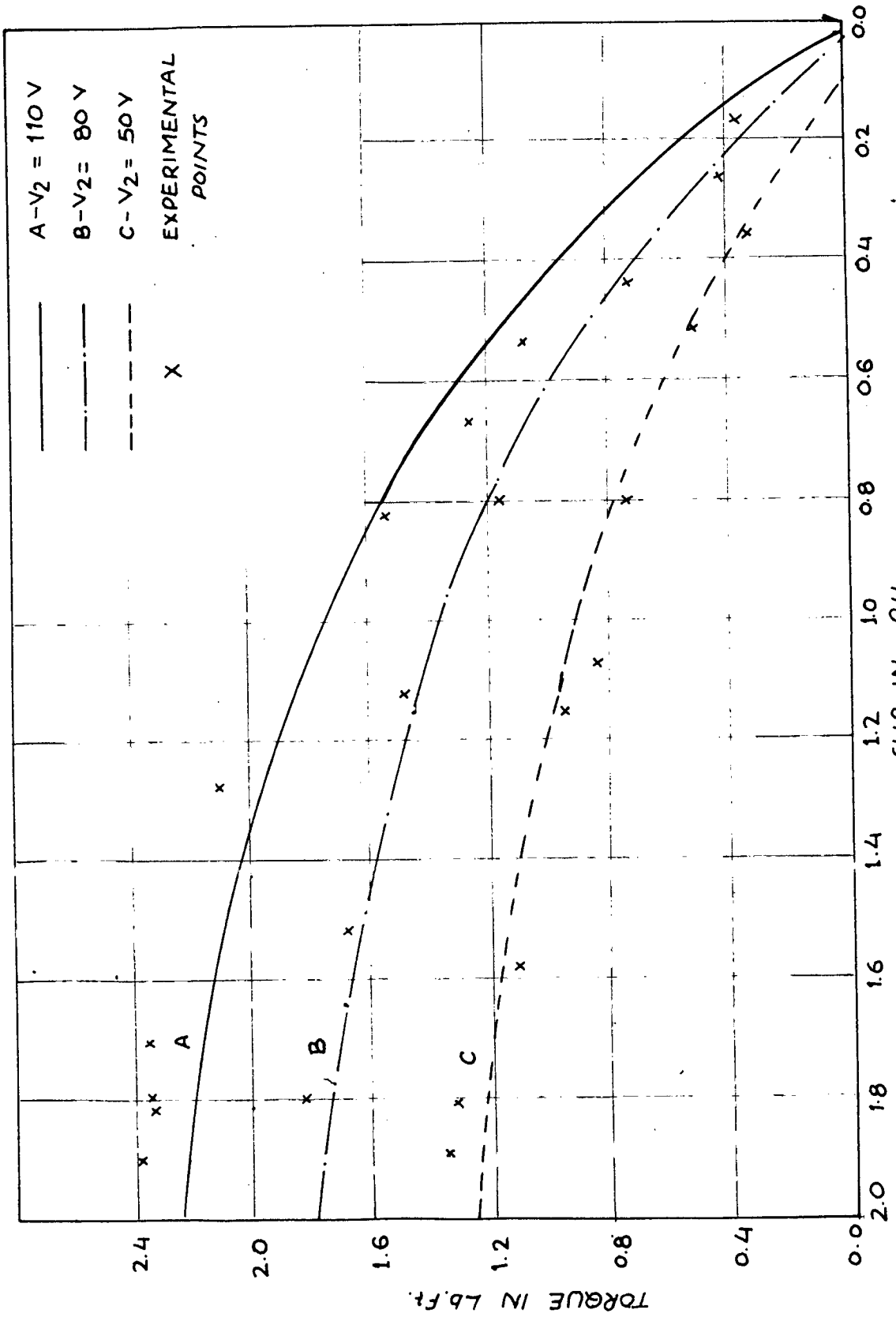
TORQUE - SPEED CURVE (COPPER THIN CUP)
 WITH UNBALANCED APPLIED VOLTAGES ($V_1 = 50$ V)

Fig. 5.1d



TORQUE - SPEED CURVE (COPPER THICK CUP)
 WITH BALANCED APPLIED VOLTAGES $V_1 = V_2 = 110$ V

FIG 52 Q



SPEED - TORQUE CURVE (COPPER THICK CUP)
 WITH UNBALANCED APPLIED VOLTAGES V=110V.

FIG-5.2b

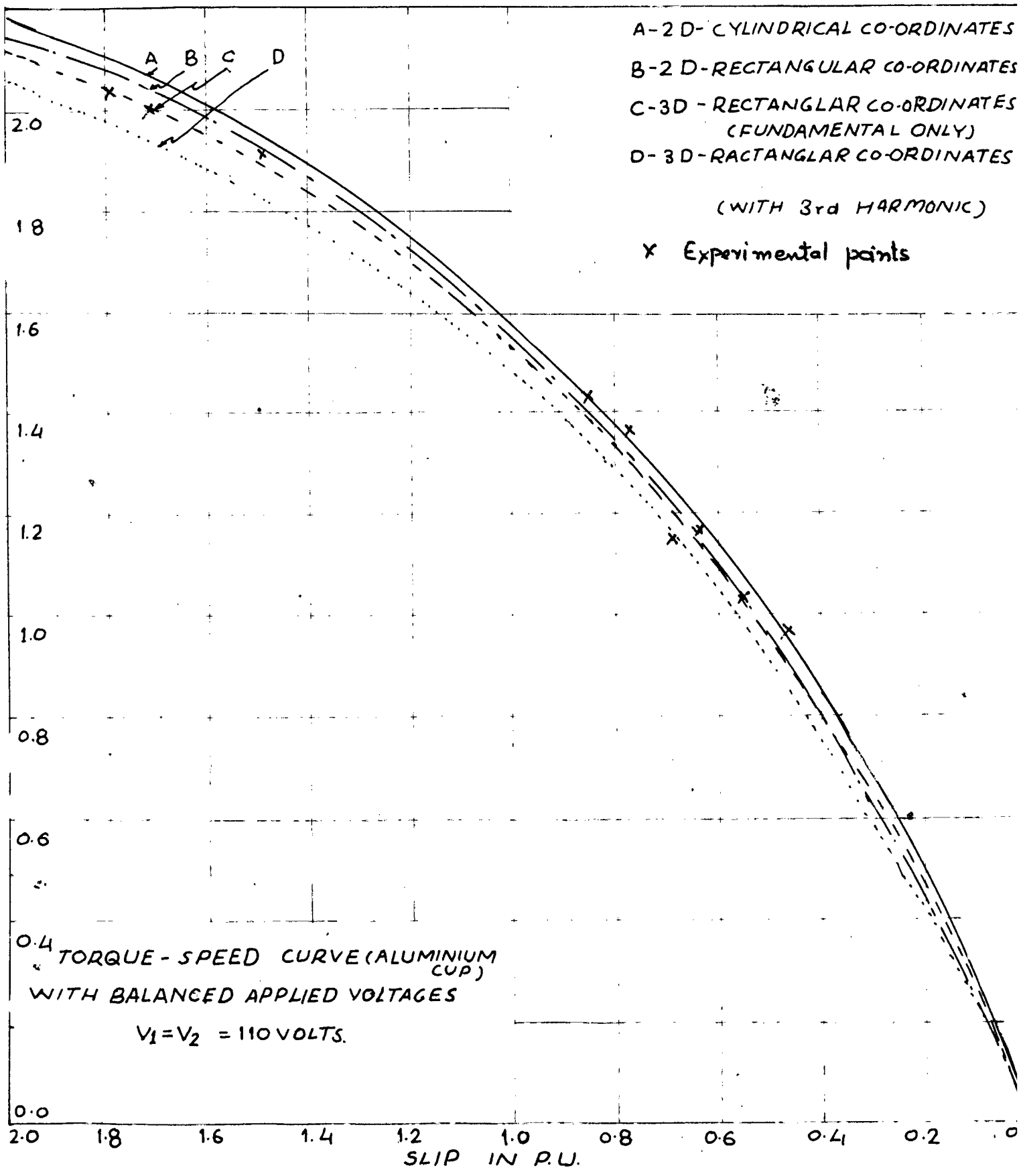
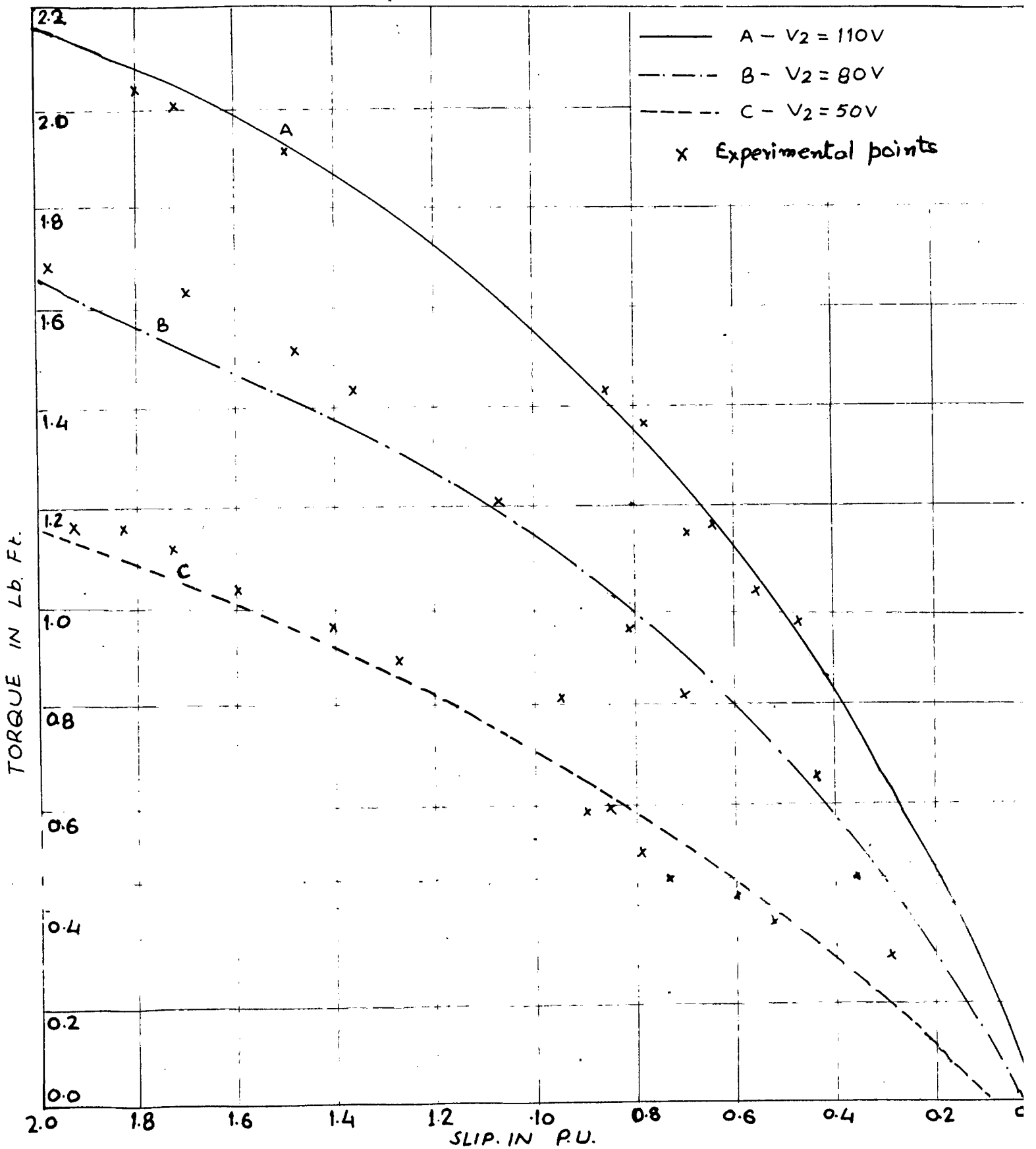
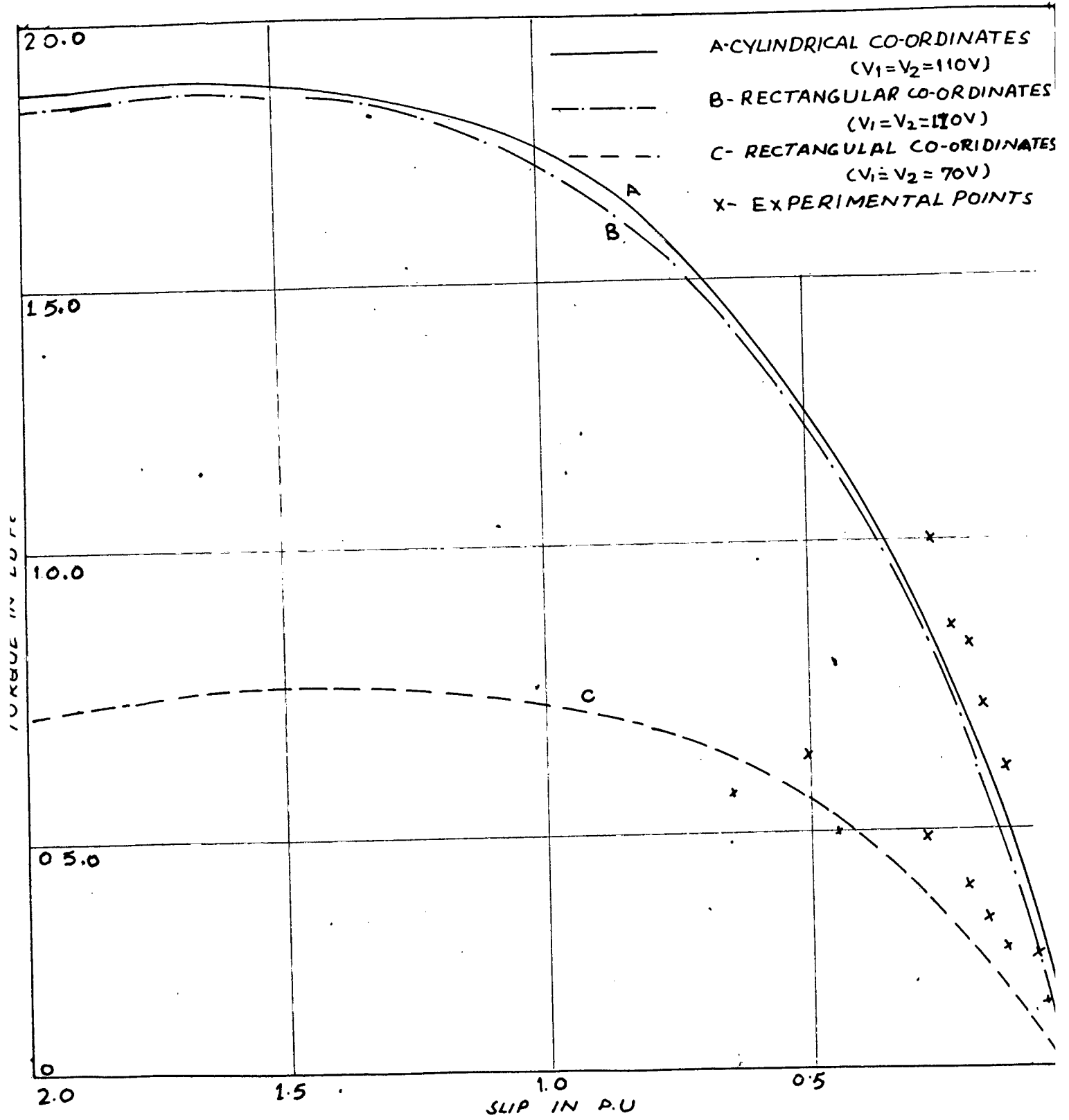


FIG. 53a



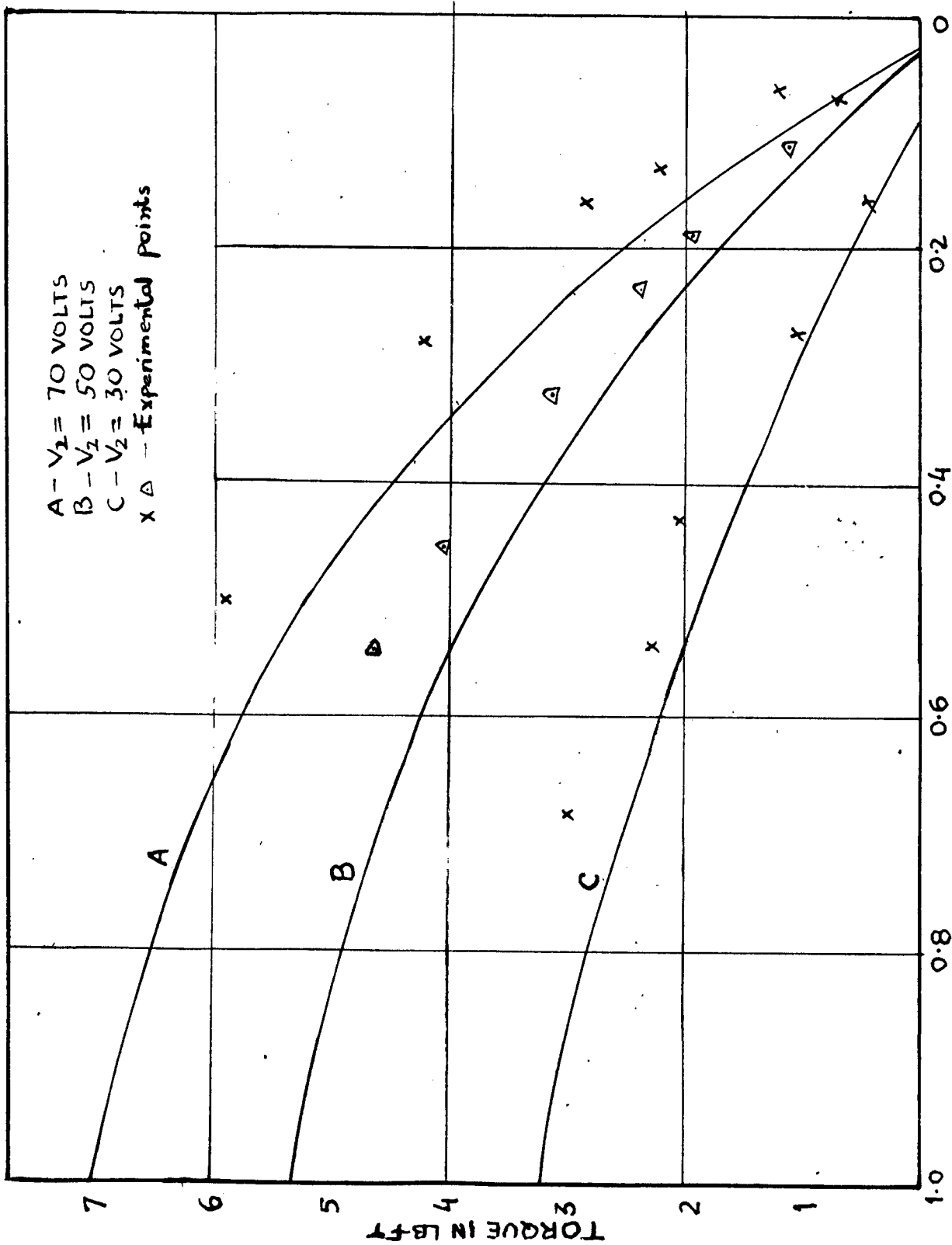
TORQUE SPEED CURVE (ALUMINIUM CUP)
 WITH UNBALANCED APPLIED VOLTAGES $V_1 = 110V$.

FIG- 5.3b

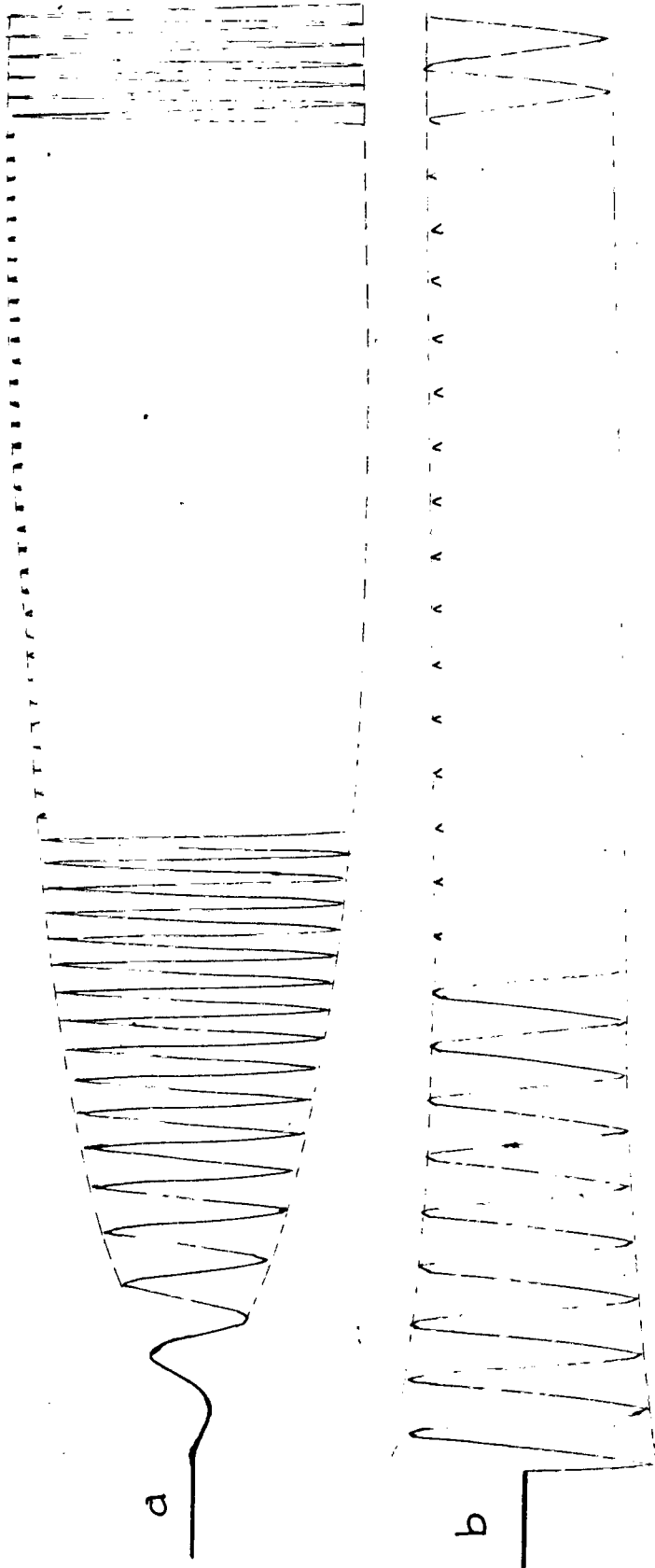


TORQUE-SPEED CURVE LARGE MOTOR WITH BALANCED APPLIED VOLTAGES

FIG 5.40



SLIP IN P. U.
 $V_1 = 70$ VOLTS
 TORQUE-SPEED CURVE (LARGE MOTOR)
 WITH UNBALANCED APPLIED VOLTAGES
 FIG. 5.4 b



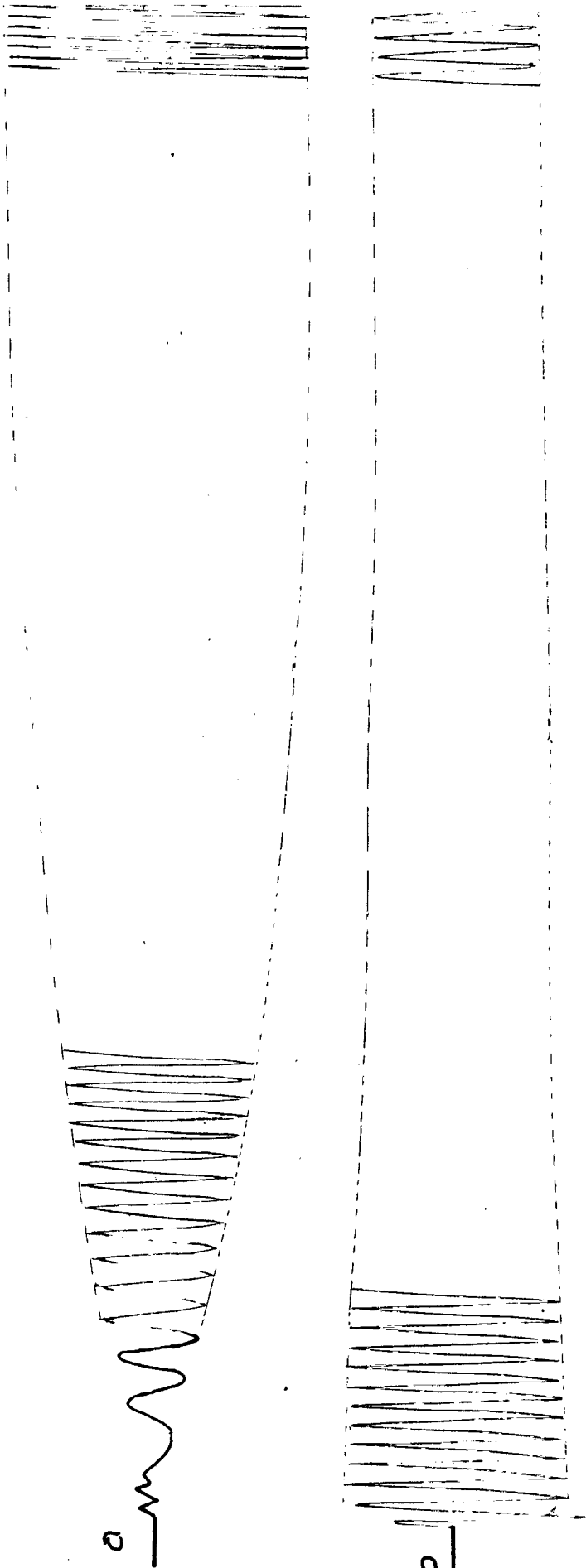
$V_1 = V_2 = 110 \text{ V}$

FINAL SPEED = 1350 R.P.M

a - SPEED TRANSIENT b - CURRENT TRANSIENT

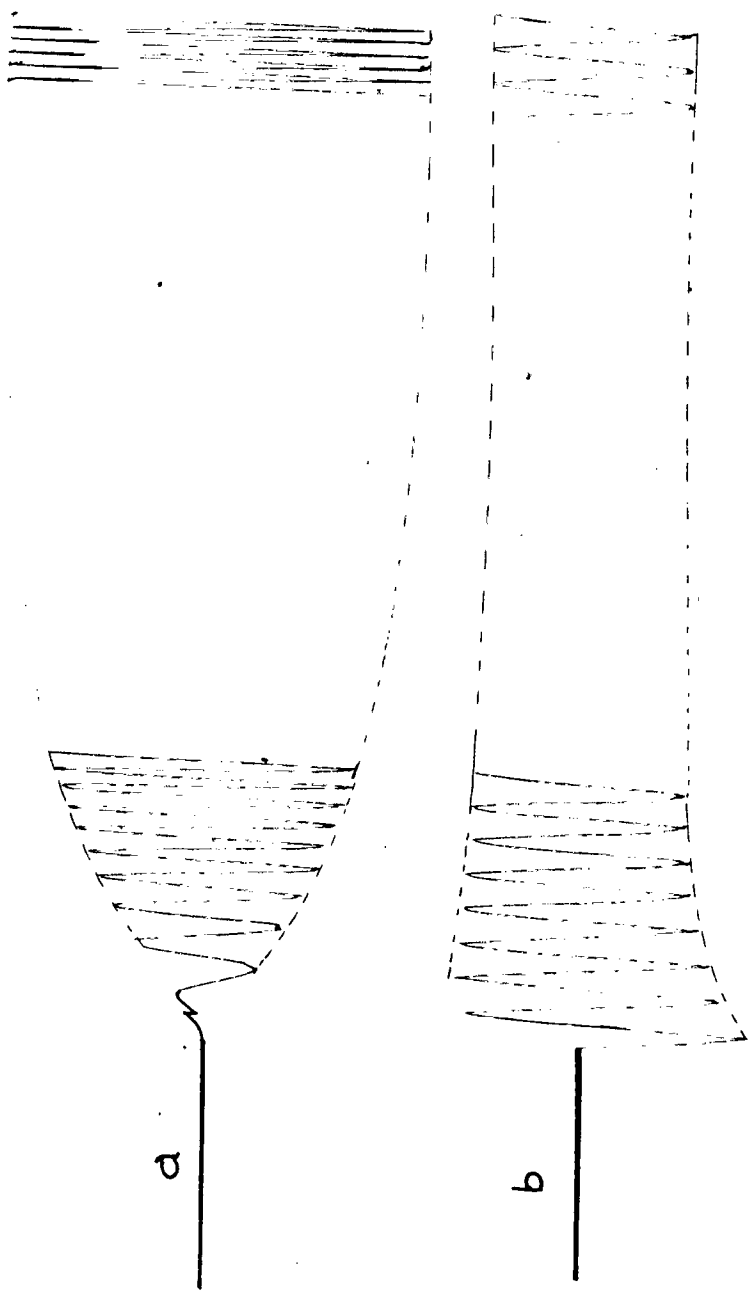
SPEED RISE (COPPER THICK CUP) AT NO LOAD

FIG. 5.5



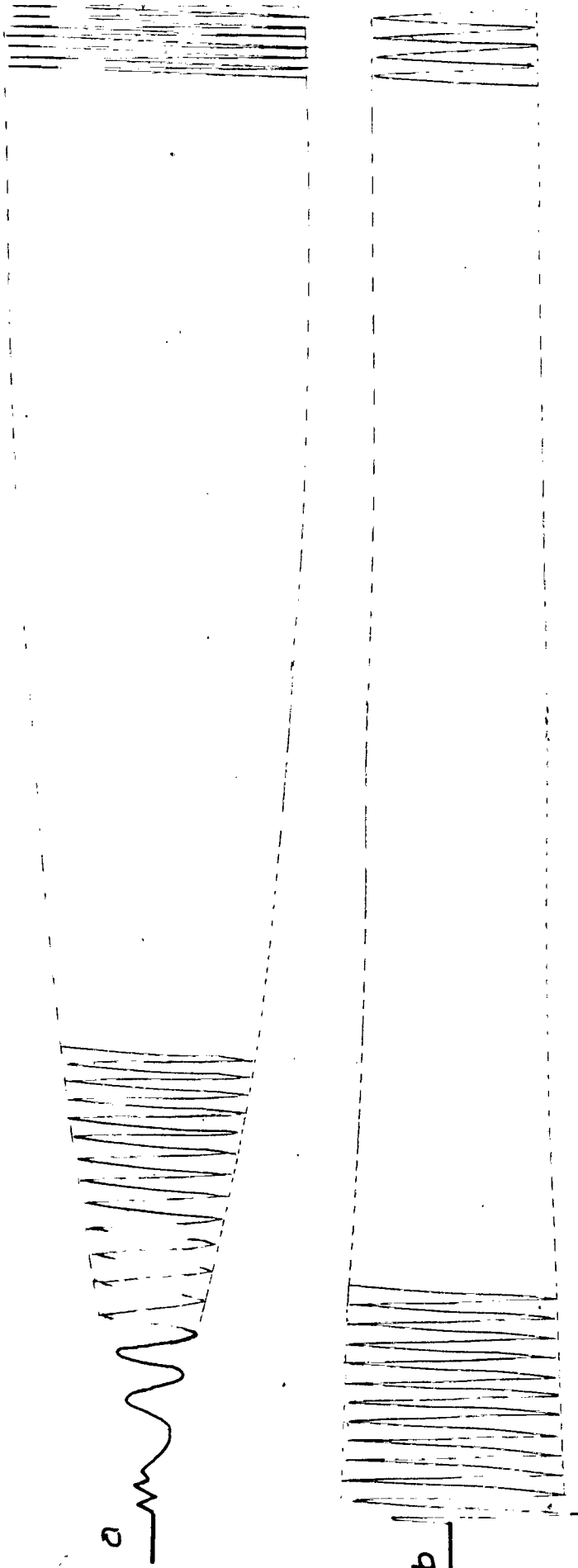
$V_1 = V_2 = 110V$
FINAL SPEED = 1000 R.P.M
a- SPEED TRANSIENT b-CURRENT TRANSIENT
SPEED RISE (COPPER THIN. CUP) AT LOAD

FIG. 5.6b



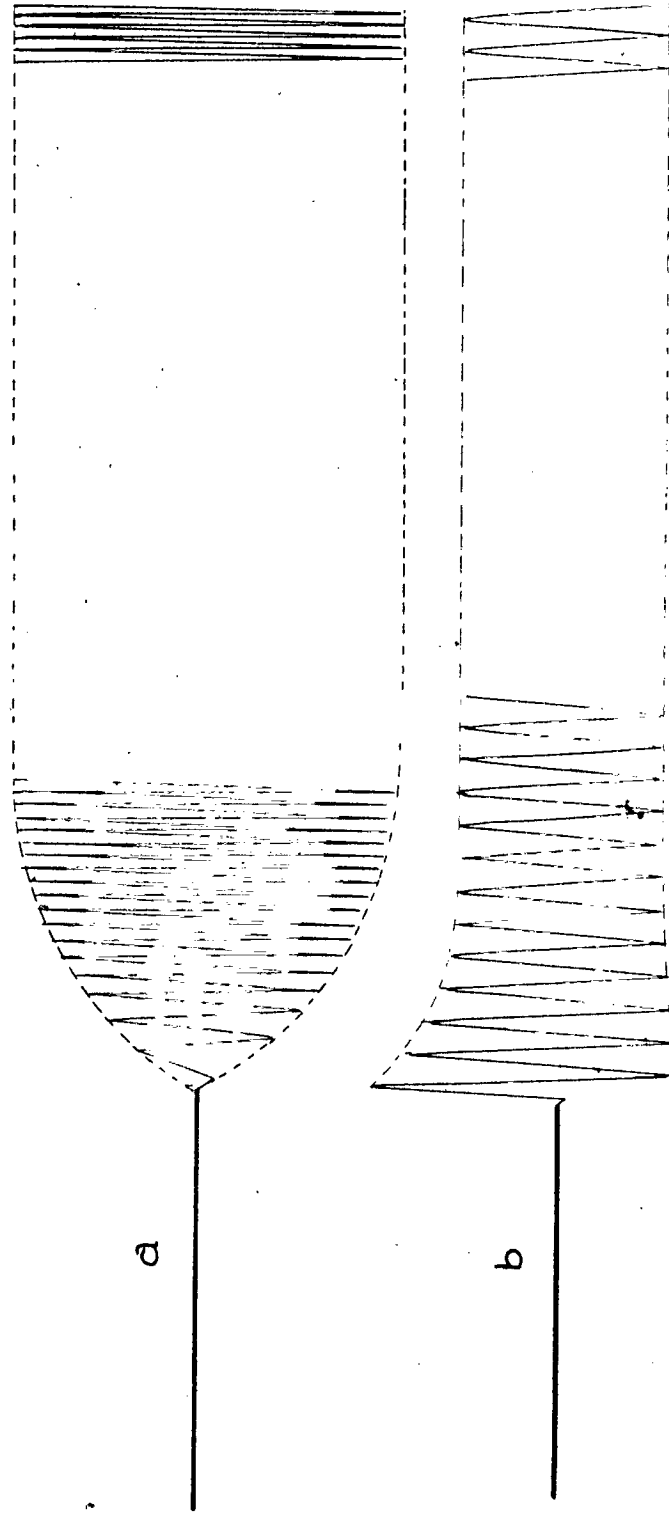
$V_1 = V_2 = 110V$
FINAL SPEED = 1360 RPM
a SPEED TRANSIENT b - CURRENT TRANSIENT
SPEED RISE (COPPER THIN CUP) AT NO LOAD

FIG. 5.6a



$V_1 = V_2 = 110V$
FINAL SPEED = 1000 R.P.M
a- SPEED TRANSIENT b-CURRENT TRANSIENT
SPEED RISE (COPPER THIN. CUP) AT LOAD

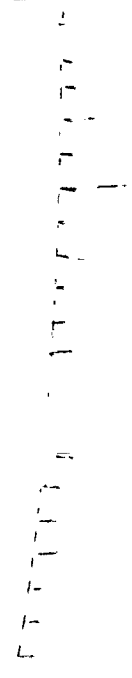
FIG. 5.6b



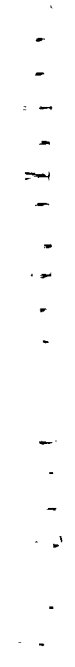
$V_1 = V_2 = 110V$
FINAL SPEED = 1380 R.P.M
a - SPEED TRANSIENT b - CURRENT TRANSIENT

SPEED RISE (ALUMINIUM CUP) AT NO LOAD

FIG. 5.7a



b



C_O_N_C_L_U_S_I_O_N_S.

In the present work it has been shown that the curvature of the machine has little effect on torque-speed^{characteristic} except at low values of slip. The results of rectangular co-ordinates are in close agreement with those of cylindrical ones. The three dimensional analysis gives better results over two dimensional analysis in a restricted region of torque speed curve. Results differ significantly at higher slips. Effect of harmonics is pronounced particularly when /^{tangentially} directed component of vector potential is taken into account.

The so-called $Z_{\text{gap-rotor}}$ the combined impedance of rotor airgaps and central core leads to equivalent circuit which is by no means simple as the rotor resistance and reactance are both functions of slip and the number of harmonics of Fourier series.

APPENDIX I

I.1. EFFECT OF NEGLECTING TANG. COMPONENT OF VECTOR POTENTIAL:

When tangential component of vector potential or rotor current is neglected we have-

$$A_{gx1} = 0$$

Now from the analysis of Chapter 3, section 3.1.4 from expression for voltage induced in a conductor we have-

$$v(\text{cond}) = \int_0^{\lambda} (T_z E_{gz})_{y=g_1} dz$$

Voltage induced in a phase of stator winding is-

$$V_1 = p \int_0^T v(\text{cond}) dx_1$$

Substituting the values of T_z and E_{gz} in expression for $v(\text{cond})$ and carrying out the required integrations one obtains-

$$V_1 = \frac{3}{8} \frac{p \lambda \mu \omega T}{\beta_n \lambda^2} \left(\frac{4}{n\pi} \right)^2 B_1^2 a_5 \cdot I_1 \sin \omega t$$

Dividing by $I_1 \sin \omega t$,

$$Z_{\text{gap-rotor}} = \frac{6p \lambda \omega \mu \cdot B_1^2}{\beta_n \pi n^2 a} \cdot a_5 \quad \dots \quad (I-1)$$

I.2. ANALYSIS BASED ON AMPERE TURN CONCEPT:

At stator surface $y = g_1$, ampere turn has sinusoidal variation in x direction and periodic variation in z direction. Therefore,

$$\Delta T = \frac{4}{n\pi} M \sum_{n=1,3..} \sin \frac{n\pi z}{\lambda} \exp(j(\omega t + \frac{x\pi}{T}))$$

where M is the peak value of ampere turn.

$$\begin{aligned} H_x &= \frac{\partial(\Delta T)}{\partial x} \\ &= \frac{\pi}{T} \sum_{n=1,3..} \frac{4}{n\pi} M \cdot \sin \frac{n\pi z}{\lambda} \exp(j(\omega t + \frac{\pi x}{T})) \end{aligned} \quad \dots \quad (I.2)$$

Flux density variation in x direction is sinusoidal and periodic in z direction. We may therefore write,

$$H_y = \sum_{n=1,3,\dots} \frac{A_n}{n\pi} \frac{B_0}{\mu_0} \exp(j(\omega t + \frac{\pi x}{l})) \left(\sin \frac{n\pi z}{\lambda} \right) \dots (1.3)$$

From (3.25) and (3.10) of Chapter 3,

$$H_x = \sum_{n=1,3,\dots} \frac{\beta_n}{\mu_0} \exp(j(\omega t + \frac{\pi x}{l})) (A_{n1} \sinh \beta_n z_1 + B_{n1} \cosh \beta_n z_1) \sin \frac{n\pi z}{\lambda} \dots (1.4)$$

$$\text{and } H_y = -\sum_{n=1,3,\dots} j \frac{1}{\mu_0} \frac{\pi}{l} \exp(j(\omega t + \frac{\pi x}{l})) (A_{n1} \cosh \beta_n z_1 + B_{n1} \sinh \beta_n z_1) \dots (1.5)$$

For equations (1.2), (1.4) and (1.3), (1.5) to be identical one must have-

$$\frac{\pi}{l} \cdot \frac{A_n}{n\pi} \pi = \frac{\beta_n}{\mu_0} (A_{n1} \sinh \beta_n z_1 + B_{n1} \cosh \beta_n z_1)$$

$$\frac{A_n}{n} = \frac{-j}{\mu_0} \frac{\pi}{l} (A_{n1} \cosh \beta_n z_1 + B_{n1} \sinh \beta_n z_1)$$

From the c-

$$\pi = j \frac{\beta_n}{\mu_0} \frac{\omega a^2}{\pi^2 c_5}$$

Therefore,

$$Z_{\text{cup-rotor}} = \frac{\text{volts } \times \text{turns}}{\pi} \\ \text{Low voltage} = \frac{4.442 \frac{f}{\text{ph}} K}{p} D.$$

$$\text{Therefore, } Z_{\text{cup-rotor}} = \frac{4.442 \frac{f^2}{\text{ph}} K \pi D \lambda \mu_0 a^2 a_s}{p \beta_n} \dots (1.6)$$

Now since $\beta_n^2 = \frac{\mu^2 \pi^2}{\lambda^2} + a^2$ a comparison of (3.44) and (1.6)

indicates that these two expressions will be identical if $\beta_n^2 = a^2$.

So neglecting tangential component of vector potential will give

accurate results only if $\beta_n \approx 0$ that means the term $\frac{n\pi}{\lambda}$ is small. This would in turn mean that λ , the machine length in a direction must be large for this assumption to hold. So the consideration of ^{tangential} component of vector potential is necessary if machine length is not very large.

APPENDIX II

SOME RELATIONS OF BESSEL FUNCTIONS

Bessel function of 1st kind and order n is given by-

$$J_n(z) = \sum_{m=0}^{\infty} \frac{(-1)^m (z/2)^{n+2m}}{\Gamma(m) \Gamma(n+m)} \dots \quad (1)$$

where z may be real or complex.

Bessel function of 2nd kind and order n is given by-

$$Y_n(z) = 2J_n(z) \left(\log_e \frac{z}{2} + \psi \right) - \sum_{m=0}^{\infty} \frac{(-1)^m (z/2)^{n+2m}}{\Gamma(m) \Gamma(n+m)} \left(\sum_{r=1}^{n+m} \frac{1}{r} + \sum_{r=1}^n \frac{1}{r} \right) - \sum_{m=0}^{n-1} \frac{(z/2)^{-n+2m} \Gamma(n-m-1)}{\Gamma(m)} \dots \quad (2)$$

where $\psi =$ Euler's constant. $= 0.5772157$

When argument is purely imaginary it is convenient to use modified Bessel functions which are defined by-

$$I_n(x) = \sum_{m=0}^{\infty} \frac{(x/2)^{n+2m}}{\Gamma(m) \Gamma(n+m)} \dots \quad (3)$$

$$K_n(x) = \frac{1}{2} \sum_{m=0}^{\infty} \frac{(-1)^m \Gamma(n-m-1) (x/2)^{-n+2m}}{\Gamma(m)} +$$

$$\frac{(x/2)^{n+2m}}{\Gamma(m) \Gamma(n+m)} \left\{ \log_e \frac{x}{2} + \psi - \frac{1}{2} \sum_{r=1}^m \frac{1}{r} - \frac{1}{2} \sum_{r=1}^{n+m} \frac{1}{r} \right\} \dots \quad (4)$$

Expressions (1) to (4) above are slowly convergent and are not suited for computation of functions of large arguments. Use is made of asymptotic expansions when arguments are large. They are defined as under-

$$I_n(x) = \frac{e^x}{(2x)^{\frac{1}{2}}} \sum_{m=0}^{\infty} \frac{(-1)^m (n, m)}{(2x)^m} +$$

$$\frac{e^{-x + j(n+\frac{1}{2})\pi}}{(2\pi x)^{\frac{1}{2}}} \sum_{m=0}^{\infty} \frac{(n, m)}{(2x)^m}$$

$$\text{where, } (n, m) = \frac{(4n^2-1^2)(4n^2-3^2)\dots(4n^2-(2m-1)^2)}{2^{2m} \underline{m}}$$

$$K_n(x) = \left(\frac{\pi}{2x}\right)^{\frac{1}{2}} \exp(-x) \cdot \left[1 + \frac{4n^2-1^2}{\underline{1}(8x)} + \frac{(4n^2-1^2)(4n^2-3^2)}{\underline{2}(8x)^2} + \right.$$

$$\left. \frac{(4n^2-1^2)(4n^2-3^2)(4n^2-5^2)}{\underline{3}(8x)^3} + \dots \right]$$

$$J_n(z) \sim \left(\frac{2}{\pi z}\right)^{\frac{1}{2}} \left\{ \left[1 - \frac{(4n^2-1^2)(4n^2-3^2)}{\underline{2}(8z)^2} + \right. \right.$$

$$\left. \frac{(4n^2-1^2)(4n^2-3^2)(4n^2-5^2)(4n^2-7^2)}{\underline{4}(8z)^4} - \dots \right\} \cos \varphi -$$

$$\left\{ \frac{4n^2-1^2}{\underline{1}(8z)} - \frac{(4n^2-1^2)(4n^2-3^2)(4n^2-5^2)}{\underline{3}(8z)^3} + \dots \right\}$$

$$\sin \varphi.$$

$$Y_n(z) \sim \left(\frac{2}{\pi z}\right)^{\frac{1}{2}} \left[\left\{ 1 - \frac{(4n^2-1^2)(4n^2-3^2)}{\underline{2}(8z)^2} + \right. \right.$$

$$\left. \frac{(4n^2-1^2)(4n^2-3^2)(4n^2-5^2)(4n^2-7^2)}{\underline{4}(8z)^4} - \dots \right\} \sin \varphi +$$

$$\left. \left\{ \frac{4n^2-1^2}{\underline{1}(8z)} - \frac{(4n^2-1^2)(4n^2-3^2)(4n^2-5^2)}{\underline{3}(8z)^3} + \dots \right\} \cos \varphi \right]$$

$$\text{where } \varphi = z - (n+\frac{1}{2})\frac{\pi}{2}$$

and $|z| \gg 1$, $|z| \gg |n|$

$$I'_\nu(x) = \frac{1}{2} (I_{\nu+1}(x) + I_{\nu-1}(x))$$

$$K'_\nu(x) = -\frac{1}{2} (K_{\nu+1}(x) + K_{\nu-1}(x))$$

$$J'_\nu(z) = \frac{1}{2} (J_{\nu-1}(z) - J_{\nu+1}(z))$$

$$Y'_\nu(z) = \frac{1}{2} (Y_{\nu-1}(z) - Y_{\nu+1}(z))$$

APPENDIX-IIIDESIGN CONSIDERATIONS

The design of servomotors differs from that of conventional motors in several aspects. This is because of requirements demanded in control system applications. These are quick response, ability to operate over a wide range of speed, linearity of torque speed characteristic, maximum torque at stall, ability to withstand the continuously applied reference phase voltage and not to single phase when error signal is zero.

In order to achieve low inertia which is necessary for low motor time constant a small diameter rotor is employed while at the same time employing as large a rotor as practicable to maintain the torque capability. To further reduce the inertia the thickness of cup should be small. This results in a quick response motor but at the same time leads to high rotor resistance and, therefore, limiting the output of the machine. In higher rating machines a thick cup must be used. This not only increases the acceleration time constant but results in an excessive input current which may not be easy to handle and at the same time the linearity of the torque speed characteristic is worsened. This is because of an excessively large airgap which is created by the thick cup which forms a part of the airgap. The length of the airgap may be several times higher than that in conventional induction machines. One more trouble that may arise because of sufficiently thick cup is that the torque speed characteristic may become drooping at higher slips and thereby limiting the stability region of the motor. In a balanced design an agreement must therefore, ^{be} attained between the thickness of cup, maximum power output and the linearity of the torque speed characteristic. A compromise is also necessary

in the degree of linearity that can be achieved. The straighter the torque speed curve, the less stall torque available, which means smaller initial accelerations. In most of the servomotors a compromise is made so that the maximum torque occurs between slips of 1.5 and 2 per unit. This gives reasonable linearity and stalled torque.

It is essential that such a motor does not 'single-phase'. That is, when the control field voltage is reduced to zero, it should not develop a positive torque at any speed. This requirement is met if the torque speed characteristic peaks at a negative speed. It has been shown⁽²⁰⁾ that a motor will not single-phase provided that the secondary resistance is smaller than the airgap reactance.

In order to keep the exciting current low the clearances between the moving cup and the stationary parts should be limited only by mechanical considerations, allowance being made for uneven heating and gyroscopic torques.

The temperature rise is much pronounced in servomotors. This is because of the continuous excitation of one phase and large exciting current. This may be a problem in large motors. Unlike ordinary rotary machines the temperature cannot be lowered by increased machine dimensions the specific load is determined by the time constant condition as well as the capacity. Consequently this problem must be solved by heat resistant and insulating materials and cooling systems. For large capacities a separately driven blower may be necessary to cool the machine. In still bigger sizes it is necessary to add a circuit which will lower excitation voltage at the zero position to restrain the temperature rise.

DESIGN:

The design of test motor was carried out by the conventional method of induction motor design technique except for rotor. Some of the points earlier discussed were taken into consideration to modify the conventional motor design. The stator winding was designed to carry the extra excitation current required by the large airgap. Ratio of length to diameter was kept high in order to reduce the moment of inertia of rotor.

Design of cup rotor was done somewhat arbitrarily. Torque speed curve was predetermined and a value of cup thickness selected so as to meet the requirement of power output and linearity of torque speed curve.

ABSOLUTE DRAG DATA - 3/4 h.p. motor.Stator frame:

No. of slots = 24
 Tooth pitch = 0.49"
 Pole pitch = 2.95"
 Average dia = 3.75"
 stack length = 3"

Stator windings:

No. of coils = 24, frequency = 50 c/s.
 coil pitch = 6, No. of poles = 4
 No. of turns per coil = 40
 No. of turns in series per phase = 552
 Each coil of 2 strands of 24 AWG super enameled copper wire.
 Double layer two phase winding.
 Stator resistance at 20°C = 8.0 ohms
 Stator resistance at 50°C = 10.3 ohms
 Stator leakage reactance = 7.5 ohms

Rotor core-

14 mil transformer iron sheet laminations.

Rotor cup-

1. Copper mixed with 10% tin. Resistivity of rotor material at 60°C = 4.5×10^{-8} ohm meter.

2. Commercial aluminium. Resistivity at 60°C = 3.58×10^{-8} ohm meter.

Dimensions are as per fig. 1.

PRINCIPAL DESIGN DATA - 3 h.p. motor:Stator frame:

No. of slots = 36
 Tooth pitch = .48"
 Mean slot opening = .125"
 Average diameter = 5.5"
 Stack length = 6.5"

Stator windings:

No. of coils = 36
 Coil pitch = 4, 5 wound so as to obtain minimum of unbalance.

No. of poles = 4
 Frequency = 50 c/s

No. of turns per coil = 9

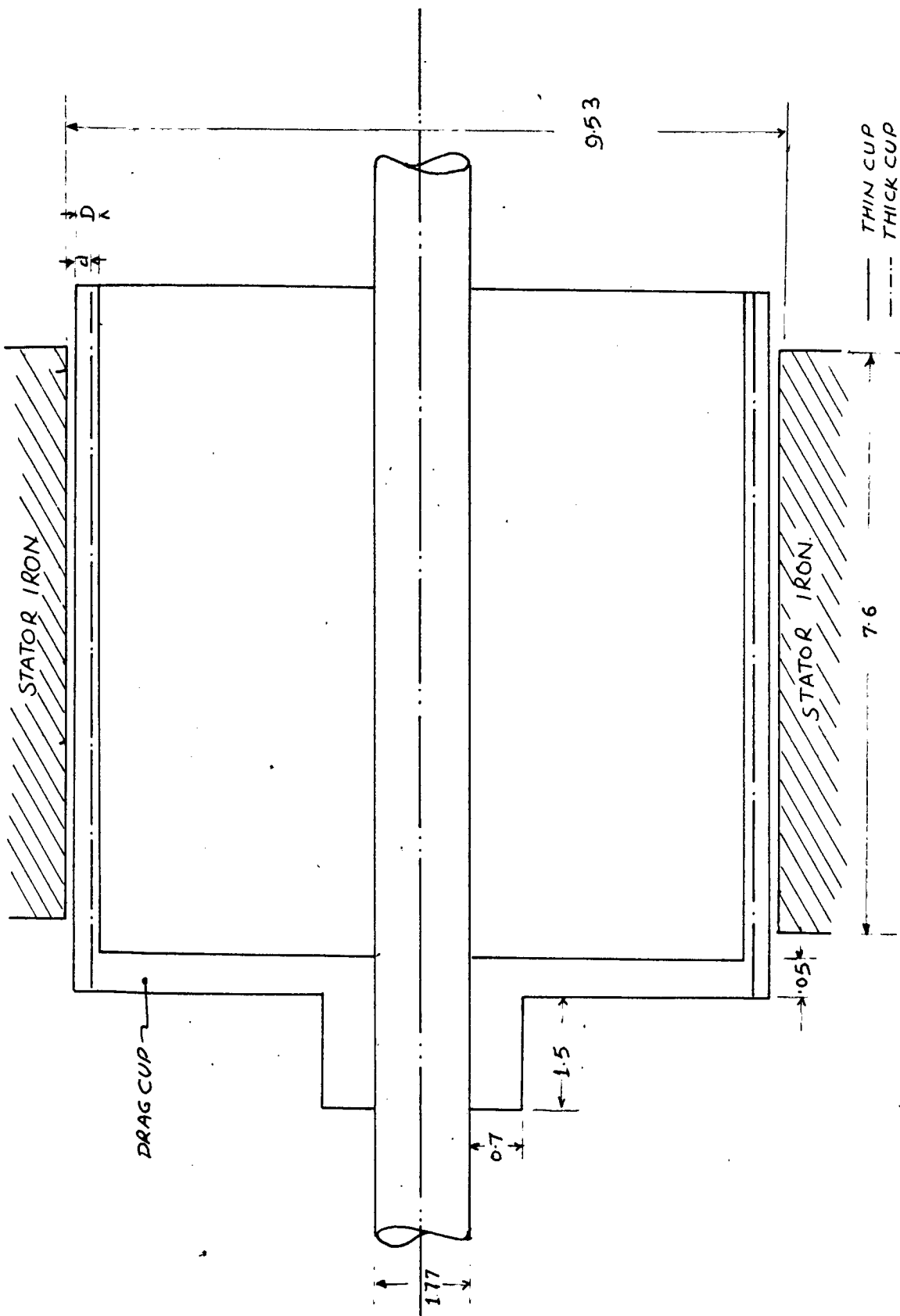
Each coil of 4 strands, 3 strands of 18 SWG and 1 strand of 19 SWG super-enamelled copper wire.

Double layer two phase winding.

No. of turns in series per phase = 162.

Stator resistance at 50°C = 10 ohms

Stator leakage reactance = 1.5 ohms



ALL DIMENSIONS ARE IN
CENTIMETERS
SCALE 4:3

ROTOR OF SMALL MOTOR

FIG - 1

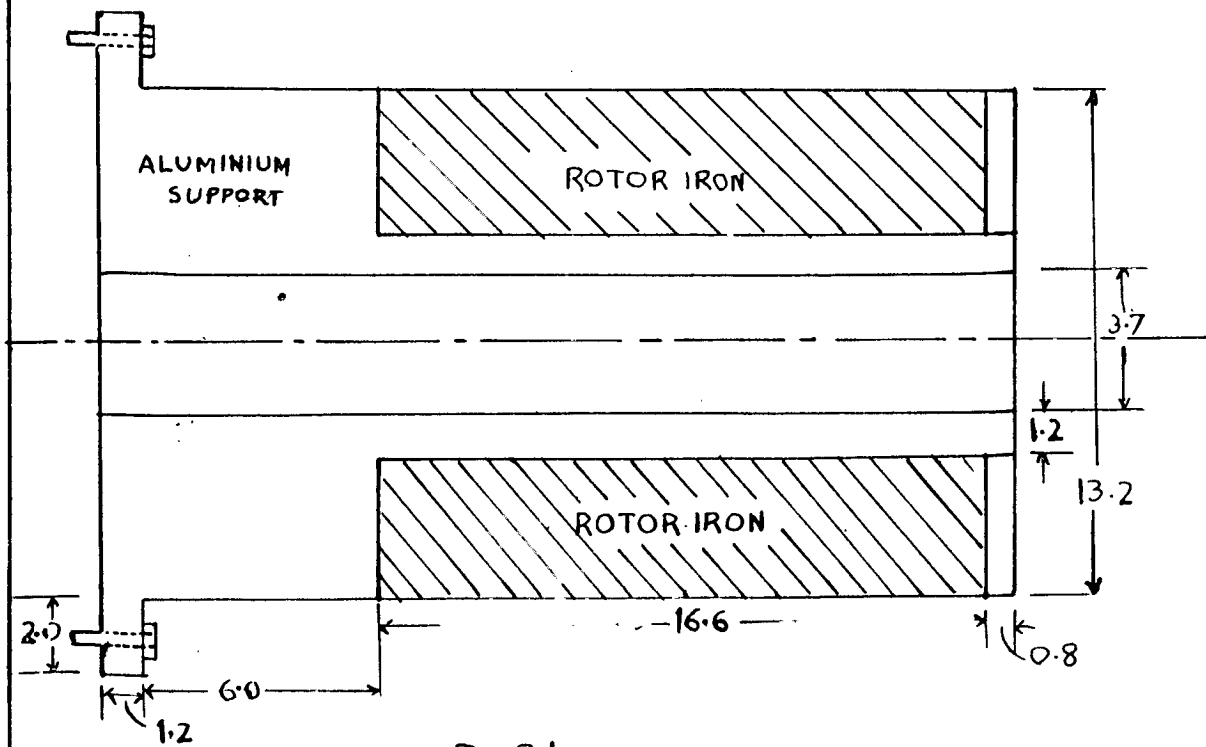
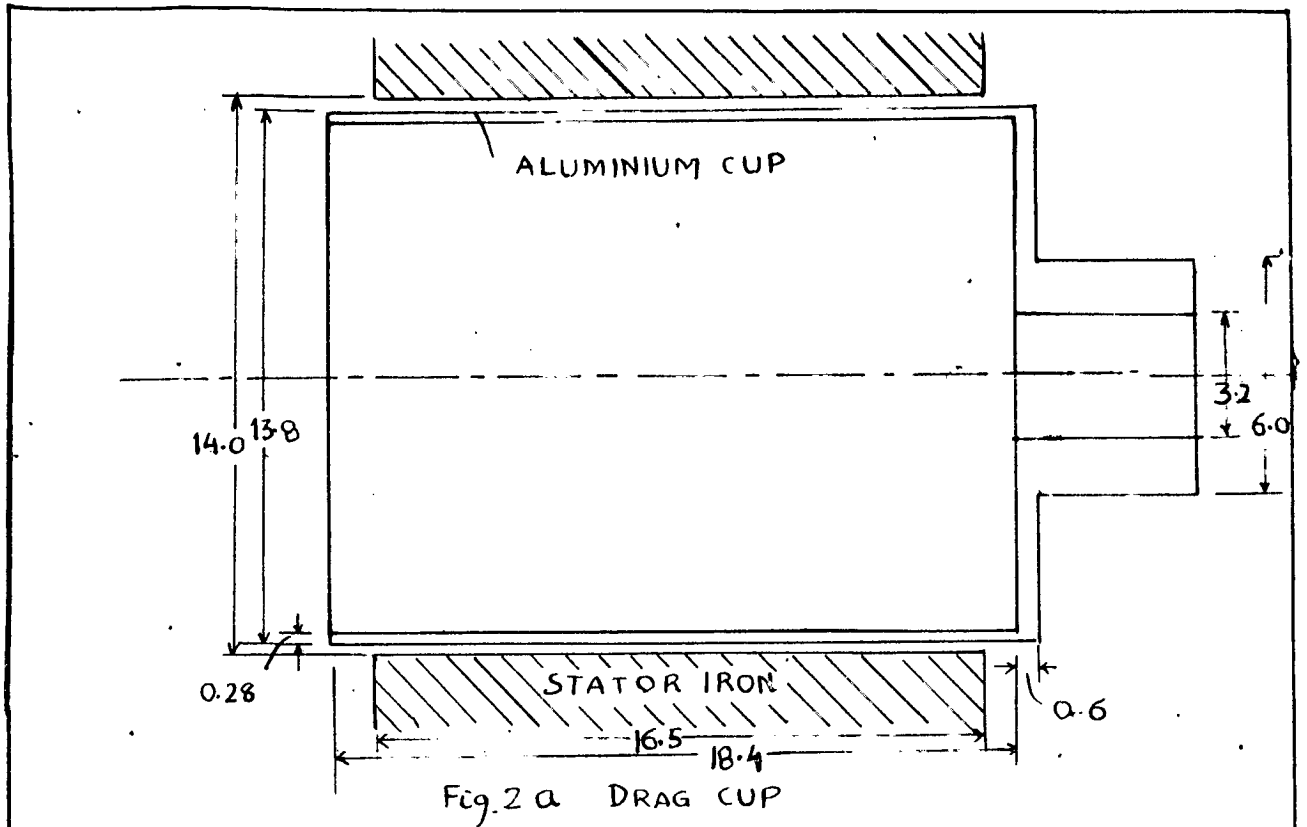


Fig 2 b CENTRAL CORE

ROTOR OF LARGE MOTOR

FIG. 2.
 (ALL DIMENSIONS ARE IN CM.)
 SCALE - HALF FULL SIZE

Rotor core:

18 mil transformer iron sheet laminations.

Rotor cup:

Commercial aluminium.

Resistivity at 60°C = 3.58×10^{-8} ohm meter

Dimensions are as per fig.2.

LIST OF NOTATIONS

QUANTITIES IN INPUT LIST

AMU=PERMEABILITY OF FREE SPACE
 OMEGA=ANGULAR FREQUENCY
 G1,G2,G3=MACHINE DIMENSIONS IN Y AXIS IN METERS
 DIA=DIAMETER IN INCHES AT STATOR BORE
 BO=AIRGAP FLUX DENSITY IN WEBERS PER METRE AT DESIGNED VOLTAGE
 VOLTS =DESIGN TERMINAL VOLTAGE
 TURNS =NO. OF TURNS PER PHASE
 SYN=SYNCHRONOUS SPEED
 R1=STATOR RESISTANCE
 X1=STATOR LEAKAGE REACTANCE
 RHO=RESISTIVITY OF CUP MATERIAL IN OHM METERS
 I1,I2,I3,XY ARE QUANTITIES USED TO CONTROL THE VARIATION OF SLIP
 (I1=REFERENCE PHASE VOLTAGE
 V2=CONTROL PHASE VOLTAGE
 I=NO. OF TIMES THE OPERATION IS REPEATED WITH DIFFERENT VALUES OF V1 AND V2

QUANTITIES IN OUTPUT LIST

ZFQ=EQUIVALENT IMPEDANCE OF ROTOR, ROTOR CORE AND AIRGAPS
 $Z^2 + JX^2 = ZEQ$
 TORQ=TORQUE IN LB-FT

```

C C TORQUE SPEED CURVE BY TWO DIMENSIONAL ANALYSIS IN RECT. CO-OR. Z
  DIMENSION SLIP(20),R2(20),X2(20),ZEQ(20)
  PUNCH101
101 FORMAT(5X50HTORQUE SPEED CHARACTERISTIC 2D CASE RECT. COORD.)
  READ30,L
30  FORMAT(I2)
  LL=1
312  READ50,AMU,OMEGA,G1,G2,G3
50   FORMAT(E20.5,4F10.0)
  READ20,DIA,BO,VOLTS,URNS,POLES,SYN
20   FORMAT(7F10.0)
  READ60,R1,X1
60   FORMAT(2F10.0)
  READ40,RHO
40   FORMAT(F20.5)
  READ102,I1,I2,I3,XY
102  FORMAT(3I2,F10.0)
  PUNCH300,RHO,G1,G2,G3
300  FORMAT(/4HRHO=E12.5,4X3HG1=F7.5,4X3HG2=F7.5,4X3HG3=F7.5)
  PY=3.14159
  TA=PY*DIA*.0254/POLES
  BRB=180./PY$ CCC=1./BRB$ DDD=2.*7.04/1500.
  AK=PY/TA$ TAK=AK+AK$ A=EXPF(TAK*G3)$ A1=-1./A
  B=EXPF(TAK*G2)$ G=EXPF(TAK*G1)$ X=AK*AK
  DOI1=I1,I2,I3$ AI=I$ SLIP(I)=AI/XY $ S=SLIP(I)
  Y=S*OMEGA*AMU/RHO$ RH=(X*X+Y*Y)**.25
  FI=ATANF(Y/X)$ FIH=FI/2.$ AMR=RH*COSF(FIH)$ AMI=RH*SINF(FIH)
  AA=2.*G2*AMR$ BB=2.*G2*AMI$ CC=EXPF(AA)$ CR=CC*COSF(BB)
  CI=CC*SINF(BB)$ DD=2.*G1*AMR$ EE=2.*G1*AMI$ GG=EXPF(DD)
  DR=GG*COSF(EE)$ DI=GG*SINF(EE)$ AB=A1*B$ F=(1.-AB)/(1.+AB)
  E=F*AMR$ H=F*AMI$ O=E-AK$ P=(E+AK)*CR-CI*H$ Q=CI*(F+AK)+CR*H
  R=P*P+Q*Q$ AR2=(O*P+H*Q)/R$ AI2=(H*P-O*Q)/R$ T=AR2*DR-AI2*DI
  U=AR2*DI+AI2*DR$ V=1.+T$ W=1.-T$ USQ=U*U$ Z=W*W+USQ
  PR=(V*W-USQ)/Z$ PI=(V*U+U*W)/Z$ HH=AK*PR$ OO=AK*PI$ PP=AMR-HH
  QQ=AMI-OO$ RR=(AMR+HH)*G$ SS=(AMI+OO)*G$ TT=RR*RR+SS*SS

```

```

AR3=(PP*RR+QQ*SS)/TT$ AI3=(QQ*RR-PP*SS)/TT$ UU=1.+AR3$ VV=1.-AR3
XX=AI3*AI3$ WW=VV*VV+XX$ YY=(UU*VV-XX)/WW$ ZZ=(UU*AI3+VV*AI3)/WW
AAA=BO/(AMU*AK)$ ATR=YY*AAA$ ATI=ZZ*AAA
ANGLE=ATANF(ATI/ATR)*BBB$ AT=SQRTF(ATR*ATR+ATI*ATI)
ZEQ(I)=VOLTS*URNS/AT$ ANG=(90.-ABSF(ANGLE))*CCC
R2(I)=ZEQ(I)*COSF(ANG)$ X2(I)=ZEQ(I)*SINF(ANG)
1 CONTINUE
31 READ150,V1,V2 $ VF=(V1+V2)/2.$ VB=(V1-V2)/2.
150 FORMAT(2F10.0)
PUNCH250,V1,V2
250 FORMAT(3HV1=F5.0,5X3HV2=F5.0)
VFSQ=VF*VF$ VBSQ=VB*VB
DO2I=I1,I2,I3 $ CFSQ=VFSQ/((R1+R2(I))**2+(X1+X2(I))**2)
IF(VB)22,21,22
21 CBSQ=0.$ GOTO51
22 CBSQ=VBSQ/((R1+R2(I2-I))**2+(X1+X2(I2-I))**2)
51 TF=CFSQ*R2(I)$ TB=CBSQ*R2(I2-I)$ TORQ=(TF-TB)*DDD
2 PUNCH200,SLIP(I),ZEQ(I),R2(I),X2(I),TORQ
200 FORMAT(F4.1,4F17.3)
LL=LL+1
IF(LL-L)31,31,32
32 STOP
END

```

SAMPLE DATA

```

1
.12566E-05314. .00055 .00365 .00395
5.5 .36 220. 162. 4.
1. 1.5
3.58E-08
120 1 10.
110. 110.

```

PROGRAM NO. 2

QUANTITIES IN INPUT LIST

NPH=NO. OF PHASES

Y1=STATOR RESISTANCE

Z1=STATOR LEAKAGE REACTANCE

V1=STATOR TERMINAL VOLTAGE

SPFED=SYNCHRONOUS SPEED IN RPM

M1,M2,M3,AAM ARE THE FACTORS TO CONTROL THE VARIATION OF SLIP

R1,R2,R3,R4 ARE RADII AT STATOR BORE, ROTOR OUTER, ROTOR INNER AND ROTOR CORE I

URNS=NO. OF TURNS IN SERIES PER PHASE

VOLT=DESIGN TERMINAL VOLTAGE

BO=AIRGAP FLUX DENSITY IN WB/METRE SQUARE AT DESIGNED VOLTAGE

QUANTITIES IN OUTPUT LIST

S=SLIP IN PER UNIT

ANGLE=ANGLE OF AMPERE-TURN

ZFE=EQUIV. IMPEDANCE OF ROTOR, AIRGAPS, ROTOR CORE

Y2+JZ2=ZEQ

```

C C TORQUE SPEED CURVE BY TWO DIMENSIONAL ANALYSIS IN CYLIN. COORD. Z
READ 150,NPH
150 FORMAT (I1)

```



```

READ 100,Y1,Z1,V1,SPEED
100 FORMAT(4F10.0)
  READ 50,M1,M2,M3,AAM
 50 FORMAT (3I2,F10.0)
  READ 10,OMEGA,AMEW,RHO,R1,R2,R3,R4
10 FORMAT(F6.0,E13.6,E11.4,4F10.0)
  RFAD20,URNS,VOLT,BO
20 FORMAT(3F10.0)
  RT=SQRTF(2.)$ TRT=2.*RT$ RTO=RT*OMEGA$ RTSIX=RT*6.$ RRR=R3
  TG=1.154432$ PY=3.14159$ AA=PY*RT$ DD=PY/12.$ EE=PY/2.
  AMEWA=0.125664E-05$ XA=R1**4$ XB=R3**4$ XC=R2**4$ A1=-R4**4
  XD=XB-A1$ XE=XB+A1$ OJ=180./PY$ OH=BO*R1/(2.*AMEWA)$ ANPH=NPH
  OK=VOLT*URNS$ B=OMEGA*AMEW/RHO$ DDD=ANPH*7.04/SPEED
  DO 1I=M1,M2,M3$ CONS=I$ S=CONS/AAM$ R3=RRR
C
  CALCULATION FOR A2,A4,A7,A9 BEGINS
  DO 2J=1,2$ GO TO (3,4),J
 4 R3=R2$ GO TO 3
 3 AL=SQRTF(S*B)$ AO=AL/RTO$ X=AL*R3/2.$ CC=2.*LOGF(X)$ GG=X**3
  XSQN=-X*X$ XX=XSQN$ YY=XX$ SIGN=XSQN*X
  D1=1.$ K=1$ SR1=X/RT$ SI1=SR1$ SR3=-GG/RTSIX$ SI3=-SR3
 6 AK=K$ D1=D1*AK*(AK+1.)$ Q=1.+2.*AK$ U=Q*.7854$ T=SIGN/D1
  TR=T*COSF(U)$ TI=T*SINF(U)$ SR1=SR1+TR$ SI1=SI1+TI$ K=K+1
  SIGN=SIGN*XSQN$ IF(ABSF(TR)-.000001)5,6,6
 5 IF(ABSF(TI)-.000001)7,6,6
 7 L=1$ SIGN=XSQN*GG$ D3=6.
 9 ZL=L$ D3=D3*ZL*(ZL+3.)$ R=3.+2.*ZL$ V=R*.7854$ W=SIGN/D3
  WR=W*COSF(V)$ WI=W*SINF(V)$ SR3=SR3+WR$ SI3=SI3+WI$ L=L+1
  SIGN=SIGN*XSQN$ IF(ABSF(WR)-.000001)8,9,9
 8 IF(ABSF(WI)-.000001)11,9,9
 11 A2R=AO*(SR1+SR3-SI1-SI3)$ A2I=AO*(SR1+SR3+SI1+SI3)
  A4R=AO*(SR3-SR1+SI3-SI1)$ A4I=AO*(SR1-SR3-SI1+SI3)
C
  CALCULATION FOR A3,A5,A8,A10 BEGINS
  AR=X/AAS$ AI=ARS$ KK=1$ D11=1.$ D2=CC-1.+TG$ DI=EE
 14 BK=KK$ CK=BK*EES$ D11=D11*BK*(BK+1.)$ SA=0.0
  DO 12M=1,CK$ AM=M
 12 SA=SA+1./AM$ AM1=AM+1.$ SR=SA+1./AM1$ BB=XX/D11
  BR=BB*COSF(CK)$ BI=BB*SINF(CK)$ CR=CC-SA-SB+TG$ CI=EE
  ER=BR*CR-BI*CI$ FI=BR*CI+BI*CR$ DR=DR+ER$ DI=DI+EI$ KK=KK+1
  XX=XX*XSQN
  IF(ABSF(ER)-.000001)13,14,14
 13 IF(ABSF(EI)-.000001)15,14,14
 15 FR=1./(AA*X)$ FI=-FR$ GR=AR*DR-AI*DI$ GI=AR*DI+AI*DR
  YR1=GR-FR$ YI1=GI-FI
C
  CALCULATION FOR 3RD ORDER BEGINS
  HR=-GG/AAS$ HI=-HR$ QR=(CC-1.83333+TG)/6.$ QI=DD$ D33=6.$ N=1
 19 CN=NS$ DN=CN*EES$ D33=D33*CN*(CN+3.)$ SC=0.$ SD=0.$ N3=N+3
  DO 16M=1,NS$ BM=M
 16 SC=SC+1./BM$ DO17M=1,N3$ CM=M
 17 SD=SD+1./CM$ O=YY/D33$ OR=O*COSF(DN)$ OI=O*SINF(DN)
  PR=CC-SC-SD+TG$ PI=EE$ RR=OR*PR-OI*PI$ RI=OR*PI+OI*PR
  YY=YY*XSQN
  QR=QR+RR$ QI=QI+RI$ N=N+1$ IF(ABSF(RR)-.000001)18,19,19
 18 IF(ABSF(RI)-.000001)21,19,19
 21 PP=RT/GG$ QQ=1./(X*RT)$ ZZ=X/TRT$ SR=(-PP+QQ+ZZ)/PY
  SI=(-PP-QQ+ZZ)/PY$ ZR=HR*QR-HI*QI$ ZI=HR*QI+HI*QR$ YR3=ZR-SR
  YI3=ZI-SI$ A3R=AO*(YR1+YR3-YI1-YI3)$ A3I=AO*(YR1+YR3+YI1+YI3)
  A5R=AO*(YR3-YR1+YI3-YI1)$ A5I=AO*(YR1-YR3-YI1+YI3)$ GO TO(23,22),J

```

```

23 B2R=A2R$ B2I=A2I$ B4R=A4R$ B4I=A4I$ B3R=A3R$ B5R=A5R$ B5I=A5I
   B3I=A3I$ GOTO2
22 A7R=A2R$ A7I=A2I$ A9R=A4R$ A9I=A4I$ A8R=A3R$ A8I=A3I$ A10R=A5R
   A10I=A5I
   2 CONTINUE
C   CALCULATON FOR A6,A11 BEGINS
   UR=B5R*XD-B3R*XE$ UI=B5I*XD-B3I*XE$ DDR=B2R*XE-B4R*XD
   DDI=B2I*XE-B4I*XD$ XG=DDR*DDR+DDI*DDI$ A6R=(UR*DDR+UI*DDI)/XG
   A6I=(UI*DDR-UR*DDI)/XG$ XH=A6R*A7R-A6I*A7I+A8R
   XI=A6I*A7R+A6R*A7I+A8I$ XJ=A6R*A9R-A6I*A9I+A10R
   XK=A6I*A9R+A6R*A9I+A10I$ XL=XH+XJ$ XM=XI+XK$ XN=XJ-XH$ XO=XK-XI
   XP=XN*XN+XO*XO$ XQ=XC*XP$ A11R=(XL*XN+XM*XO)/XQ
   A11I=(XM*XN-XL*XO)/XQ$ OA=A11R*XA$ OB=OA+1.$ OC=A11I*XA
   OD=OA-1.$ OE=OD*OD+OC*OC$ OF=(OB*OD+OC*OC)/OE
   OG=(OC*OD-OB*OC)/OE$ ATR=OG*OH$ ATI=-OF*OH
   AT=SQRTF(ATR*ATR+ATI*ATI)$ OI=ATI/ATR$ ANGLE=ATANF(OI)*OJ
   ZEQ=OK/AT$ ANG=(90.-ABSF(ANGLE))/OJ$ Y2=ZEQ*COSF(ANG)
   Z2=ZEQ*SINF(ANG)$ C1=V1/((Y1+Y2)*(Y1+Y2)+(Z1+Z2)*(Z1+Z2))
   C1R=C1*(Y1+Y2)$ C1I=-C1*(Z1+Z2)$ SDR=C1R*Y1-C1I*Z1
   SDI=C1I*Y1+C1R*Z1$ EOR=V1-SDR$ EOI=-SDI$ PIN=EOR*C1R+EOI*C1I
   TORQ=DDD*PIN
1   PUNCH 200,S,ANGLE,ZEQ,Y2,Z2,TORQ
200 FORMAT(F5.1,F11.2,4F14.3)
   STOP $ END

```

SAMPLE DATA

```

2
1.      1.5      220.      1500.
120 1 10.
314.   0.125664E-05 0.3580E-07.06985   .0693   .0662   .0659
162.   220.      .36

```

THE ABOVE PROGRAM IS ONLY FOR A MACHINE HAVING 4 POLES. FOR A MACHI
HAVING A DIFFERENT NO. OF POLES A SIMILAR PROGRAM CAN BE WRITTEN USING THE
PROPER ORDER BESSEL FUNCTIONS OR A GENERAL PROGRAM CAN BE WRITTEN ON THE LIN
OF PROGRAM NO. 4

PROGRAM NO. 3

```

C C TORQUE SPEED CURVE BY THREFF DIMENSIONAL ANALYSIS IN RECT. COOR. Z
  DIMENSION G(31),GSQ(31),S(40),BC(20),BD(20),R2(31),X2(31),X(31)
  RFAD100,B,RL,P,TPH,WDF$ READ101, R,G1,D,G2
  READ102,I1,I2,I3,N1,N2,N3,N4,AII $ READ103,R1,X1,SYN,V1
100  FORMAT(7F10.0)
101  FORMAT(F20.5,5F10.0)
102  FORMAT(7I2,5F10.0)
103  FORMAT(7F 10.0)
C SFNSE SWITCH 1 ON WILL PUNCH THE VALUES OF R2(N) AND X2(N) FOR VARIOU:
C VALUES OF N
  W=314.
  U=.12566E-05
  PY=3.14159265$ B=B*.0254$ RL=RL*.0254$ DD=2.*7.04*V1*V1/SYN
  C=U*W/R$ B1BAR=TPH*WDF/B$ A=P/B$ ASQ=A*A$ E=PY*PY/(RL*RL)
  CC=6.*P*RL*W*U*B1BAR*B1BAR*P /(PY*ASQ*A)
500  DO1N=N1,N2,N3$ SQN=N*N$ GSQ(N)=ASQ+SQN*E$ G(N)=SQRTF(GSQ(N))
  GN=G(N)$ F=EXPF(GN*G2)$ H=1./F$ A1=(F+H)/(F-H)$ O=EXPF(GN*D)
  Q=1./O$ T=(O-Q)*.5$ V=(Q+O)*.5$ X(N)=GN*(-A1*T+V)/(A1*V-T)
501  BA=EXPF(GN*G1)$ BB=1./BA$ BC(N)=.5*(BA-BB)$ BD(N)=.5*(BA+BB)
  CONTINUE
  DO2I=I1,I2,I3$ AI=I$ S(I)=AI/AII$ Y=C*S(I)
  DO3N=N1,N2,N3$ GSON=GSQ(N)$ GN=G(N)$ SQN=N*N $ Z=X(N)
  CALL SQRT(GSQN,Y,ZR,ZI)$ CALL DIV(Z,0.,ZR,ZI,A21,A22)
  CALL SICO(ZR*D,ZI*D,AA,AB,AC,AD)$ CALL MUL(A21,A22,AA,AB,AE,AF)
  CALL MUL(A21,A22,AC,AD,AG,AH)
  CALL DIV(AE+AC,AF+AD,AA+AG,AB+AH,A31,A32)
  CALL DIV(GN*A31,GN*A32,ZR,ZI,A41,A42)$ BCN=BC(N)$ BDN=BD(N)
503  CALL DIV(A41*BDN+BCN,A42*BDN,A41*BCN+BDN,A42*BCN,A51,A52)
  DA=CC*GN/SQN$ DB=DA*A51$ DC=DA*A52$ CALL MUL(0.,1.,DB,DC,DE,DF)
  R2(N)=DE$ X2(N)=DF
  3 CONTINUE
  IF(SFNSE SWITCH1)11,12
11  PUNCH200,S(I),(N,R2(N),X2(N),N=N1,N2,N4)
12  RR=0.$ XX=0.$ DO4N=N1,N2,N3$ RR=RR+R2(N)
  4 XX=XX+X2(N)$ TORQ=DD*RR/((R1+RR)**2+(X1+XX)**2)
  2 PUNCH201,S(I),TORQ,RR,XX
200  FORMAT(2HS=F4.1,2X,(2HN=I2,2F7.1,2HN=I2,2F7.1,2HN=I2,2F7.1,2HN=I2,
12F7.1))
201  FORMAT(2HS=F4.1,5X7HTORQUE=F7.3,5X3HR2=F7.2,5X3HX2=F7.2)$ STOP
  END

SUBROUTINE SQRT(A,B,C,D)$ E=(A*A+B*B)**.25$ TH=ATANF(B/A)*.5
C=E*COSF(TH)$ D=E*SINF(TH)$ RETURN $ END

SUBROUTINE DIV(A,B,C,D,E,F)$ H=C*C+D*D$ E=(A*C+B*D)/H
F=(B*C-A*D)/H$ RETURN $ END

SUBROUTINE MUL(A,B,C,D,E,F)$ E=A*C-B*D$ F=A*D+B*C$ RETURN $ END

SUBROUTINE SICO(A,B,C,D,CC,DD)$ E=EXPF(A)$ F=1./E
G=(E+F)*.5$ H=(E-F)*.5$ P=SINF(B)$ Q=COSF(B)
C=Q*H$ D=P*G$ CC=Q*G$ DD=P*H$ RETURNS $ END

```

SAMPLE DATA

3.7		3.5	4.	552.	.9
		4.50000E-08.	0.00027	.00175	.0019
120	1 1 7 2 2 10.		4.		
10.3		7.5	1500.	110.	

PROGRAM NO.4

LIST OF NOTATIONS

QUANTITIES IN INPUT LIST

RHO = RESISTIVITY OF ROTOR MATERIAL IN OHM METERS
 RL = MACHINE LENGTH IN METERS
 R1, R2, R3, R4 ARE RADII AT STATOR BORE, ROTOR OUTER, ROTOR INNER
 AND ROTOR CORE RESPECTIVELY.
 RS = STATOR RESISTANCE
 XS = STATOR LEAKAGE REACTANCE
 TURNS = NO. OF TURNS IN SERIES PER PHASE
 WK = WINDING DISTRIBUTION FACTOR
 NP = NO. OF POLE PAIRS
 I1, I2, I3, ALL ARE THE FACTORS TO CONTROL THE VARIATION OF SLIP
 K1, K2, K3, K4 ARE THE FACTORS TO CONTROL VARIATION OF NO. OF
 HARMONICS UP TO WHICH CALCULATION IS DESIRED.
 DIA = AVERAGE ROTOR DIAMETER AND METERS.

QUANTITIES IN OUTPUT LIST

XA = ROTOR RESISTANCE (AN HARMONIC CONSIDERED)
 XB = ROTOR REACTANCE (N HARMONIC CONSIDERED)
 TOR(N) = TORQUE IN LB.FT. WITH AN HARMONIC CONSIDERED.

PROGRAM NO. 4

```

C C TORQUE SPEED CURVE BY THREE DIMENSIONAL IN CYLIN. CO-OR. Z
COMMON PY,ACC
DIMENSION B(31),C(31),D(31),E(31),F(31),G(31),H(31),O(31),P(31)
DIMENSION Q(31),RR(31),XR(31),PD(31),XA(31),XB(31),RC(31),RD(31),T
10R(31)
READ100,RHO,RL,R1,R2,R3,R4
100 FORMAT(E10.2,6F10.0)
READ101,RS,XS,URNS,WK,NP
101 FORMAT(5F10.0,I2)
READ102,I1,I2,I3,N1,N2,N3,N4,A11,DIA,SYN,ACC
102 FORMAT(7I2,5F10.0)
NA=NP+1
NB=NP-1
PY=3.14159265
W=314.
R1SQ=R1*R1
PYSQ=PY*PY
PN=NP
U=.12566E-05
A=W*U/RHO
B1BAR=URNS*WK/DIA
B1SQ=B1BAR*B1BAR
PA=12.*B1SQ*U*W/PY*2.*P
PB=PYSQ*R1SQ/(PN*RL)
PC=PN*RL
DDD=2.*7.04*V1*V1/SYN
501 DO1N=N1,N2,N3
AN=N
B(N)=AN*PY/RL
BN=B(N)
BR3=BN*R3
BR4=BN*R4
CALL KNX(NA,BR4,AA)
CALL KNX(NB,BR4,AB)
CALL INX(NA,BR4,AC)
CALL INX(NB,BR4,AD)
502 A1=(AA+AB)/(AC+AD)
CALL INX(NA,BR3,AE)
CALL INX(NB,BR3,AF)
CALL KNX(NA,BR3,AG)
CALL KNX(NB,BR3,AH)
CALL INX(NP,BR3,AP)
CALL KNX(NP,BR3,AQ)
C(N)=B(N)*(A1*(AE+AF)-AG-AH)/(2.*(A1*AP+AQ))
503 BR2=BN*R2
BR1=BN*R1
CALL KNX(NP,BR2,DA)
CALL KNX(NA,BR2,DB)
CALL KNX(NB,BR2,DC)
CALL INX(NA,BR2,DE)
CALL INX(NB,BR2,DF)
CALL INX(NP,BR2,DG)
D(N)=DA
E(N)=- (DB+DC)*.5
F(N)=(DE+DF)*.5

```

```

504  G(N)=DG
      CALL INX(NP,BR1,FA)
      CALL KNX(NP,BR1,EB)
      CALL INX(NA,BR1,EC)
      CALL INX(NB,BR1,ED)
      CALL KNX(NA,BR1,EE)
      CALL KNX(NB,BR1,FF)
505  H(N)=EA
      O(N)=EB
      P(N)=.5*(EC+ED)
      Q(N)=-.5*(FF+FF)
      PD(N)=PA*(PC/AN+PB*AN)/(BN*AN)
1    CONTINUE
      IF(SENSE SWITCH 1)10,11
10   PUNCH201,(PD(N),N=N1,N2,N3)
201  FORMAT(5E16.8)
11   DO2I=11,12,13
      AI=I
      S=AI/AII
      CX=A*S
      RA=0.
      RB=0.
506  DO3N=N1,N2,N3
      BN=B(N)
      AN=N
      BNSQ=BN*BN
      CALL SQRT(-BNSQ,-CX,G1,G2)
      B1=R3*G1
      B2=R3*G2
      CALL JNZ(NP,B1,B2,BA,BB)
      CALL YNZ(NA,B1,B2,BC,BD)
      CALL YNZ(NB,B1,B2,BE,BF)
      CALL JNZ(NA,B1,B2,BG,BH)
      CALL JNZ(NB,B1,B2,BI,BJ)
507  CALL YNZ(NP,B1,B2,BK,BL)
      CN=C(N)
      CALL DIV(CN,0.,G1,G2,A21,A22)
      CALL MUL(A21,A22,BA,BB,BM,BQ)
      CALL MUL(A21,A22,BK,BL,BO,BP)
      CALL DIV(BO-.5*(BE-BC),BP-.5*(BF-BD),.5*(BI-BG)-BM,.5*(BJ-BH)-BQ,A
131,A32)
      C1=R2*G1
      C2=R2*G2
      CALL JNZ(NA,C1,C2,CA,CB)
      CALL JNZ(NB,C1,C2,CC,CD)
      CALL YNZ(NA,C1,C2,CE,CF)
508  CALL YNZ(NB,C1,C2,CG,CH)
      CALL JNZ(NP,C1,C2,CI,CJ)
      CALL YNZ(NP,C1,C2,CK,CL)
      CALL MUL(A31,A32,.5*(CC-CA),.5*(CD-CB),CO,CP)
      CALL MUL(A31,A32,CI,CJ,CQ,CR)
      CALL MUL(G1/BN,G2/BN,CO+.5*(CG-CE),CP+.5*(CH-CF),CS,CT)
      CALL DIV(CS,CT,CQ+CK,CR+CL,A41,A42)
      FA=A41*D(N)-F(N)
      FB=A42*D(N)
      FC=F(N)-A41*G(N)
      FD=-A42*G(N)

```

```

509 CALL LIV(PA, PB, PC, PD, A51, A521)
A61=A51*H(N)+O(N)
A62=A52*H(N)
A71=A51*P(N)+Q(N)
A72=A52*P(N)
CALL DIV(A61, A71, A72, QA, QB)
CALL MUL(O., 1., QA, QB, QC, QD)
RR(N)=PD(N)*QC
XR(N)=PD(N)*QD
XA(N)=QC
XB(N)=QD
RA=RA+RR(N)
RB=RB+RR(N)
RC(N)=RA
RD(N)=RB
TOR(N)=DDD*RA/((RS+RA)**2+(XS+RB)**2)
3 CONTINUE
IF(SENSE SWITCH 2) 12, 20
12 PUNCH202, (XA(N), XB(N), N=N1, N2, N4)
202 FORMAT(2X, (6E13.6))
20 PUNCH200, (S, RR(N), XR(N), N=N1, N2, N4)
200 FORMAT(F4.1, 6X, (10F7.3))
2 PUNCH203, (N, TOR(N), N=N1, N2, N4)
203 FORMAT(2X, (15, F8.3, 15, F8.3, 15, F8.3, 15, F8.3, 15, F8.3, 15, F8.3))
STOP
END

```

```

SUBROUTINE INX(N, X, R)
COMMON PY, ACC
AH=N
A=.5*A
ASQ=A*A
FACT=1
DO 11=1, N
A1=1
1 FACT=FACT*A1
R=B/FACT
L=0
11 L=L+1
AL=L
B=B*ASQ
FACT=FACT*AL*(AL+AL)
TERM=B/FACT
R=R+TERM
IF(TERM<ACC) 10, 10, 11
10 RETURN
END

```

```

SUBROUTINE KUX(N, X, R)
COMMON PY, ACC
N1=N-1
IF(N1) 31, 32, 31

```

```

32  IF(X-7.)33,33,34
31  IF(X-10.)33,33,34
34  FNSQ=N*N*4
    W=FNSQ-1.
    U=W*(FNSQ-9.)
    V=U*(FNSQ-25.)
    EX=8.*X
    EXS=EX*EX
    EXC=EXS*EX
    R=SQRTF(PY/(X+X))*EXPF(-X)*(1.+W/EX+U/(EXS+EXS)+V/(6.*EXC))
    GOTO23
33  GAMA=.5772157
    AN=N
    N2=N+1
    AN1=N1
    A=.5*X
    ASQ=A*A
    B1=A**N
    B2=1./B1
    C=LOGF(A)+GAMA
    R=0.
    M=0
    SIGN1=1.
    SIGN2=(-1.)**N2
    FACTN=1.
    FACT1=1.
    FACN1=1.
    DO1I=1,N
    AI=I
1   FACTN=FACTN*AI
    FACT3=FACTN
    IF(N1)21,22,21
21  DO2I=1,N1
    AI=I
2   FACN1=FACN1*AI
22  FACT2=FACN1
    F=0.
    DO3I=1,N
    AI=I
3   F=F+1./AI
    G=0.
24  AM=M
    ANM=N+M
    IF(M)17,18,17
17  G=G+1./AM
    F=F+1./ANM
    B1=B1*ASQ
    FACT1=FACT1*AM
    FACT3=FACT3*(AN+AM)
    FACT2=1.
    NM1=N1-M
    IF(NM1)9,10,7
7   DO4I=1,NM1
    AI=I
4   FACT2=FACT2*AI
10  B2=B2*ASQ
18  P=SIGN1*B2*FACT2/FACT1

```



```

GOTO19
9   P=0.
19  D=C-.5*(G+F)
190 Q=B1*D/(FACT1*FACT3)
    M=M+1
    SIGN1=-SIGN1
    R=R+.5*P+SIGN2*Q
    IF(ABS(F(Q)-ACC)23,23,24
23  RETURN
    END

SUBROUTINE JNZ(N,X,Y,S1,S2)
COMMON PY,ACC
AN=N
D=X*X+Y*Y
Z=SQRTF(D)
ASQ=D*.25
A=.5*Z
TH=ATANF(Y/X)
IF(N-2)21,21,22
21  IF(Z-10.)23,24,24
22  IF(Z-15.)23,24,24
24  FNSQ=4*N*N
    DA=FNSQ-1.
    DB=DA*(FNSQ-9.)
    S=DB*(FNSQ-25.)
    T=S*(FNSQ-49.)
    EZ=8.*Z
    EZS=EZ*EZ
    EZC=EZS*EZ
    TTH=TH+TH
    THTH=TTH+TH
    FTH=TTH+TH
    SIR=X-(AN+.5)*PY*.5
    SII=Y
    CA=EXPF(SII)
    CB=1./CA
    COSH=(CA+CB)*.5
    SINH=(CA-CB)*.5
    COS=COSF(SIR)
    SIN=SINF(SIR)
    COSR=COS*COSH
    COSI=-SIN*SINH
    SINR=SIN*COSH
    SINI=COS*SINH
    W=SQRTF(2./(PY*Z))
    AC=.5*TH
    CR=W*COSF(AC)
    CI=-W*SINF(AC)
    U=DB/(2.*EZS)
    V=T/(24.*EZC*EZ)
    AA=DA/EZ
    AR=S/(6.*EZC)
    PR=1.-U*COSF(TTH)+V*COSF(FTH)
    PI=U*SINF(TTH)-V*SINF(FTH)
    QR=AA*COSF(TH)-AB*COSF(THTH)
    QI=-AA*SINF(TH)+AB*SINF(THTH)
    BR=PR*COSR-PI*COSI-QR*SINR+QI*SINI

```

```

BI=PI*COSR+PR*COSI-QR*SINI-QI*SINR
S1=CR*BR-CI*BI
S2=CI*BR+CR*BI
GOTO12
23  B=A**N
    PHI=TH*AN
    TTH=TH+TH
    T1=0.
    T2=0.
    S1=0.
    S2=0.
    SIGN=1.
    M=0
    FACT1=1.
    FACT2=1.
    DOLI=1,N
    AI=I
1   FACT2=FACT2*AI
11  C=SIGN*B/(FACT1*FACT2)
    T1=C*COSF(PHI)
    T2=C*SINF(PHI)
    M=M+1
    AM=M
    SIGN=-SIGN
    B=B*ASQ
    PHI=PHI+TTH
    FACT1=FACT1*AM
    FACT2=FACT2*(AN+AM)
    S1=S1+T1
    S2=S2+T2
10  IF(ABSF(T2)-ACC)10,10,11
12  IF(ABSF(T1)-ACC)12,12,11
    RETURN
    END

SUBROUTINE YNZ(N,X,Y,S1,S2)
COMMON PY,ACC
AN=N
N1=N-1
GAMA=.5772157
ASQ=.25*(X*X+Y*Y)
A=SQRTF(ASQ)
Z=A+A
TH=ATANF(Y/X)
51  IF(N-1)51,51,52
52  IF(Z-9.)23,24,24
24  IF(Z-12.)23,24,24
    FNSQ=4*N*N
    DA=FNSQ-1.
    DB=DA*(FNSQ-9.)
    S=DB*(FNSQ-25.)
    T=S*(FNSQ-49.)
    EZ=8.*Z
    EZS=EZ*EZ
    EZC=EZS*EZ
501 TTH=TH+TH
    THTH=TTH+TH
    FTH=THTH+TH

```

```

SIR=X-(AN+.5)*PY*.5
SII=Y
CA=EXPF(SII)
CB=1./CA
COSH=(CA+CB)*.5
SINH=(CA-CB)*.5
COS=COSF(SIR)
SIN=SINF(SIR)
502 COSR=COS*COSH
COSI=-SIN*SINH
SINR=SIN*COSH
SINI=COS*SINH
W=SQRTF(2./(PY*Z))
AC=.5*TH
CR=W*COSF(AC)
CI=-W*SINF(AC)
U=DB/(2.*EZS)
V=T/(24.*EZC*EZ)
503 AA=DA/EZ
AB=S/(6.*EZC)
PR=1.-U*COSF(TTH)+V*COSF(FTH)
PI=U*SINF(TTH)-V*SINF(FTH)
QR=AA*COSF(TH)-AB*COSF(THTH)
QI=-AA*SINF(TH)+AB*SINF(THTH)
BR=PR*SINR-PI*SINI+QR*COSR-QI*COSI
BI=PI*SINR+PR*SINI+QR*COSI+QI*COSR
S1=CR*BR-CI*BI
S2=CI*BR+CR*BI
GOTO120
23 B=A**N
R=1./B
PHI=TH*AN
ALPHA=-PHI
TTH=TH+TH
F1=0.
F2=0.
SE1=0.
SE2=0.
T1=0.
T2=0.
S1=0.
S2=0.
SIGN=1.
M=0
FACT1=1.
FACN1=1.
FACTN=1.
DOLI=1,N
AI=I
1 FACTN=FACTN*AI
FACT2=FACTN
C=.5*LOGF(ASQ)+GAMA
Q1=C+C
Q2=TTH

```

```

F=0.
G=0.
DO2I=1,N
AI=I
2 F=F+1./AI
IF(N1)21,22,21
21 DO3I=1,N1
AI=I
3 FACN1=FACN1*AI
22 FACT3=FACT3*AI
15 COS=COSF(PHI)
SIN=SINF(PHI)
E=SIGN*B/(FACT1*FACT2)
FT1=E*COS
FT2=E*SIN
P=E*(F+G)
ST1=P*COS
ST2=P*SIN
IF(N1-M)43,42,42
42 Q=R*FACT3/FACT1
TT1=Q*COSF(ALPHA)
TT2=Q*SINF(ALPHA)
GOTO44
43 TT1=0.
TT2=0.
44 F1=F1+FT1
F2=F2+FT2
SE1=SE1+ST1
SF2=SE2+ST2
T1=T1+TT1
T2=T2+TT2
IF(ABSF(FT1)-ACC)10,10,11
10 IF(ABSF(ST1)-ACC)13,13,11
13 IF(ABSF(FT2)-ACC)14,14,11
14 IF(ABSF(ST2)-ACC)16,16,11
16 IF(ABSF(TT1)-ACC)17,17,11
17 IF(ABSF(TT2)-ACC)12,12,11
11 M=M+1
AM=M
NM1=N-M-1
SIGN=-SIGN
B=B*ASQ
R=R*ASQ
FACT1=FACT1*AM
FACT2=FACT2*(AN+AM)
FACT3=1.
IF(N1-M)34,34,33
33 DO4I=1,NM1
AI=I
4 FACT3=FACT3*AI
34 PHI=PHI+TTH
ALPHA=ALPHA+TTH
G=G+1./AM
ANM=N+M
F=F+1./ANM
GOTO15

```

```

12  S1=(F1*Q1-F2*Q2-SE1-T1)/PY
    S2=(F1*Q2+F2*Q1-SE1-T2)/PY
120 RETURN
    END

```

```

SUBROUTINE SQRT(A,B,C,D)$  E=(A*A+B*B)**.25$ TH=ATANF(B/A)*.5
C=E*COSF(TH)$ D=E*SINF(TH)$ RETURN $ END

```

```

SUBROUTINE DIV(A,B,C,D,E,F)$ H=C*C+D*D$ E=(A*C+B*D)/H
F=(B*C-A*D)/H$ RETURN $ END

```

```

SUBROUTINE MUL(A,B,C,D,E,F)$ E=A*C-B*D$ F=A*D+B*C$ RETURN $ END

```

```

                                SAMPLE DATA
3.58E-08.115          .0476      .04735      .0445      .04435
110.      10.3      7.5      .36      552.      220.      .9      2
12010 1 3 2 210.      .0952      1500.      .00001

```

PROGRAM NO. 5

LIST OF NOTATIONS

QUANTITIES IN INPUT LIST

N=NO. OF T,S

T(I)=VARIOUS VALUES OF TIME IN SECONDS

C1=K1,C2=K2,C3=K3,C4=K4

CJ=MOMENT OF INERTIA

QUANTITIES IN OUTPUT LIST

T(J)=TIME IN SECONDS

OMEGA(J)=SPEED

```

C  C  SPEED TRANSIENT  Z
      DIMENSION T(50),OMEGA(50)
      READ100,N
100  FORMAT(I2)
      READ103,(T(I),I=1,N)
103  FORMAT(7F10.0)
104  READ101,C1,C2,C3,C4,CJ
101  FORMAT(F10.0,4E10.4)
102  FORMAT(E20.5)
      PUNCH300,C1,C2,C3,C4,CJ
300  FORMAT(3HK1=F6.3,4H K2=E11.4,4H K3=E11.4,4H K4=E11.4,3H J=E11.4)
11   C5=C2-   C4$ Q=4.*C1*C3-C5*C5
      O=SQRTF(-Q)$ P=O/(2.*CJ)$ S=2.*C3$ R=C5/S$ U=O/S
      V=LOGF((O+C5)/(O-C5))/2.$ DO1NN=1,N$ TH=-P*T(NN)+V
      IF(TH)12,13,12
13   X=1.$ GOTO14
12   X=EXPF(TH)
14   Y=1./X$ TANH=(X-Y)/(X+Y)$ OMEGA(NN)=U*TANH-R $ L=L+1
1    CONTINUE
      PUNCH200,(T(J),OMEGA(J),J=1,L)
2000 FORMAT(F6.3,F8.0,F7.3,F8.0,F6.3,F8.0,F7.3,F8.0,F6.3,F8.0)
16  GO TO 104
      END

```

SAMPLE DATA

19						
.01	.02	.03	.04	.05	.06	.07
.08	.09	.1	.12	.15	.17	.2
.3	.5	.7	.8	1.		
1.48	-.146E-3	-.1E-8	.6E-4	.4E-5		

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